

“It from Qubit” Postdocs and studentships

To be held at:

[Vijay Balasubramanian](#) (UPenn)

[Patrick Hayden](#) (Stanford)

[Alexei Kitaev](#)(Caltech)

[Don Marolf](#) (UCSB)

[Jonathan Oppenheim](#) (University College London)

[Mark Van Raamsdonk](#) (UBC)

Other members of the collaboration include:

[Horacio Casini](#) (Bariloche)

[Daniel Harlow](#) (MIT)

[Alex Maloney](#) (McGill)

[Juan Maldacena](#) (IAS)

[Rob Myers](#) (Perimeter)

[Scott Aaronson](#) (UT Austin)

[Dorit Aharonov](#) (Jerusalem)

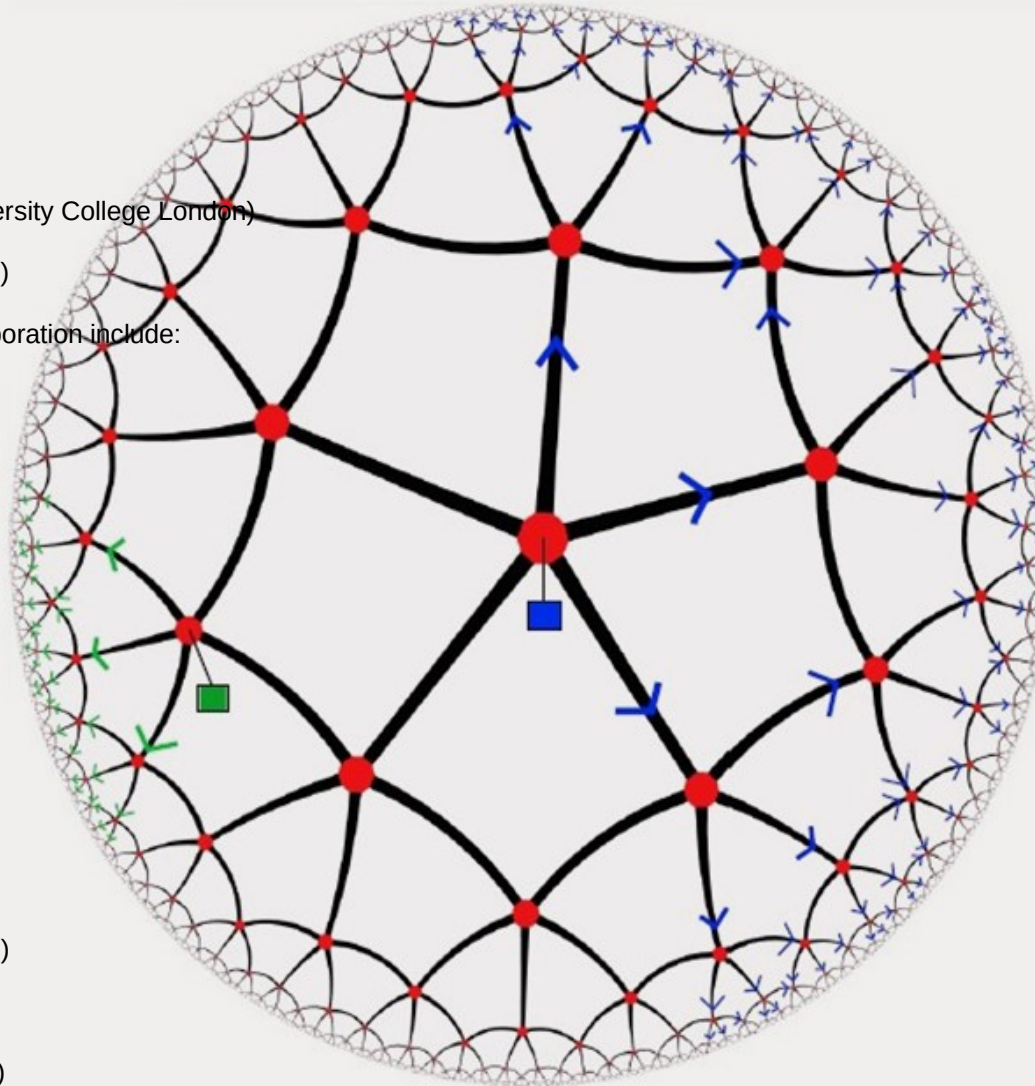
[Brian Swingle](#) (Maryland)

[Tadashi Takayanagi](#) (Kyoto)

[Matthew Headrick](#) (Brandeis)

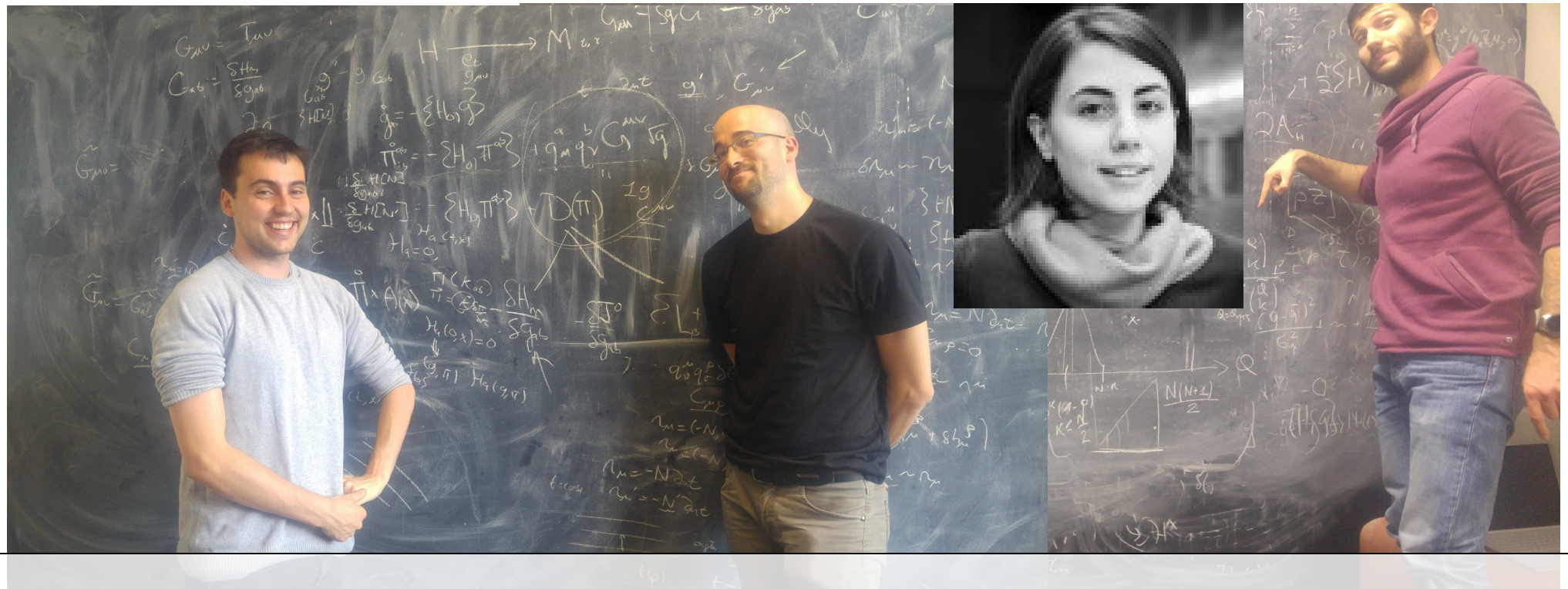
[John Preskill](#) (Caltech)

[Leonard Susskind](#) (Stanford)



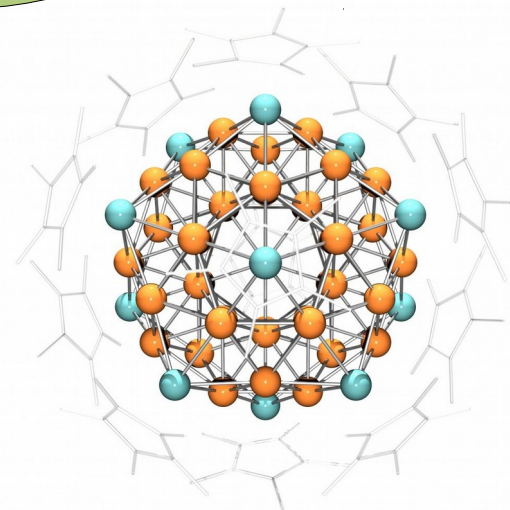
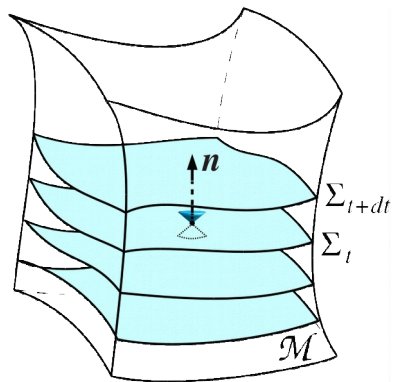
A post-quantum theory of classical gravity?

- arXiv:1811.03116, ...
- w/ Camps, Soda, Sparaciari, Weller-Davies

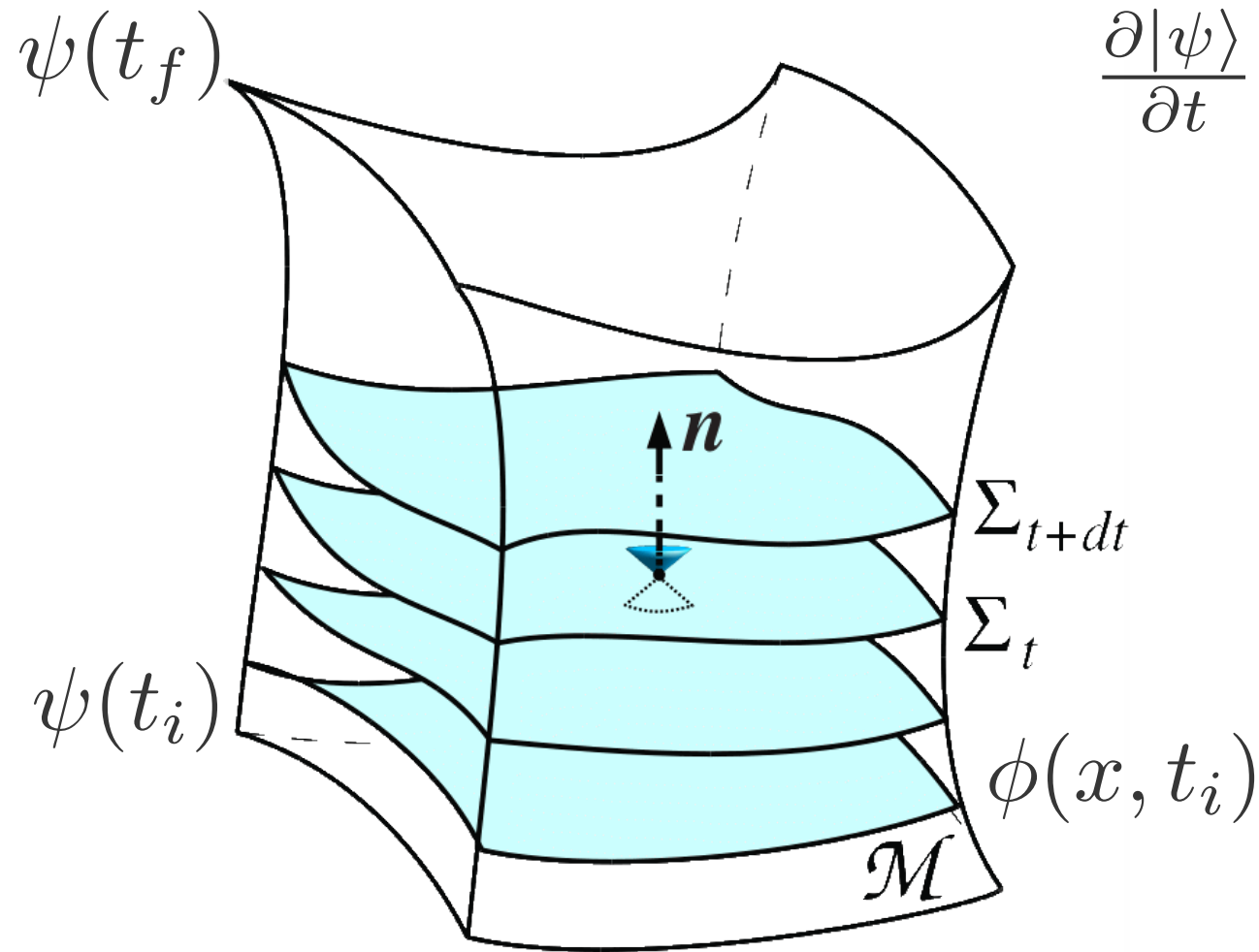


General Relativity & Quantum theory

$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

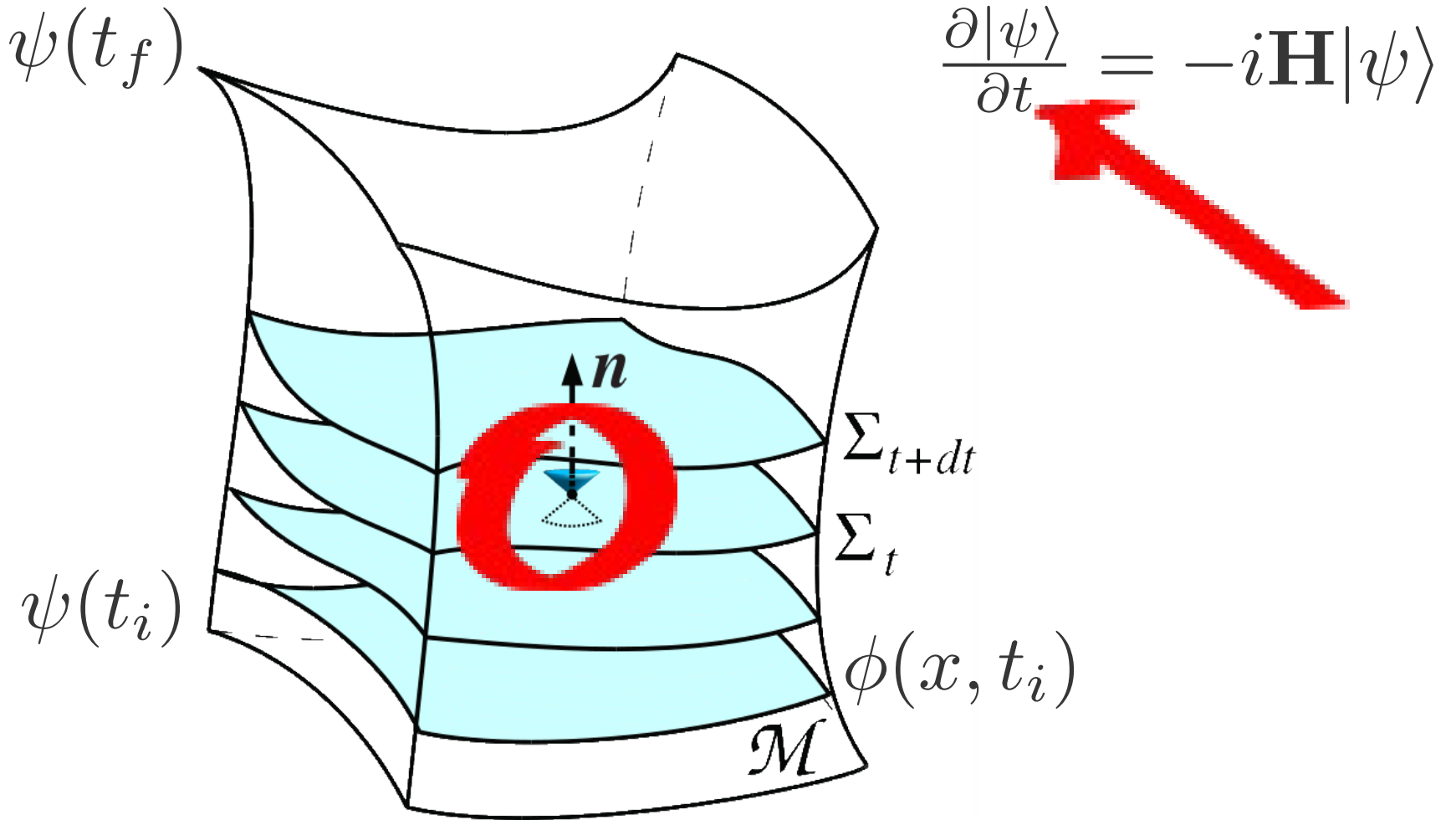


Quantum theory has a future

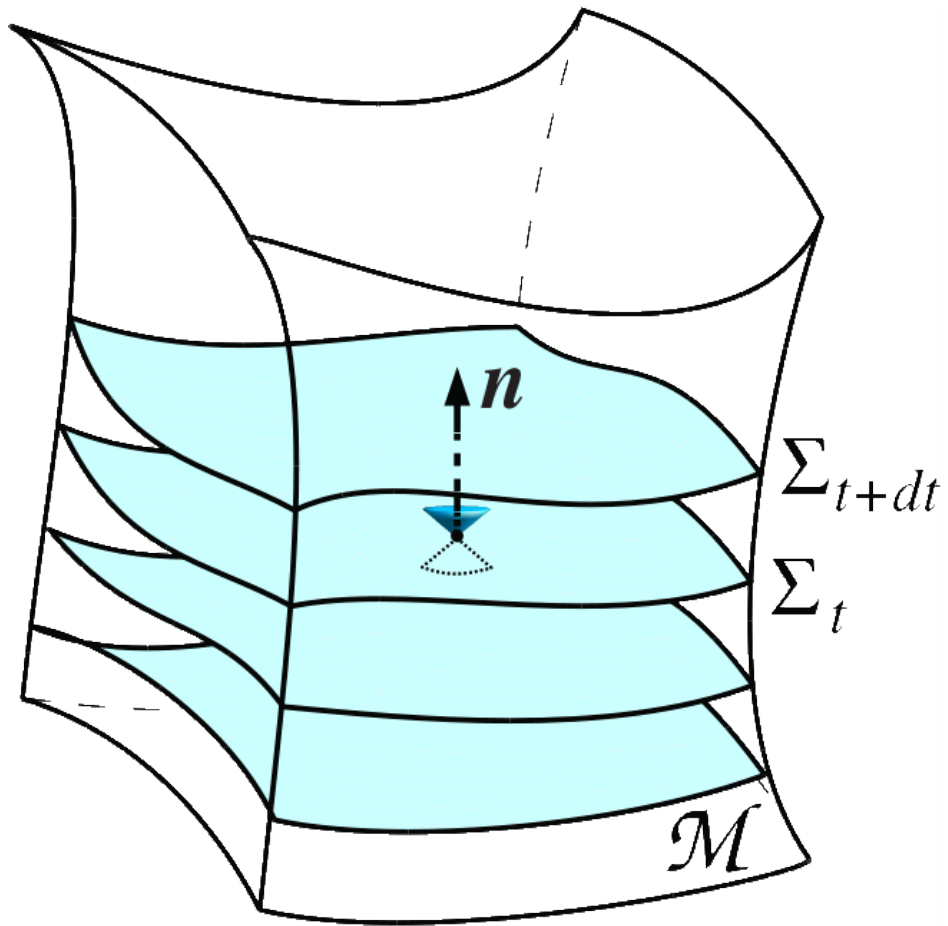


$$\frac{\partial |\psi\rangle}{\partial t} = -i\mathbf{H}|\psi\rangle$$

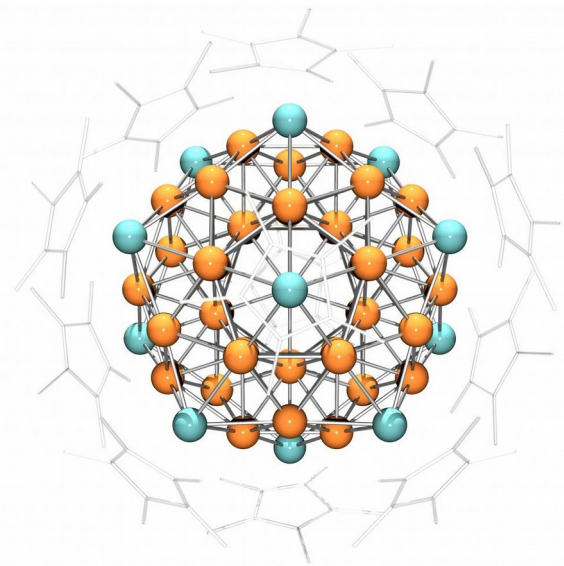
Quantum gravity, not so much...



Should we quantise space-time?



$$g_{ab}(x, t)$$



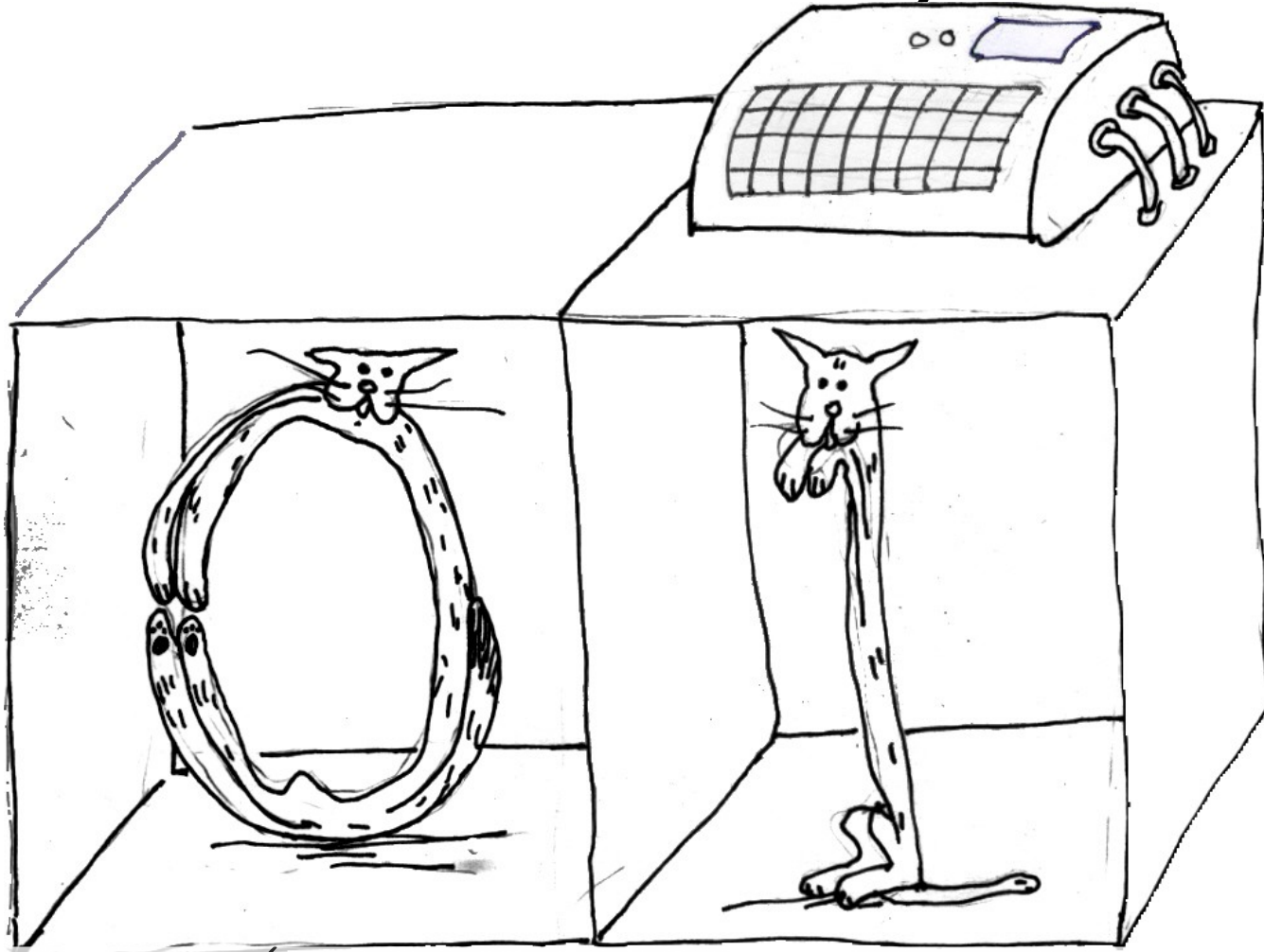
$$\phi(x, t)$$

Is space-time classical?

I have no idea. But....

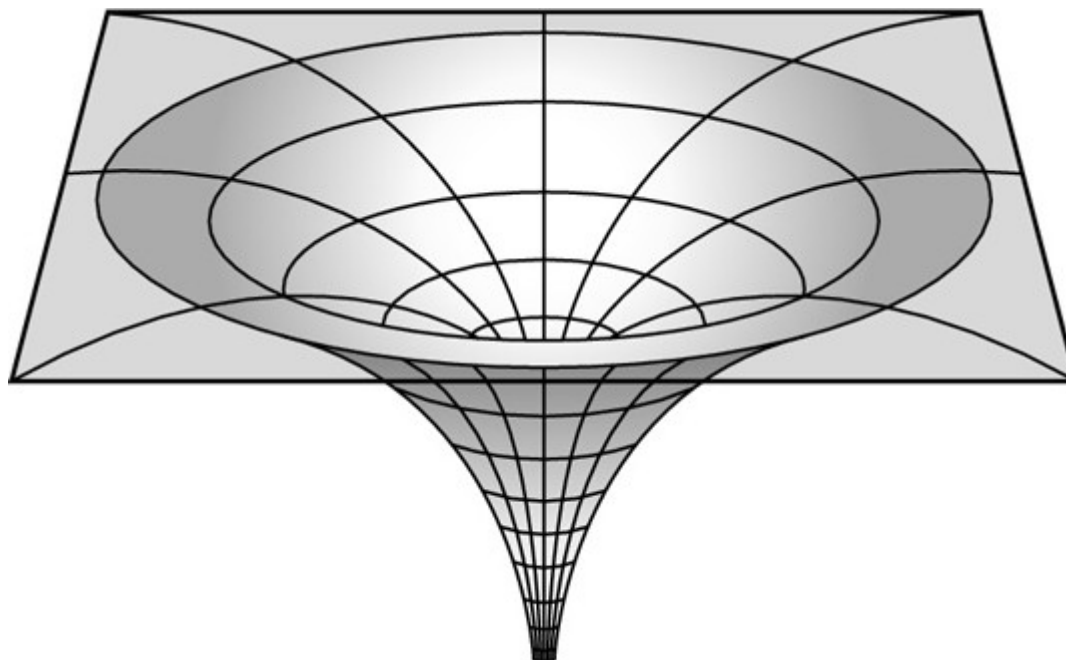
- It can be
- It would necessarily cause “collapse of the wave-function”
- It would necessarily lead to information destruction
- Effective theory?

Does it make any sense to speak of the emergence of classicality without classical systems?



$$|\psi\rangle = \frac{1}{\sqrt{2}} \left(|0000000\dots\rangle + |1111111\dots\rangle \right)$$

Black hole information problem 2.0



AMPS: If information is preserved we must break the equivalence principle (a firewall).

BPS: If information is destroyed we must sacrifice locality or energy conservation.

Hawking (1976)

Almheiri, Marolf, Polchinski, Sully (2012)

Banks, Peskin, Susskind (1984)

Outline

- Can we couple quantum and classical systems?
- What is the general form of this dynamics?
- Application to General Relativity/QFT ▽

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Can we?

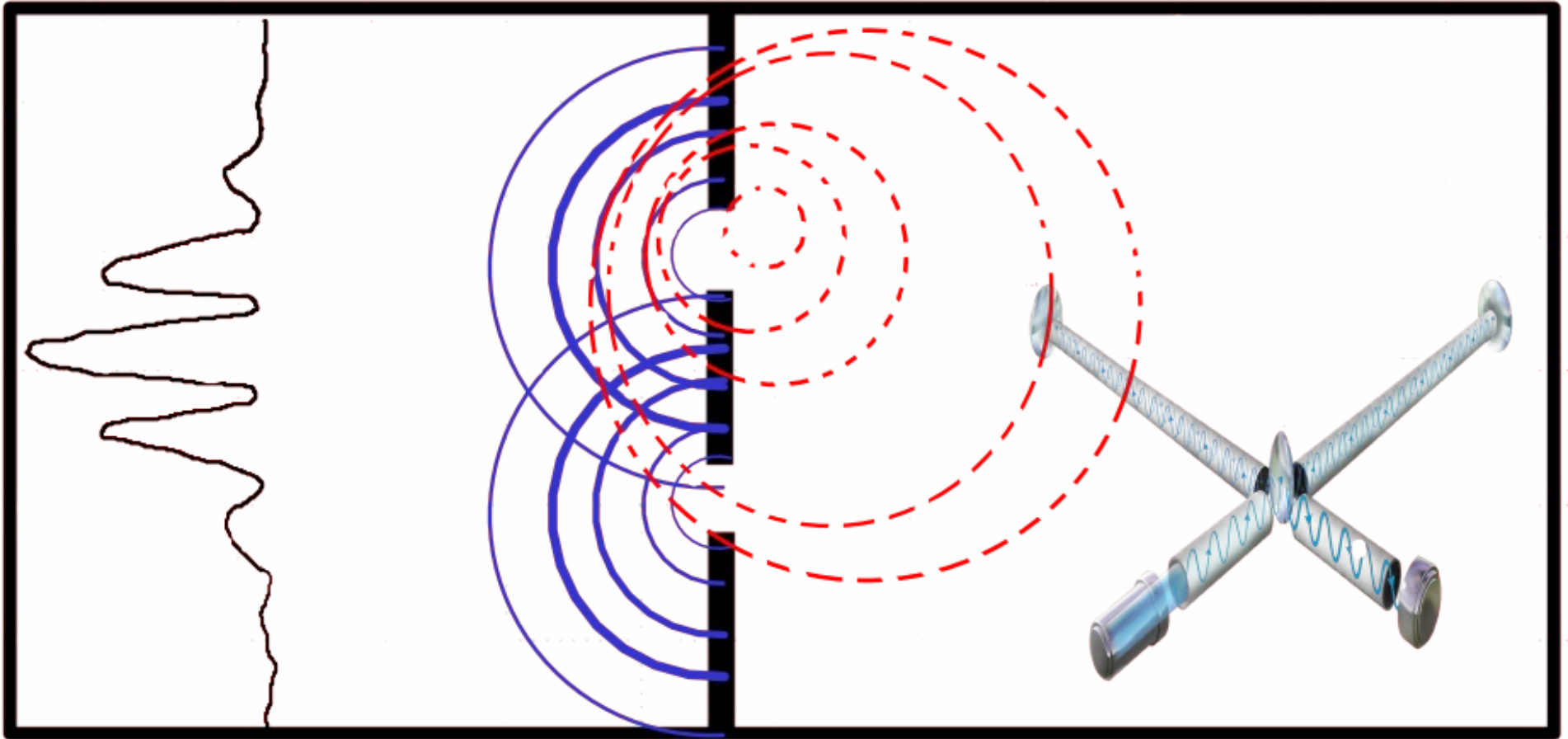
$$\frac{1}{\sqrt{2}} |\text{coin}\rangle |\text{L}\rangle + \frac{1}{\sqrt{2}} |\text{coin}\rangle |\text{R}\rangle$$

$$\sigma_c = \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}$$

$$\sigma_p = \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}$$

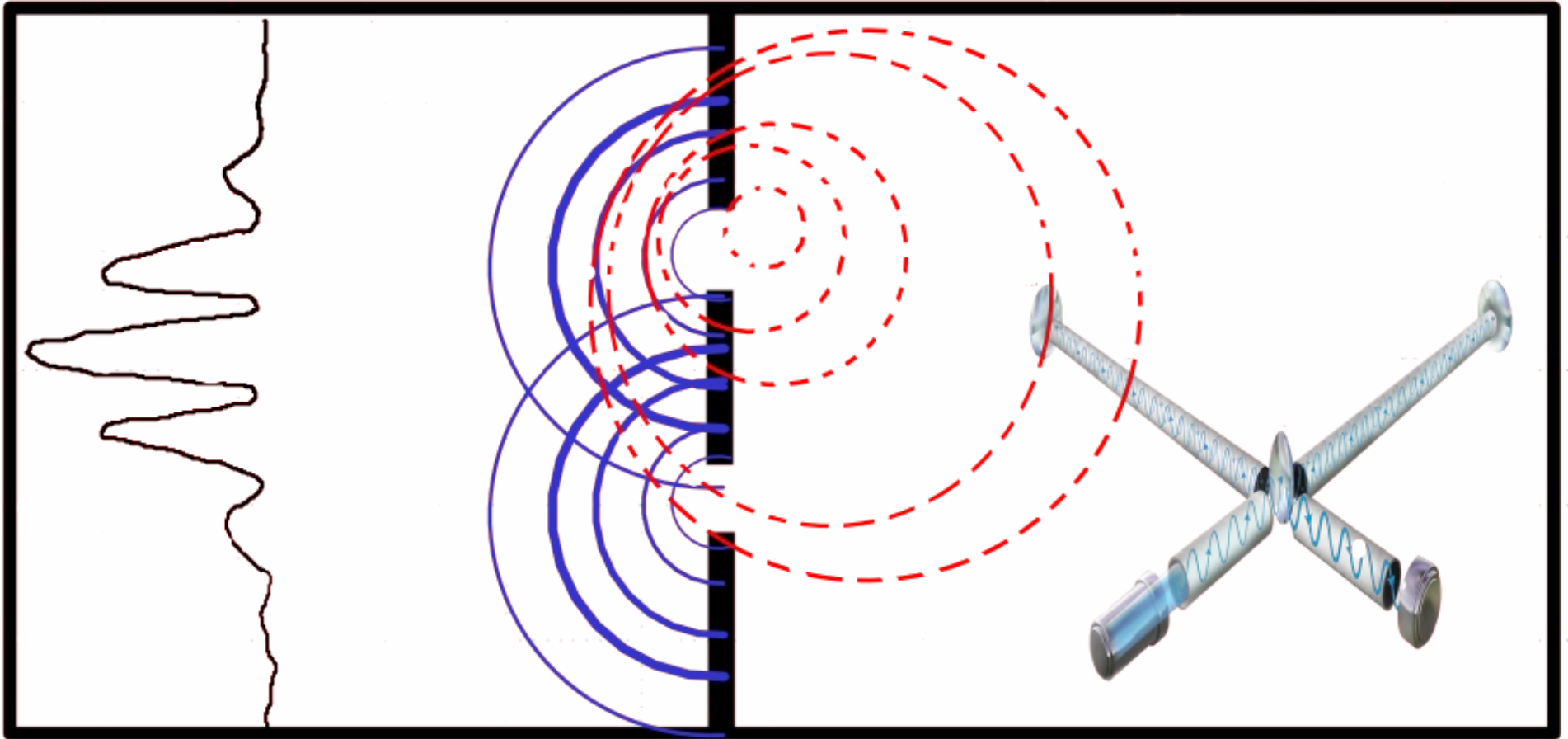
Feynman, Chapel Hill Conference (1957)
Eppley and Hannah (1977)
Marletto and Vedral (2017)

Can we?



Feynman, Chapel Hill Conference (1957)
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Marletto and Vedral (2017)

Can we?



$$\frac{1}{\sqrt{2}} |E_L\rangle |\text{Moon L}\rangle + \frac{1}{\sqrt{2}} |E_R\rangle |\text{Moon R}\rangle$$

Can we?

$$\frac{1}{\sqrt{2}} |E_L\rangle |\text{Moon}_L\rangle + \frac{1}{\sqrt{2}} |E_R\rangle |\text{Moon}_R\rangle$$

$$\alpha = \langle E_L | E_R \rangle$$

$$\sigma_p = \begin{pmatrix} \frac{1}{2} & \alpha^* \\ \alpha & \frac{1}{2} \end{pmatrix}$$

Can we?

Classical gravity is incompatible with quantum theory unless the coupling is **stochastic**.

Feynman, Chapel Hill Conference (1957)
Eppley and Hannah (1977)
Marletto and Vedral (2017)

Can we?

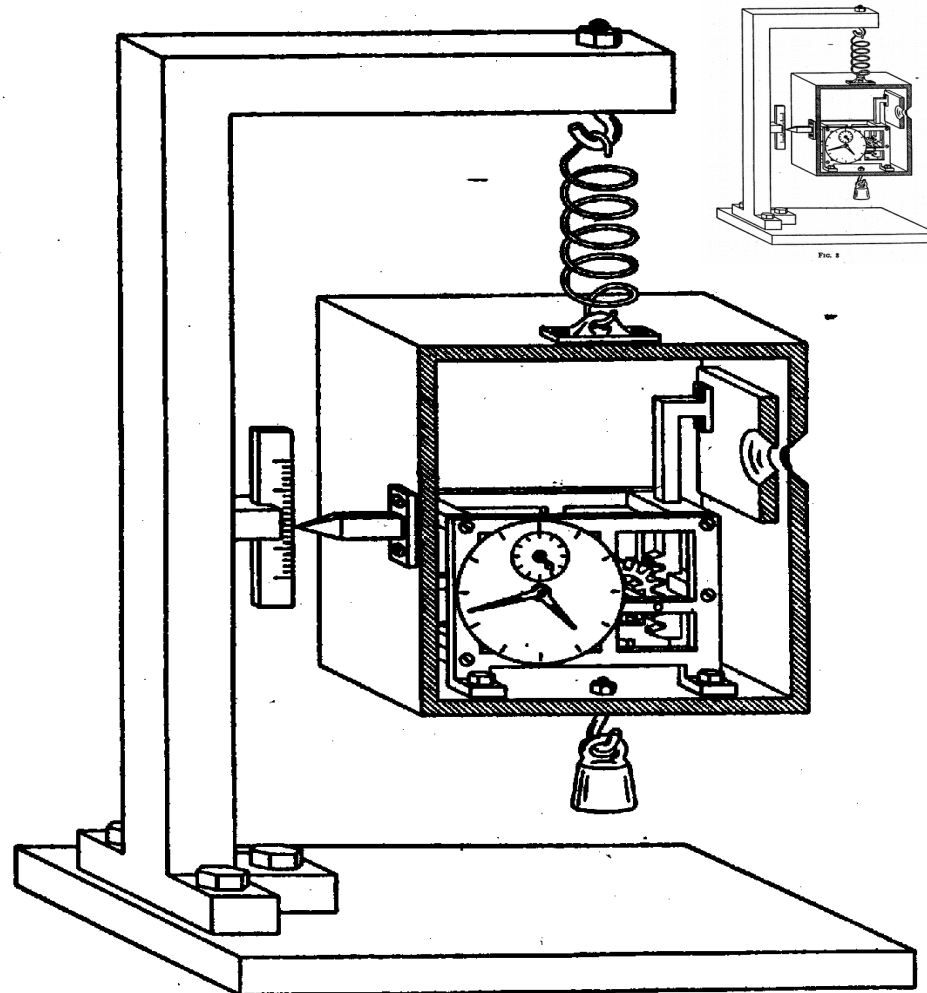
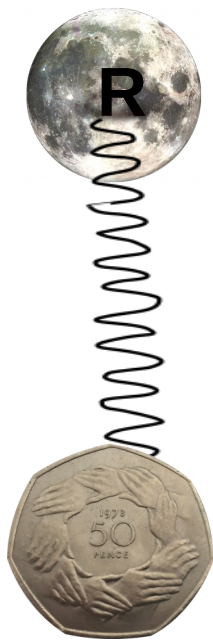
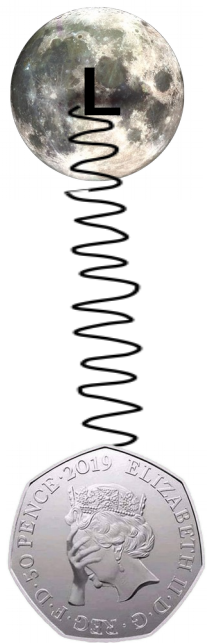


FIG. 8

Semi-classical gravity?

$$G_{\mu\nu} = \frac{8\pi G}{c^4} \langle \mathbf{T}_{\mu\nu} \rangle$$

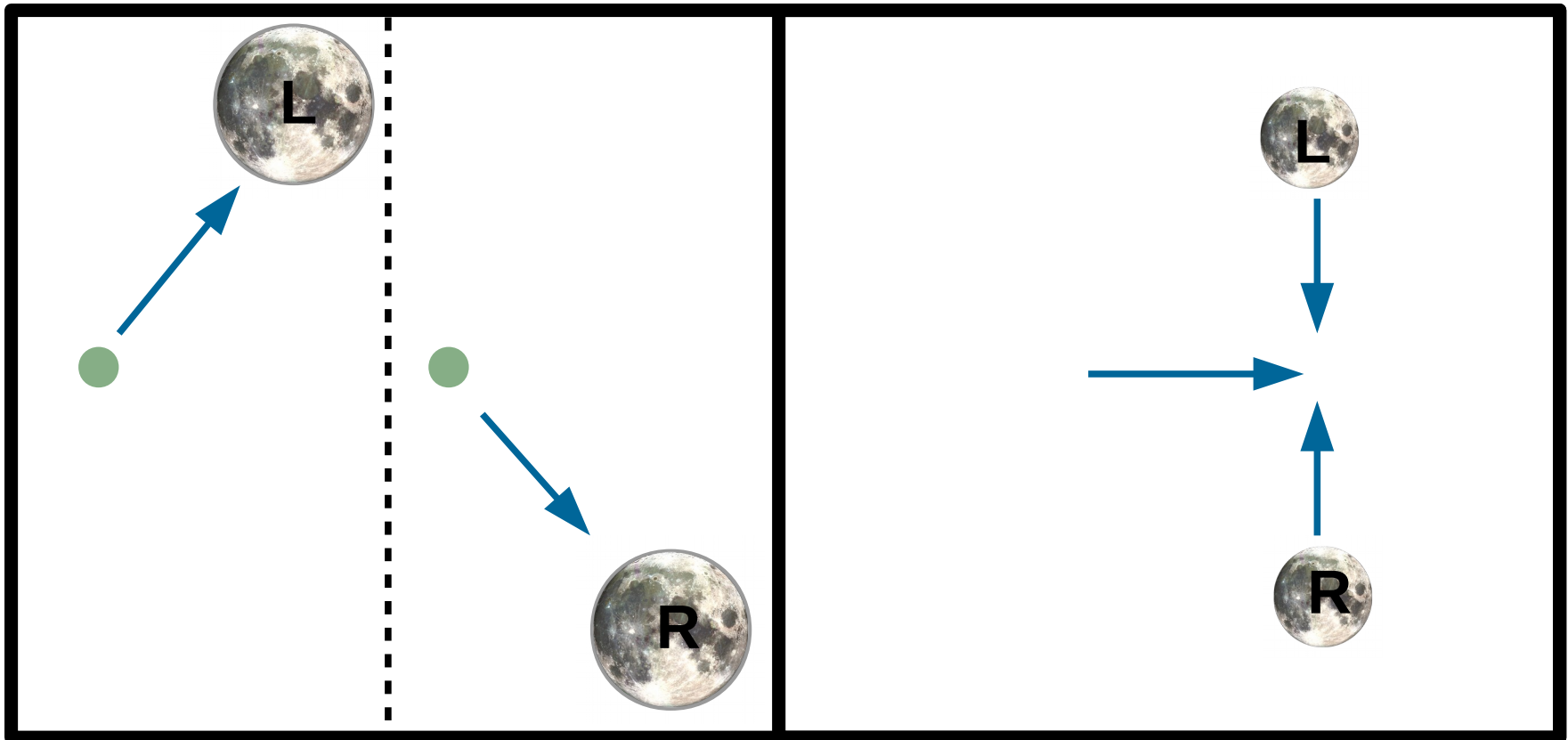


$$\langle \mathbf{T}_{\mu\nu} \rangle$$

$$\frac{1}{\sqrt{2}} \left| \begin{array}{c} \text{coin} \\ \text{L} \end{array} \right\rangle + \frac{1}{\sqrt{2}} \left| \begin{array}{c} \text{coin} \\ \text{R} \end{array} \right\rangle$$

Semi-classical gravity?

$$G_{\mu\nu} = \frac{8\pi G}{c^4} \langle \mathbf{T}_{\mu\nu} \rangle$$

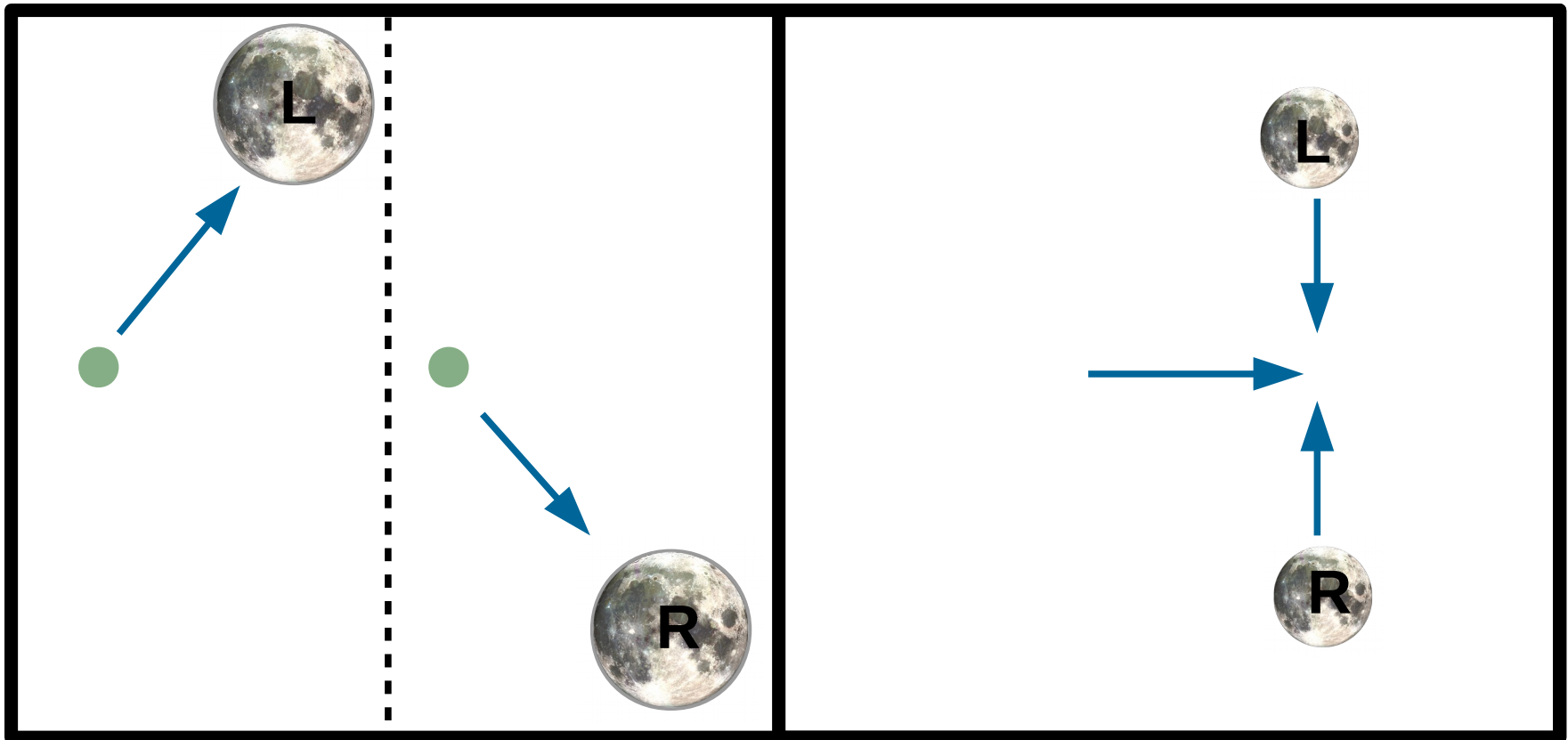


This or that

Not this!

Semi-classical gravity?

$$\cancel{G_{\mu\nu} = \frac{8\pi G}{c^4} \langle T_{\mu\nu} \rangle}$$



This or that

Not this!

Outline

- Can we couple quantum and classical systems?
- What is the general form of this dynamics?
- Application to General Relativity/QFT ▽

What kind of dynamics is allowed?

σ Quantum states are represented by a density matrix (positive, trace one) in a Hilbert space

$\rho(q, p)$ Classical states are a probability density (positive, integrate to one) over phase space

$\varrho(q, p)$ Hybrid state-space consists of a Hilbert space at each point in phase space. Hybrid states are positive, $\int dqdp \text{tr} \varrho(q, p) = 1$

What kind of dynamics is allowed?

$\rho(q, p)$ Hybrid state-space consists of a Hilbert space at each point in phase space. Hybrid states are positive, $\int dq dp \text{tr} \rho(q, p) = 1$

e.g. quantum qubit

$$\sigma = \begin{pmatrix} p_0 & \alpha \\ \alpha^* & p_1 \end{pmatrix}$$

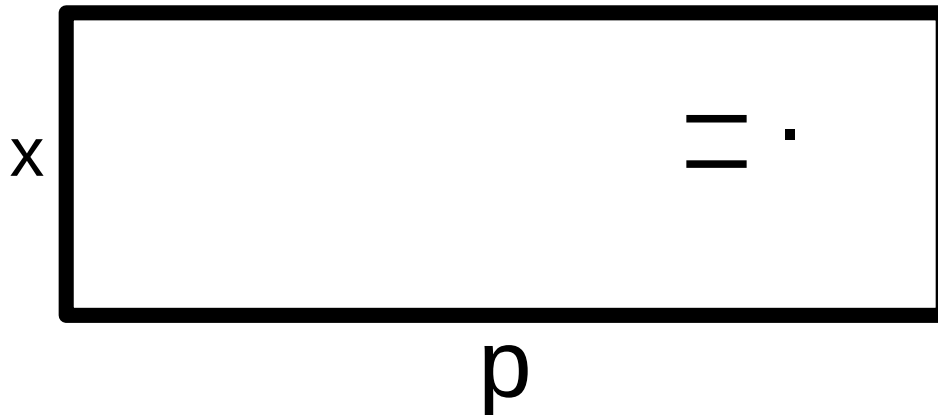


What kind of dynamics is allowed?

$\rho(q, p)$ Hybrid state-space consists of a Hilbert space at each point in phase space. Hybrid states are positive, $\int dq dp \text{tr} \rho(q, p) = 1$

e.g. hybrid qubit which has a classical x, p

$$\rho(x, p) = \begin{pmatrix} p_0(x, p) & \alpha(x, p) \\ \alpha^*(x, p) & p_1(x, p) \end{pmatrix}$$



What kind of dynamics is allowed?

It must be linear in ρ, σ or $\varrho(q, p)$

$$|\text{heads}\rangle\langle\text{heads}| \otimes \mathcal{L}(\rho_L)$$

$$|\text{tails}\rangle\langle\text{tails}| \otimes \mathcal{L}(\rho_R)$$

$$\frac{1}{2}\mathcal{L}(\rho_L) + \frac{1}{2}\mathcal{L}(\rho_R)$$

=

$$\frac{1}{\sqrt{2}}|\text{heads}\rangle\langle\text{heads}|_{\mathbf{L}} + \frac{1}{\sqrt{2}}|\text{tails}\rangle\langle\text{tails}|_{\mathbf{R}}$$

$$\mathcal{L}\left(\frac{1}{2}\rho_L + \frac{1}{2}\rho_R\right)$$

What kind of dynamics is allowed?

It must preserve the state space

$$p(z), \quad p(z) \geq 0, \quad \sum_z p(z) = 1$$

$$p(z; t) = \sum_{z'} P(z|z') p(z'; 0)$$

$$P(z|z') \geq 0, \quad \sum_z P(z|z') = 1$$

Classical Probability theory

What kind of dynamics is allowed?

It must preserve the state space

$$\sigma, \sigma \geq 0, \text{tr} \sigma = 1$$

$$\sigma(t) = \sum_{\mu} K_{\mu} \sigma(0) K_{\mu}^{\dagger}$$

$$K_{\mu}^{\dagger} K_{\mu} = \mathbb{1}$$

Kraus (87)

Quantum theory

What kind of dynamics is allowed?

It must preserve the state space

$$\rho(z) \geq 0, \quad \sum_z \text{tr} \rho(z) = 1$$

$$\rho(z; t) = \sum_{z', \mu} K_\mu(z|z') \rho(z', 0) K_\mu(z|z')^\dagger$$

$$\sum_{z', \mu} K_\mu(z|z')^\dagger K_\mu(z|z') = \mathbb{1}$$

Classical-quantum theory

What kind of dynamics is allowed?

Q $\sigma(t) = \sum_{\mu} K_{\mu} \sigma(0) K_{\mu}^{\dagger} \quad K_{\mu}^{\dagger} K_{\mu} = \mathbb{1}$

C $p(z; t) = \sum_{z'} P(z|z') p(z'; 0) \quad P(z|z') \geq 0, \sum_z P(z|z') = 1$

CQ $\rho(z; t) = \sum_{z', \mu} K_{\mu}(z|z') \rho(z', 0) K_{\mu}(z|z')^{\dagger} \quad \sum_{z', \mu} K_{\mu}(z|z')^{\dagger} K_{\mu}(z|z') = \mathbb{1}$

What kind of dynamics is allowed?

$$\frac{\partial \sigma}{\partial t} = -i[\mathbf{H}, \sigma] + \sum_{\alpha} W^{\alpha} \mathbf{L}_{\alpha} \sigma \mathbf{L}_{\alpha}^{\dagger} - \frac{1}{2} W^{\alpha} \{ \mathbf{L}_{\alpha}^{\dagger} \mathbf{L}_{\alpha}, \sigma \}_+$$

GKLS (76)

$$\frac{p(z)}{\partial t} = \{H(z), p(z)\} + \sum_{z'} \left(W(z|z') p(z') - W(z) p(z) \right)$$

$$\frac{\partial \varrho(z)}{\partial t} = -i[\mathbf{H}(z), \varrho] + \{H_c, \varrho\} + \sum_{z', \alpha} W^{\alpha}(z|z') \mathbf{L}_{\alpha} \varrho(z) \mathbf{L}_{\alpha}^{\dagger} - \frac{1}{2} W^{\alpha}(z) \{ \mathbf{L}_{\alpha}^{\dagger} \mathbf{L}_{\alpha}, \varrho(z) \}$$

Assumption: memoryless (Markovian)

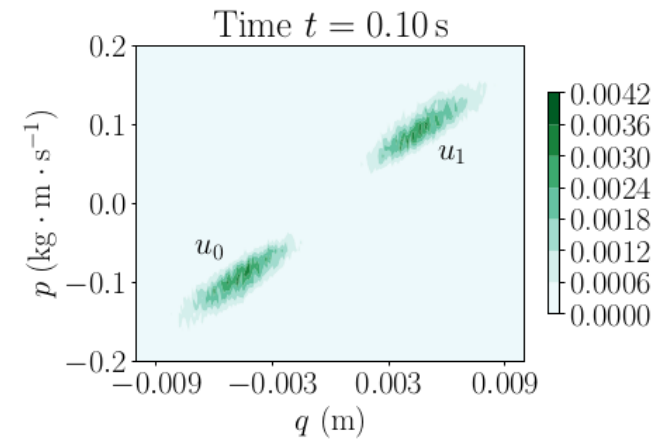
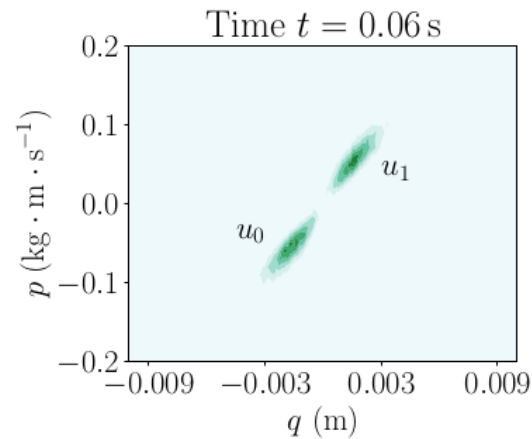
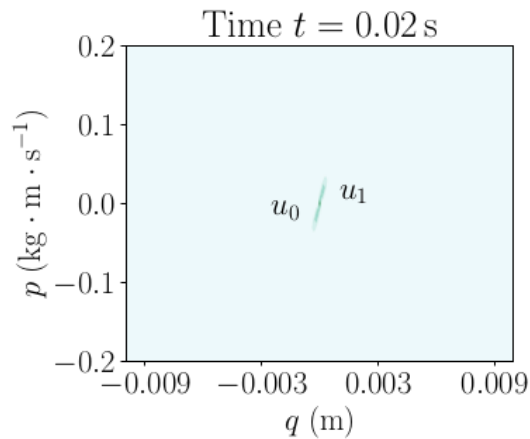
Blanchard, Jadczyk (93)

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A qubit with classical q,p

$$\frac{\partial \rho(q, p)}{\partial t} = \dots + \frac{\omega}{\tau} |1\rangle\langle 1| \rho(p + \tau B) |1\rangle\langle 1| + \frac{\omega}{\tau} |0\rangle\langle 0| \rho(p - \tau B) |0\rangle\langle 0| - \frac{\omega \rho}{\tau}$$



$$\rho(0) = \delta(x)\delta(p) |+\rangle\langle +|$$

- Quantum state collapses to 0,1
- Discreet jump in momentum
- Trade-off: coherence rate vs dispersion

Quantum field theory in a classical space-time

1) Non-commuting Hamiltonian dynamics ∇

2) $\phi(x), \pi_\phi(x) \quad g_{ab}, \pi^{ab}, N(x), N^a$

3) Dynamical ADM master equation

4) "Gauge invariance" imposes constraints $C^\mu(\varrho(g, \pi)) \approx 0,$

5) BPS \longrightarrow constraint conservation

Summary

Assume space-time is classical:

- This necessarily causes collapse of the wavefunction.
- Necessarily have stochastic and discrete classical jumps. Black holes destroy information.
- Experimentally testable
- Either fundamental, or classical limit of quantum gravity