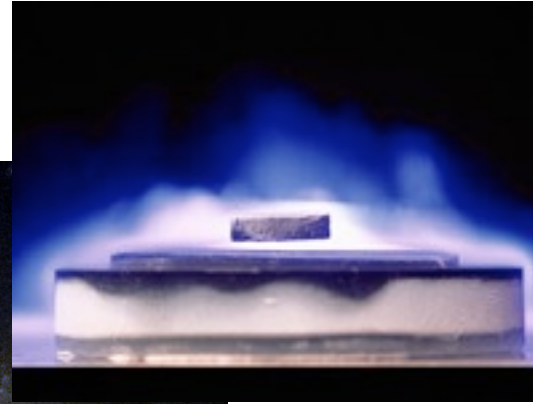




Qu-Transitions

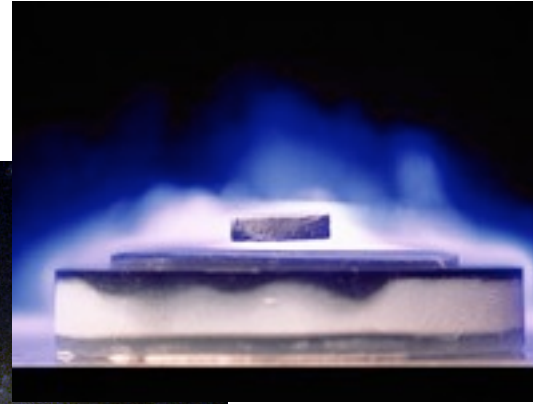


P. Coleman
(CMT, Rutgers)

University Toronto
Toronto, Sept 17
2009.

Qu-Transitions

"Phase transitions in the quantum era"



P. Coleman
(CMT, Rutgers)

University Toronto
Toronto, Sept 17
2009.

- The Qu-era: from classical to quantum.
- Classical criticality
- Heavy Fermion Quantum Criticality.
- New Approaches and Ideas
- Avoided Criticality

1758 in Paris: 72 years after “Principia”

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Classical revolution is still in full sway.

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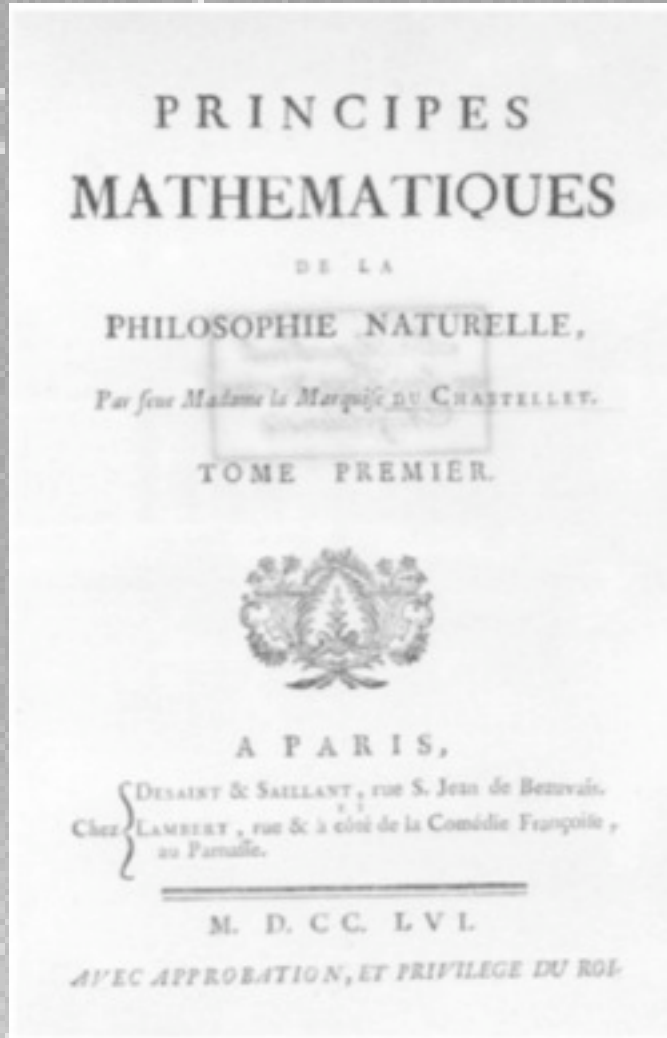
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PRINCIPES
MATHÉMATIQUES

DE LA
PHILOSOPHIE NATURELLE,
Par feu Madame la Marquise DU CHÂTELET.
TOME PREMIER.



A PARIS,
DESAINY & SAILLANT, rue S. Jean de Beauvais,
Chez LAMBERT, rue & à côté de la Comédie Française,
au Paroisse.

M. D. C. C. L. V. I.

AVEC APPROBATION, ET PRIVILEGE DU ROI.



Marquise Emilie du Châtelet
(1707-1749)

Mathematical Physicist:
Translator and interpreter of Principia.

"ce beau probleme astronomico-geometrique"

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Newtons Momentum

$$\sum_i m_i \vec{v}_i$$

Leibniz' "vis vivre"

$$\sum_i m_i (v_i)^2$$

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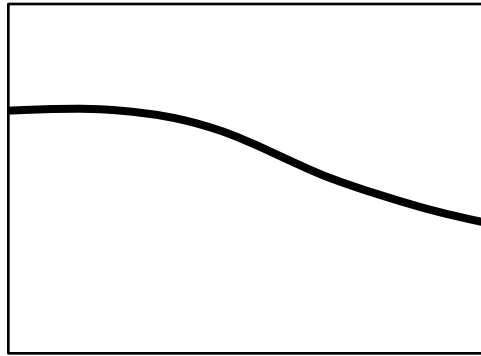
Resolution of the controversy (and the missing
factor of a half) required a further 60-80 years.

108 years after Planck, many
surprises later, the quantum era is in full sway.

108 years after Planck, many surprises later, the quantum era is in full sway.



χ

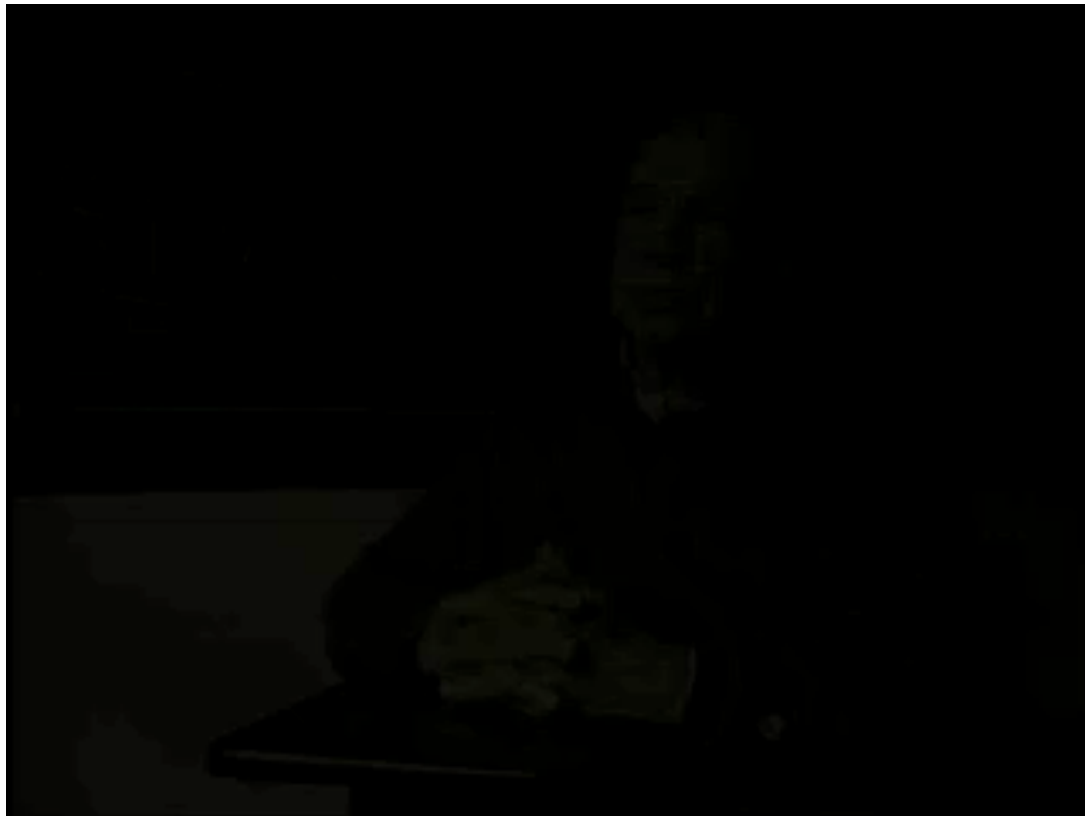


T

“With a heavy heart, I have been converted to the idea that Fermi -Dirac, not Einstein-Bose is the correct statistics. I wish to write a short note on its application to paramagnetism.”

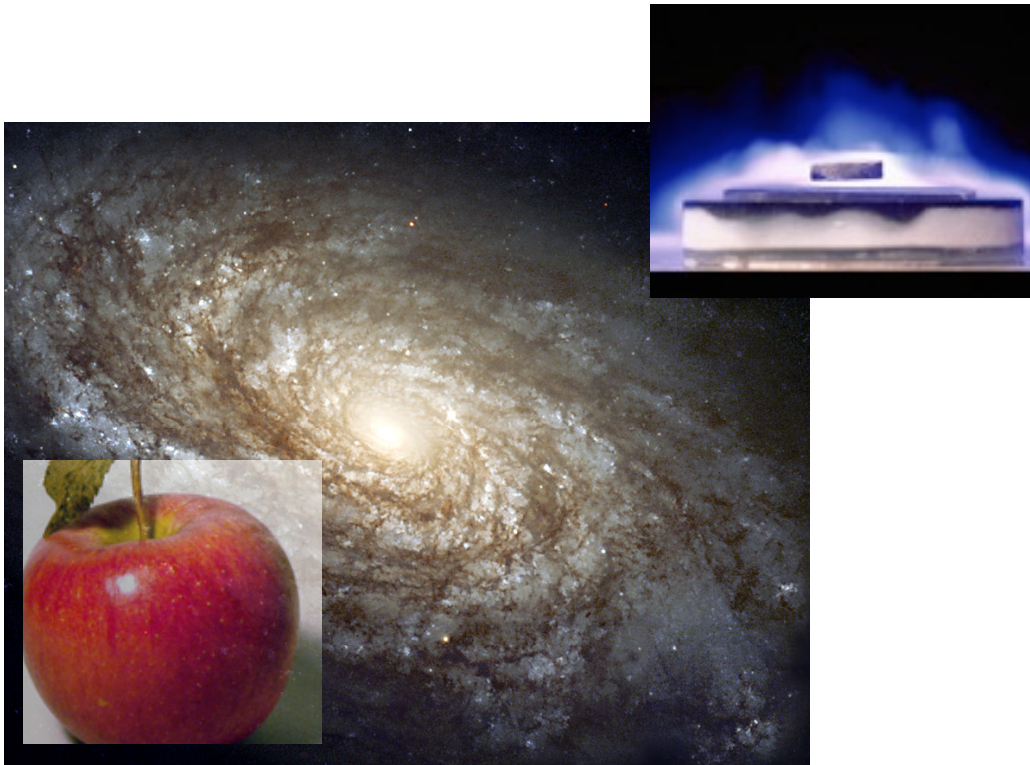
W. Pauli, in letter to Schrödinger, Dec 1926.

108 years after Planck, many
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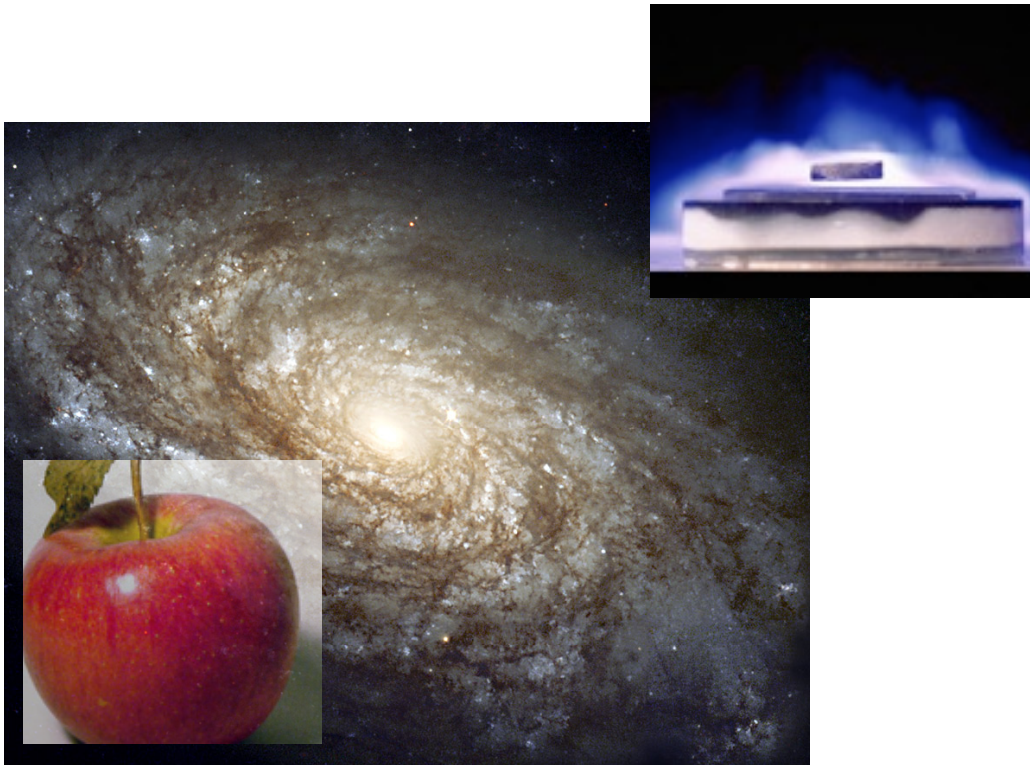


David Pines in
musicofthequantum.rutgers.edu

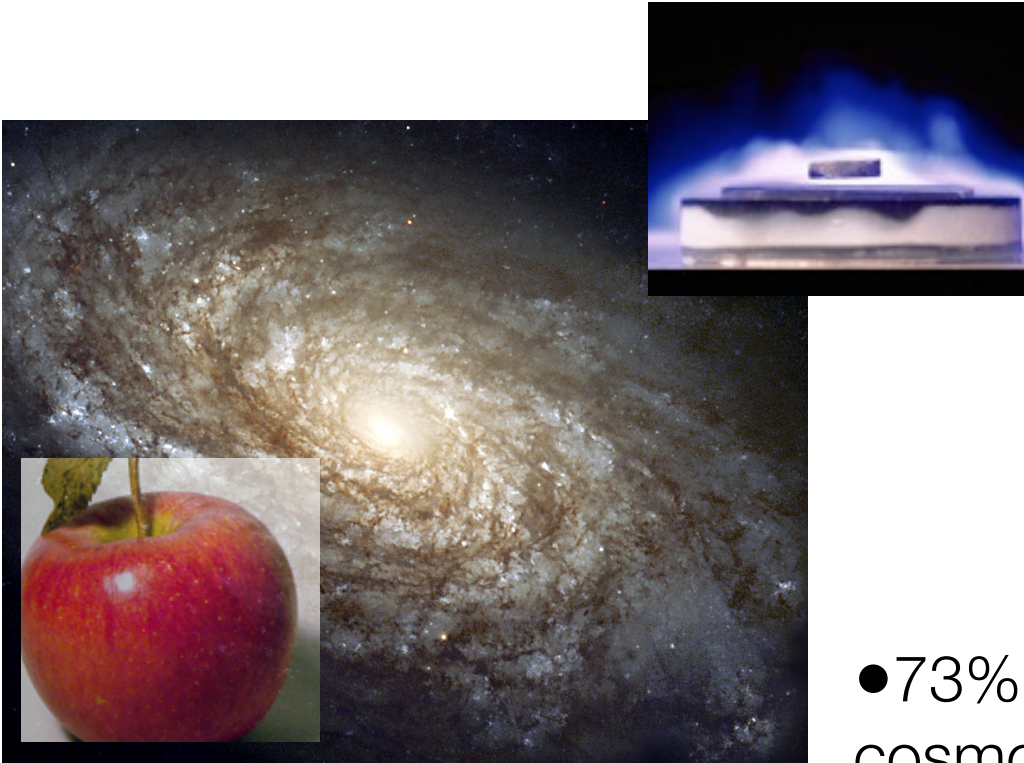
Quantum zero point fluctuations:



Quantum zero point fluctuations:
major unsolved problem of the quantum era.

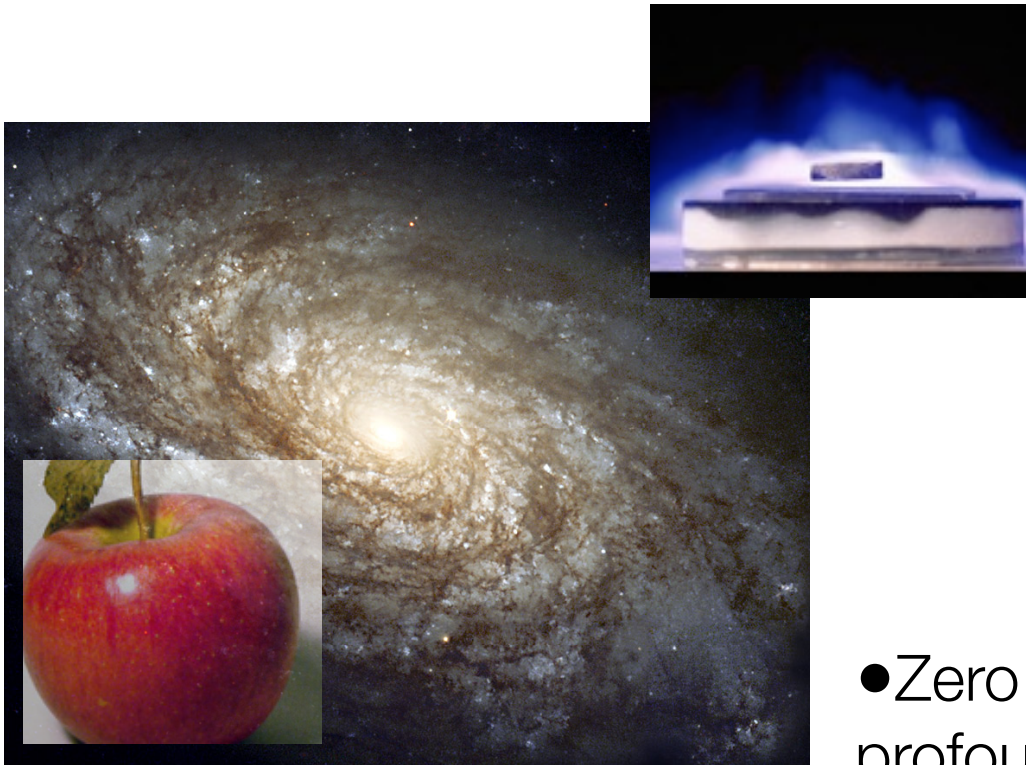


Quantum zero point fluctuations: major unsolved problem of the quantum era.



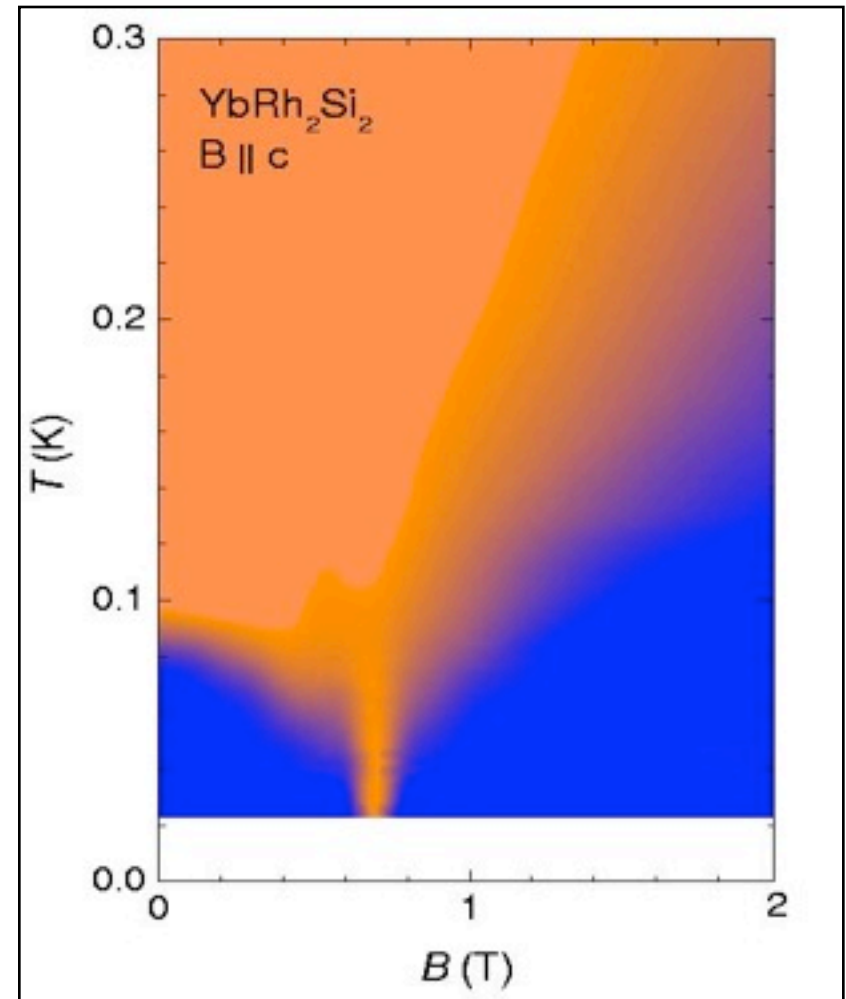
- 73% of the mass of the cosmos is “Dark Energy”: an unidentified form of zero point energy, causing the expansion to accelerate.

Quantum zero point fluctuations: major unsolved problem of the quantum era.

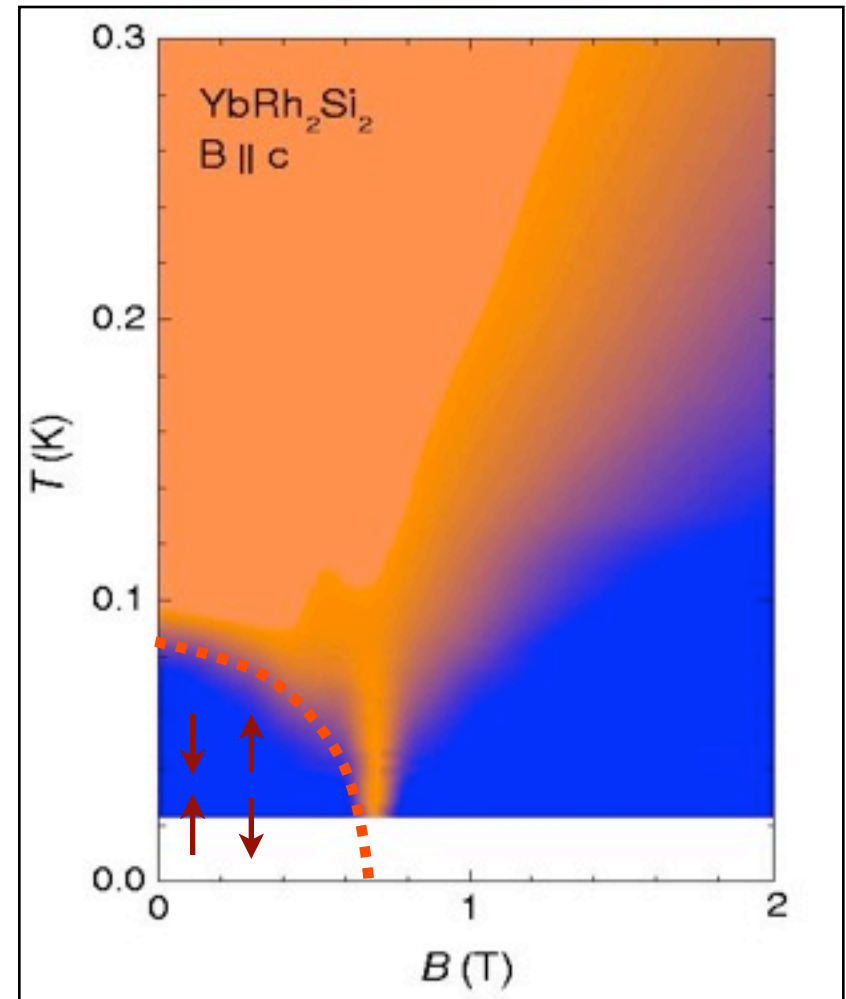


- Zero point fluctuations profoundly transform matter, endowing it with marked tendency to develop new forms of order.

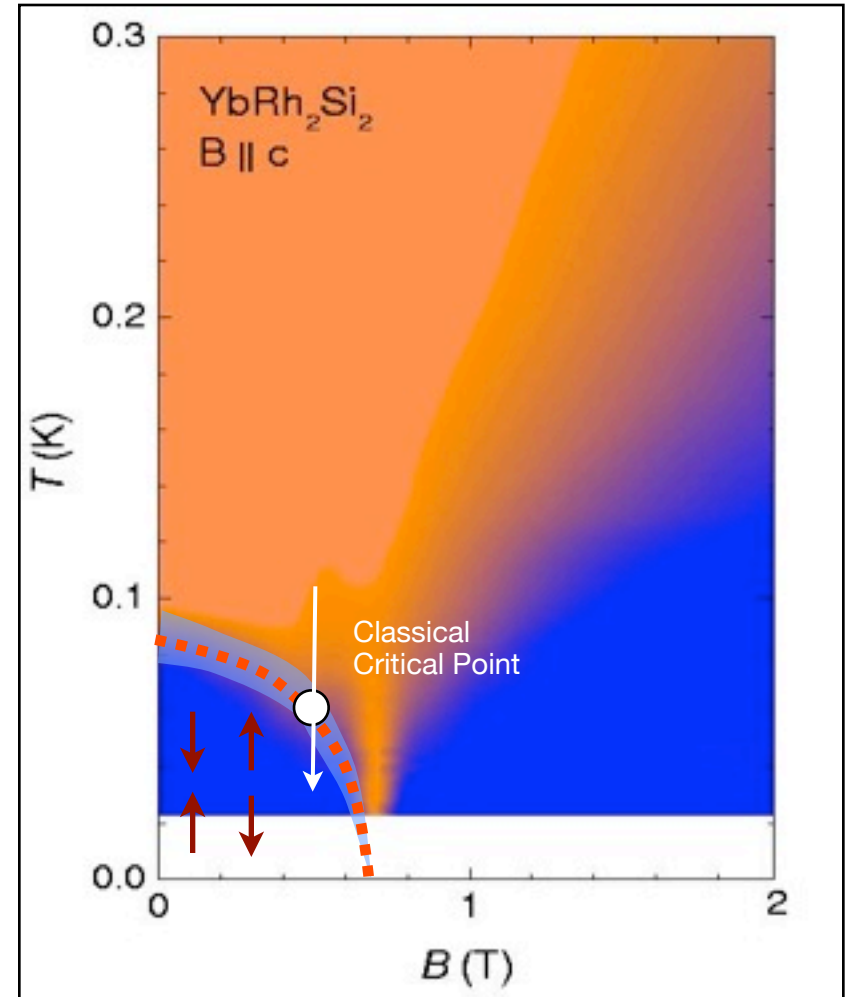
Classical Criticality



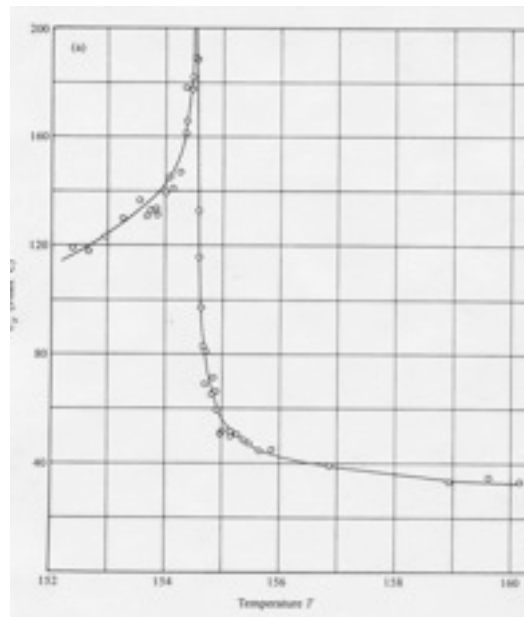
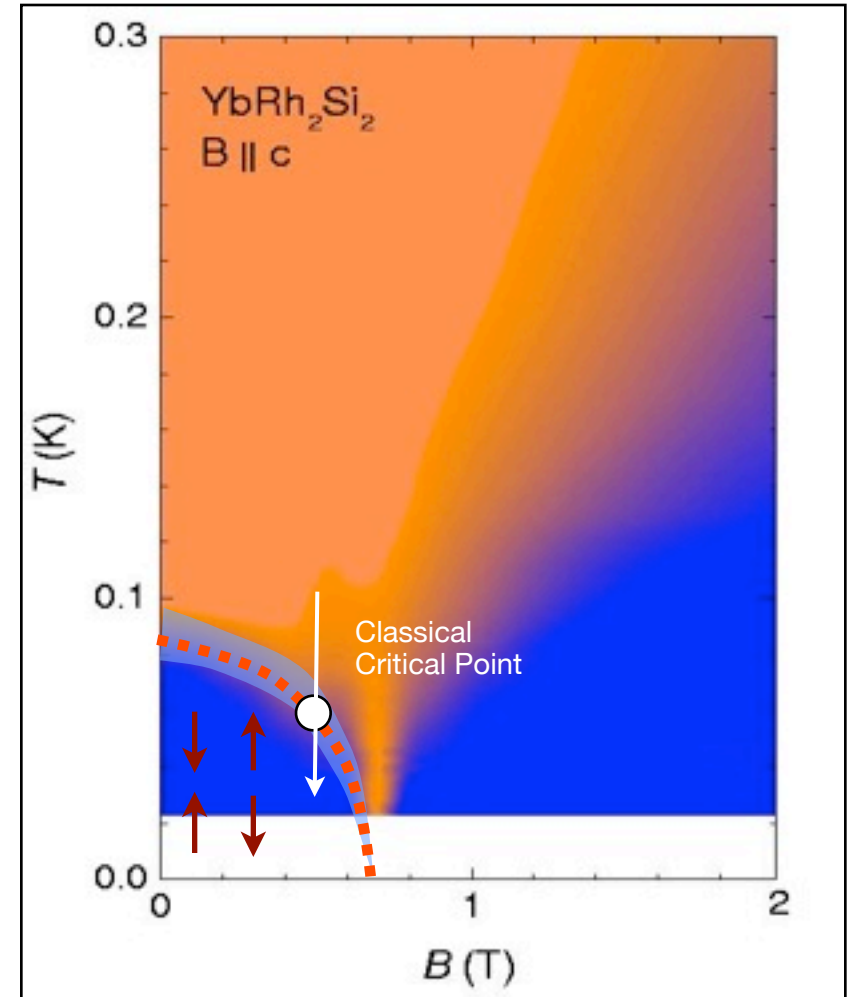
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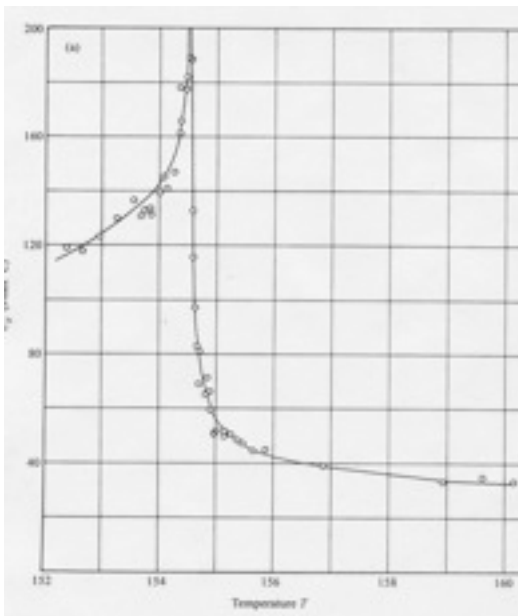
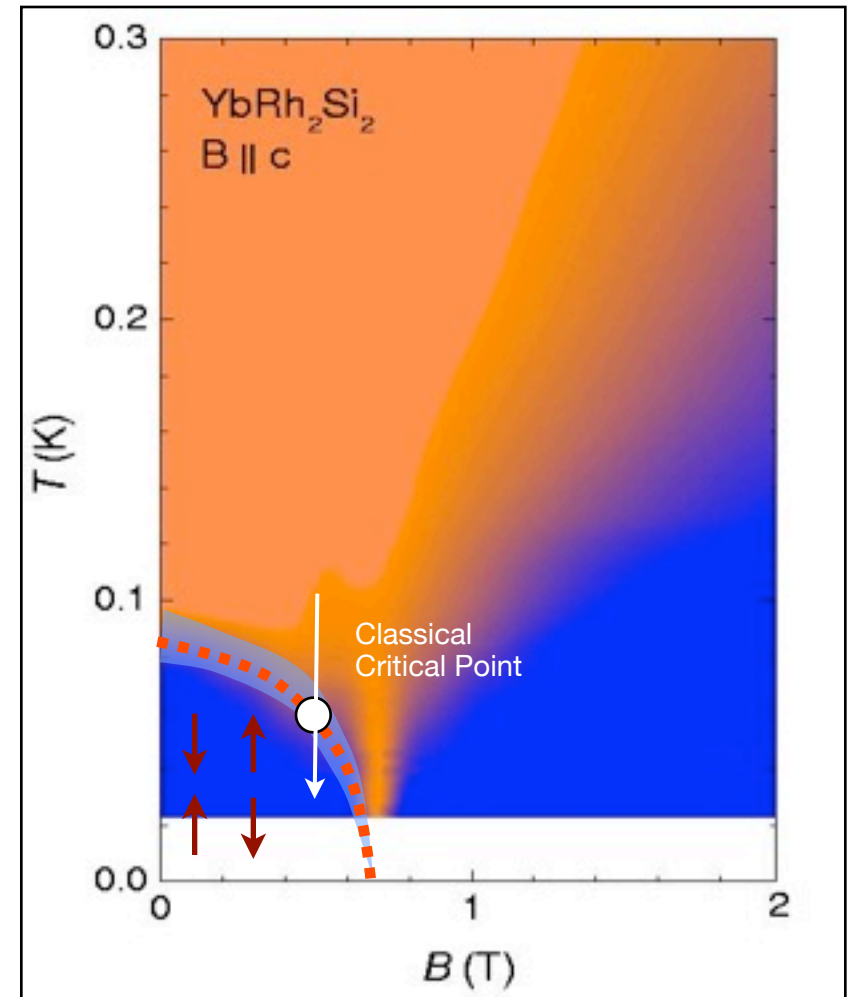
Oxygen. (Voronel et al 1963).

Classical Criticality

Michael Fisher



"New insights into physics often come from revisiting areas once thought to be closed." Michael Fisher.



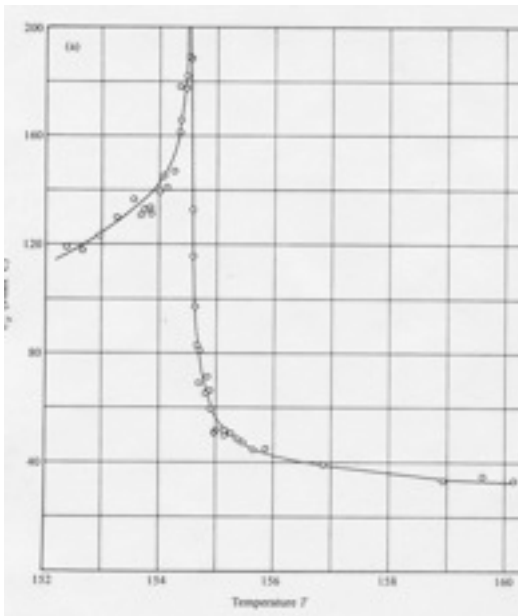
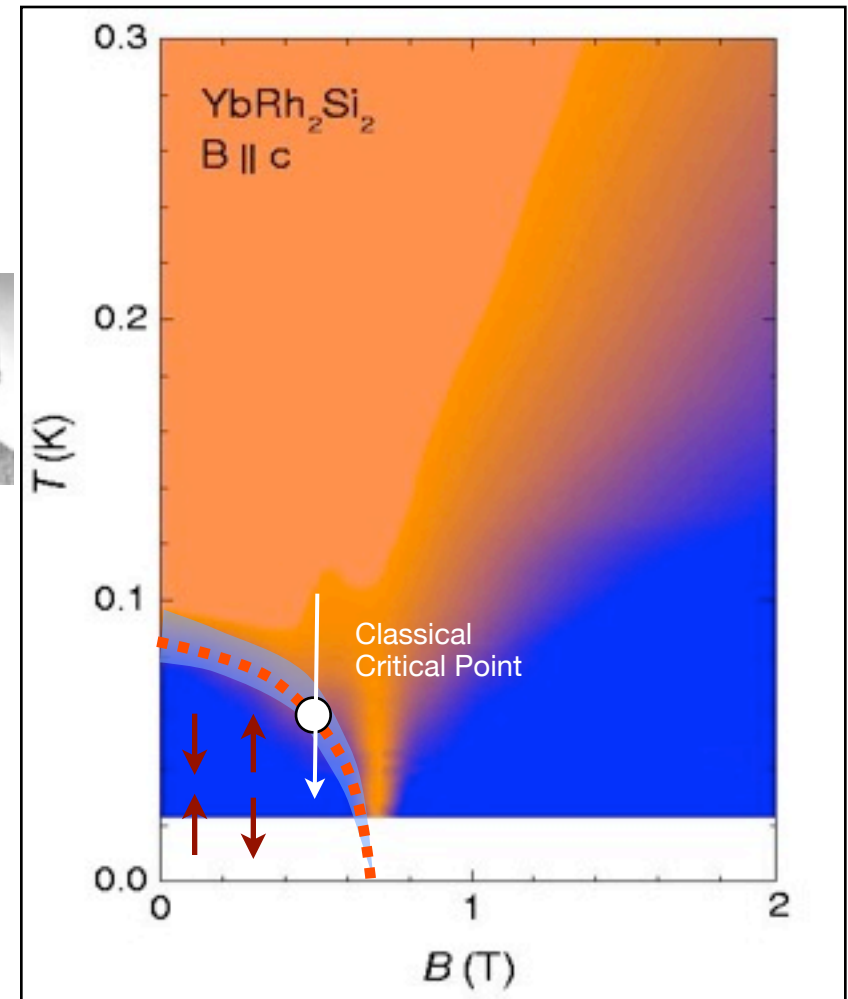
Oxygen. (Voronel et al 1963).

Classical Criticality

Michael Fisher Leo Kadanoff Ben Widom Anatoly Larkin Ken Wilson



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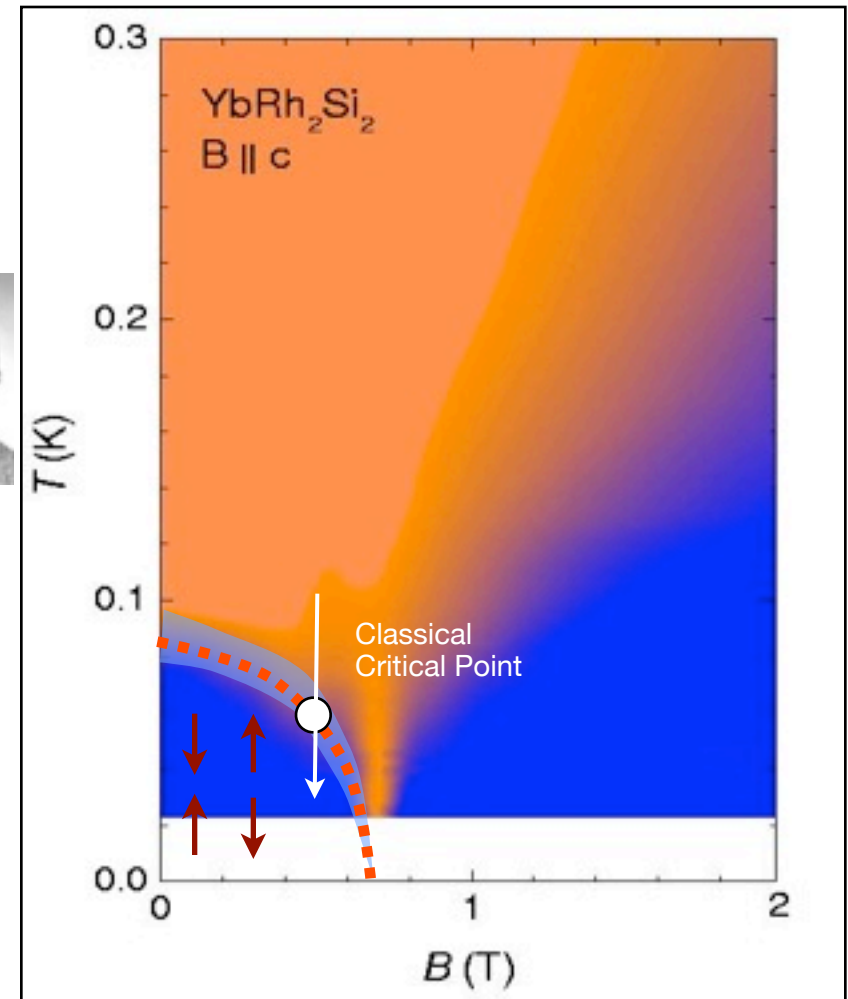
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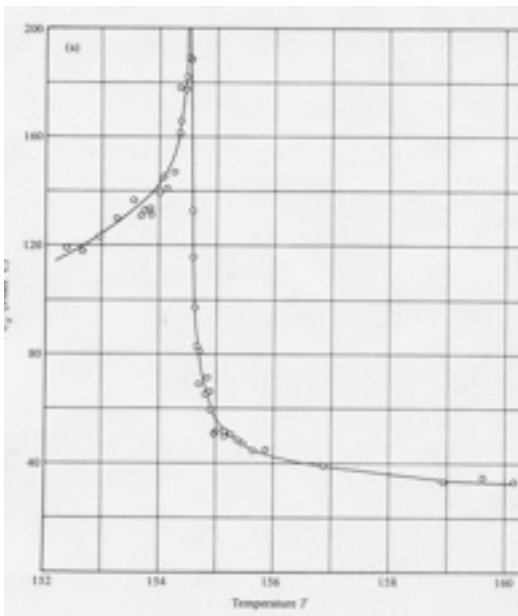


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“20th Century Revolution”

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Classical Criticality

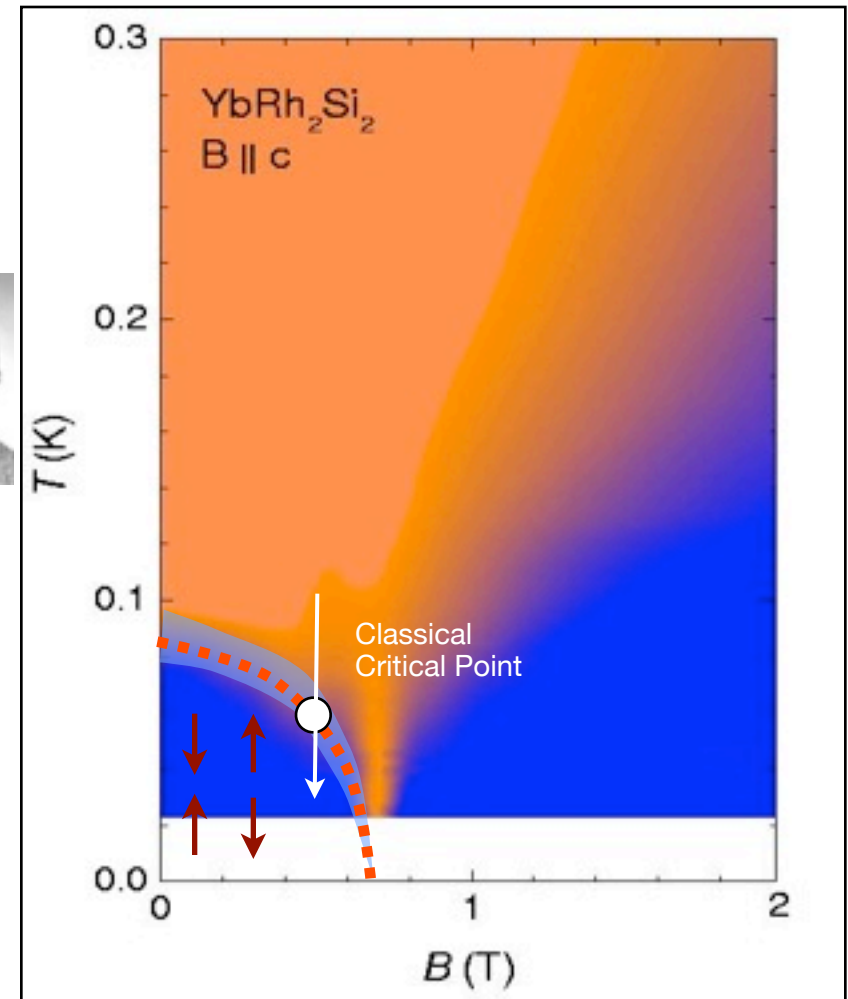
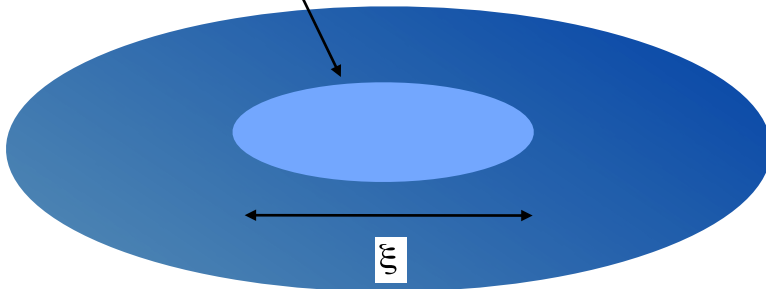
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$$\langle \psi(x)\psi(0) \rangle \sim \frac{1}{x^{d-2+\eta}}$$

Critical matter



“20th Century Revolution”

Classical Criticality

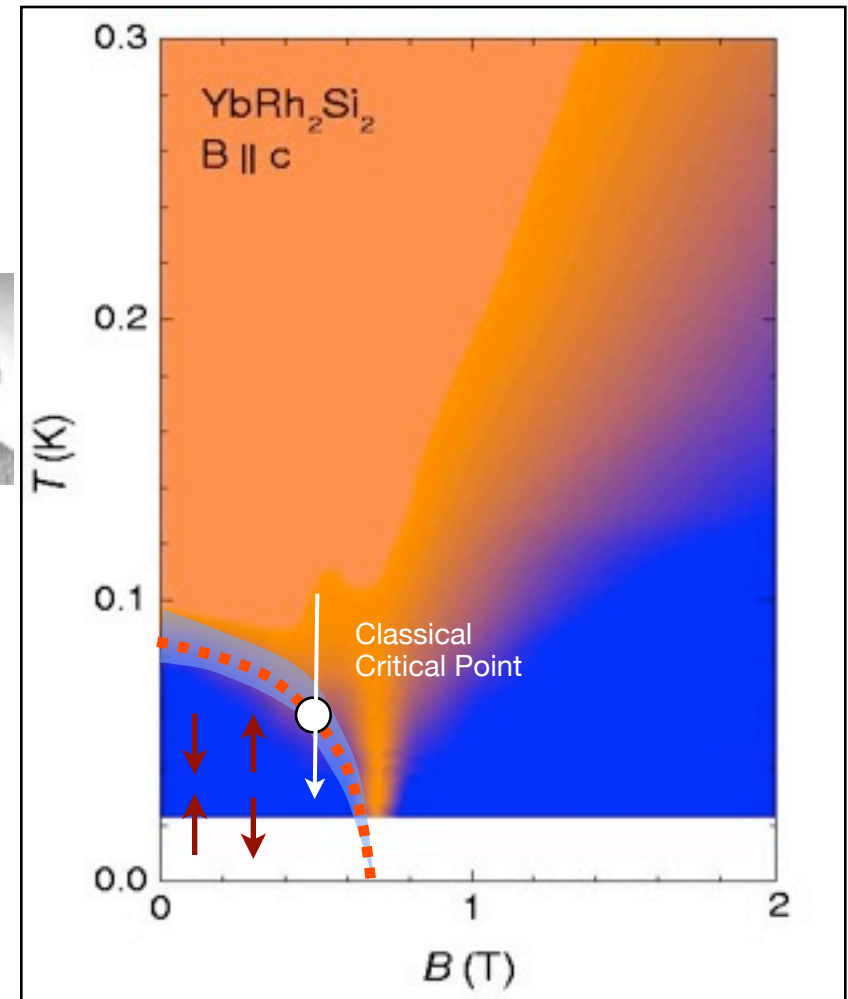
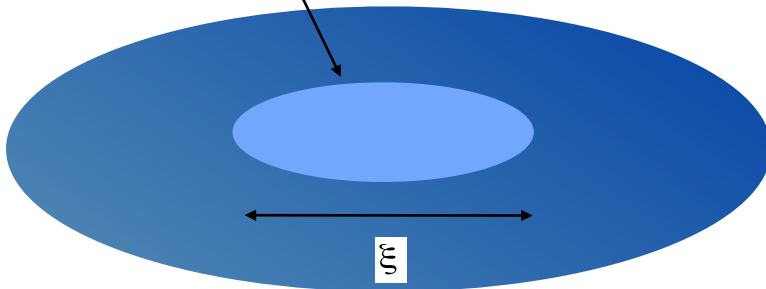
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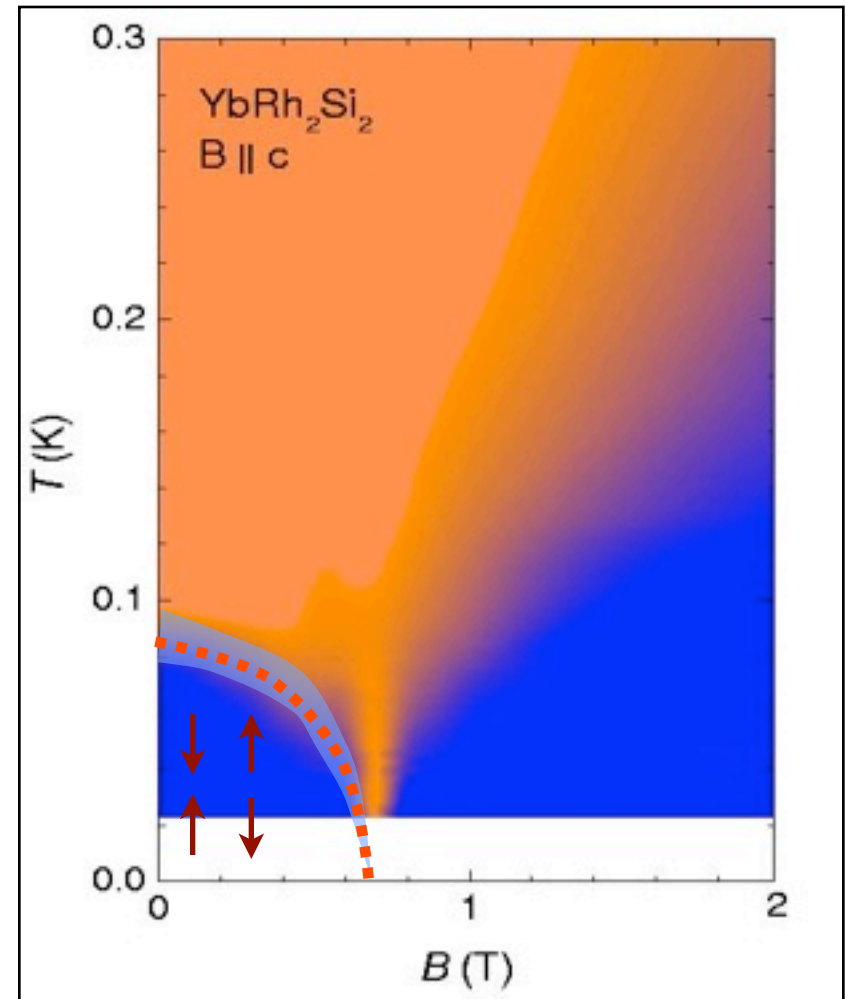
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Critical matter - universal



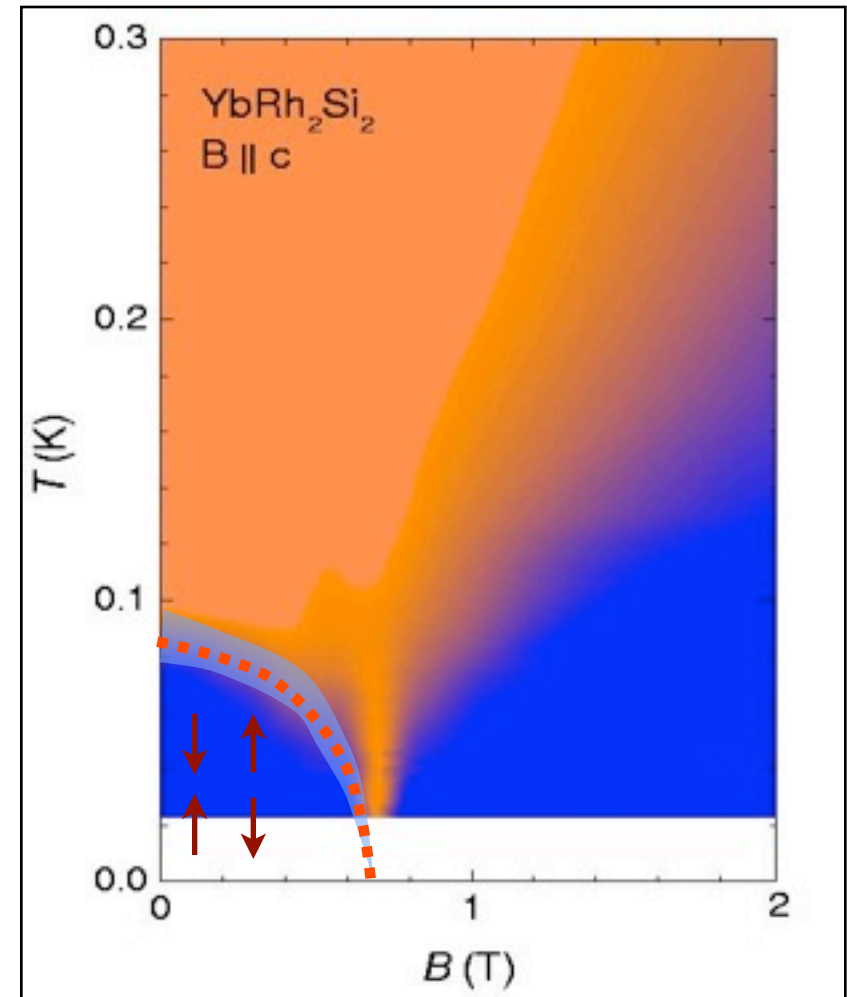
“20th Century Revolution”

Quantum Phase-Transition



Quantum Phase-Transition

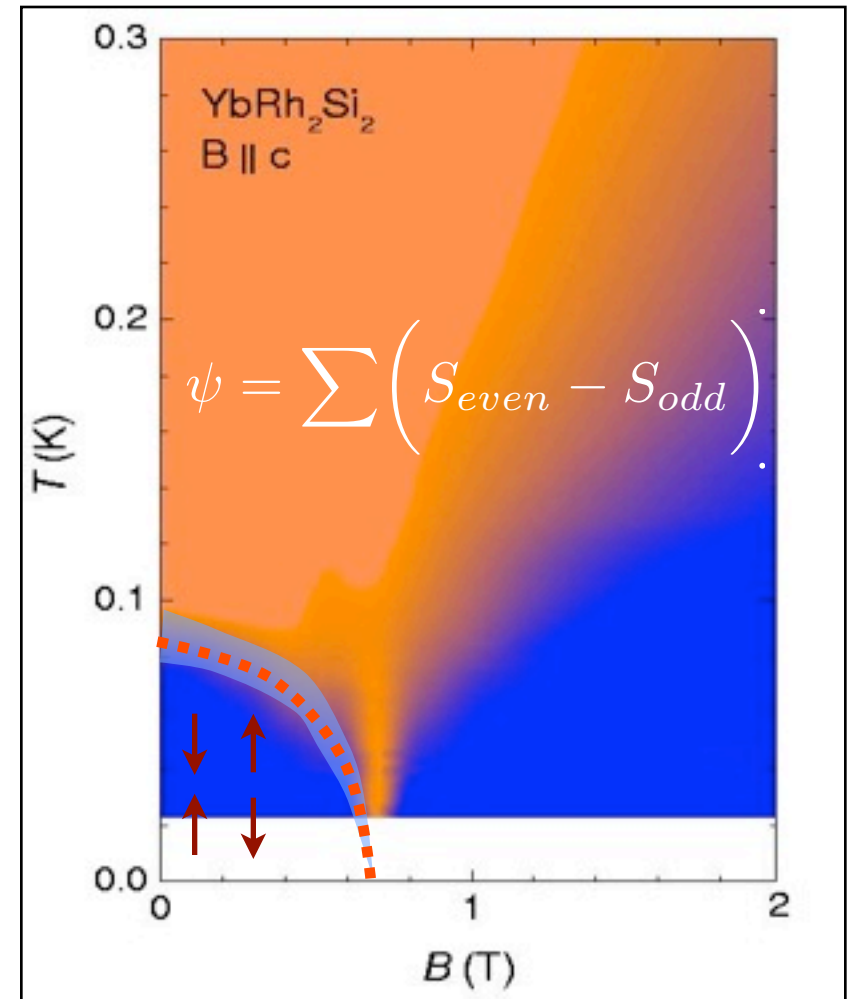
Phase transition
driven by zero point motion.



Quantum Phase-Transition

Phase transition
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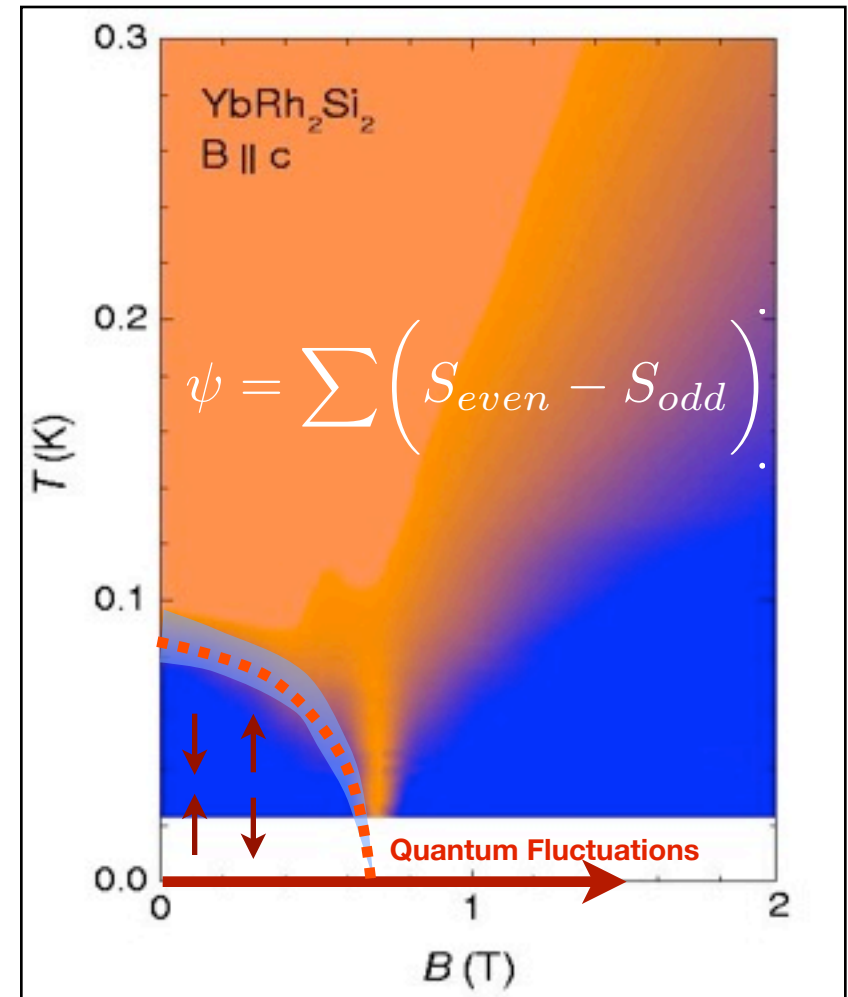
$$[H, \psi] \neq 0$$



Quantum Phase-Transition

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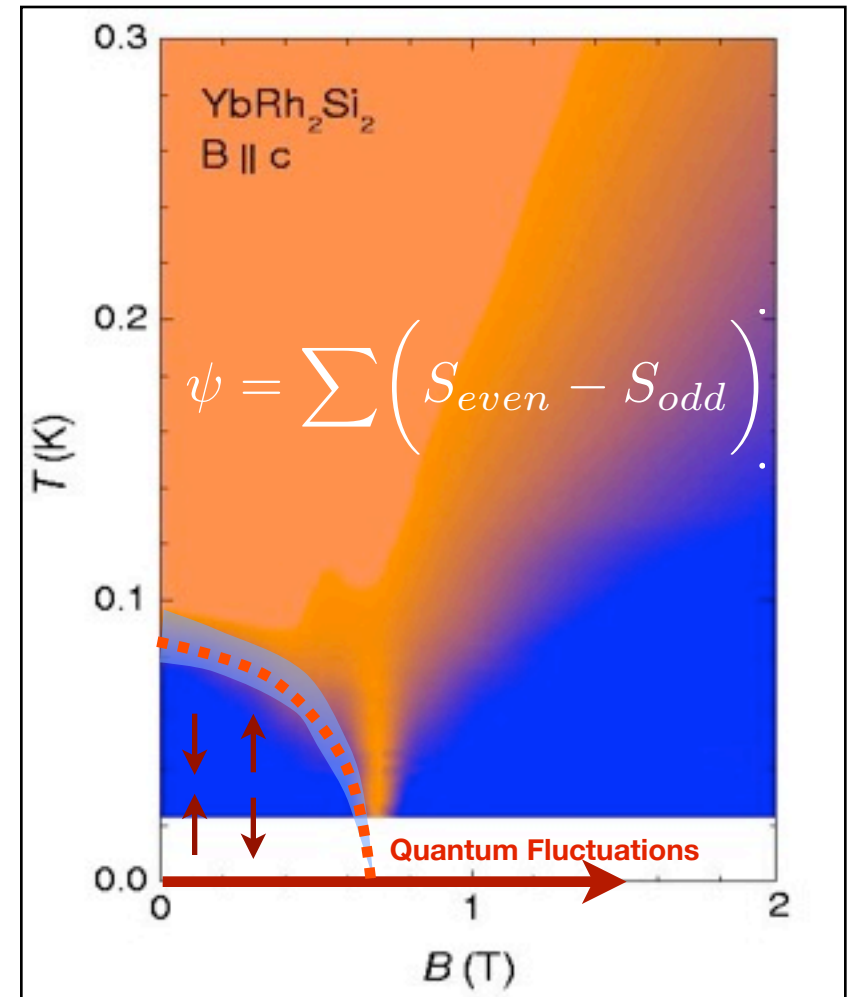
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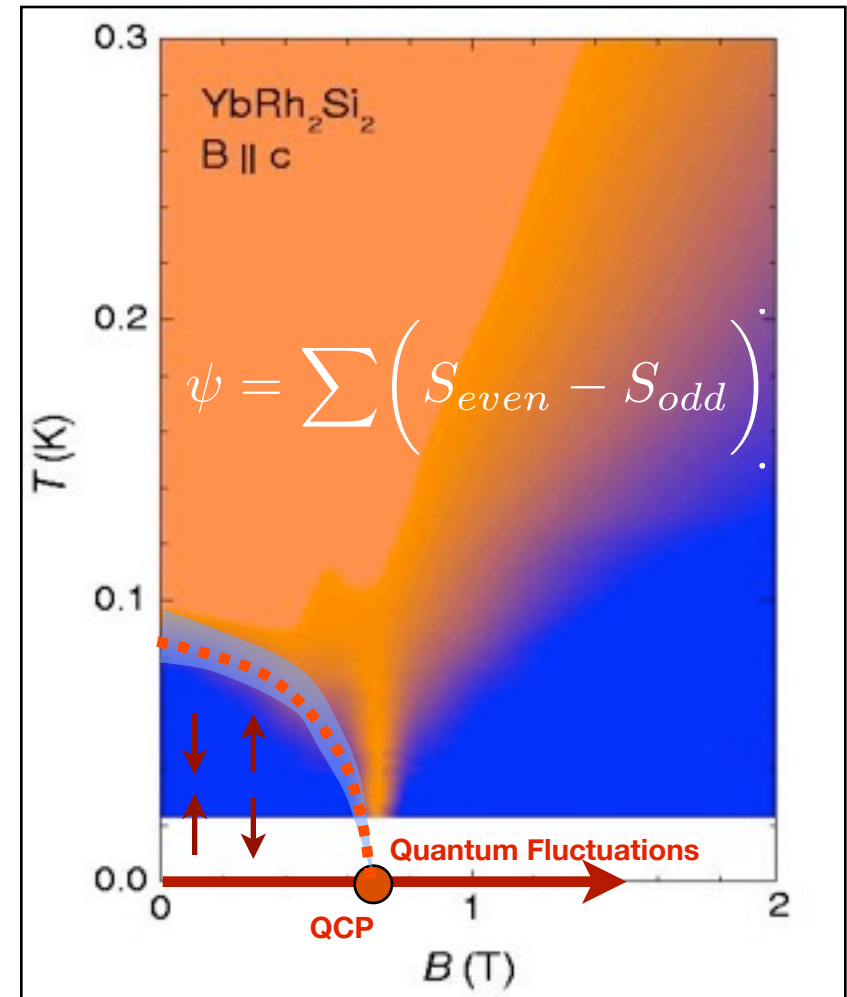


What happens when the time
and length scale of zero point
quantum fluctuations diverges?

Quantum Phase-Transition

Phase transition
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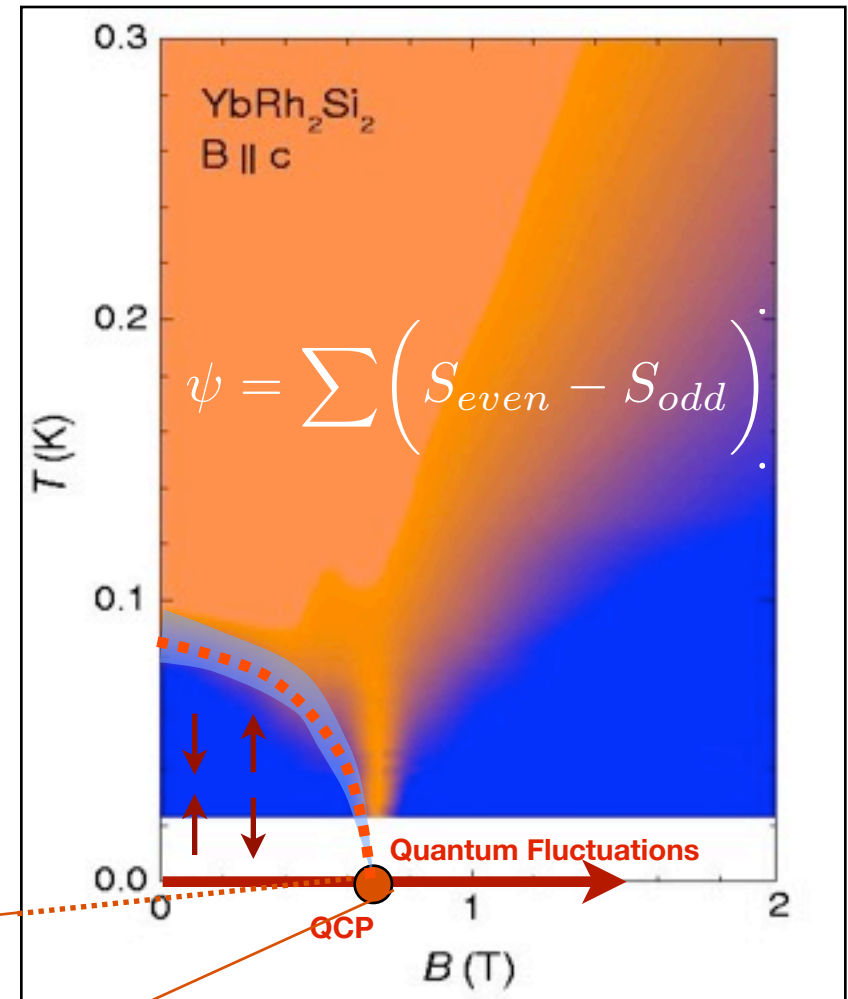
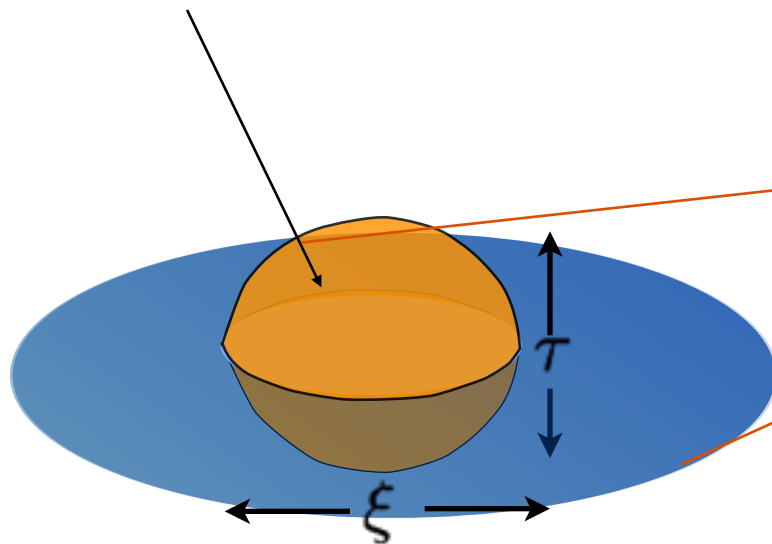
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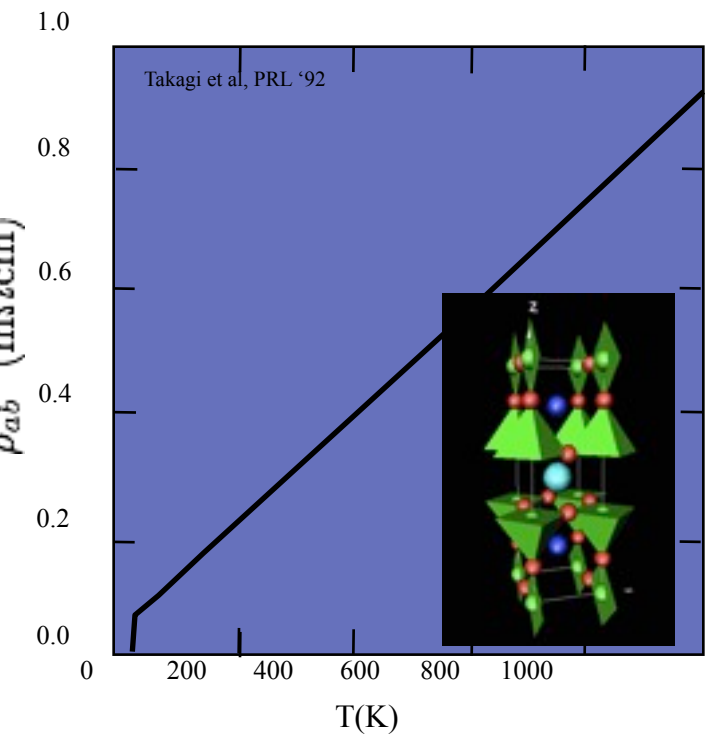
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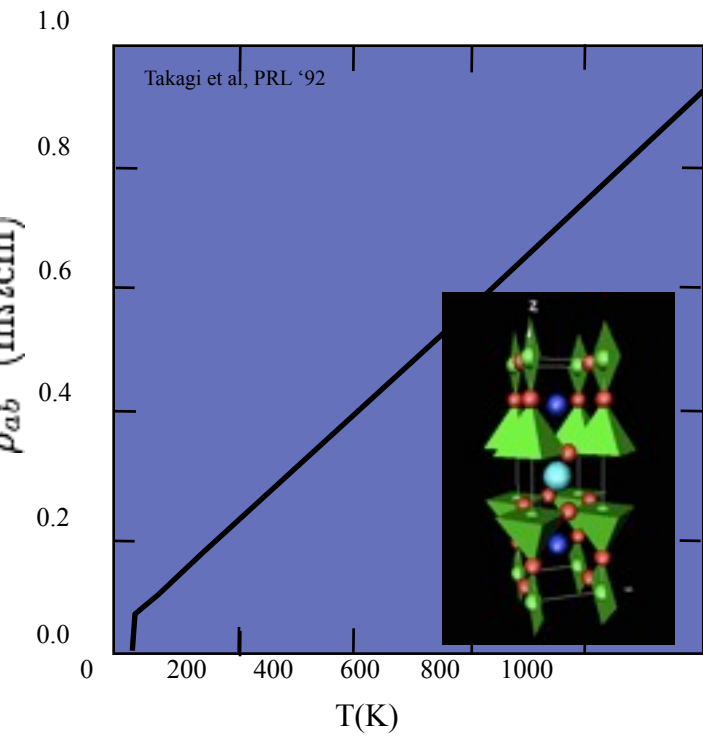
Quantum Critical matter



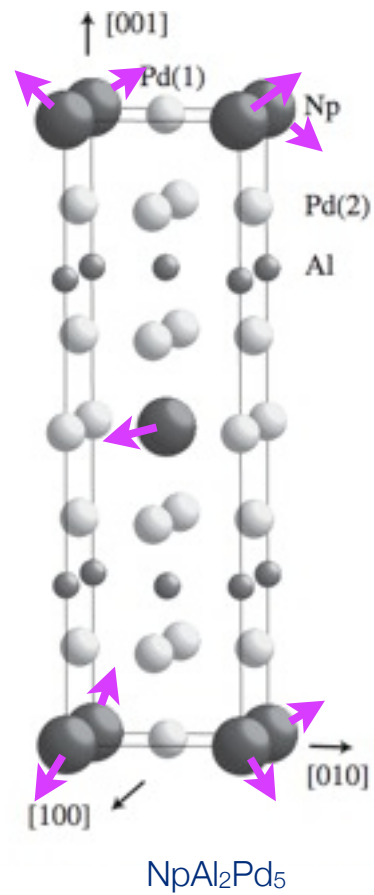
What happens when the time and length scale of zero point quantum fluctuations diverges?



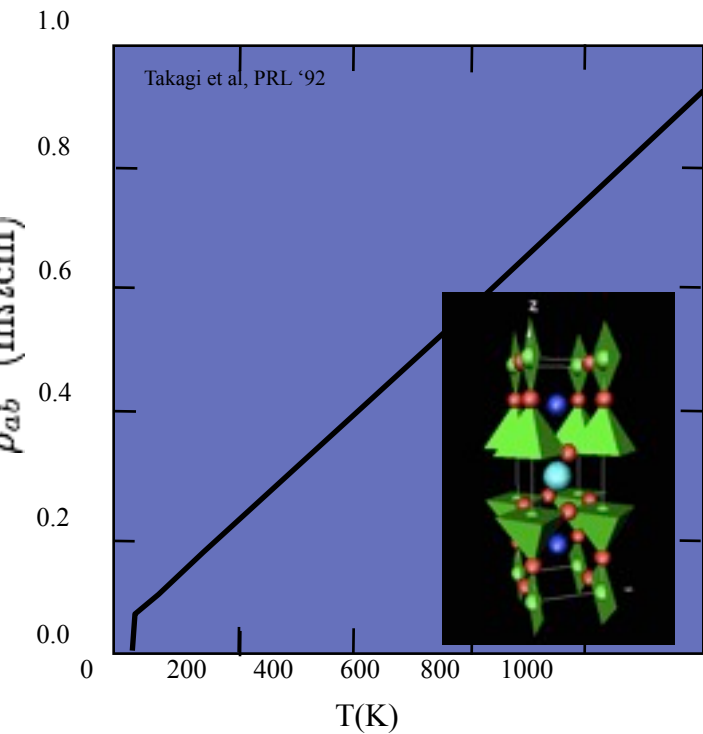
Cuprates
 $T_c=11-92\text{K}$



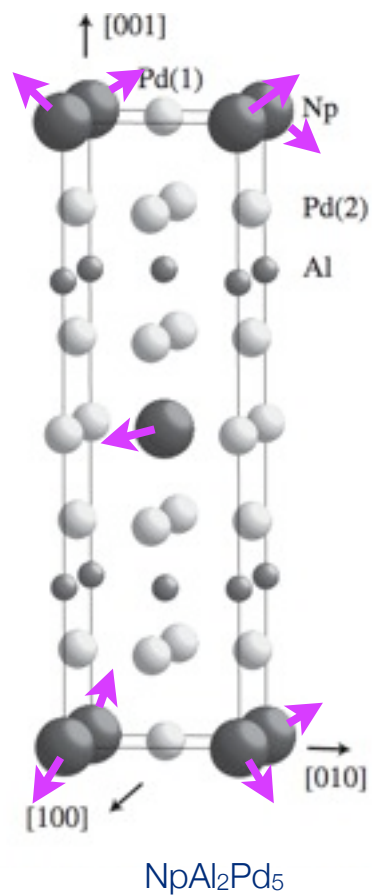
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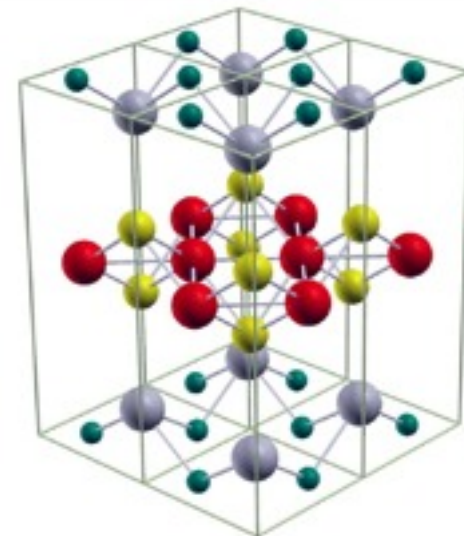
HF 115s
 $T_c = 0.2 - 18.5 \text{ K}$



Cuprates
 $T_c = 11 - 92 \text{ K}$

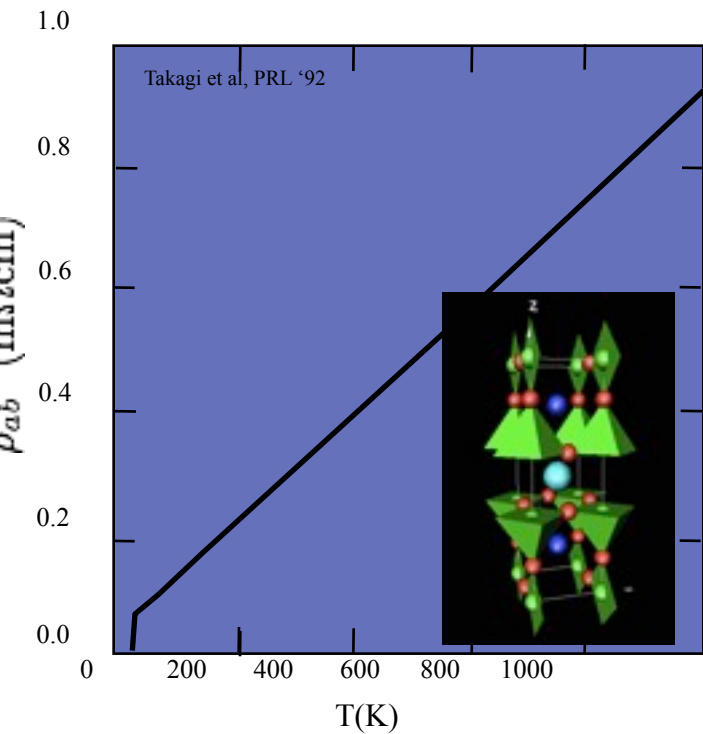


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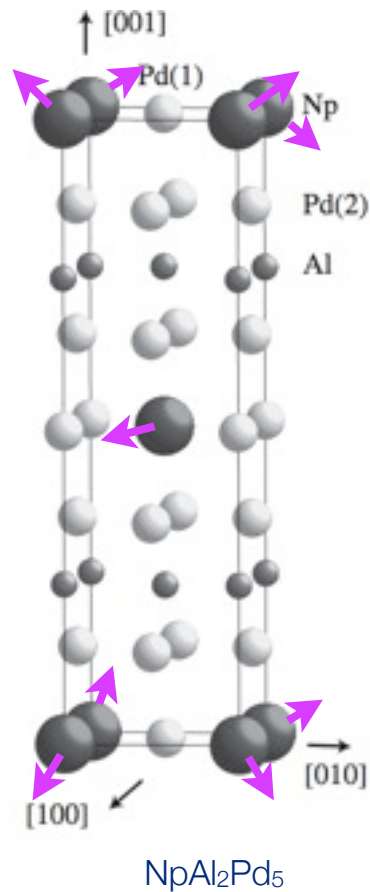


$\text{PrF}_x\text{O}_{1-x}\text{FeAs}$
 Z.A. Ren et.al, Beijing, (08)

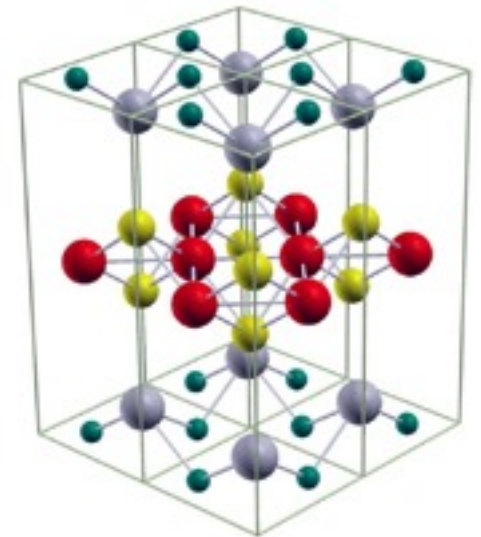
Iron based SC
 $T_c = 6 - 53 \text{ ++ ? K}$



Cuprates
 $T_c = 11 - 92 \text{ K}$



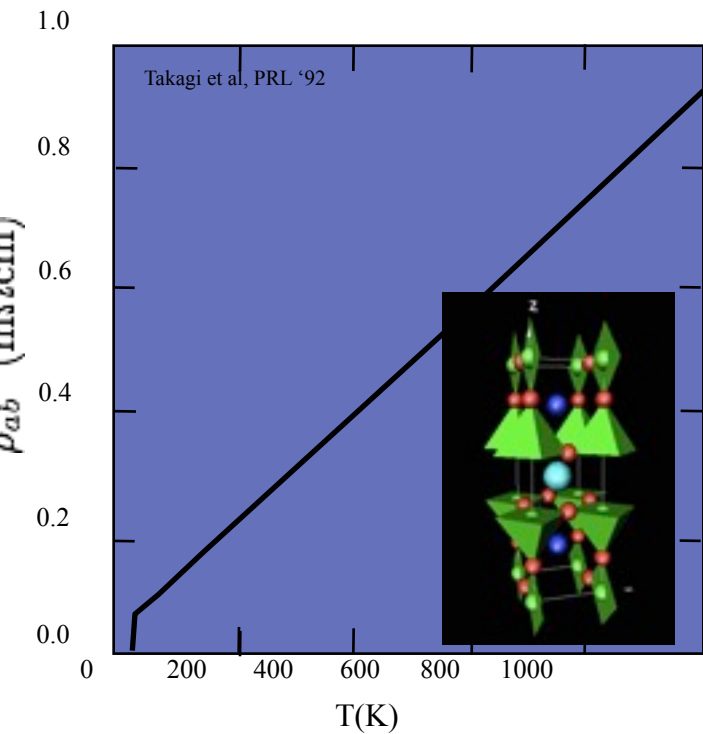
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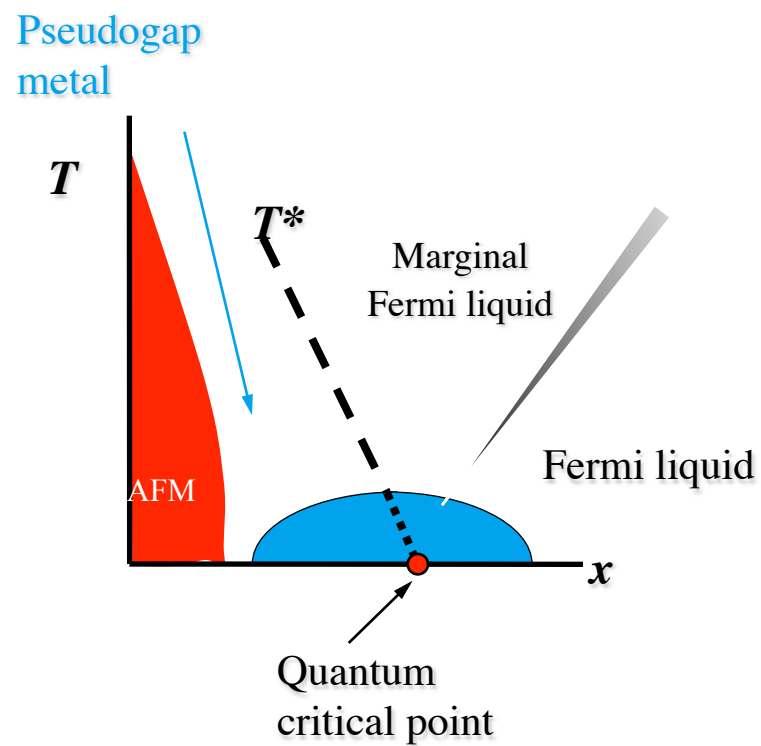
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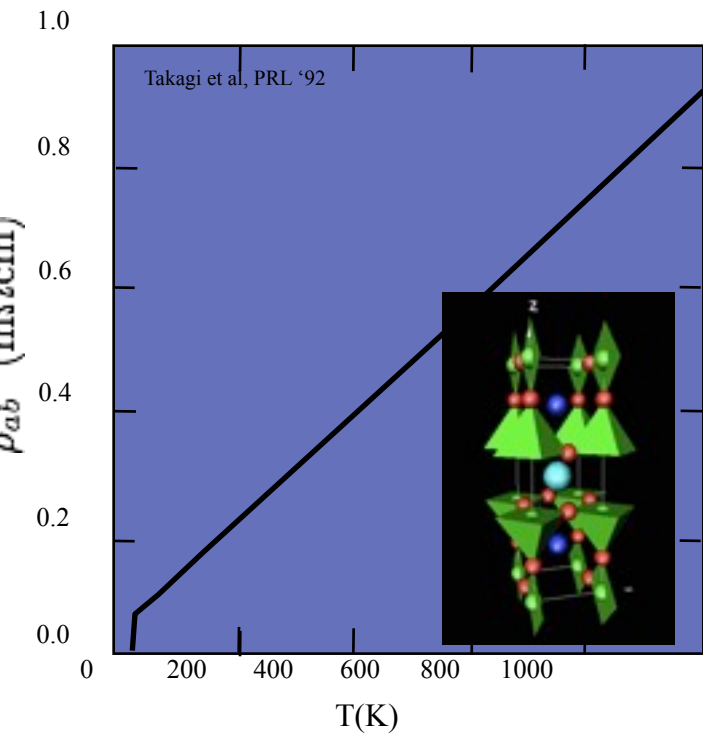
Iron based SC
 $T_c = 6 - 53 \text{ ++ ? K}$

What kind of electronic fluid yields high T_c superconductivity?

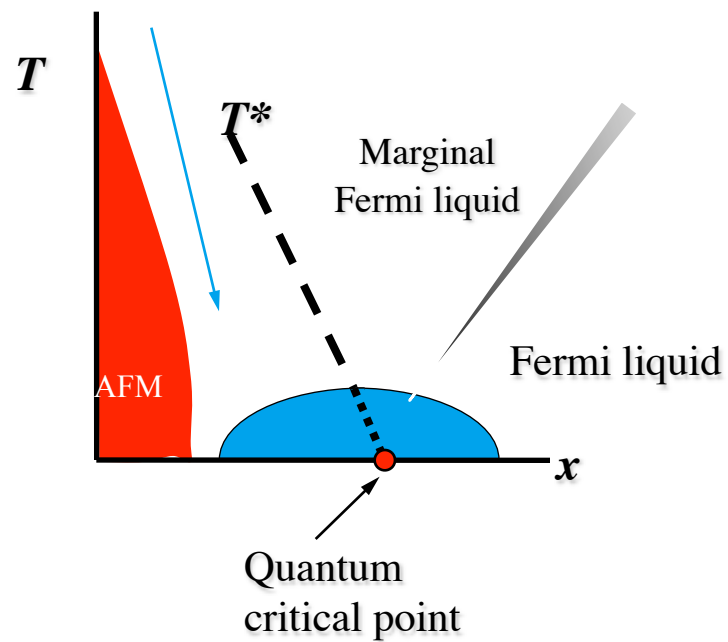


Cuprates
 $T_c = 11-92\text{K}$



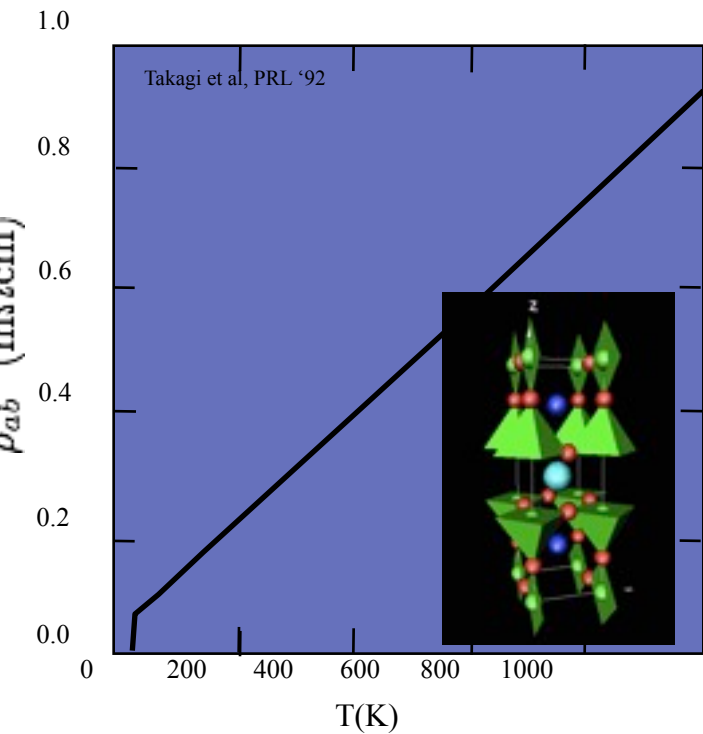


Pseudogap metal

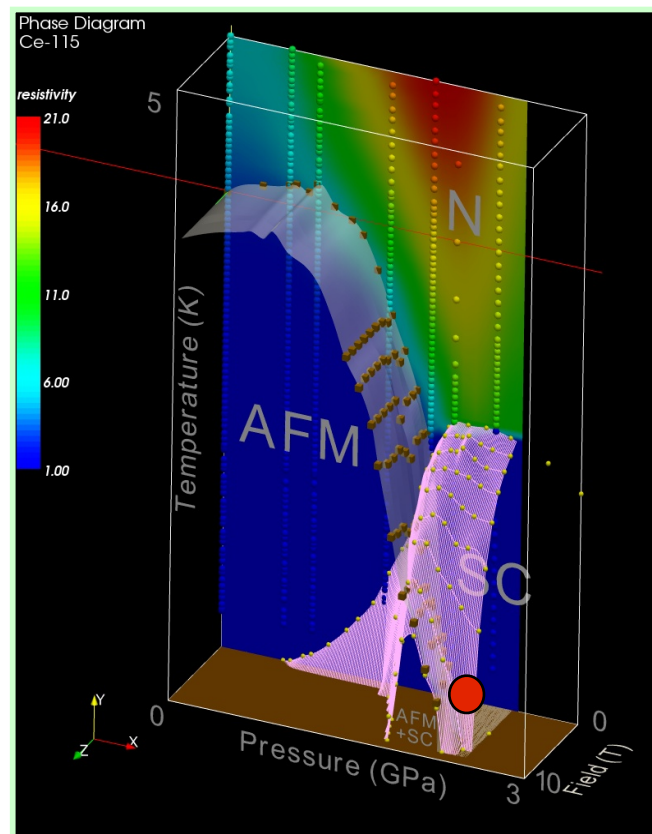


Cuprates
 $T_c=11-92K$

“Avoided criticality”



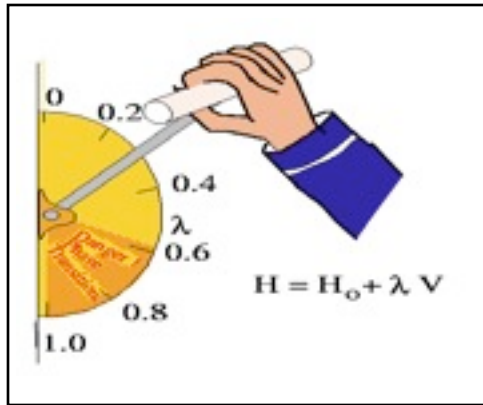
CeRhIn₅



Cuprates
T_c=11-92K

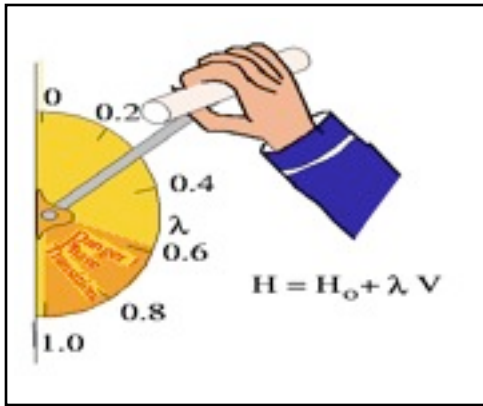
“Avoided criticality”

Tuson Park, (2007).



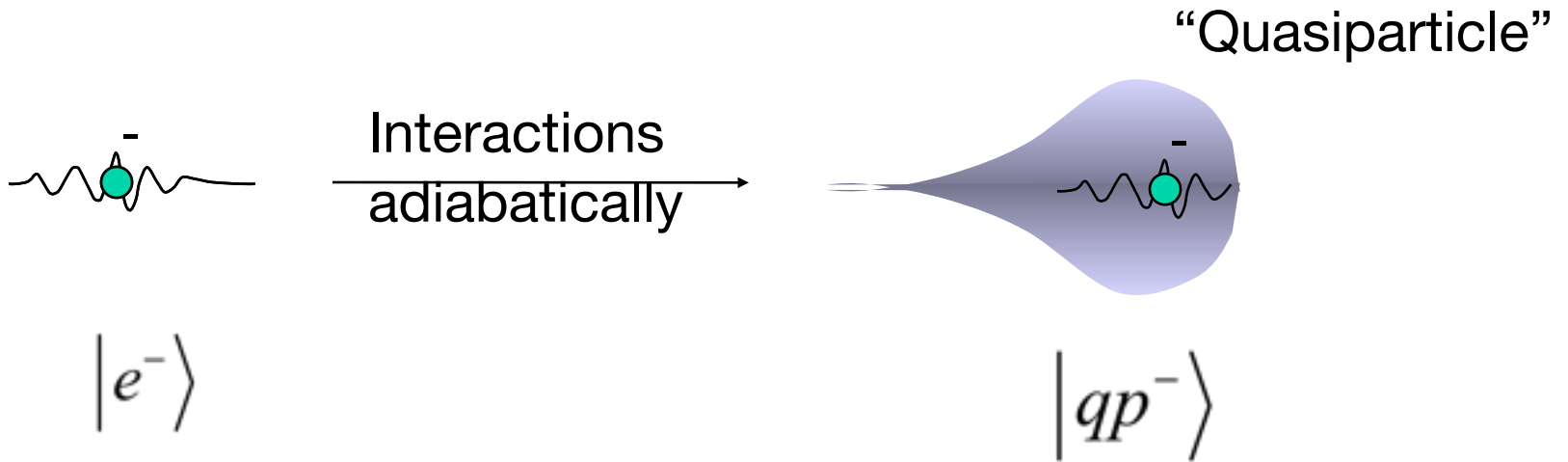
Landau: interactions can be turned on adiabatically, preserving the excitation spectrum.

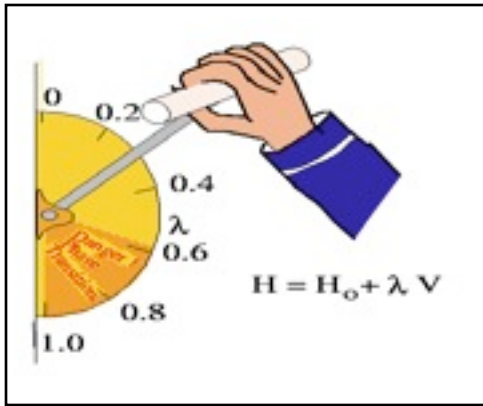
Landau, JETP 3, 920 (1957)



Landau: interactions can be turned on adiabatically, preserving the excitation spectrum.

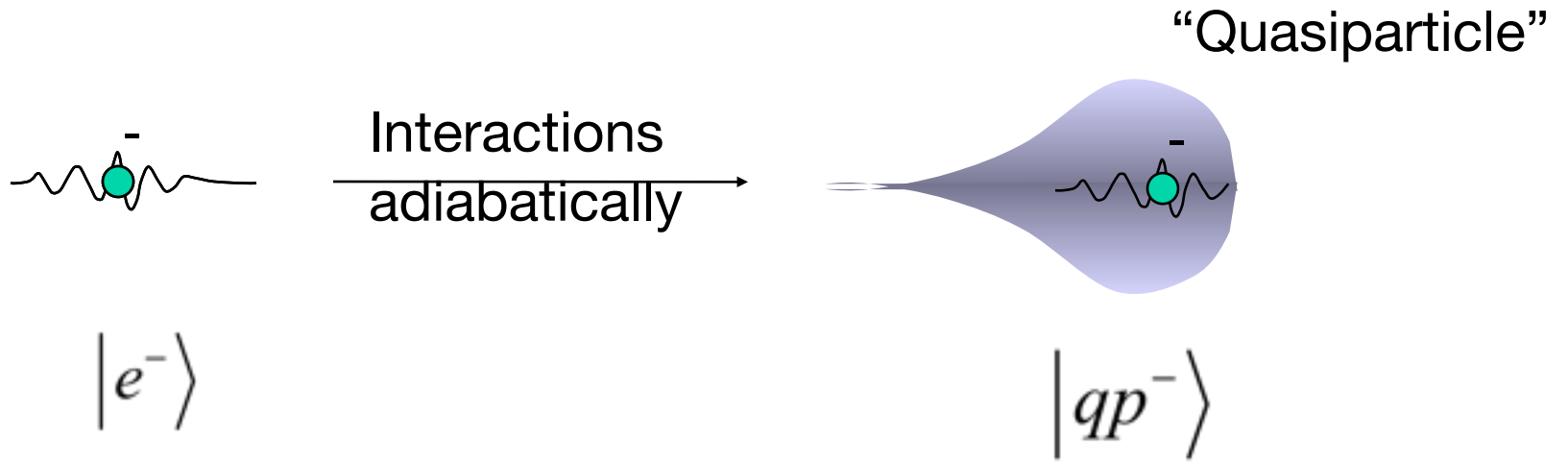
Landau, JETP 3, 920 (1957)



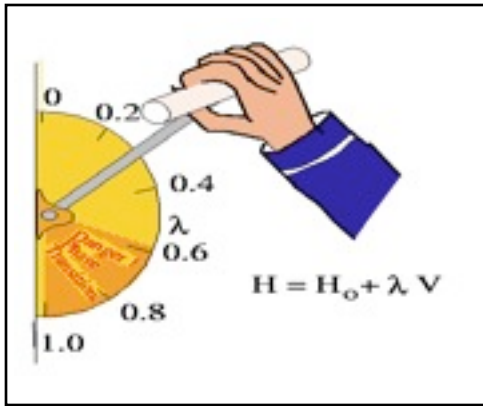


Landau: interactions can be turned on adiabatically, preserving the excitation spectrum.

Landau, JETP 3, 920 (1957)



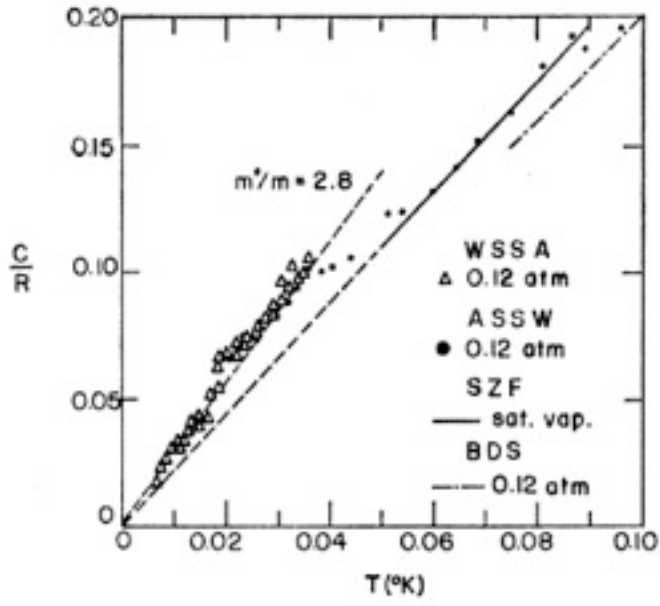
$$\frac{m^*}{m} = \frac{N(0)^*}{N(0)} = 1 + \frac{F_1^s}{3}$$



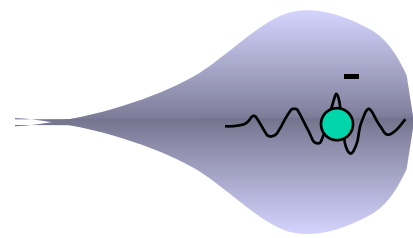
Landau: interactions can be turned on adiabatically, preserving the excitation spectrum.

Landau, JETP 3, 920 (1957)

He-3 (1950/60s)
(Fairbanks, many others)



“Quasiparticle”



$$|qp^-\rangle$$

$$\frac{m^*}{m} = \frac{N(0)^*}{N(0)} = 1 + \frac{F_1^s}{3}$$

Landau's Question.



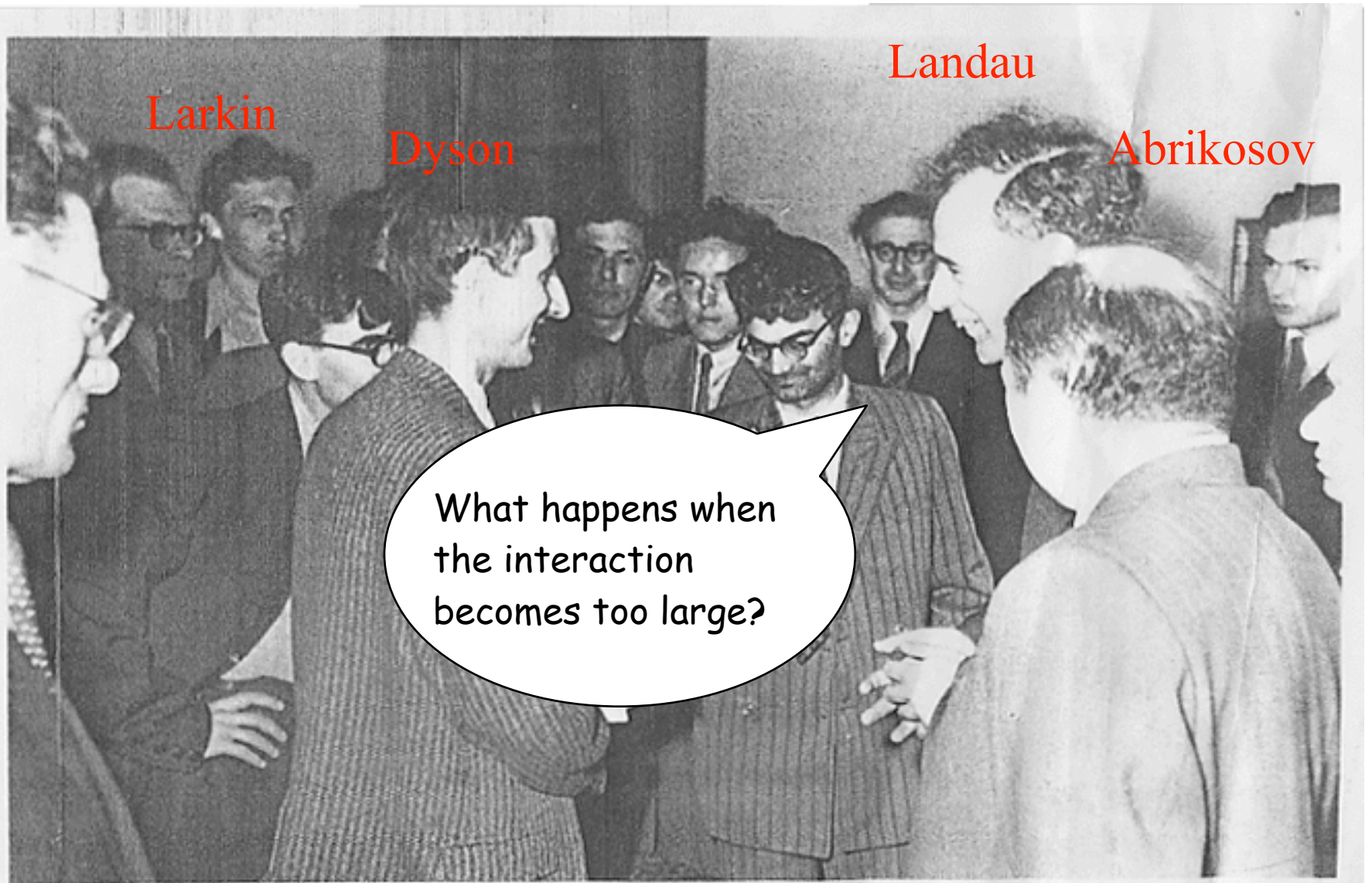
20. Moscow, 1956. Freeman Dyson (front, left), talking with I. Pomeranchuk and Lev Landau.

Landau's Question.



20. Moscow, 1956. Freeman Dyson (front, left), talking with I. Pomeranchuk and Lev Landau.

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What happens when the interaction becomes too large?



Landau 1936



“Electrons order”

What happens when the interaction becomes too large?



What happens when the interaction becomes too large?

Landau 1936

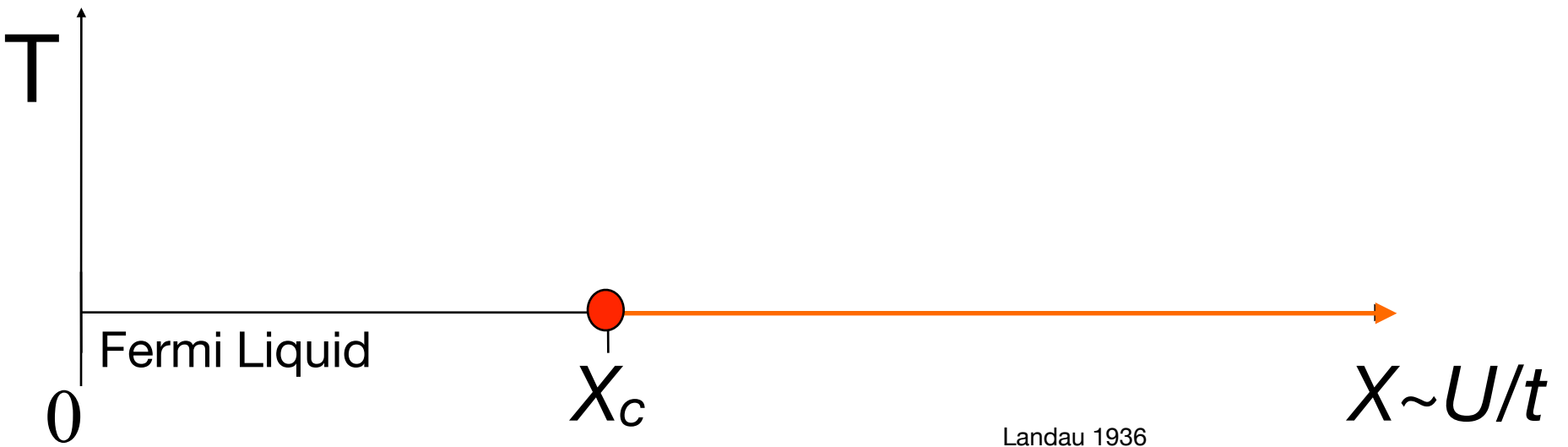


“Electrons order”

Mott 1947



“Electrons localize”



What happens when the interaction becomes too large?

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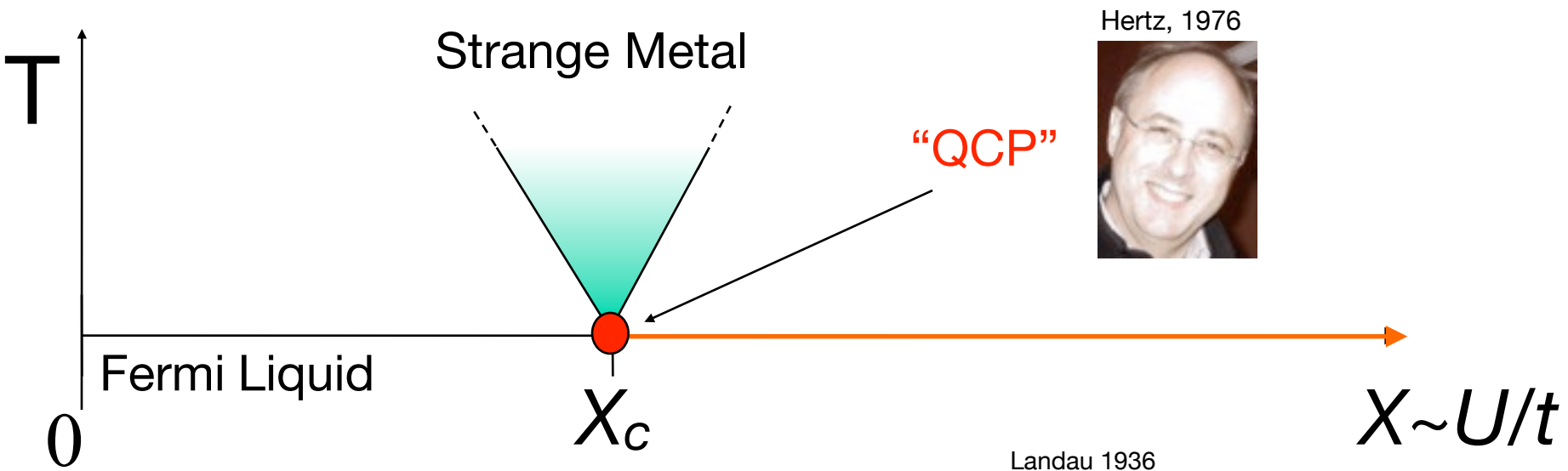


“Electrons localize”

Anderson 1961

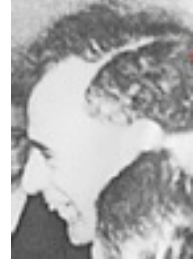


“Local moments form”



What happens when the interaction becomes too large?

Landau 1936



“Electrons order”

Mott 1947



“Electrons localize”

Anderson 1961



“Local moments form”

Kondo effect

Kondo effect

(a digression)

Kondo effect

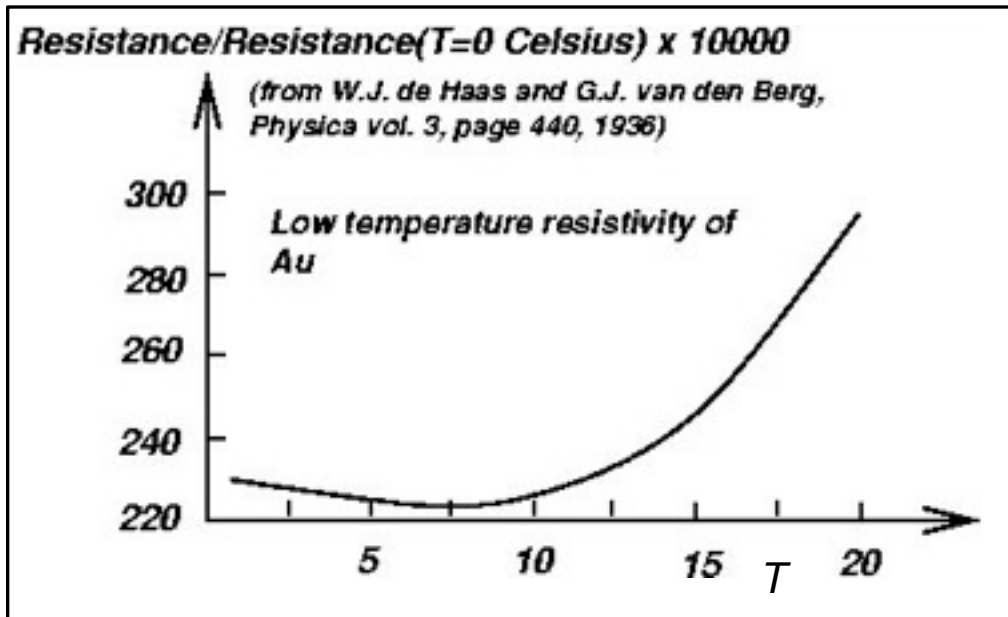
(a digression)



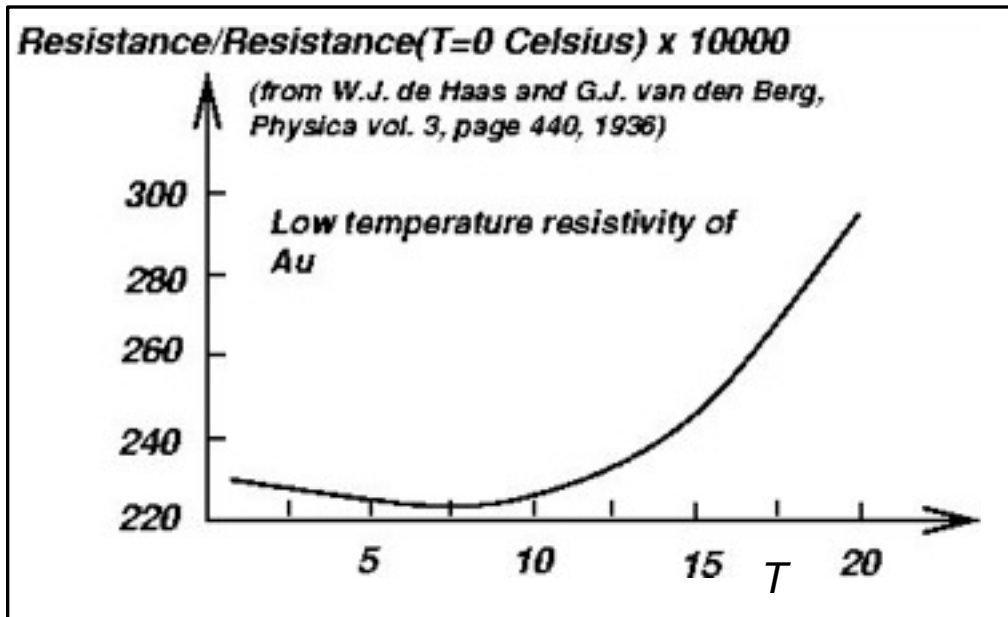
“Kondo”

Kondo KHR-2 HV, the robot that plays soccer, fights with other bots and dances salsa

Kondo effect (a digression)

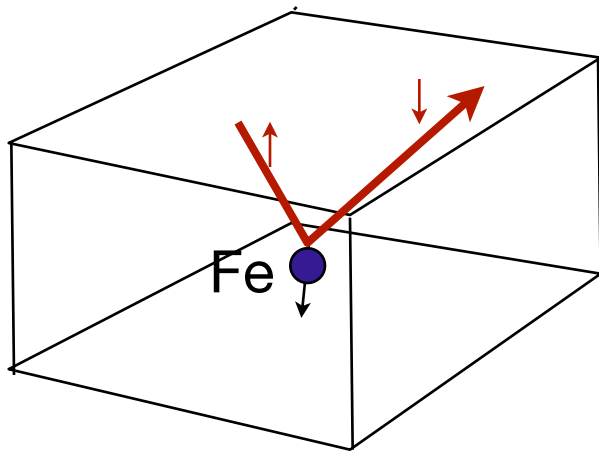


Kondo effect (a digression)



“A 75 year odyssey”

Kondo effect (a digression)



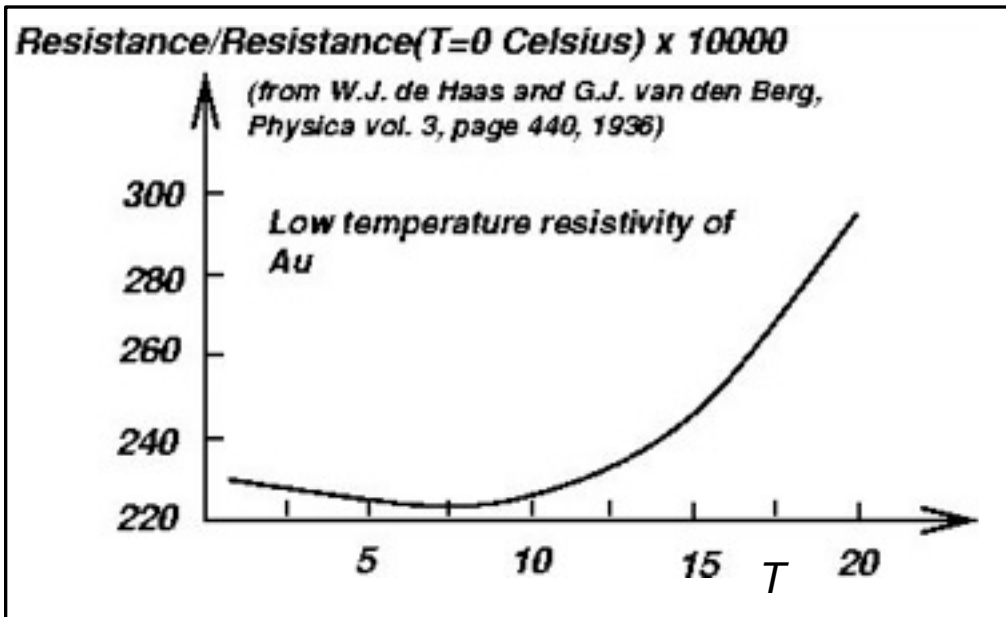
$$H = \sum \epsilon_k c_{k\sigma}^\dagger c_{k\sigma} + J(\psi^\dagger \vec{\sigma} \psi) \cdot \vec{S}$$

Kondo (1962)



$$T_K = D\sqrt{J\rho} \exp\left[-\frac{1}{2J\rho}\right]$$

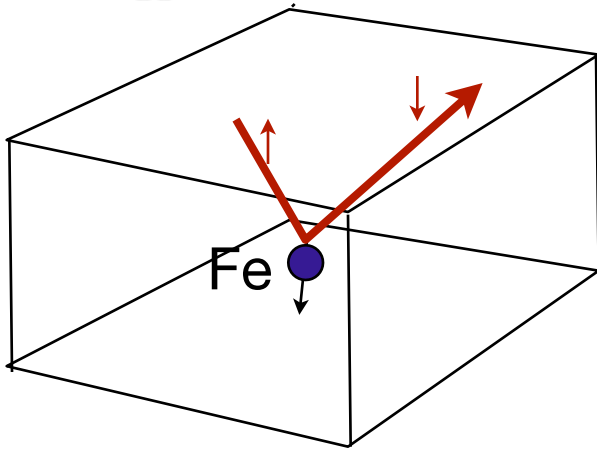
“Kondo Temperature”



“A 75 year odyssey”

Kondo effect (a digression)

$$T \gg T_K$$



Spins asymptotically free

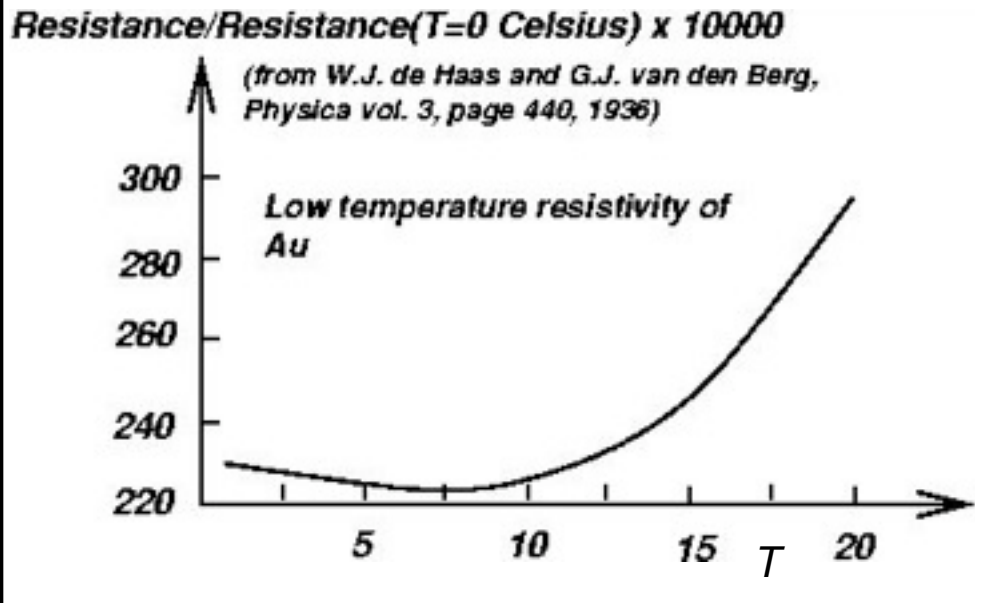
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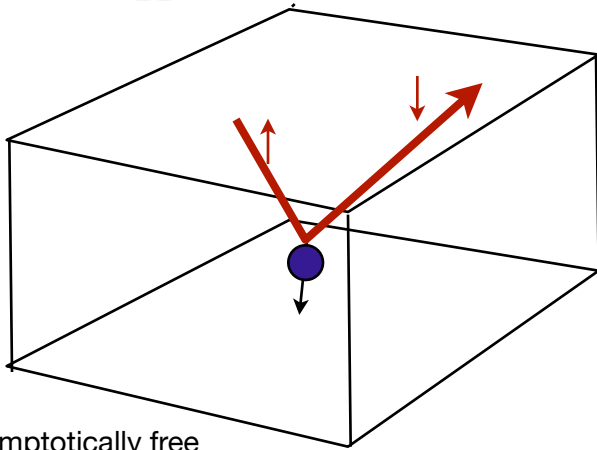
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Kondo (1962)

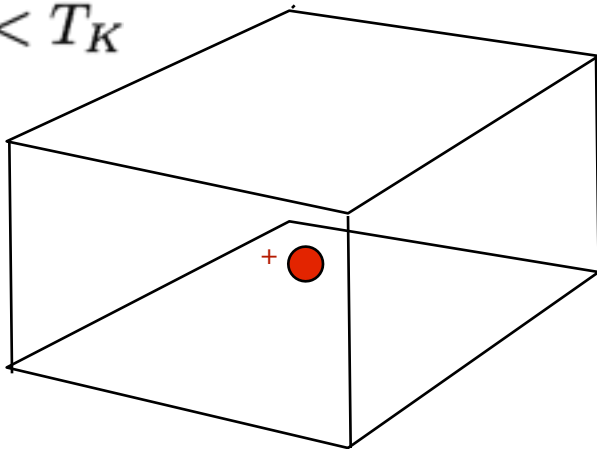


+

$$T_K = D\sqrt{J\rho} \exp\left[-\frac{1}{2J\rho}\right]$$

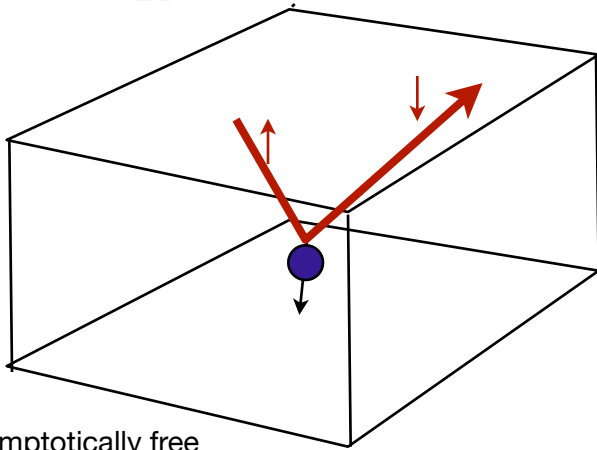
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Kondo (1962)



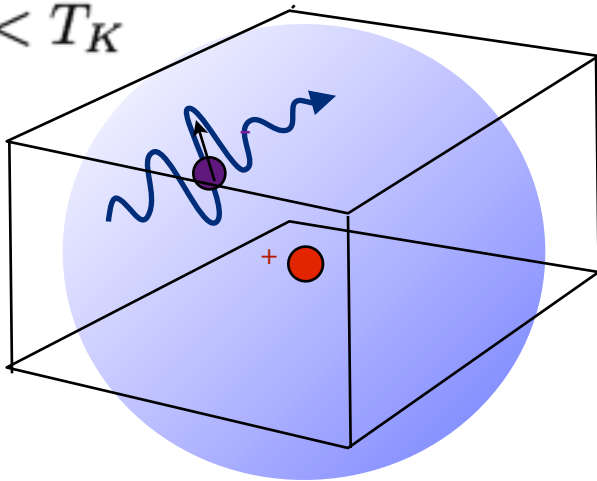
+

Spins asymptotically free

$$T_K = D\sqrt{J\rho} \exp\left[-\frac{1}{2J\rho}\right]$$

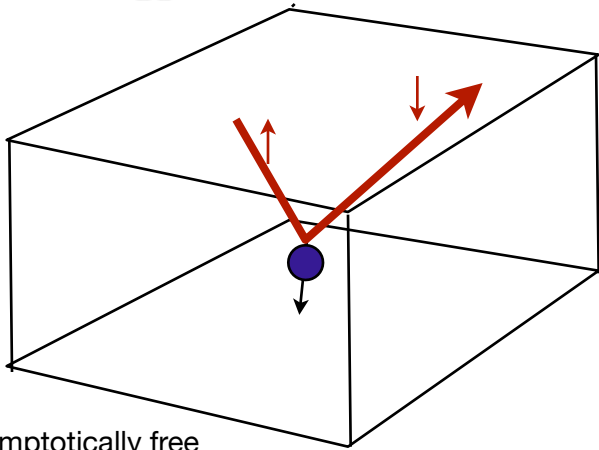
“Kondo Temperature”

$$T \ll T_K$$



Kondo effect (a digression)

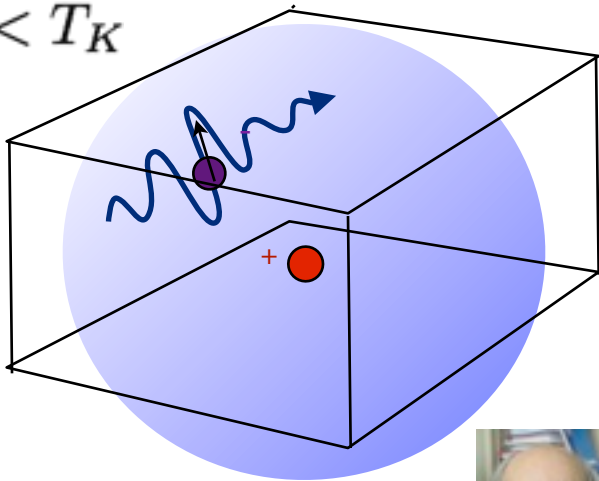
$$T \gg T_K$$



Spins asymptotically free

Spins absorbed into singlet ground-state

$$T \ll T_K$$



“Nozieres Local Fermi liquid”
(Nozieres 76)



$$H = \sum \epsilon_k c_{k\sigma}^\dagger c_{k\sigma} + J(\psi^\dagger \vec{\sigma} \psi) \cdot \vec{S}$$

Kondo (1962)



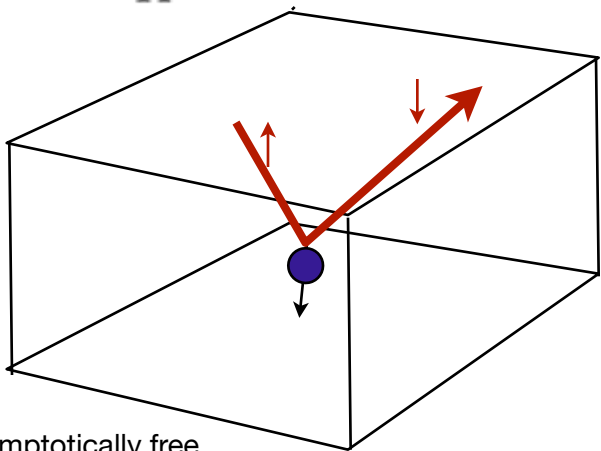
+

$$T_K = D\sqrt{J\rho} \exp\left[-\frac{1}{2J\rho}\right]$$

“Kondo Temperature”

Kondo effect (a digression)

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Kondo (1962)



+

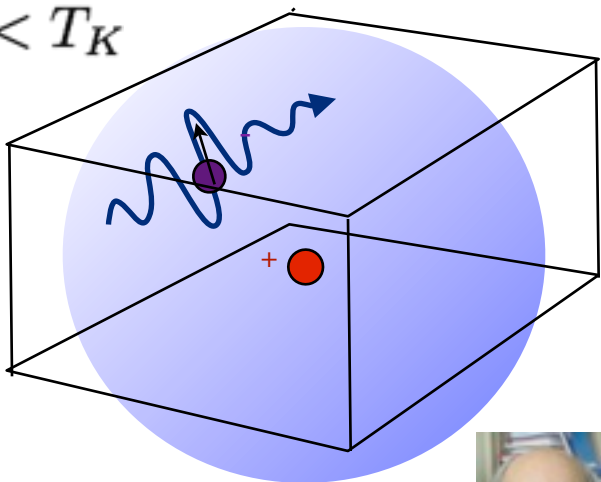
Spins asymptotically free

Spins absorbed into singlet ground-state

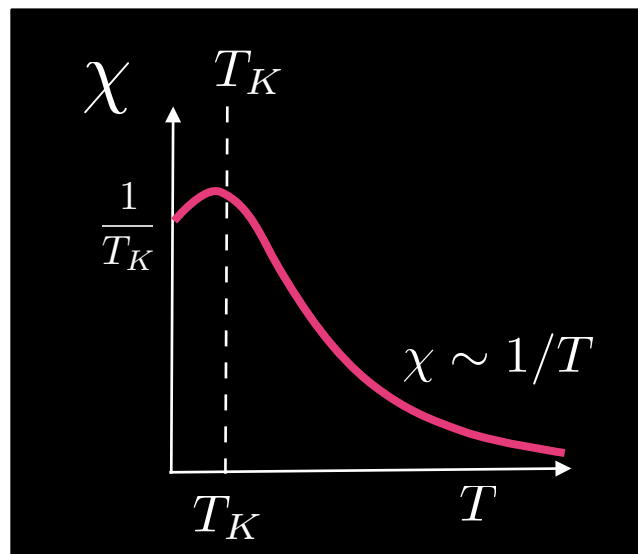
$$T_K = D\sqrt{J\rho} \exp\left[-\frac{1}{2J\rho}\right]$$

“Kondo Temperature”

$$T \ll T_K$$

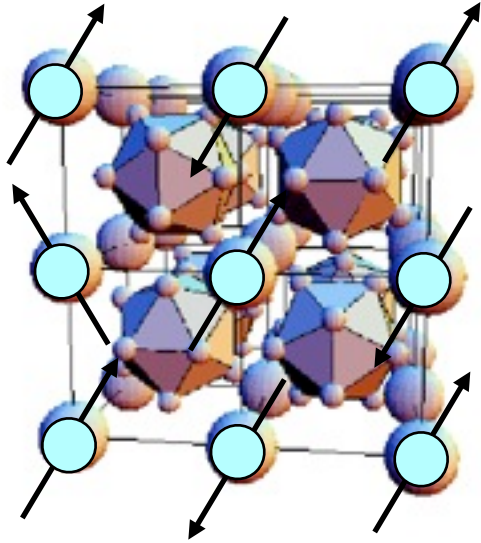


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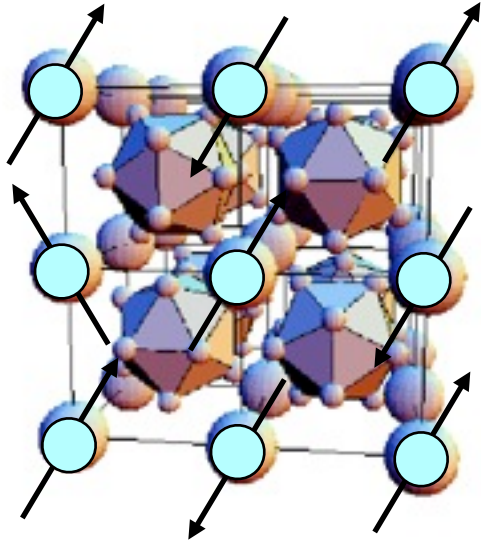
Heavy Fermion Metals

[Review: cond-mat/0612006](#)

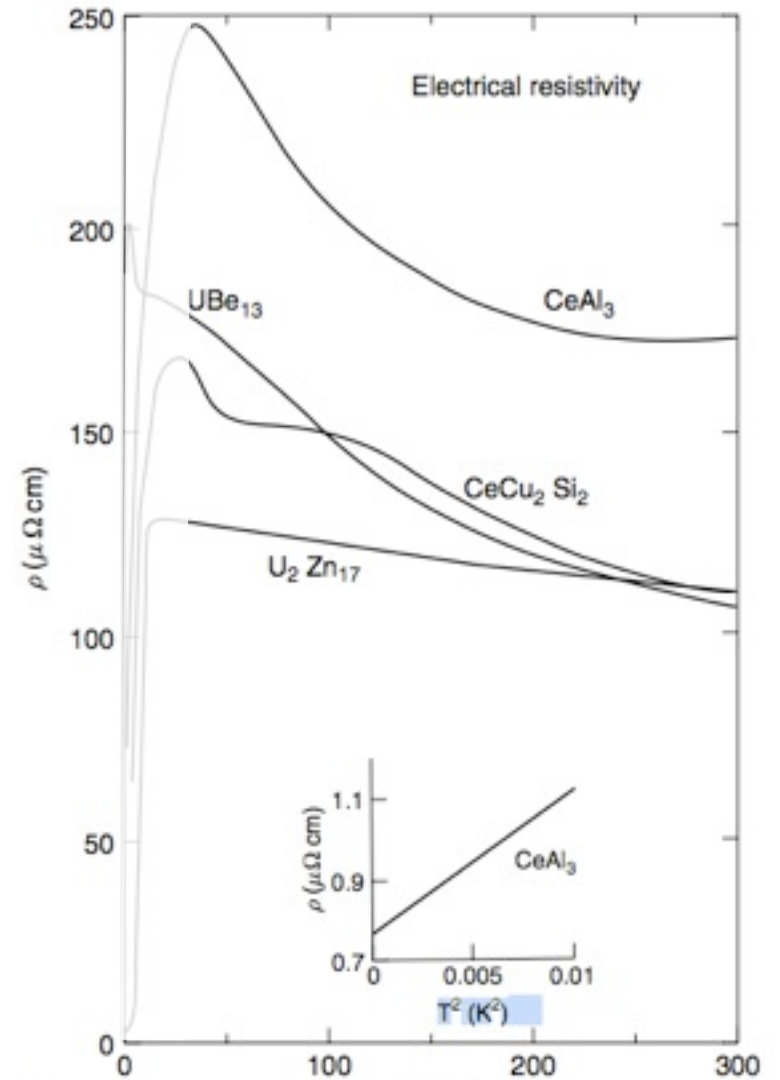


Heavy Fermion Metals

[Review: cond-mat/0612006](https://arxiv.org/abs/cond-mat/0612006)

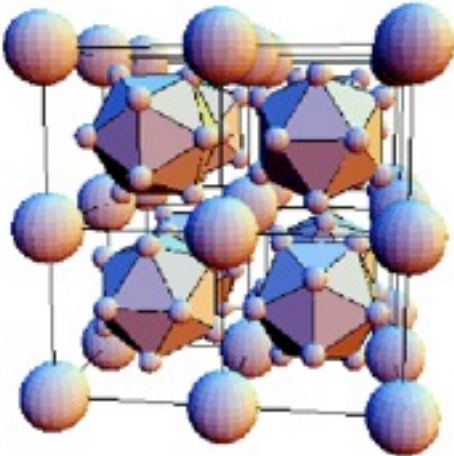


UBe_{13}



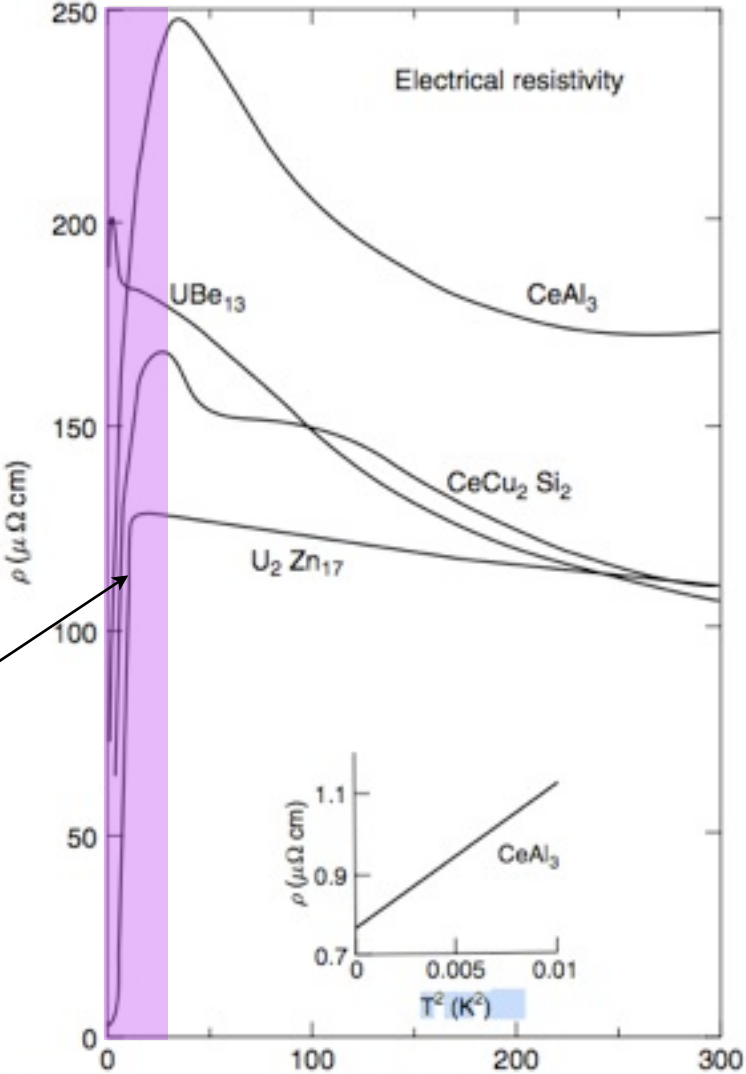
Heavy Fermion Metals

[Review: cond-mat/0612006](http://cond-mat/0612006)



UBe₁₃

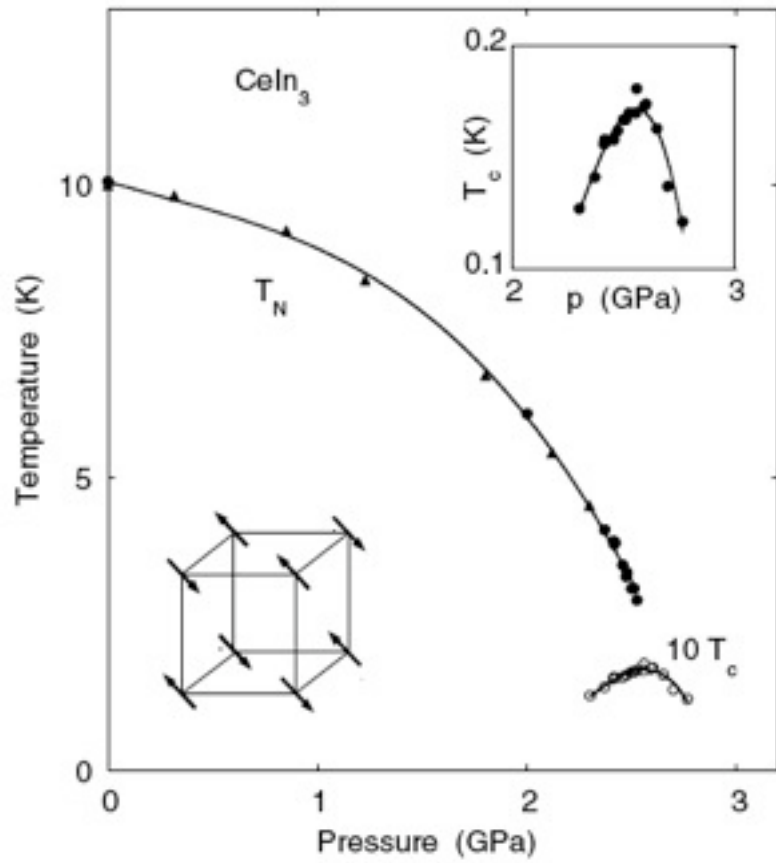
Coherent
Heavy Fermi Liquid



Heavy Electron Quantum Criticality.

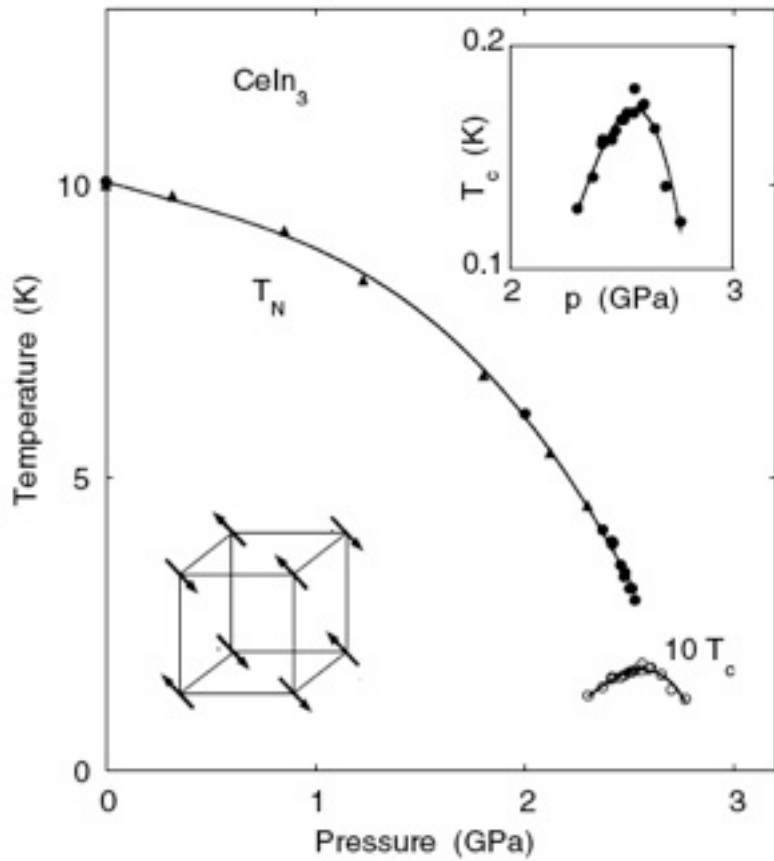
Experiments

“Avoided Criticality”



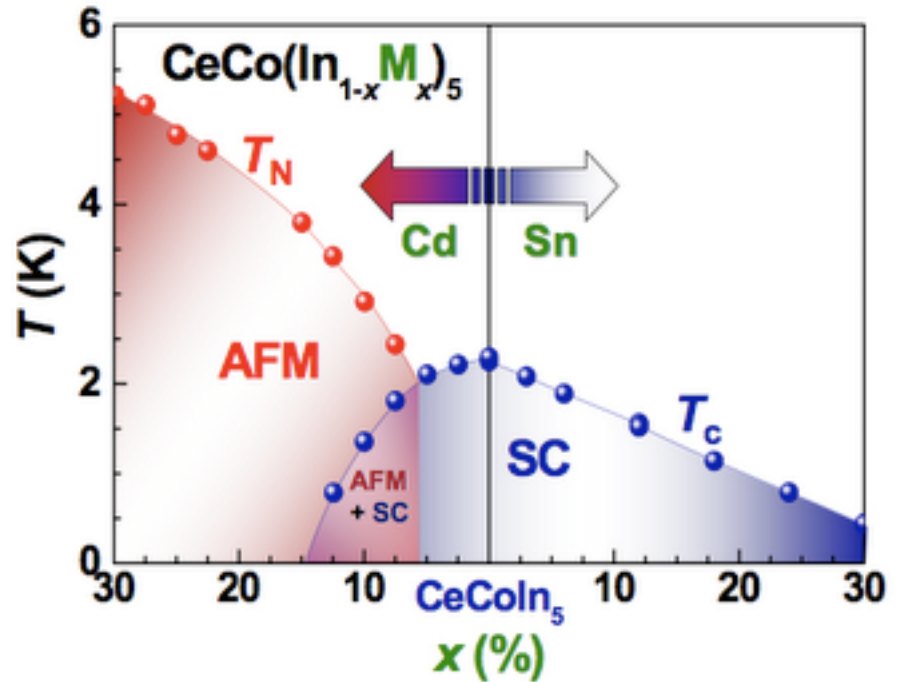
Mathur et al, Nature 394, 39 (1998)

“Avoided Criticality”



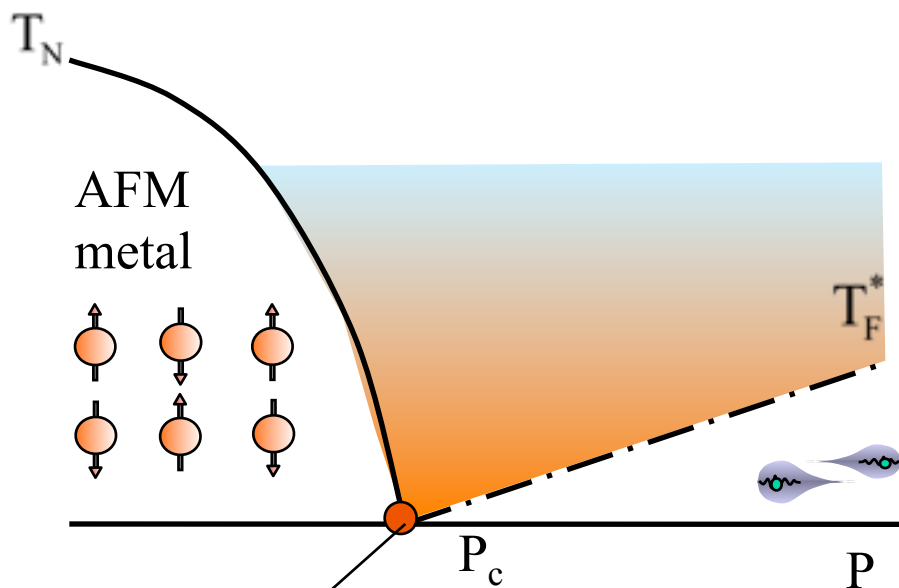
Mathur et al, Nature 394, 39 (1998)

N. Curro, PRL 99, 146402 (2007)



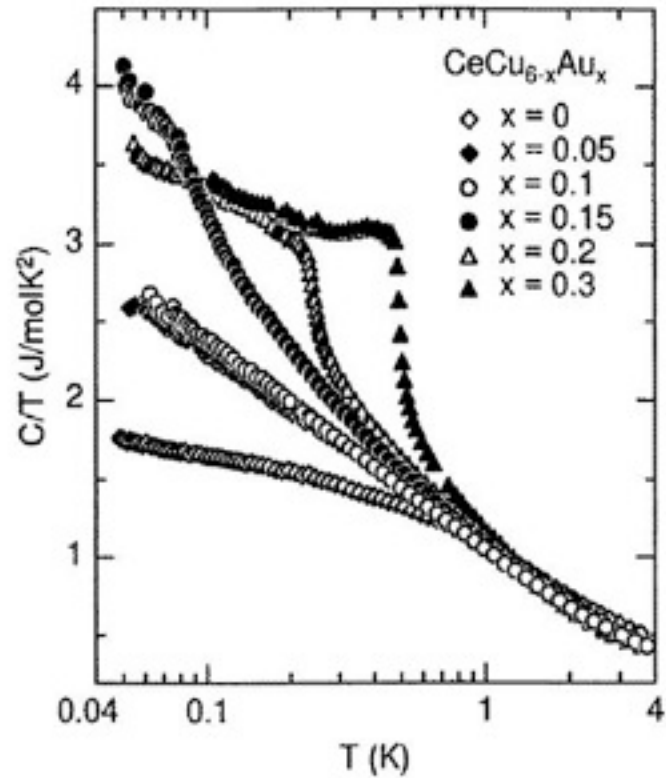
Quantum Criticality: divergent specific heat capacity

Heavy Fermion
Materials



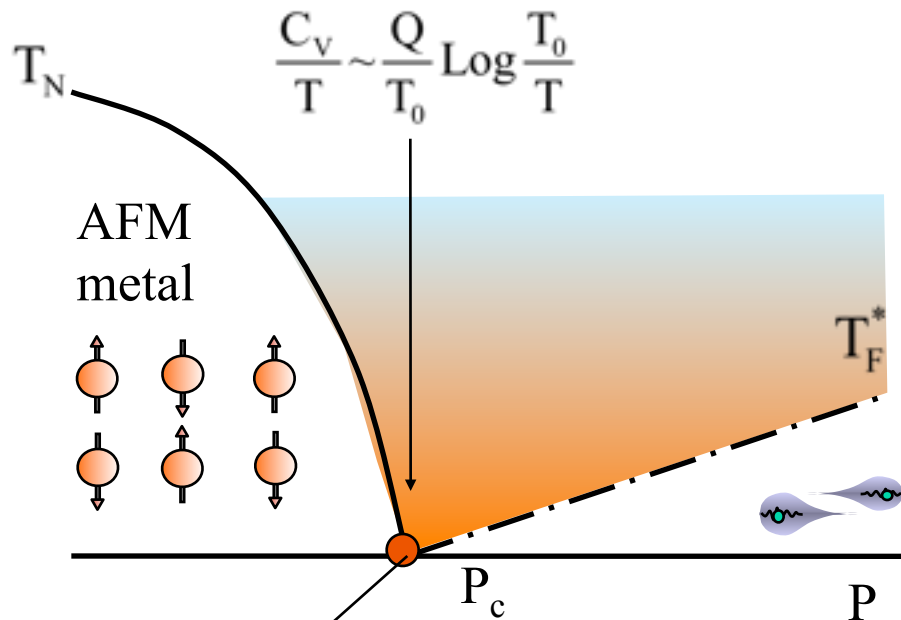
Quantum Critical
Point

H. Von Lohneyson (1996)



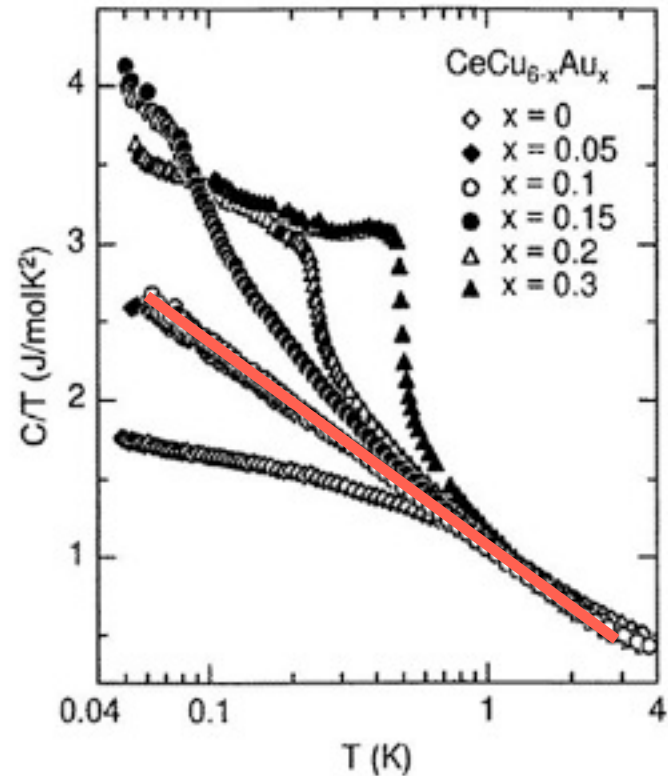
Quantum Criticality: divergent specific heat capacity

Heavy Fermion
Materials



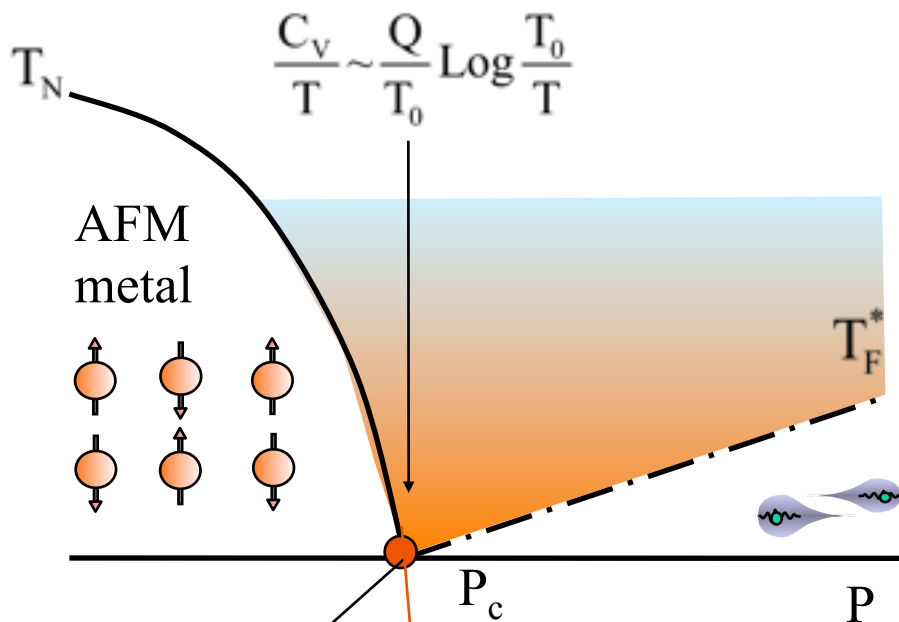
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Quantum Criticality: divergent specific heat capacity

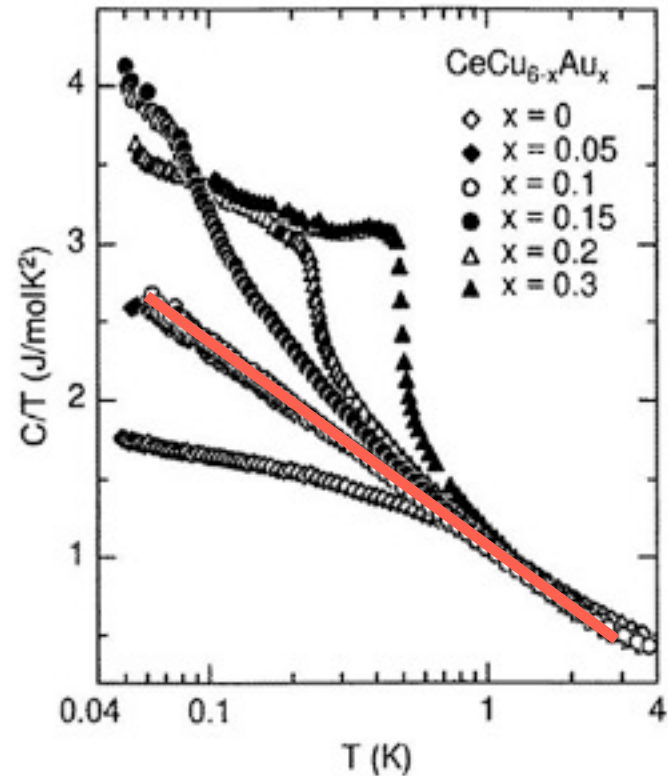
Heavy Fermion
Materials



Quantum Critical
Point

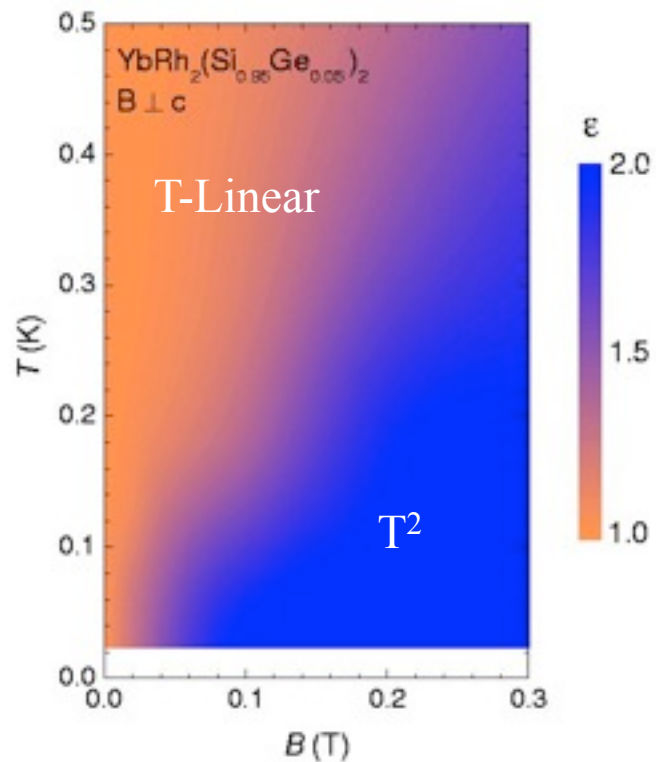
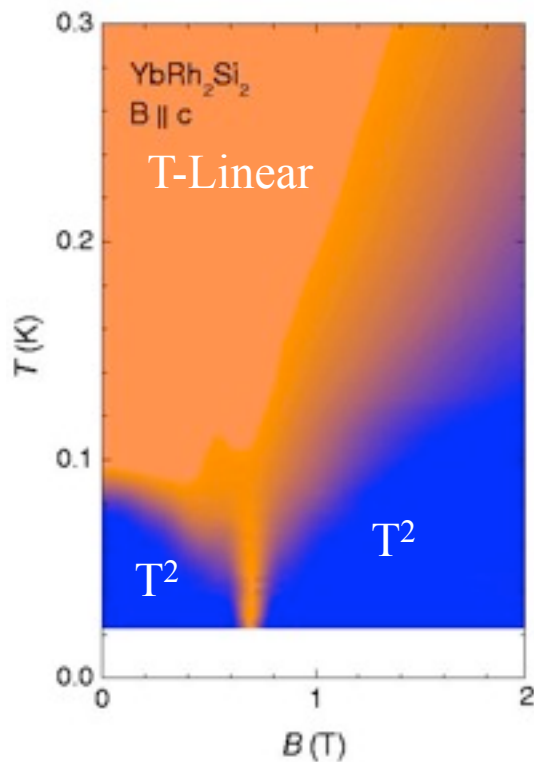
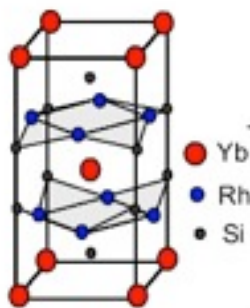
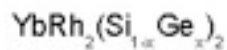
$$\frac{m^*}{m} \rightarrow \infty$$

H. Von Lohneyson (1996)

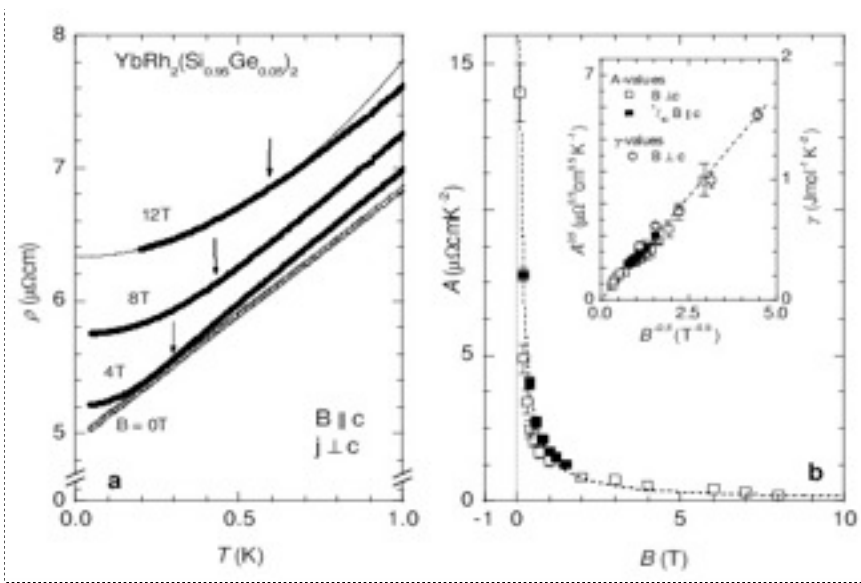


Divergence of Interaction and Effective Mass

Gegenwart et al (2002)



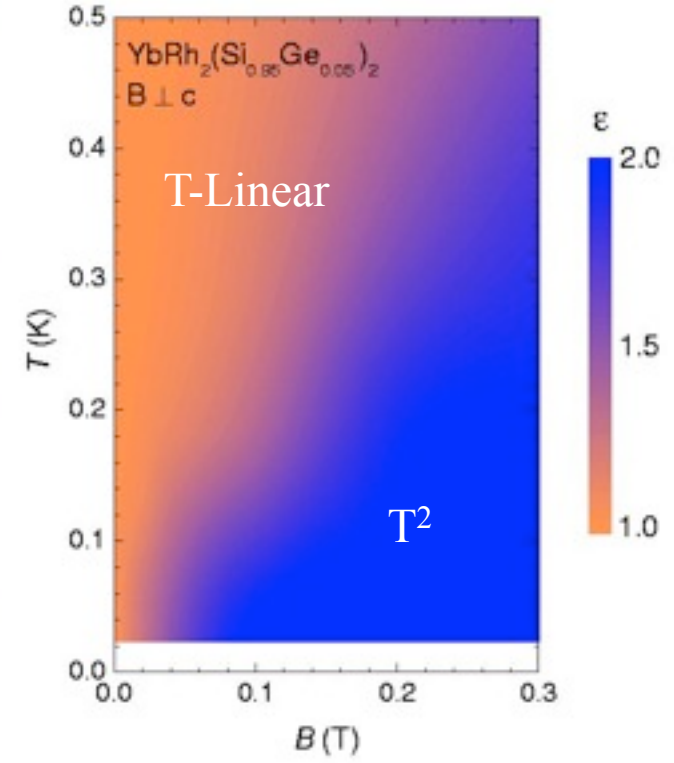
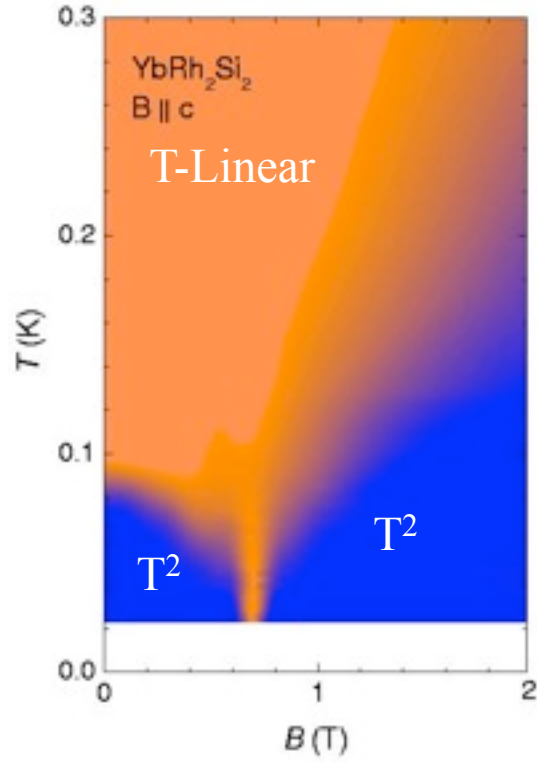
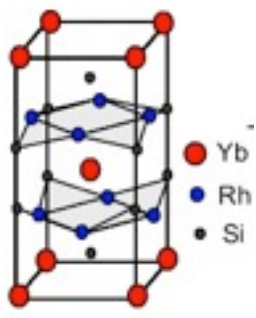
Divergence of Interaction and Effective Mass



$$\rho = AT^2 + \rho_0$$

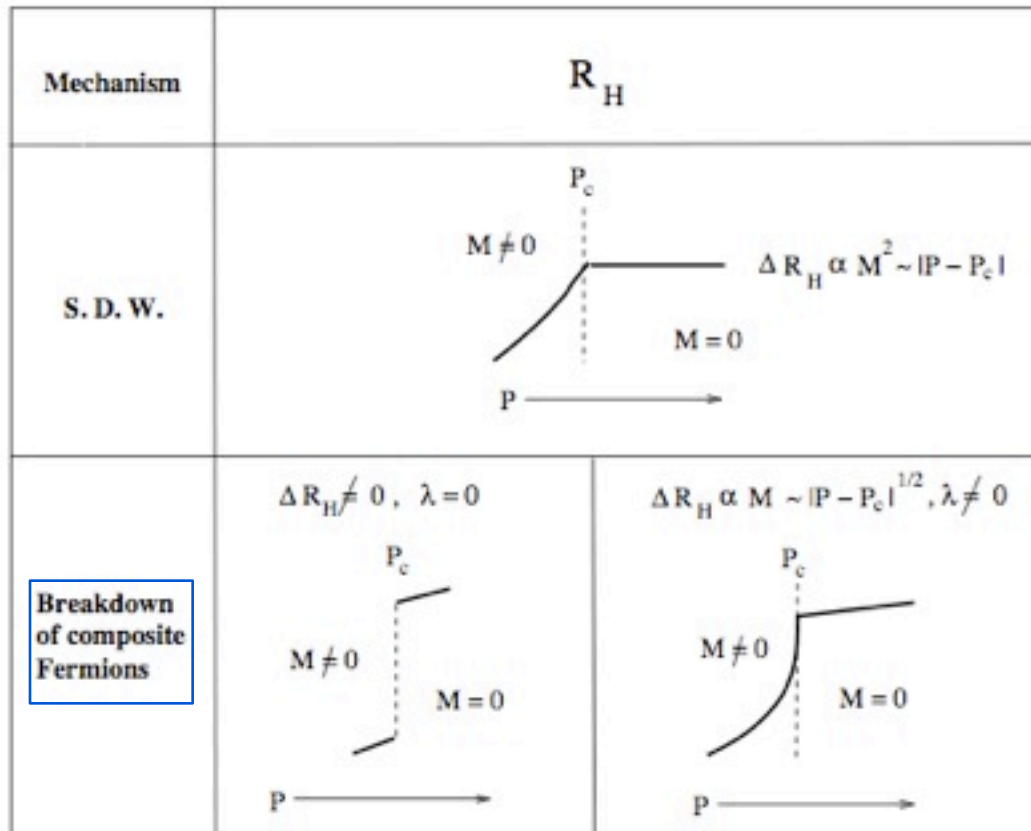
$$A \propto \frac{1}{T_F^2} \propto \frac{1}{B - B_c}$$

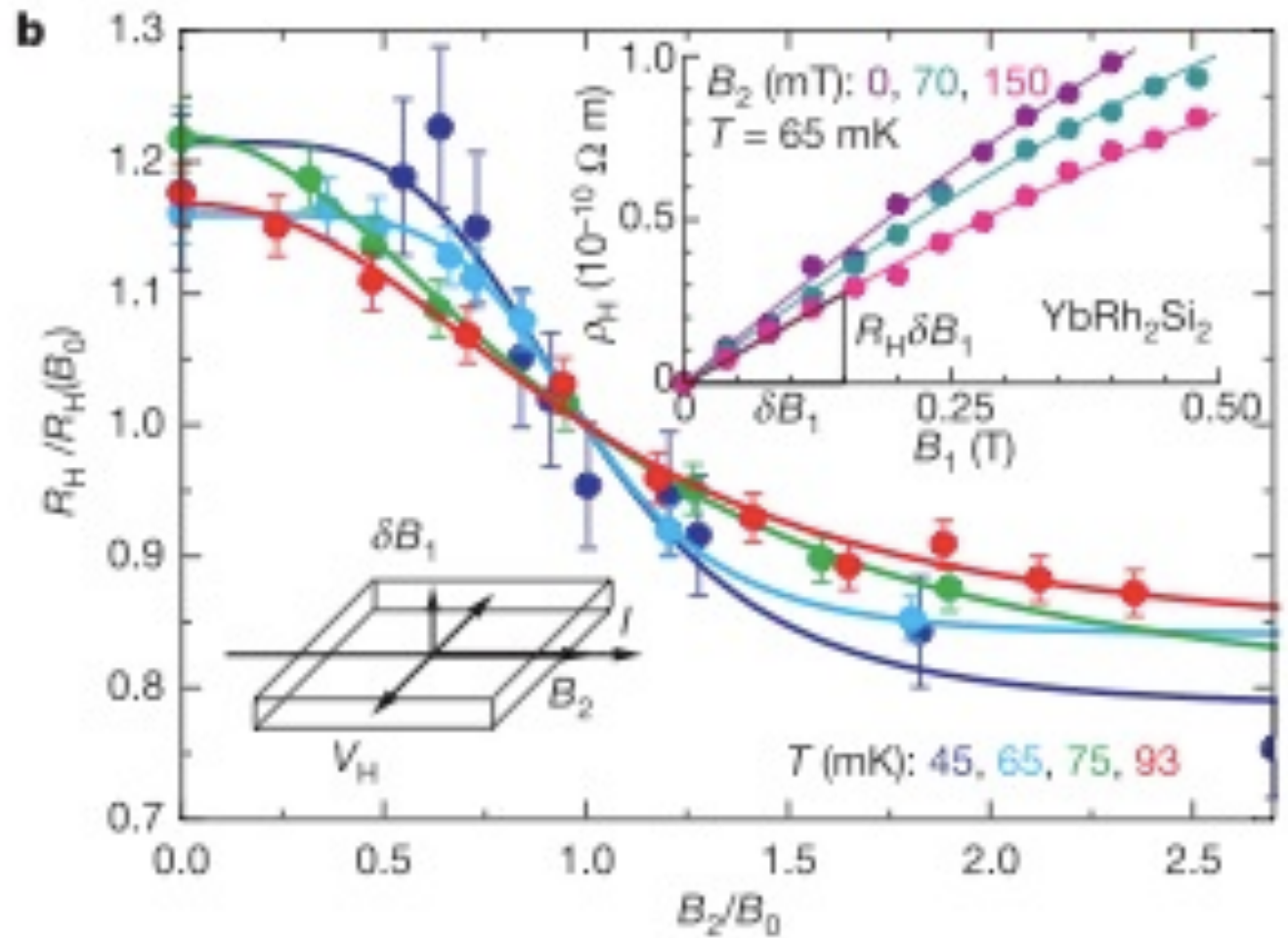
Gegenwart et al (2002)

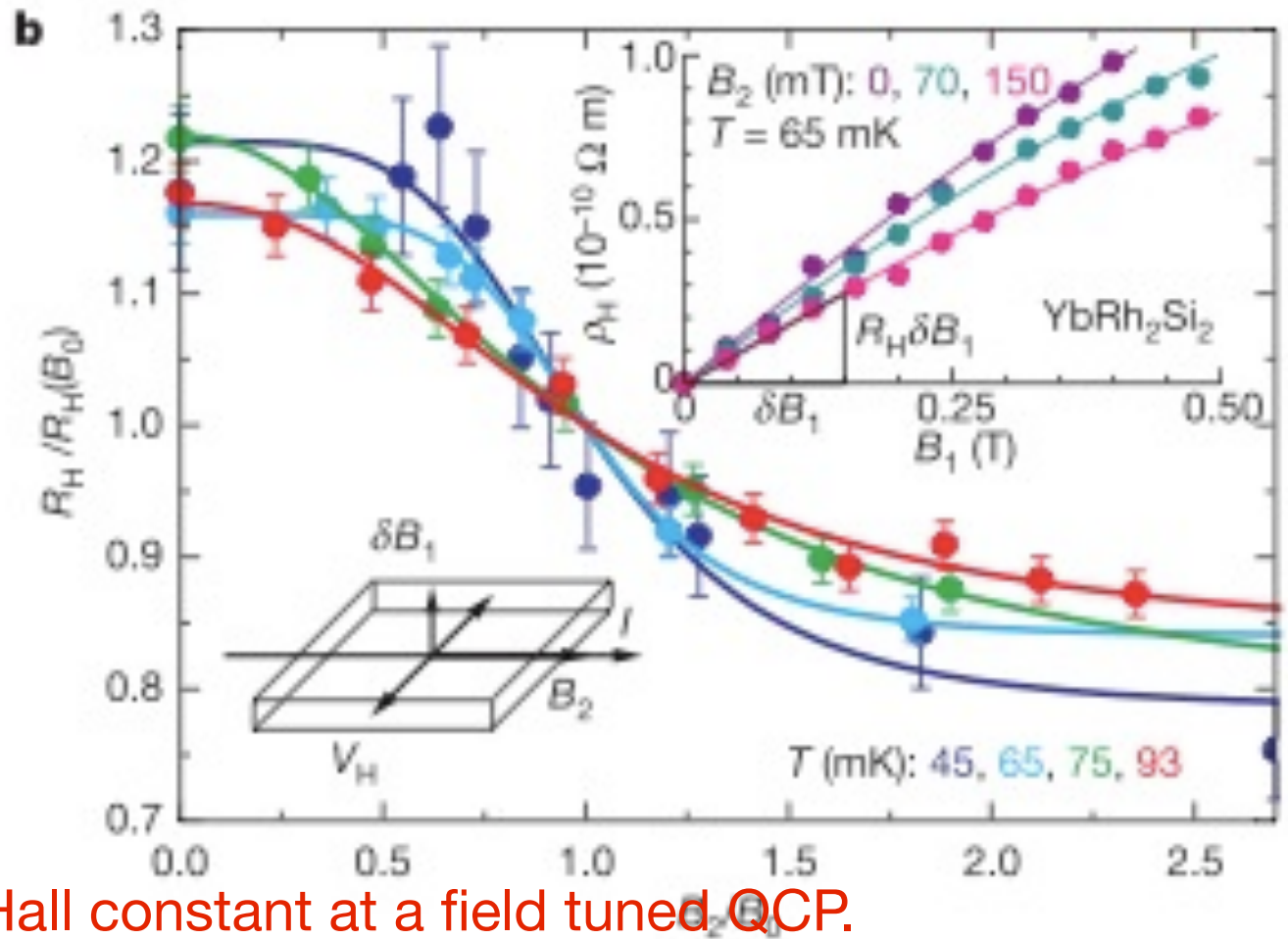


“How do fermions get heavy and die?” PC, Pepin, Si and Ramazashvili, J. Cond Matt. ,13}, R723 (2001).

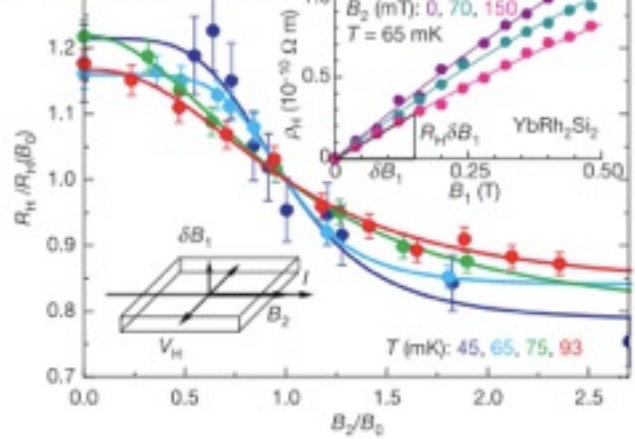
anticipated an abrupt change in FS when a composite heavy electron undergoes a Kondo “breakdown”.



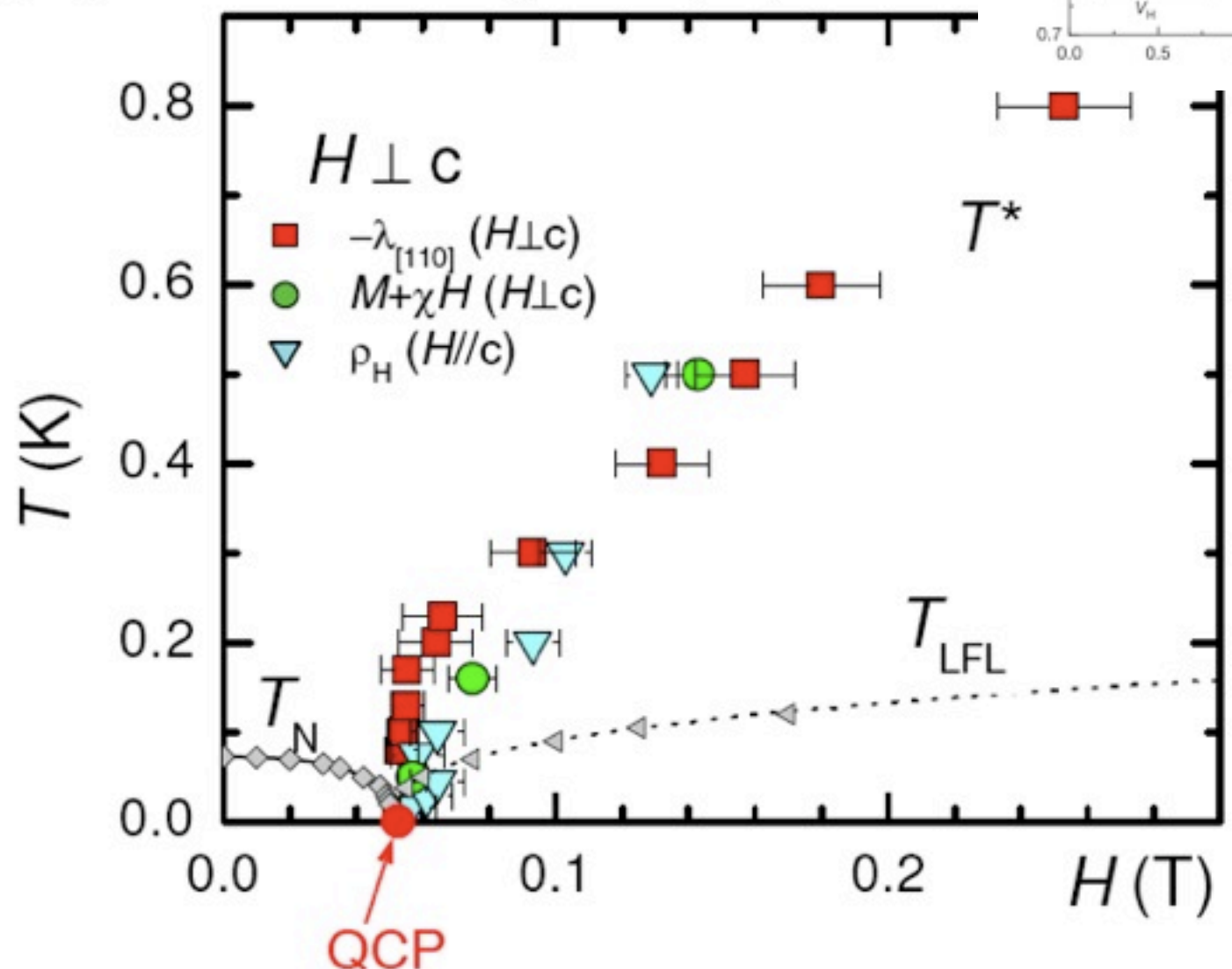
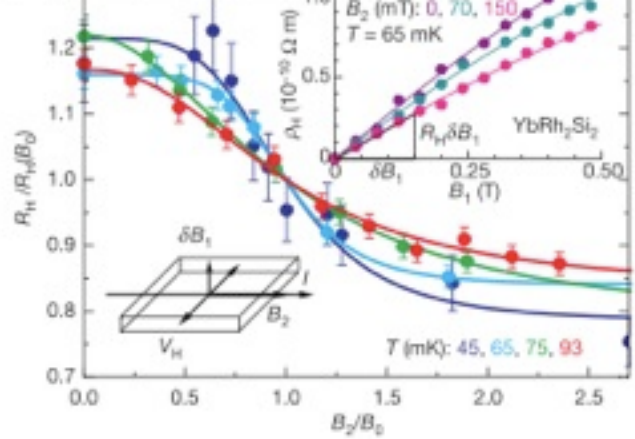




Jump in the Hall constant at a field tuned QCP.



T_{Hall} represents a new energy scale (T^*)

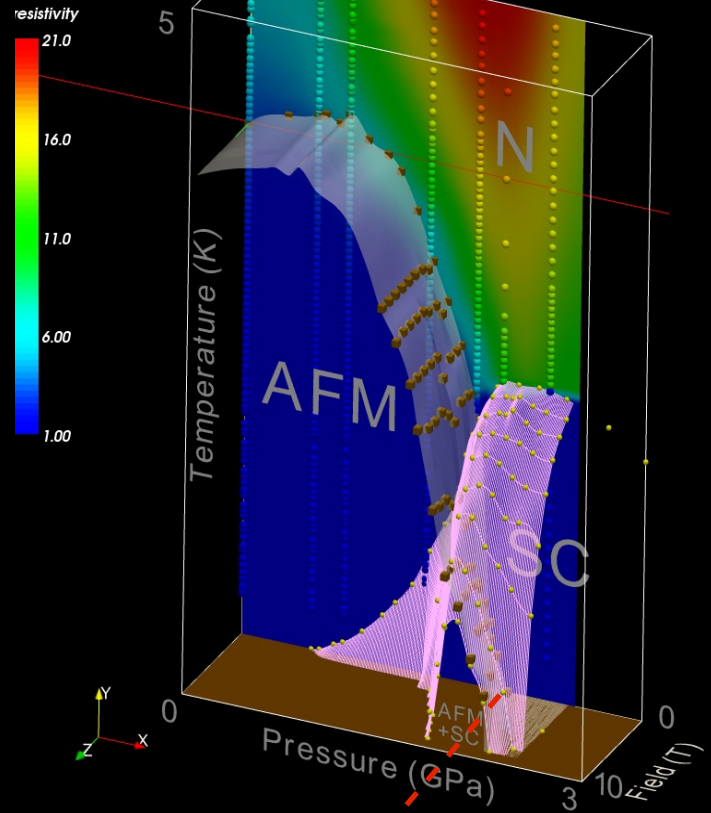


Reconstruction of the Fermi Surface and mass divergence

CeRhIn₅

Tuson Park, (2007).

Phase Diagram
Ce-115



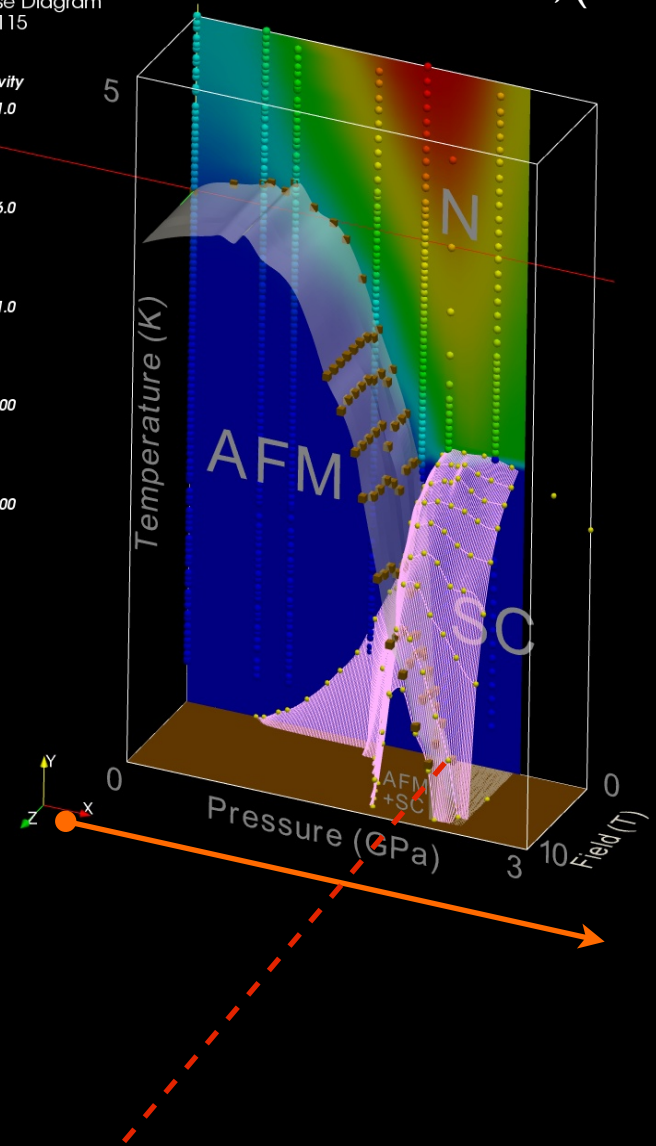
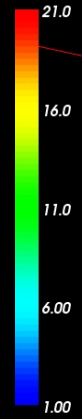
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CeRhIn₅

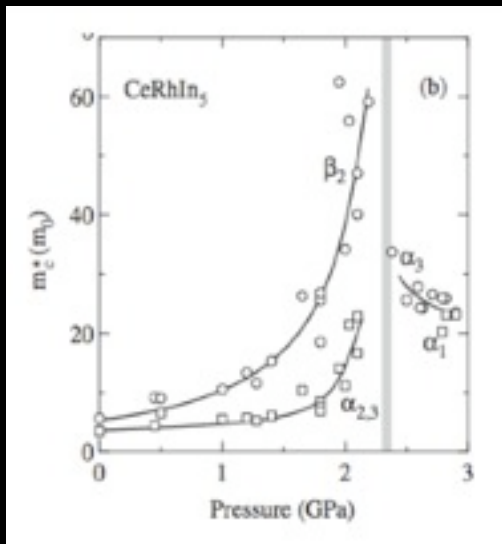
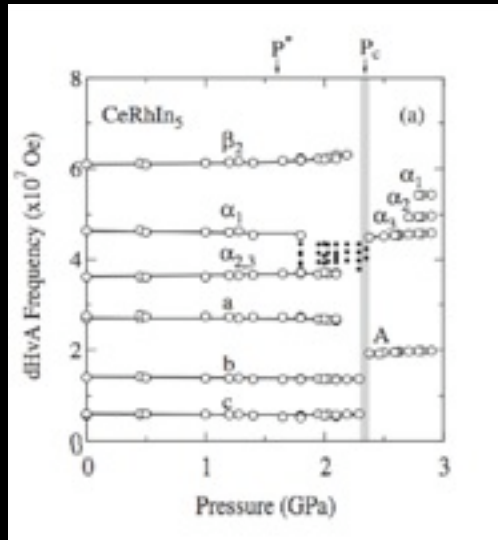
Tuson Park, (2007).

Phase Diagram
Ce-115

resistivity



Reconstruction of the Fermi Surface and mass divergence

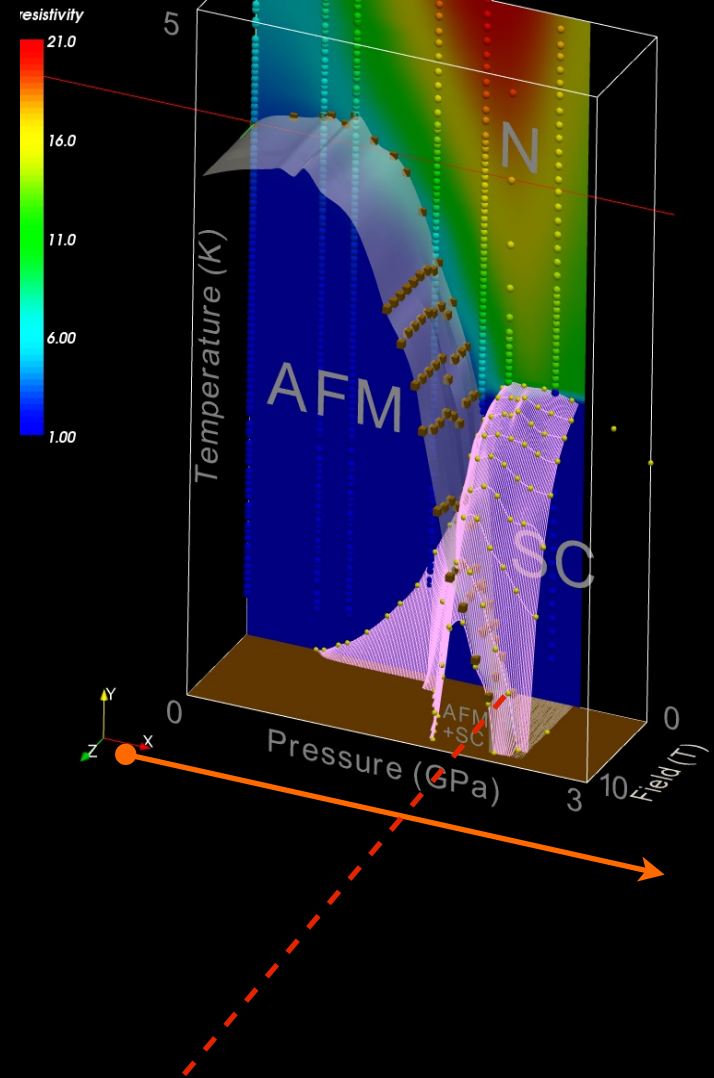


Shimuzu et al (2006)

CeRhIn₅

Phase Diagram
Ce-115

Tuson Park, (2007).

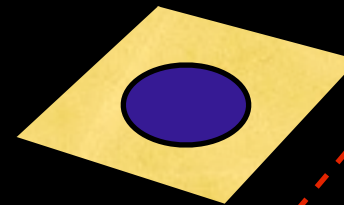
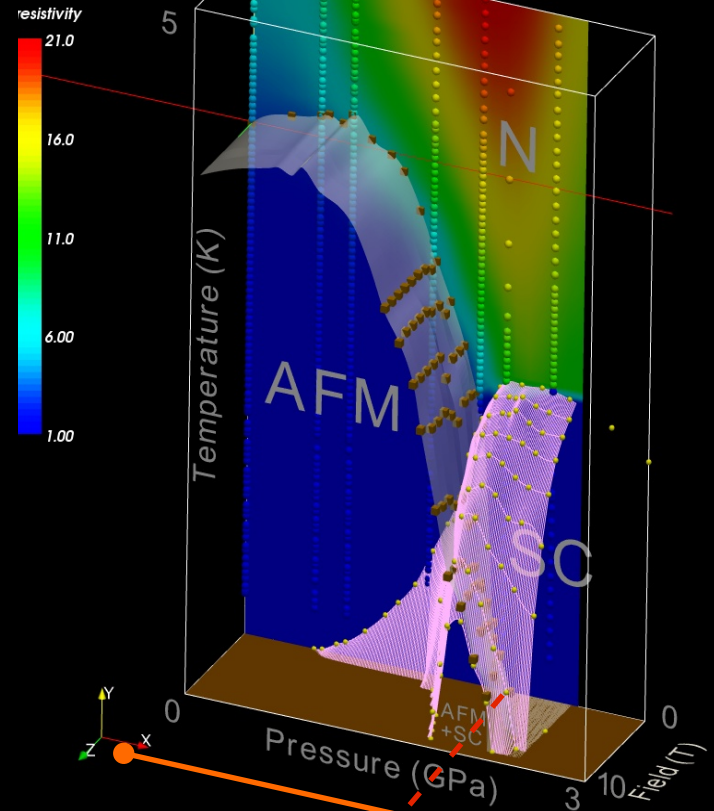
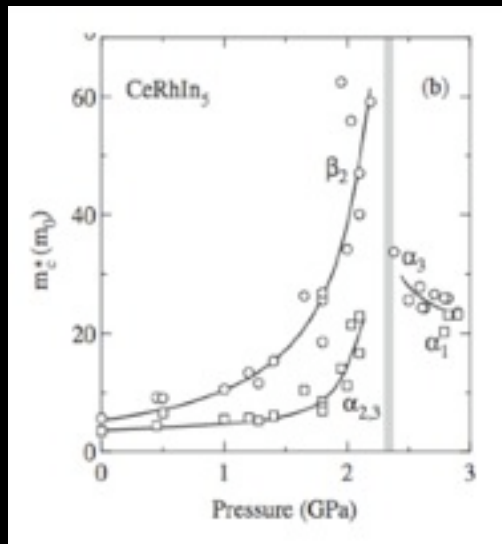
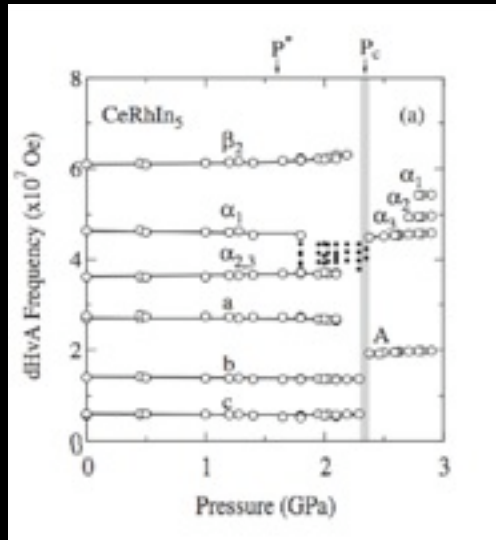


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Tuson Park, (2007).

Phase Diagram
Ce-115



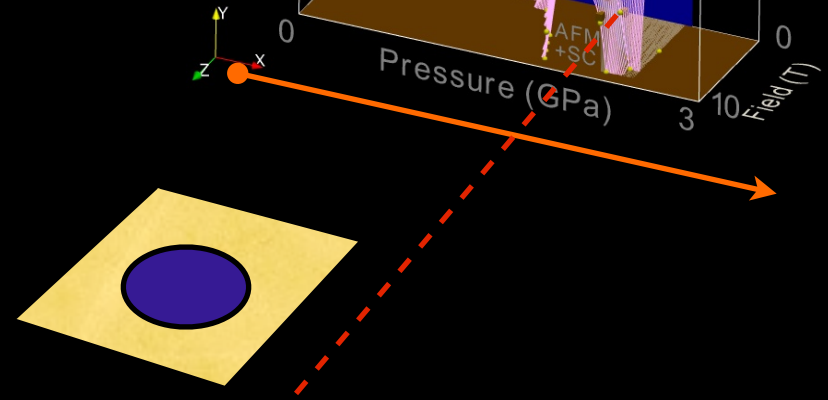
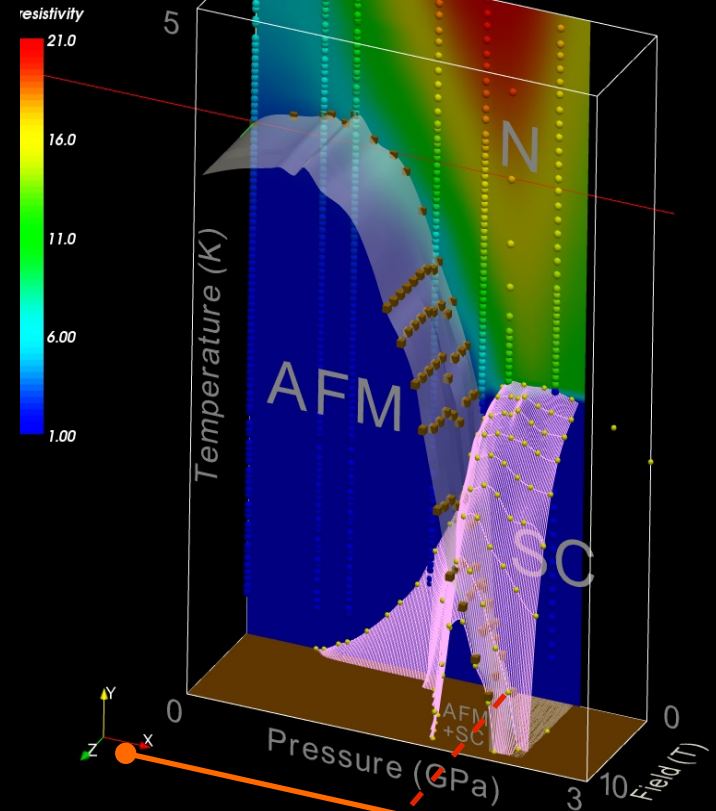
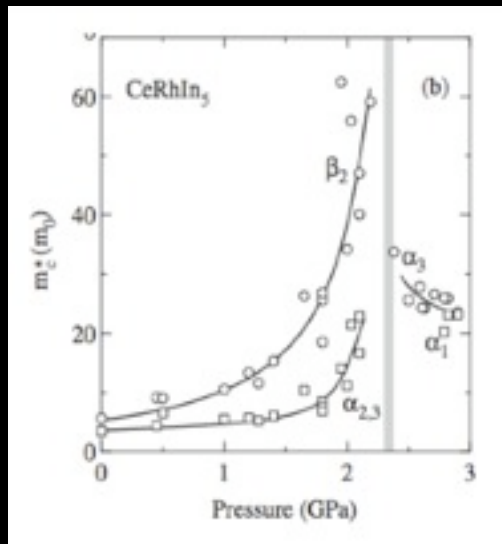
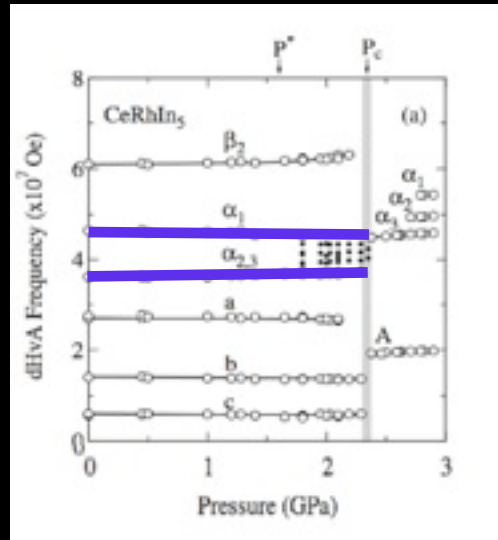
Shimuzu et al (2006)

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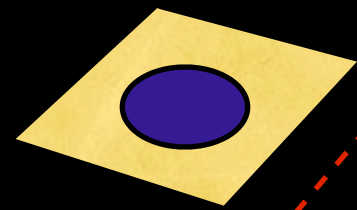
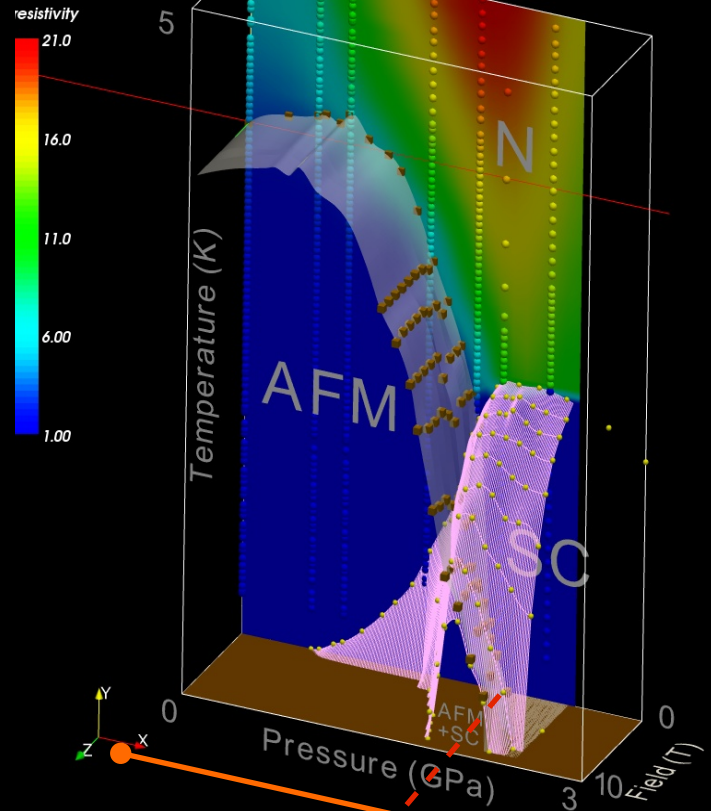
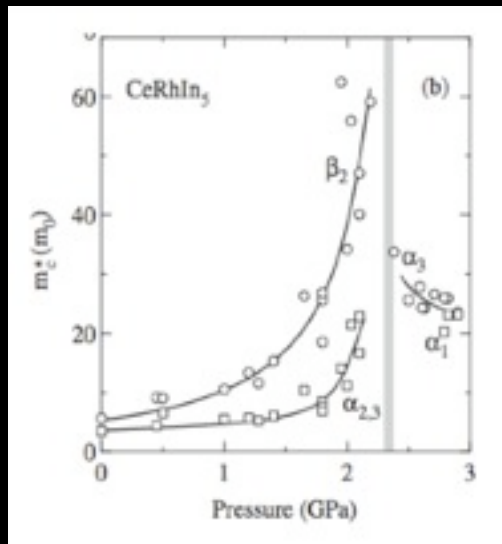
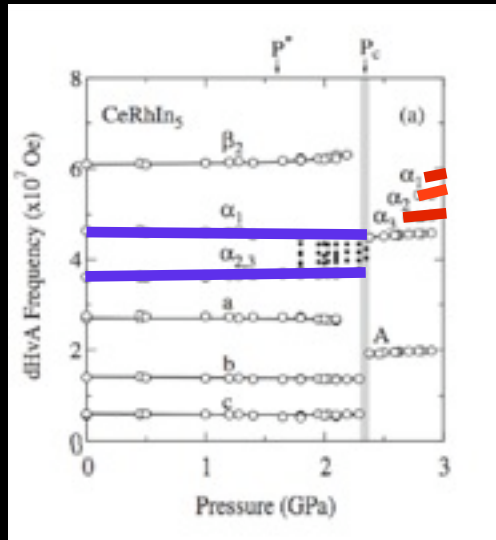
Shimuzu et al (2006)

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Phase Diagram
Ce-115



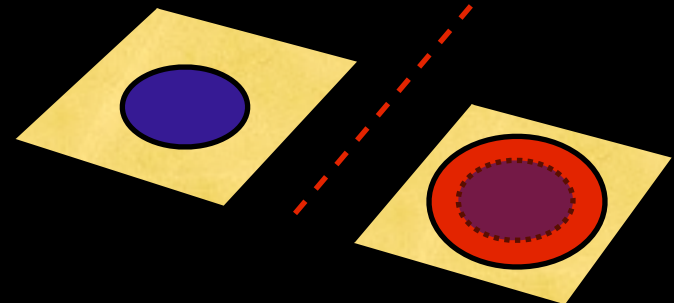
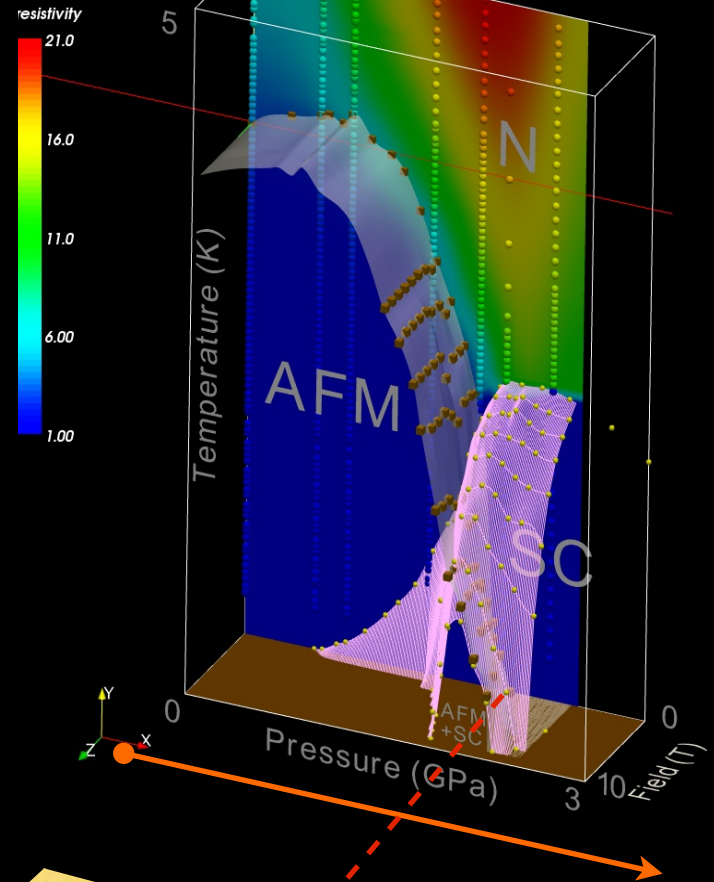
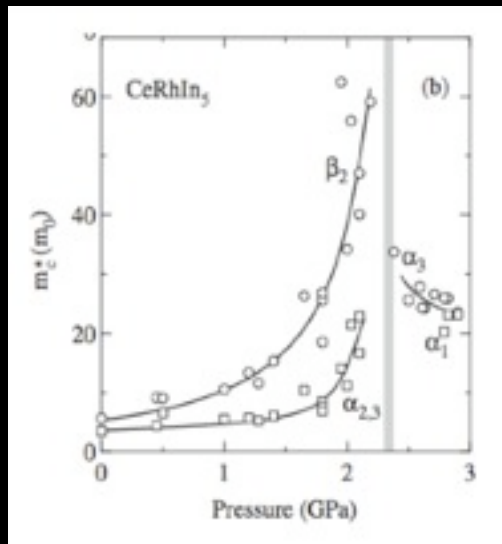
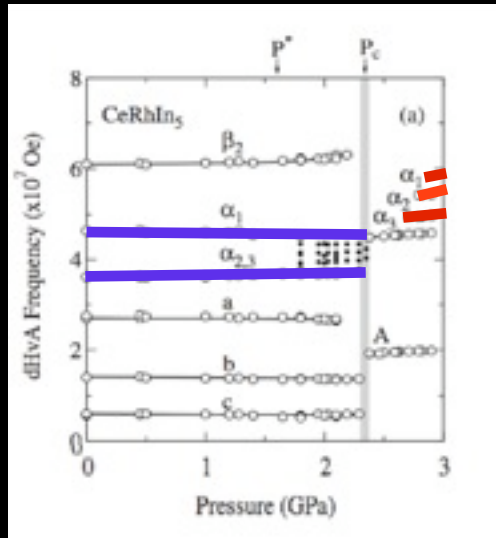
Shimuzu et al (2006)

Reconstruction of the Fermi Surface and mass divergence

CeRhIn₅

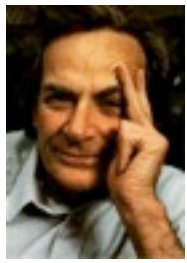
Tuson Park, (2007).

Phase Diagram
Ce-115



Shimuzu et al (2006)

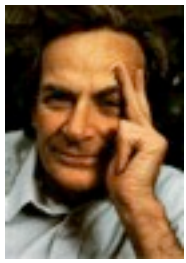
Black Hole in the Phase Diagram.



Feynman

Hertz

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$



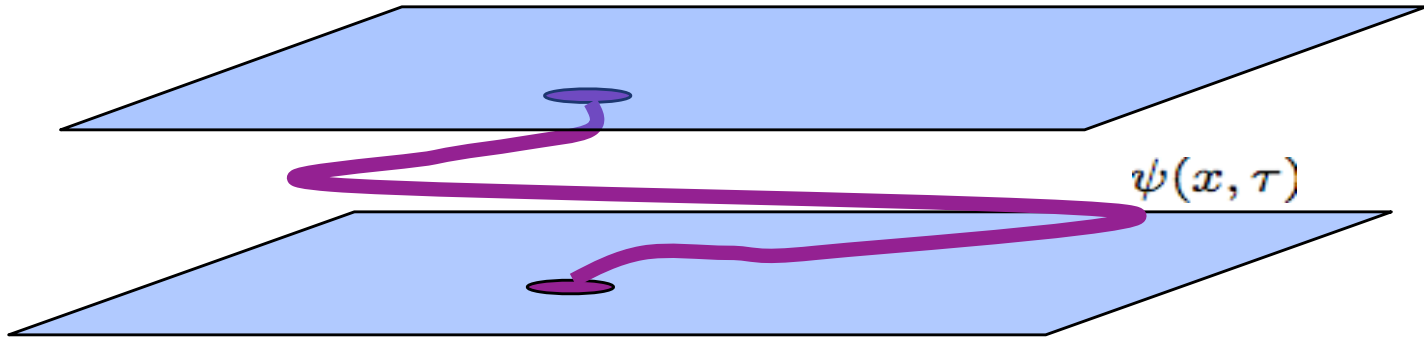
Feynman

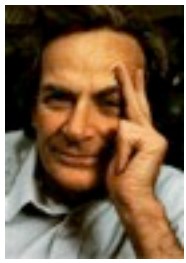


Hertz

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$

$$\frac{\hbar}{k_B T}$$





Feynman



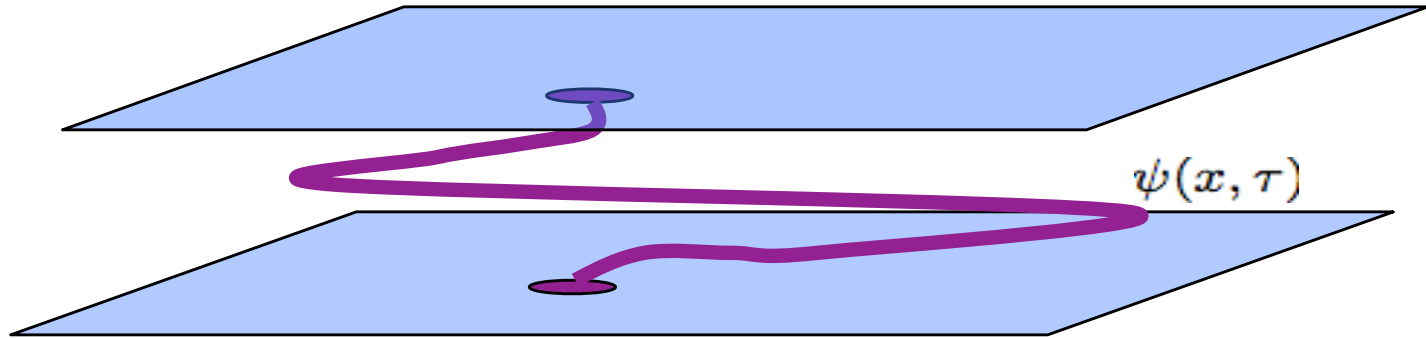
Hertz

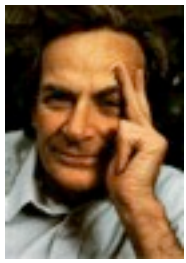
Temperature:
Boundary condition
in time.

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$

Sachdev (1999), Continentino (2001), Palova (2009).

$$\frac{\hbar}{k_B T}$$





Feynman



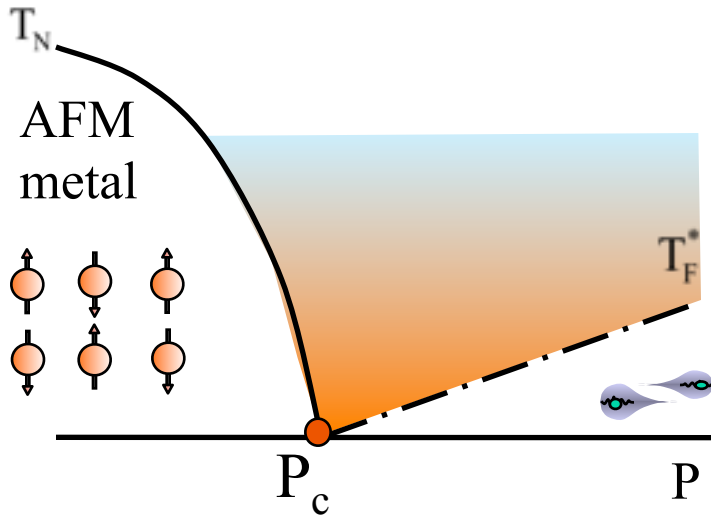
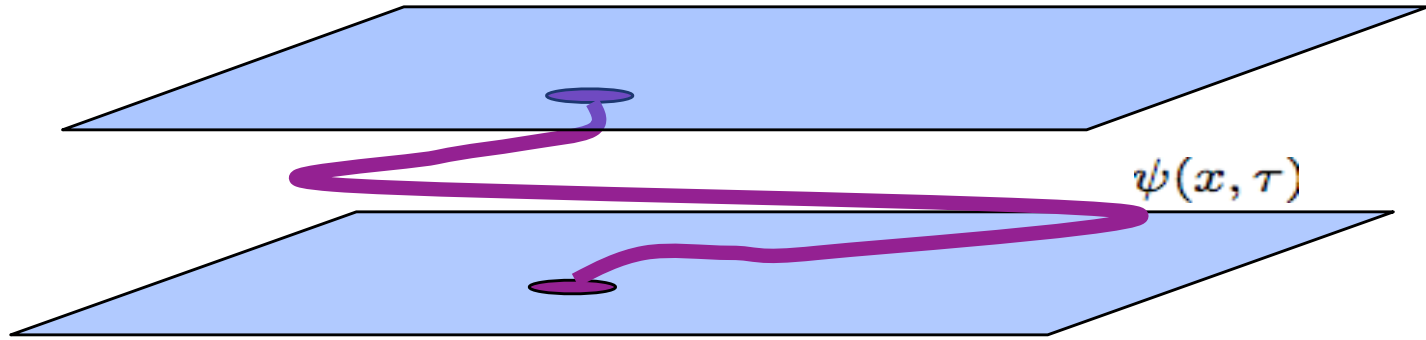
Hertz

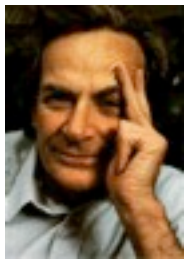
Temperature:
Boundary condition
in time.

$$Z = \sum_{\text{Histories}} \exp \left[- \int_0^{1/T} L[\psi(x, \tau)] d\tau \right]$$

Sachdev (1999), Continentino (2001), Palova (2009).

$$\frac{\hbar}{k_B T}$$





Feynman

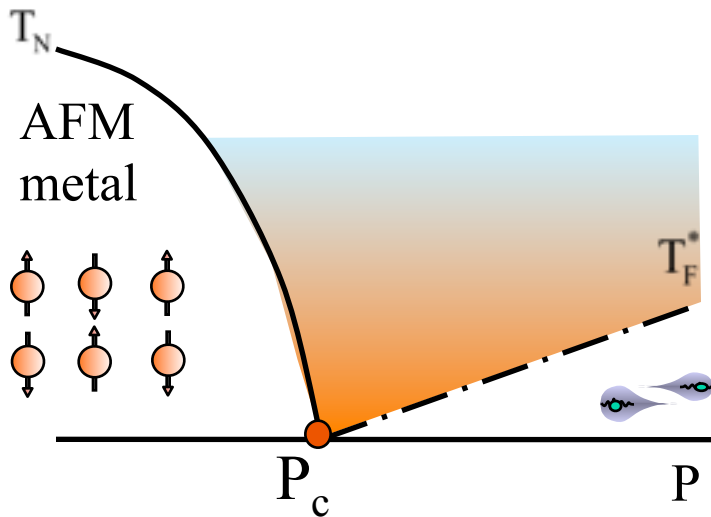
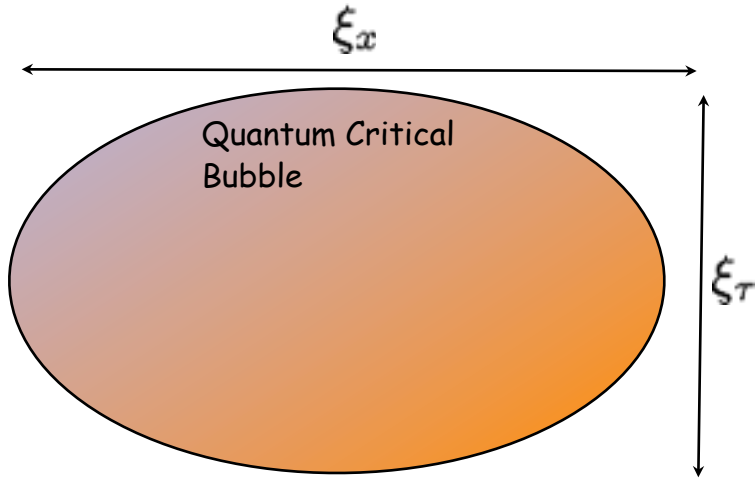


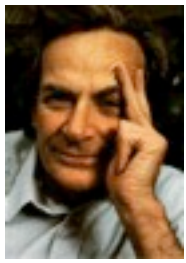
Hertz

Temperature: Boundary condition in time.

Sachdev (1999), Continentino (2001), Palova (2009).

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Feynman

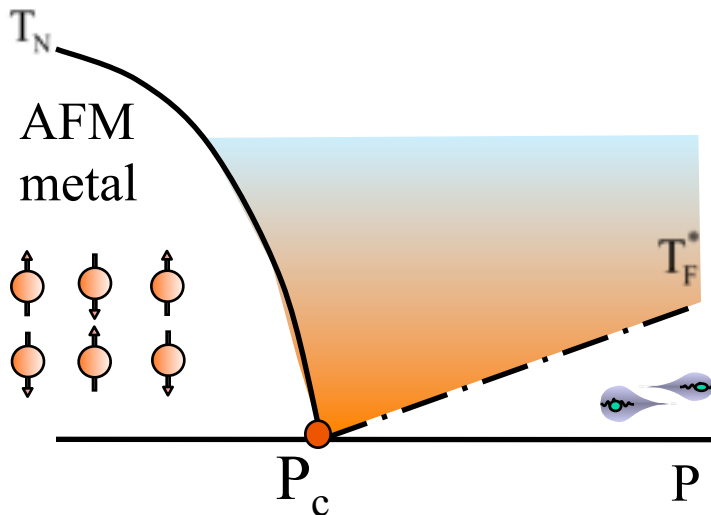
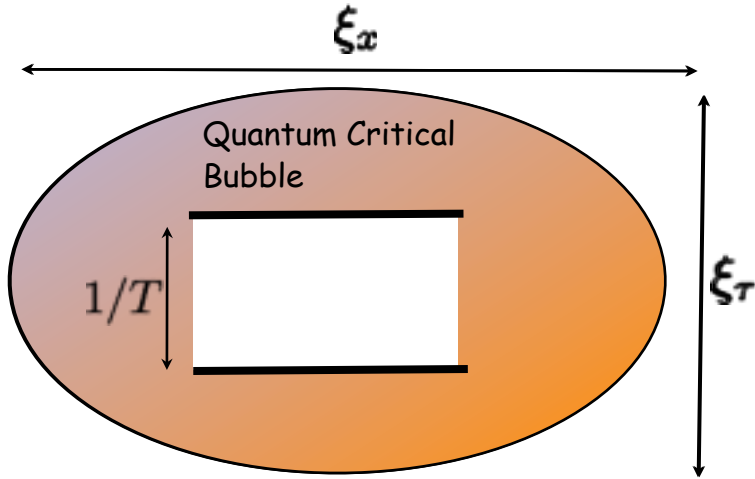


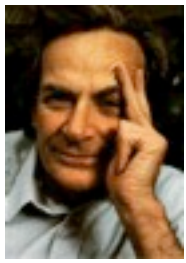
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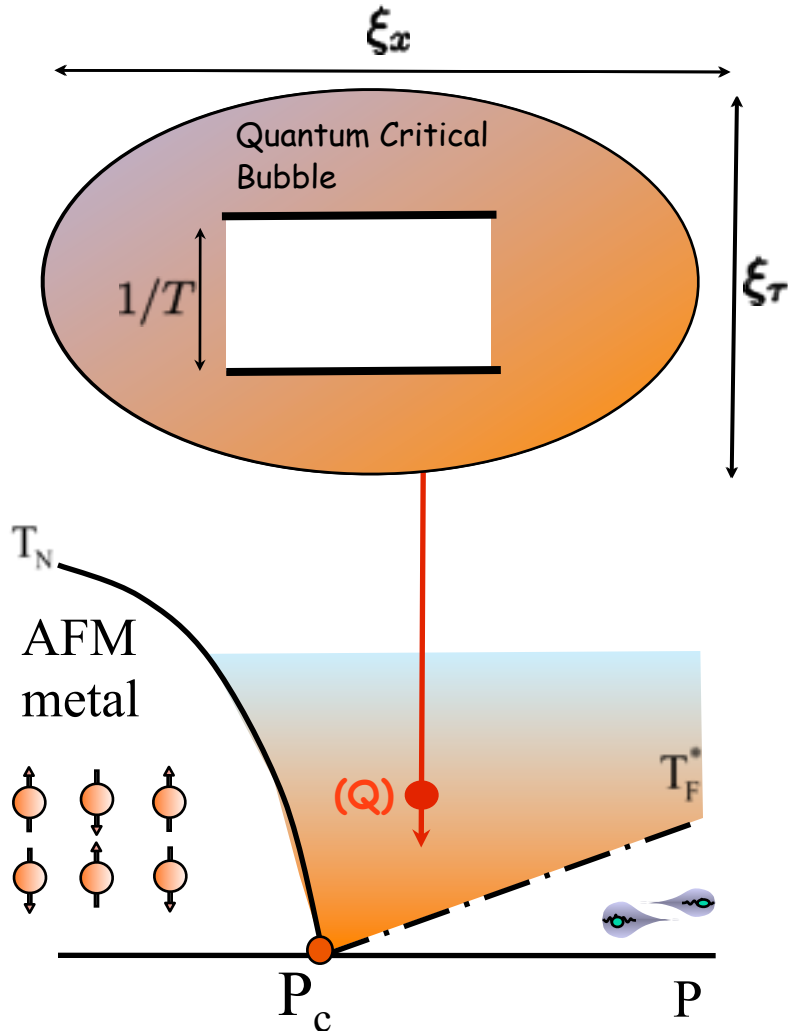
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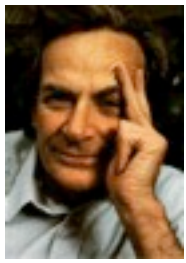
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(Q) Quantum critical region:
interior of correlation bubble.





Feynman

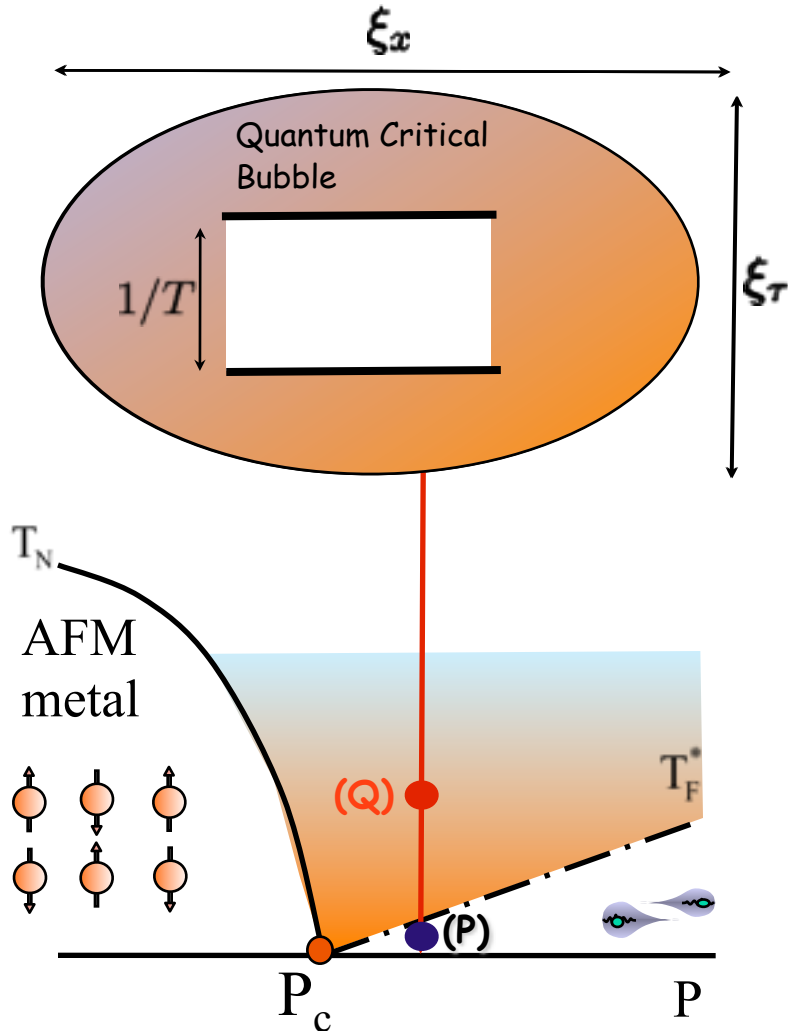


Hertz

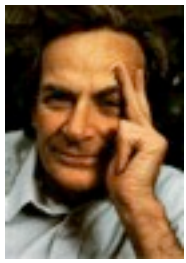
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- (Q)** Quantum critical region:
interior of correlation bubble.
- (P)** Paramagnet: probes
exterior of correlation bubble



Feynman



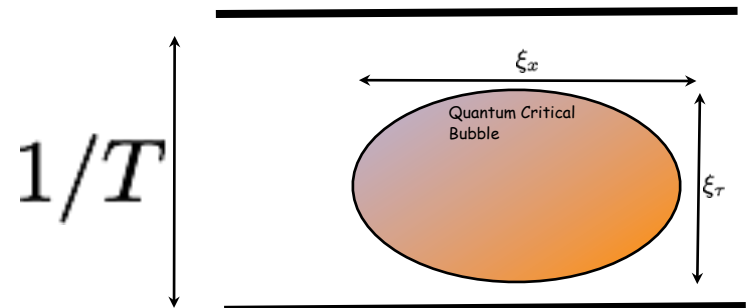
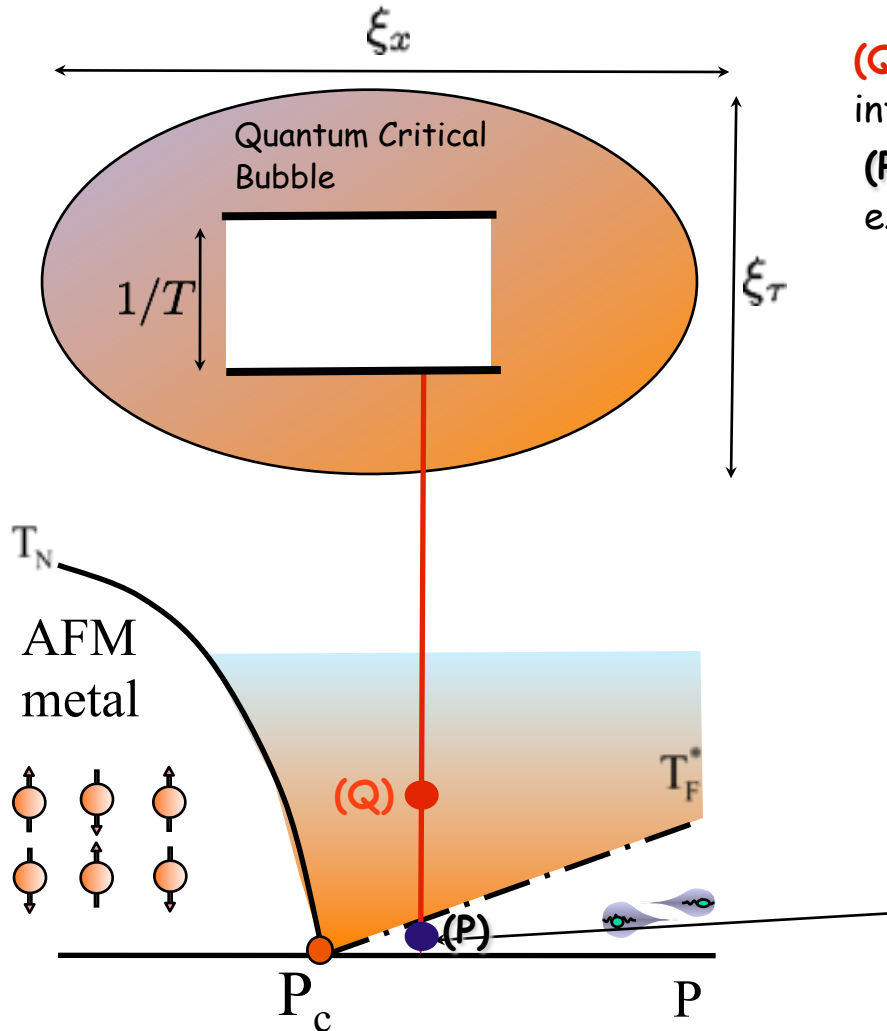
Hertz

Temperature: Boundary condition in time.

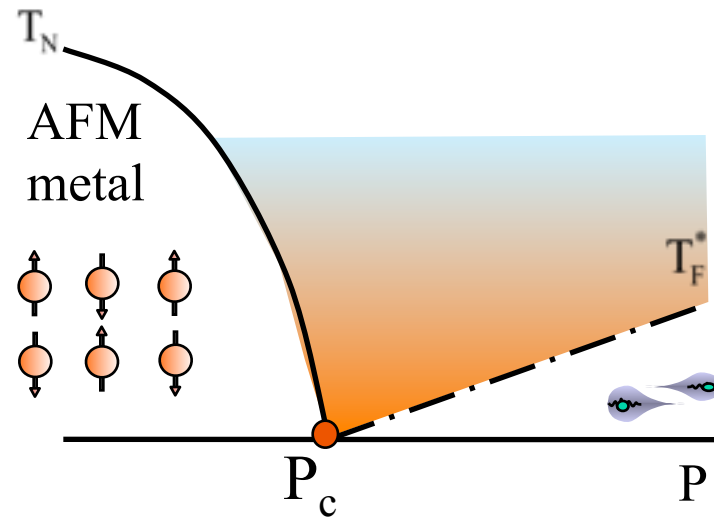
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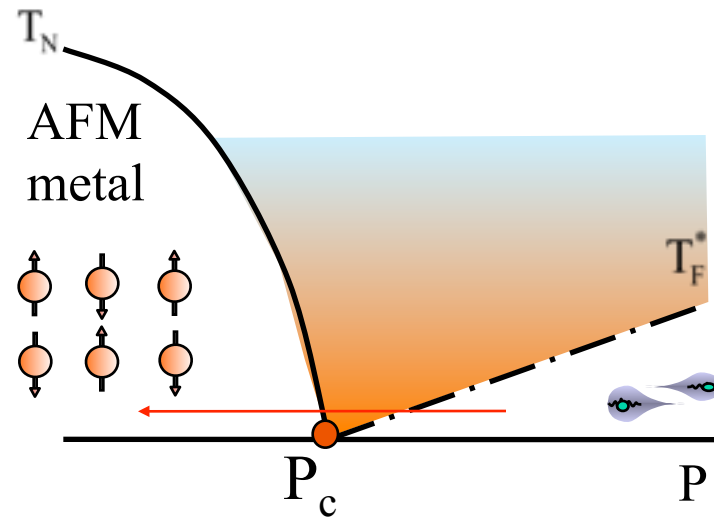
- (Q)** Quantum critical region: interior of correlation bubble.
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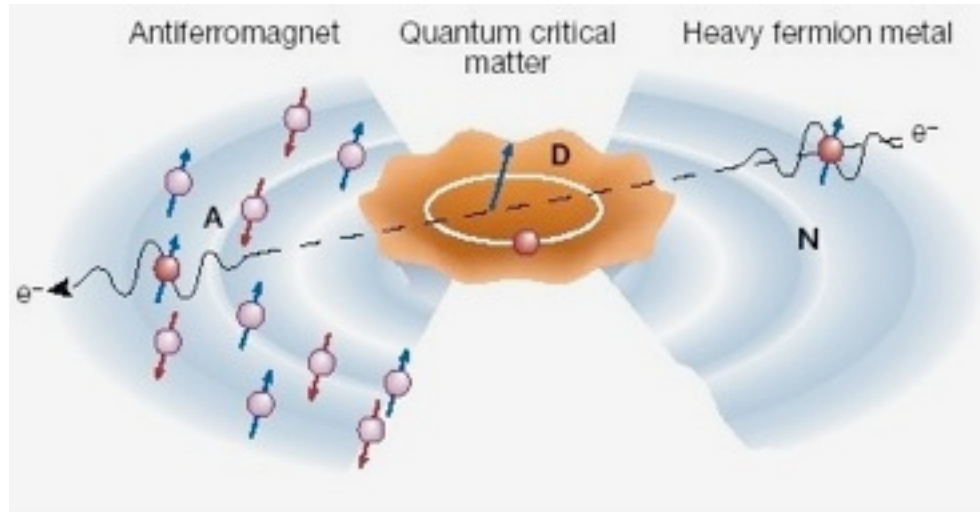
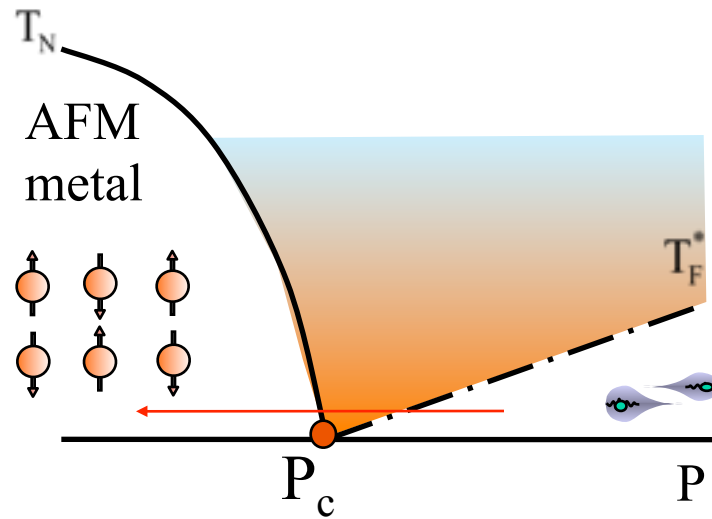
"Black Hole in the Phase Diagram".



"Black Hole in the Phase Diagram".

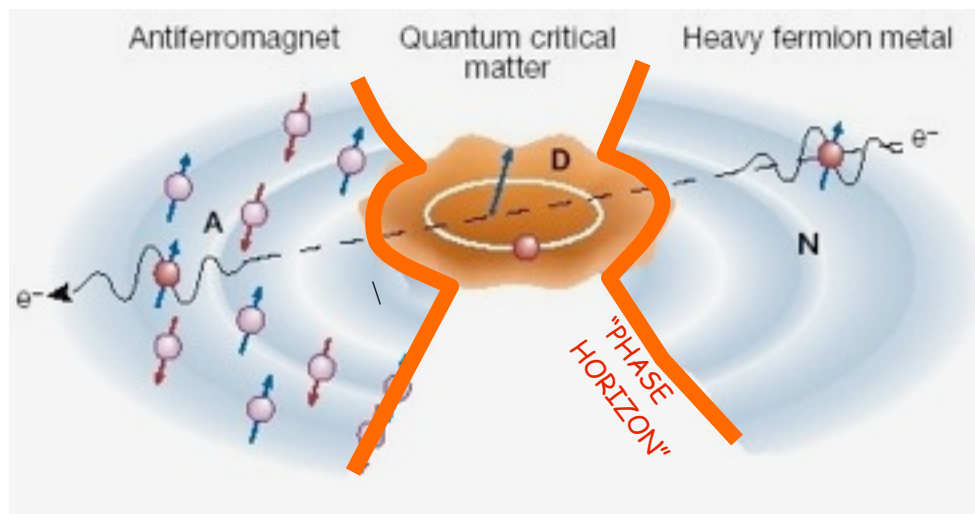
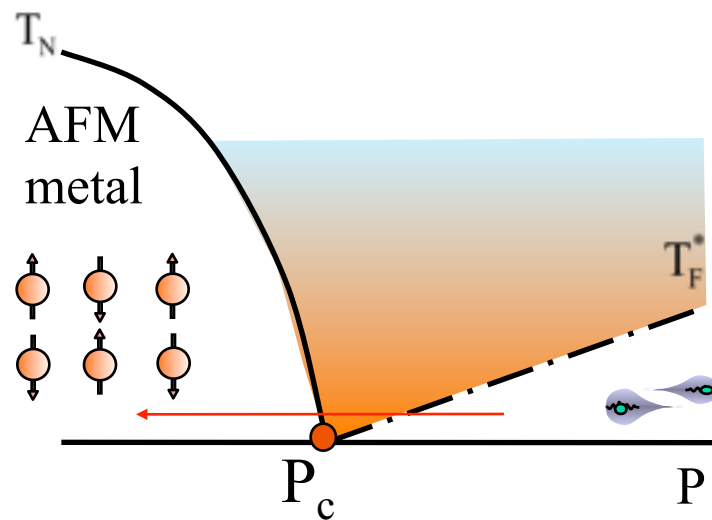


"Black Hole in the Phase Diagram".



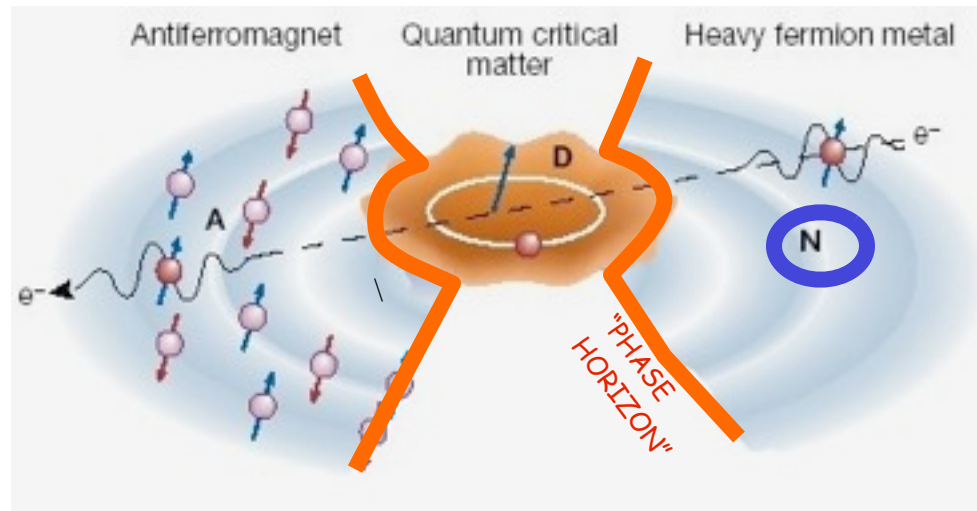
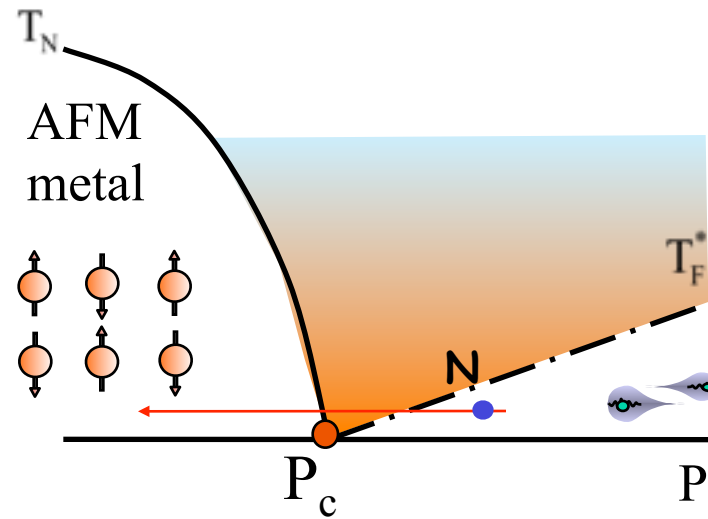
P.C and A. Schofield, Nature (2005)

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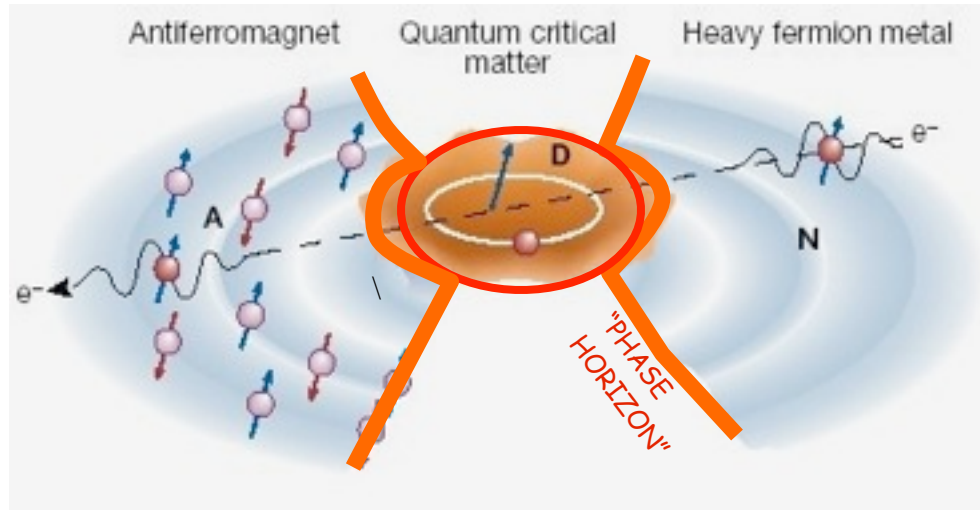
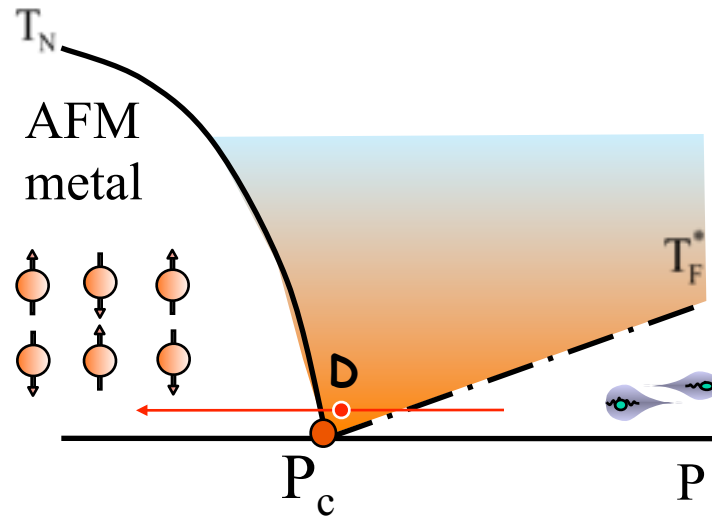
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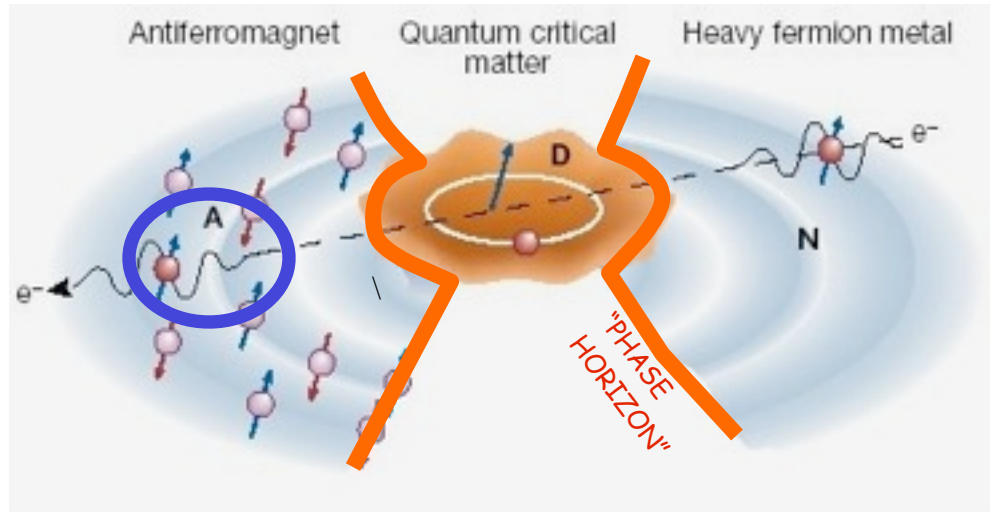
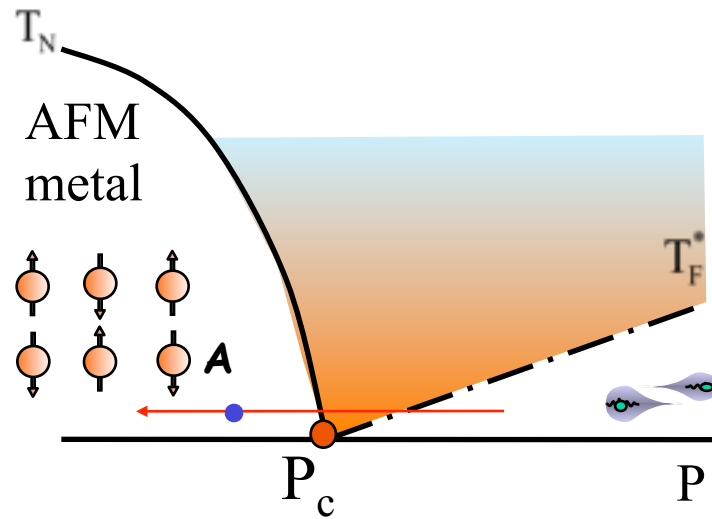
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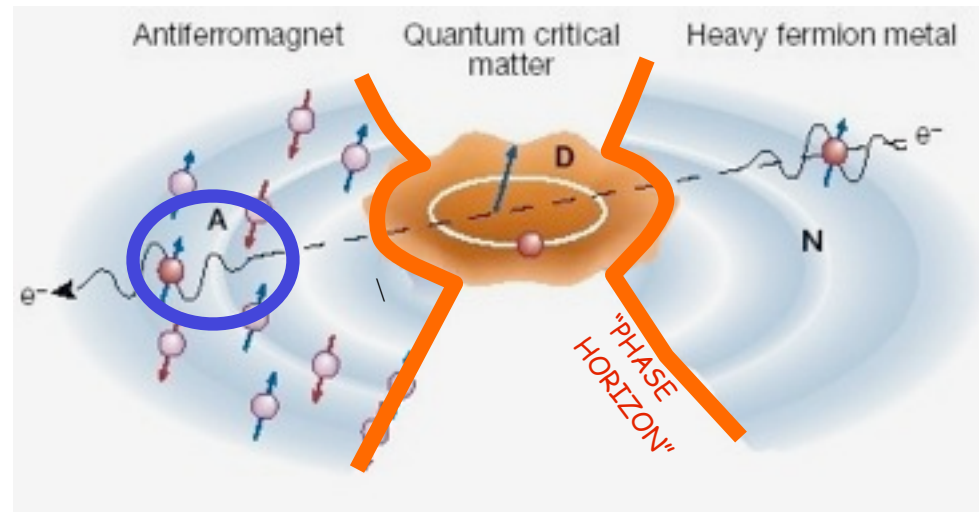
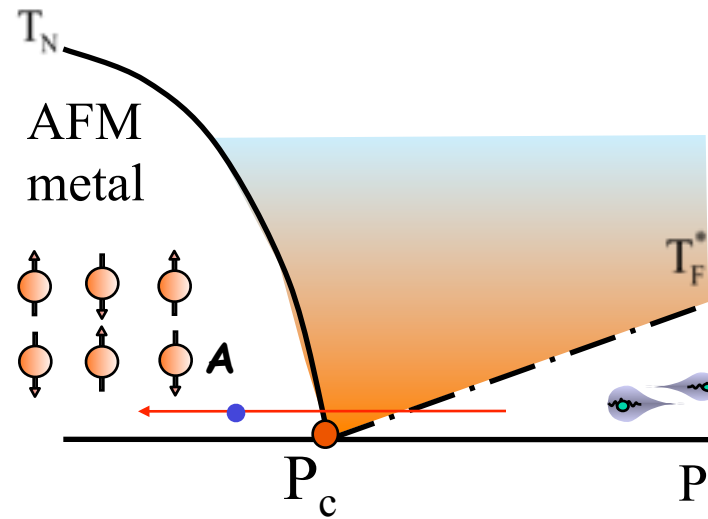
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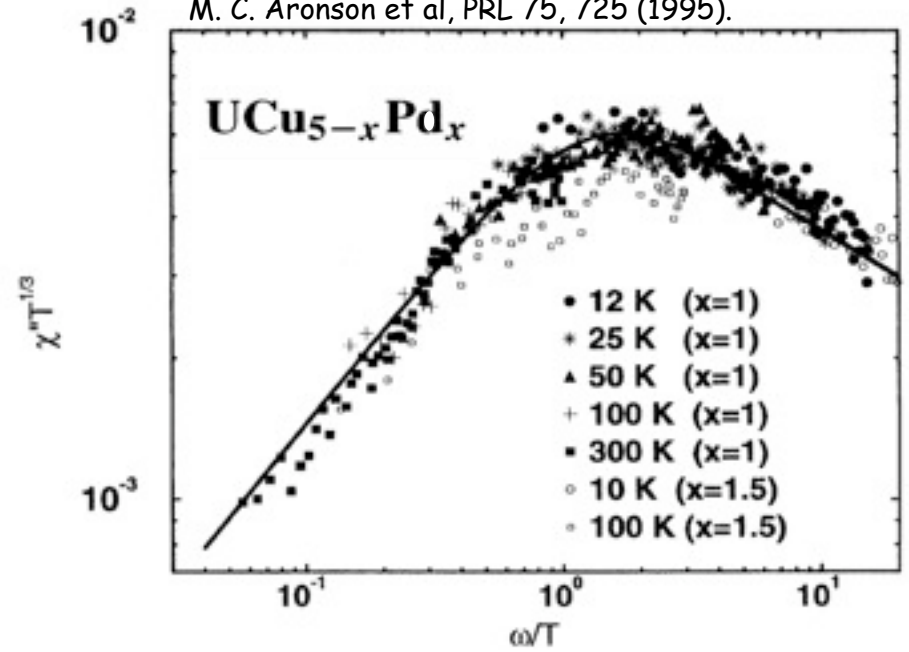
John Hertz: Critical droplet is **Quantum** if $\hbar\omega(q)|_{q=\xi^{-1}} \gg k_B T$

E/T Scaling:



Meigan Aronson

M. C. Aronson et al, PRL 75, 725 (1995).



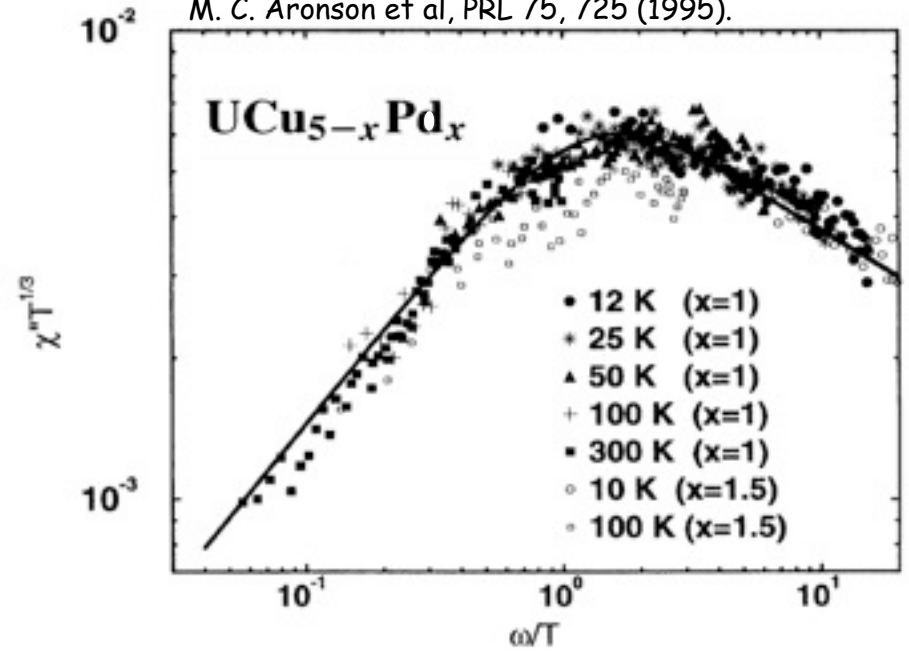
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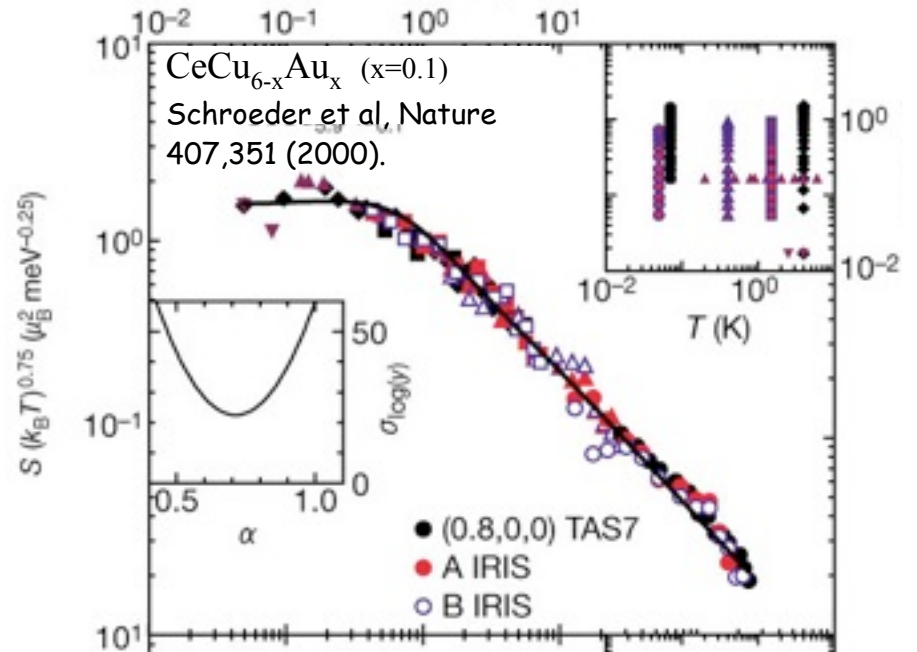
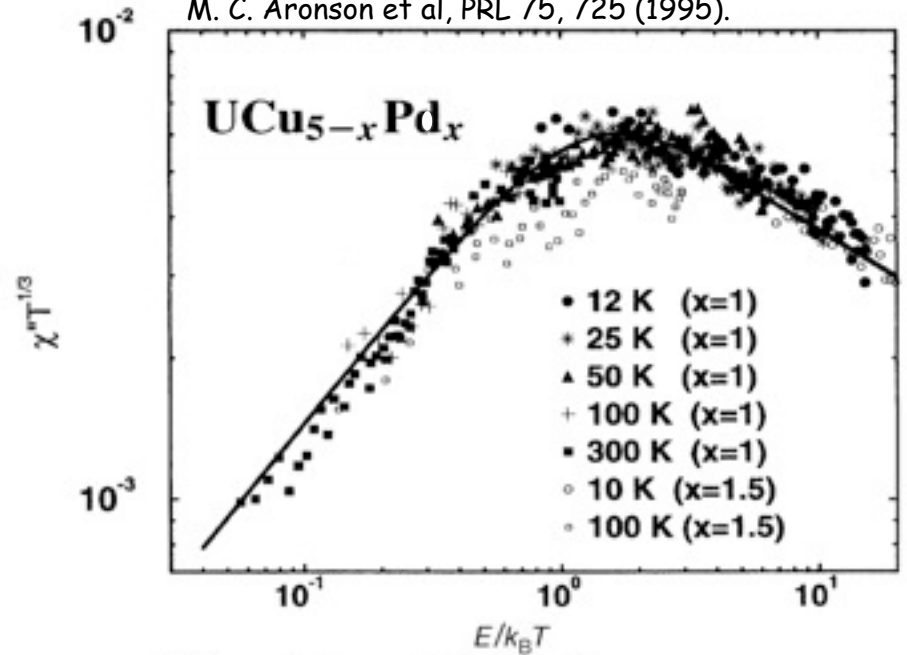
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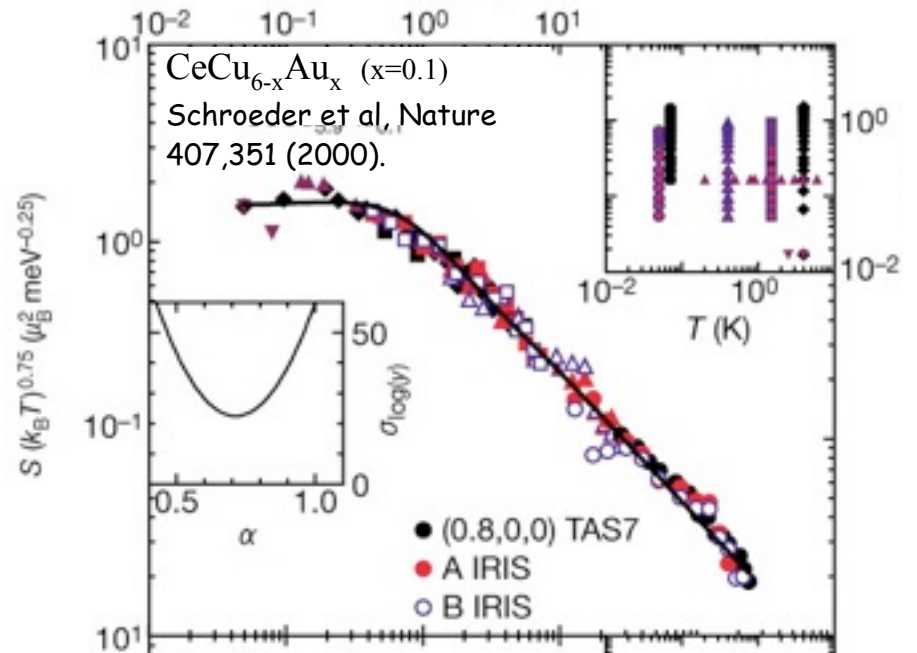
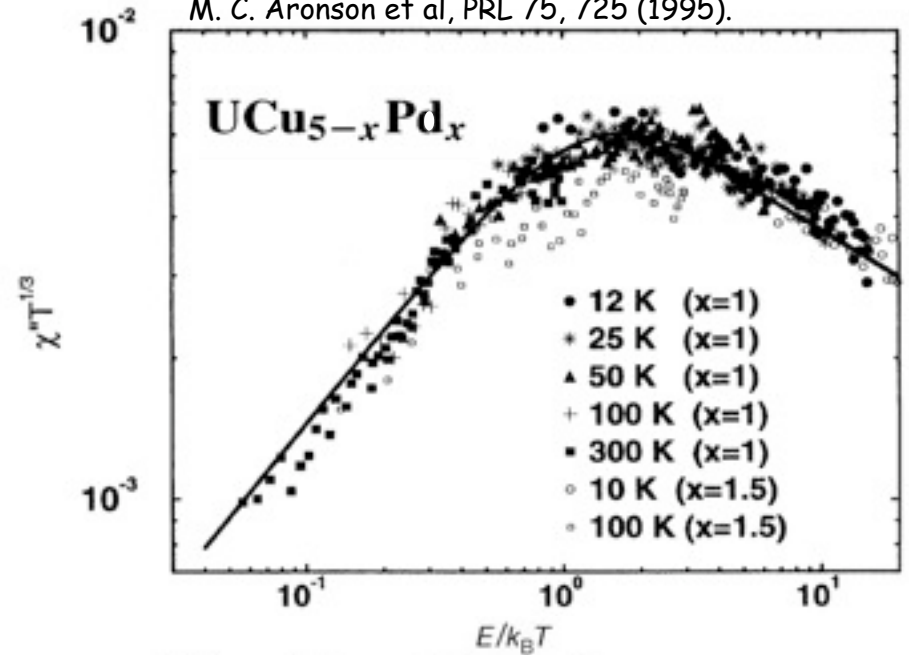


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Physics Below the upper
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M. C. Aronson et al, PRL 75, 725 (1995).



The Standard Model

Standard Model: Quantum SDW?



Doniach



Schrieffer



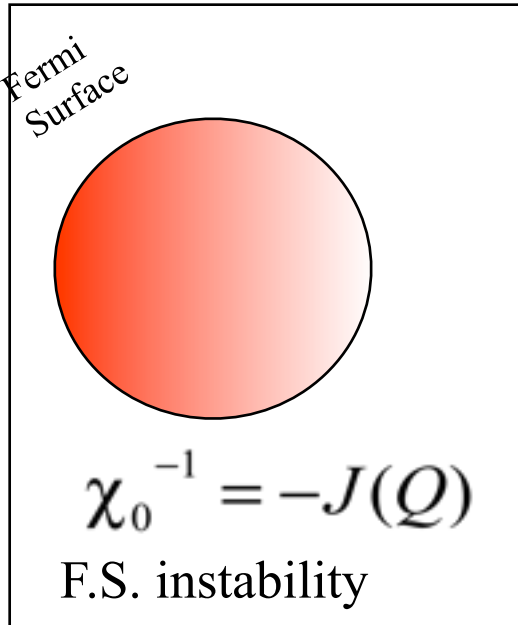
Hertz



Millis

- Moriya, Doniach, Schrieffer (60s)
- Hertz (76)
- Millis (93)

$$d_{eff} = d + z$$



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Schrieffer



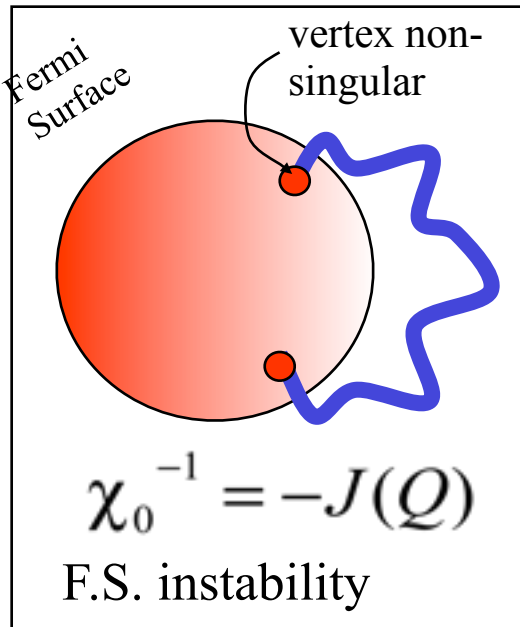
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$$\chi^{-1}(q, \omega) \propto (\xi^{-2} + (q - Q)^2 - i\omega / \Gamma)$$

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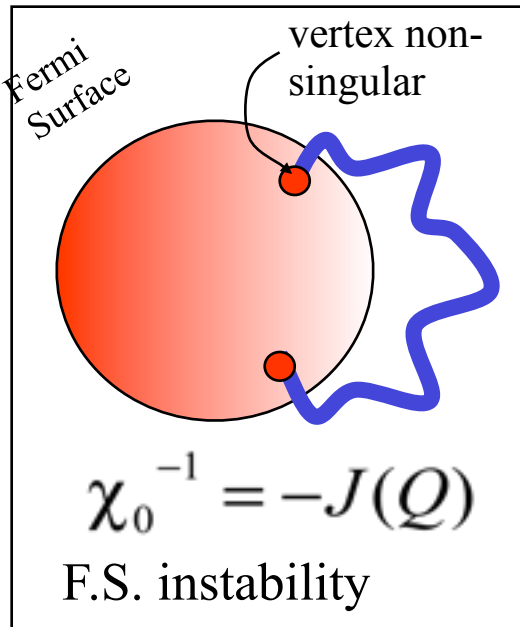
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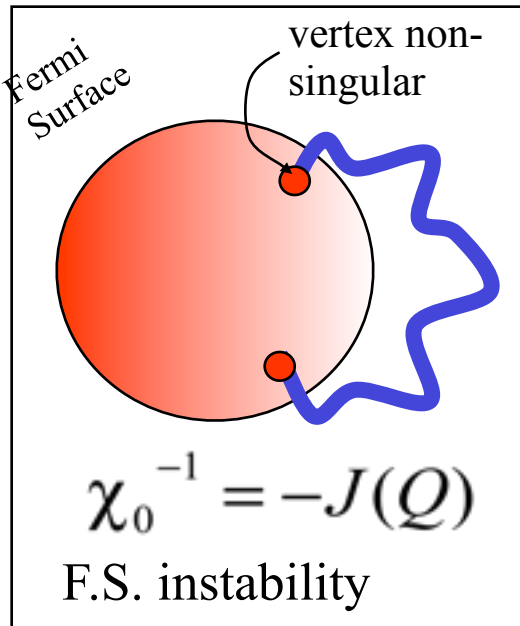
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Time counts as $z = 2$ scaling dimensions

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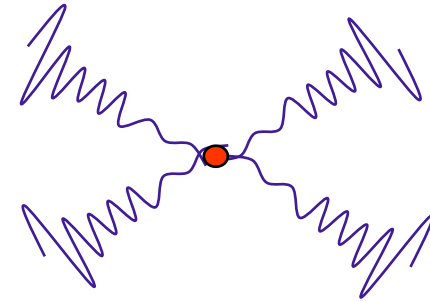
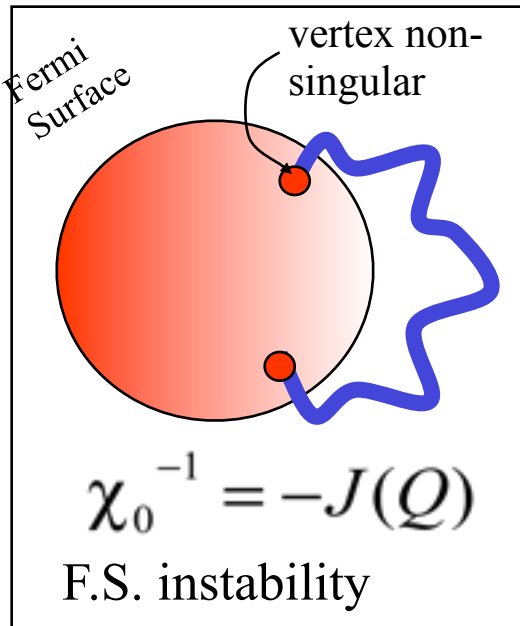
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$$d_{\text{eff}} = d + z$$



If $d + z = d + 2 > 4$:
 ϕ^4 terms “irrelevant”
 Critical modes are Gaussian.
 T is not the only energy scale.

$$\chi^{-1}(q, \omega) \propto (\xi^{-2} + (q - Q)^2 - i\omega / \Gamma)$$

$$\tau^{-1} \propto \xi^{-2}$$

Time counts as $z = 2$ scaling dimensions

New Ideas:

Break up of the electron.

$$H = \sum_k \epsilon_k c_{k\sigma}^\dagger c_{k\sigma} + J \sum_j (\psi_j^\dagger \vec{\sigma} \psi_j) \cdot \vec{S}_j$$

"THE BATTLEGROUND"

Kondo Lattice Model

(Kasuya, 1951)

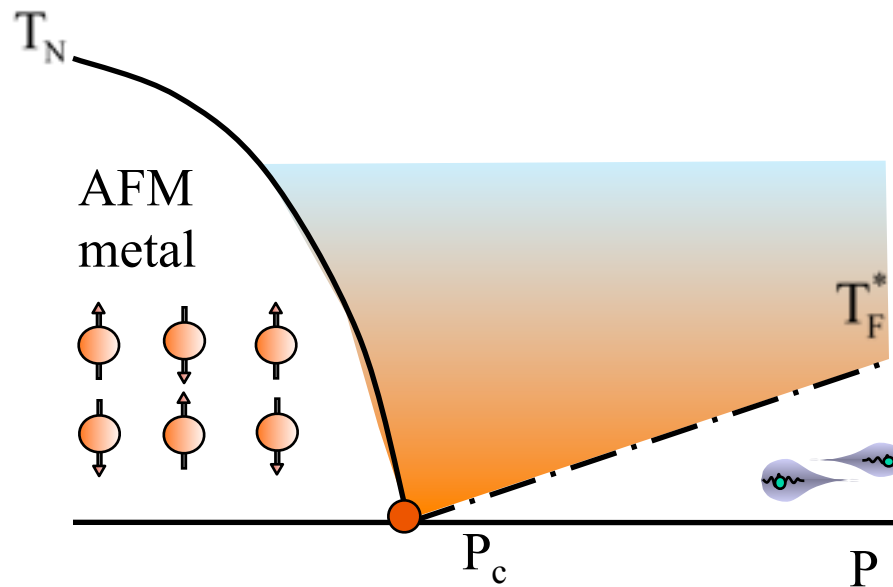
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Kondo Lattice Model

(Kasuya, 1951)

Heavy Fermion
Materials



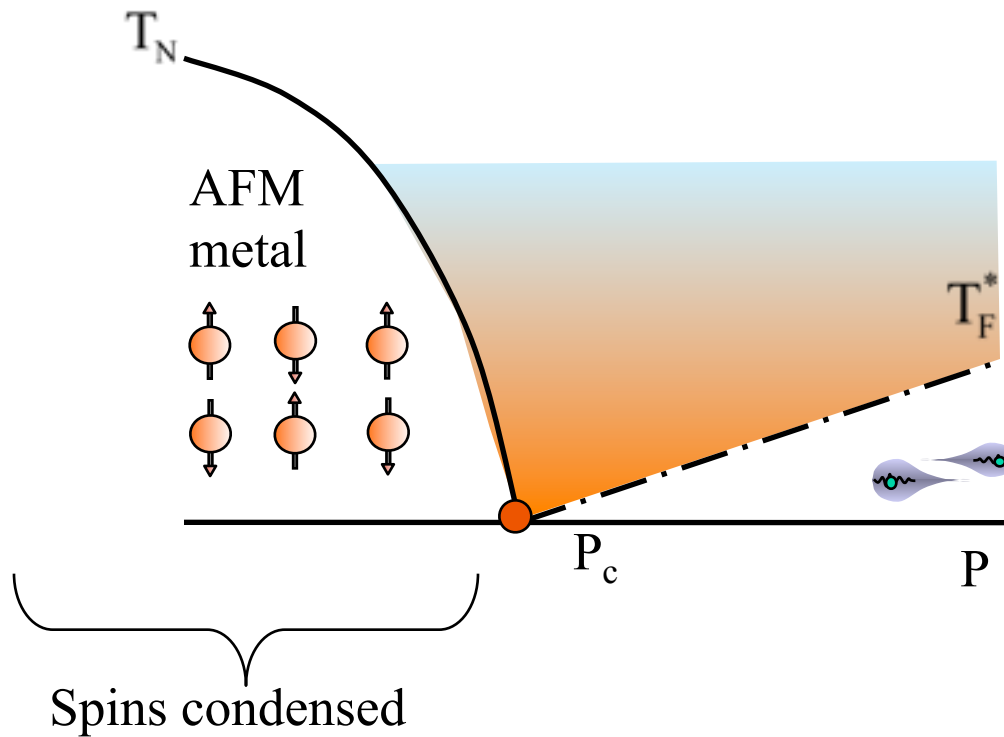
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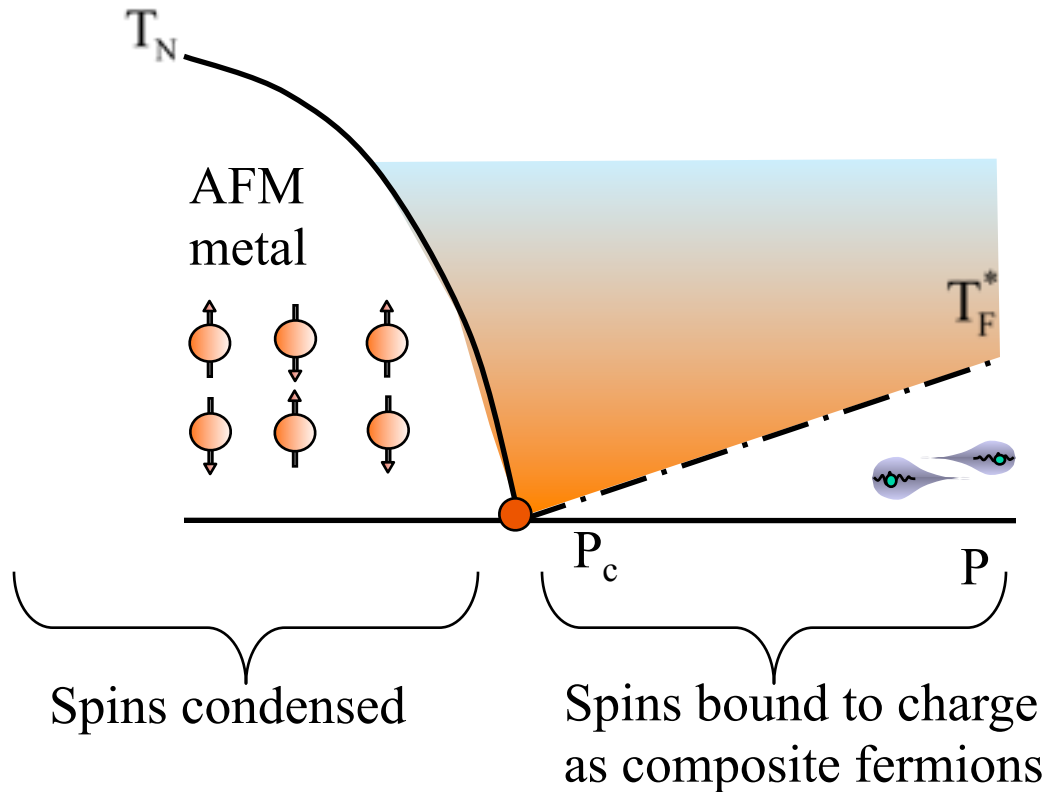
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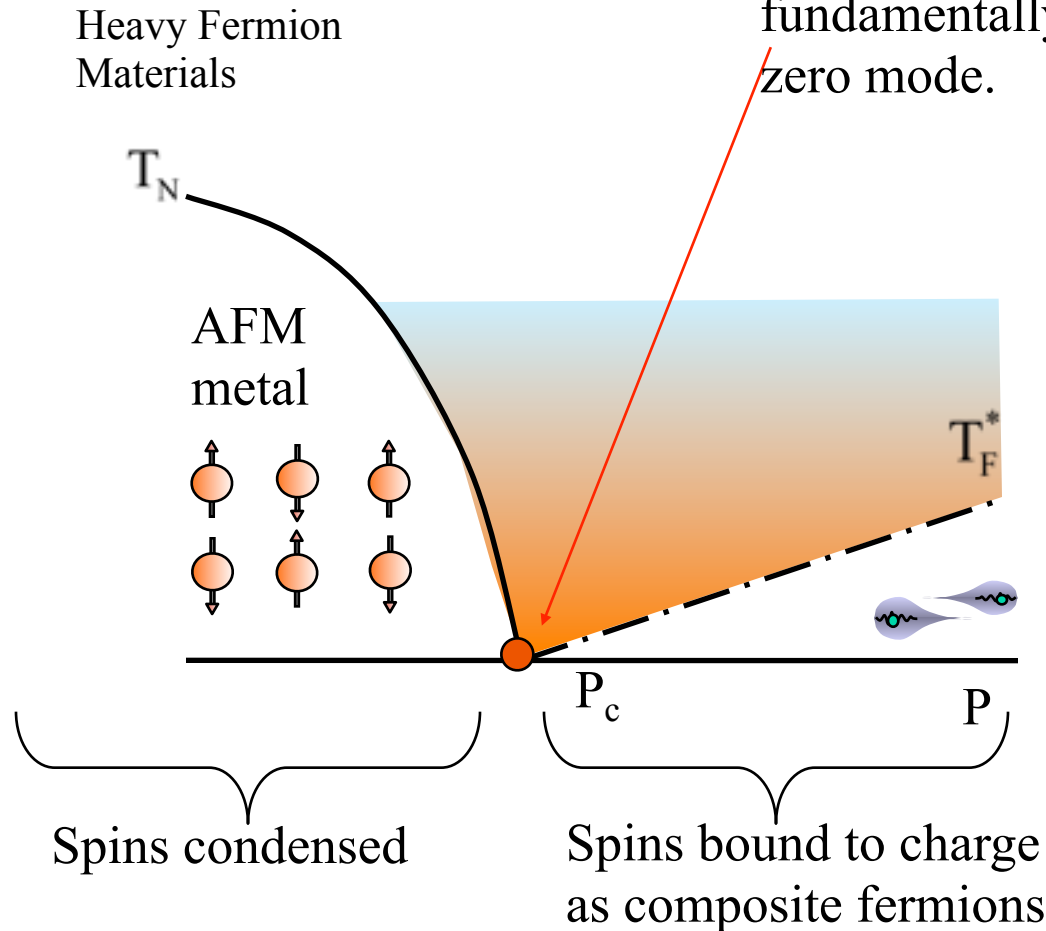
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"THE BATTLEGROUND"

Kondo Lattice Model

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Deconfinement of spin:
fundamentally new kind of
zero mode.

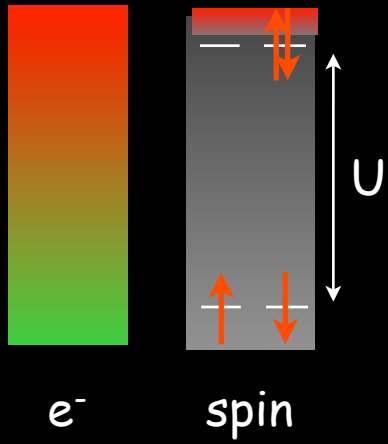


New Methods

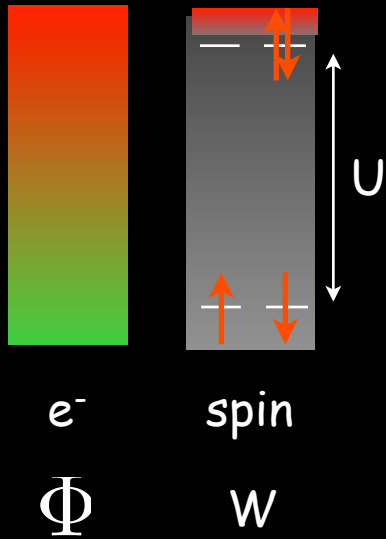


e^-

New Methods



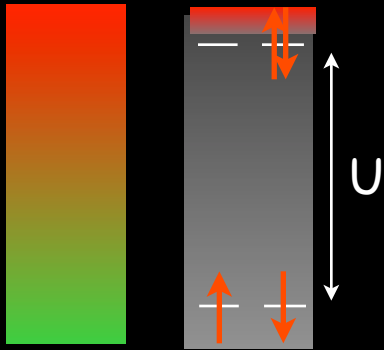
New Methods



Elimination of States
implies *Gauge Fields*.

(Read Newns, PC, Millis *Lee*... 80's)

New Methods



e^-

spin

Φ

W

$$\vec{S} = b^\dagger_\alpha \left(\frac{\vec{\sigma}}{2} \right)_{\alpha\beta} b_\beta$$

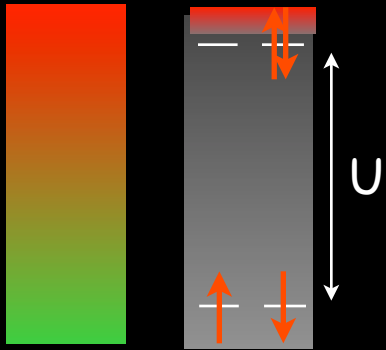
J. Schwinger '55

$$U(1) : b_j \rightarrow e^{i\theta_j} b_j$$

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e^-

spin

Φ

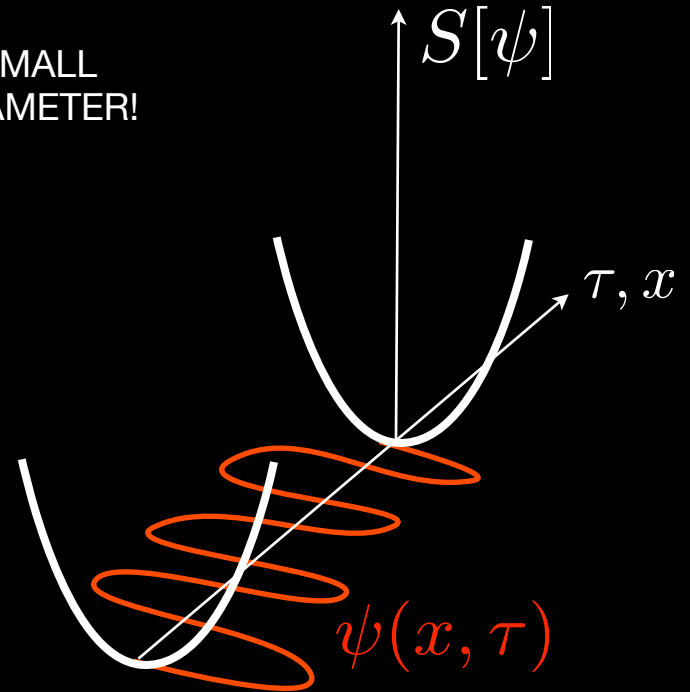
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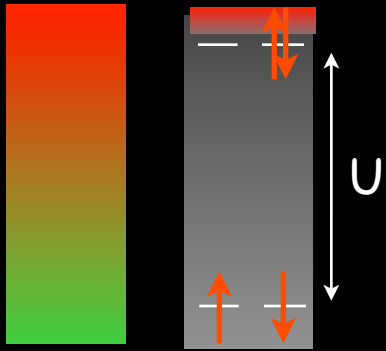
NO SMALL
PARAMETER!



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New Methods



e^-

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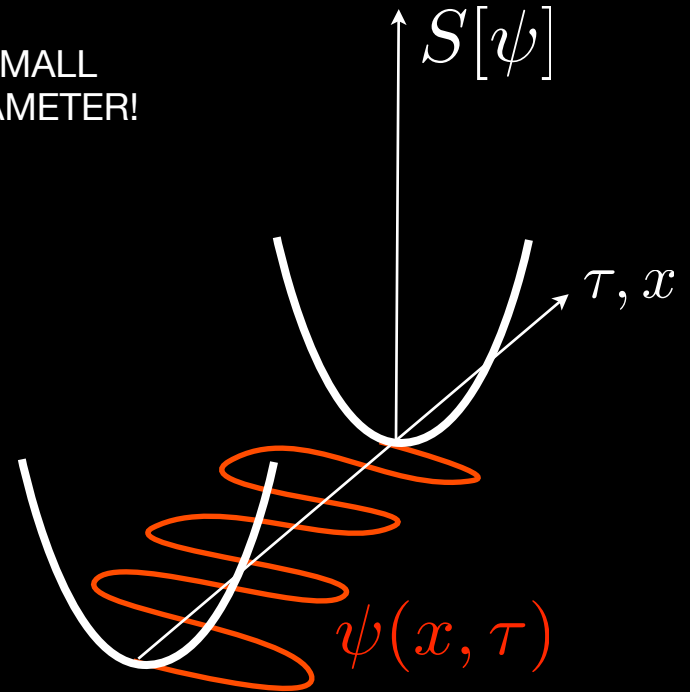
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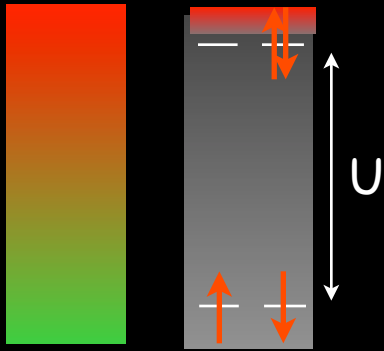


Elimination of States
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(Read Newns, PC, Millis Lee... 80's)

Large N : family of models with “N” spin components, which retain the key physics and can be solved in the large N limit.

New Methods



e^- spin

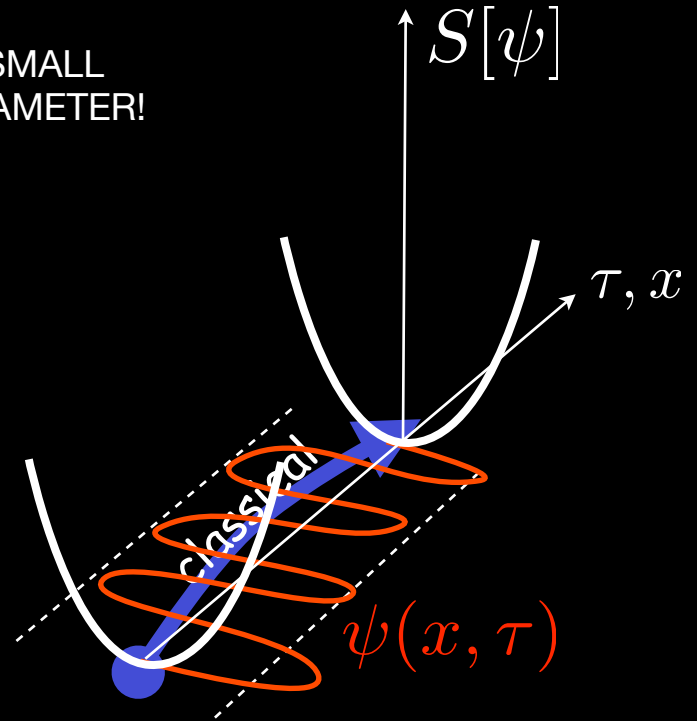
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J. Schwinger '55

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NO SMALL
PARAMETER!



$$\frac{1}{N} \sim \hbar_{eff}$$

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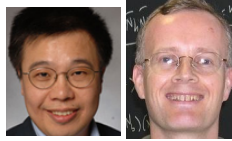
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New Ideas

New Ideas

Si, Ingersent



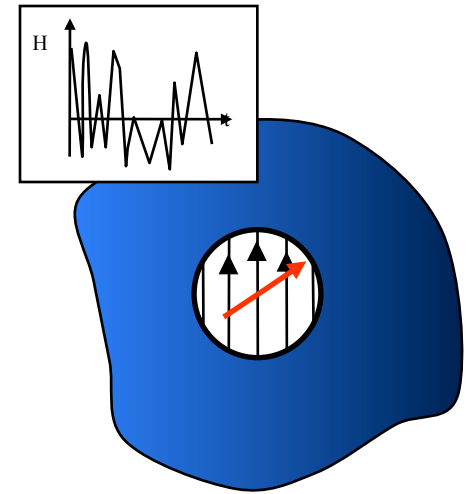
- **Local quantum criticality**

(Si, Ingersent, Smith, Rabello, Nature 2001):

Spin is the critical mode,

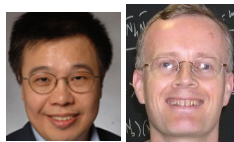
Fluctuations critical in time.

Requires a two dimensional spin fluid



New Ideas

Si, Ingersent

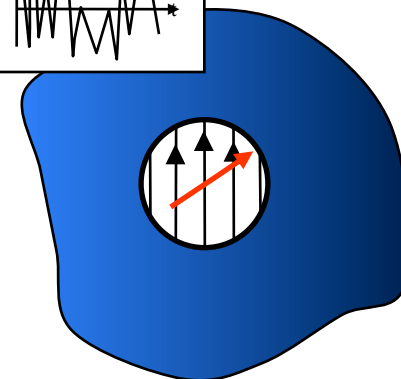
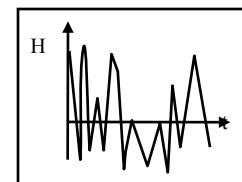


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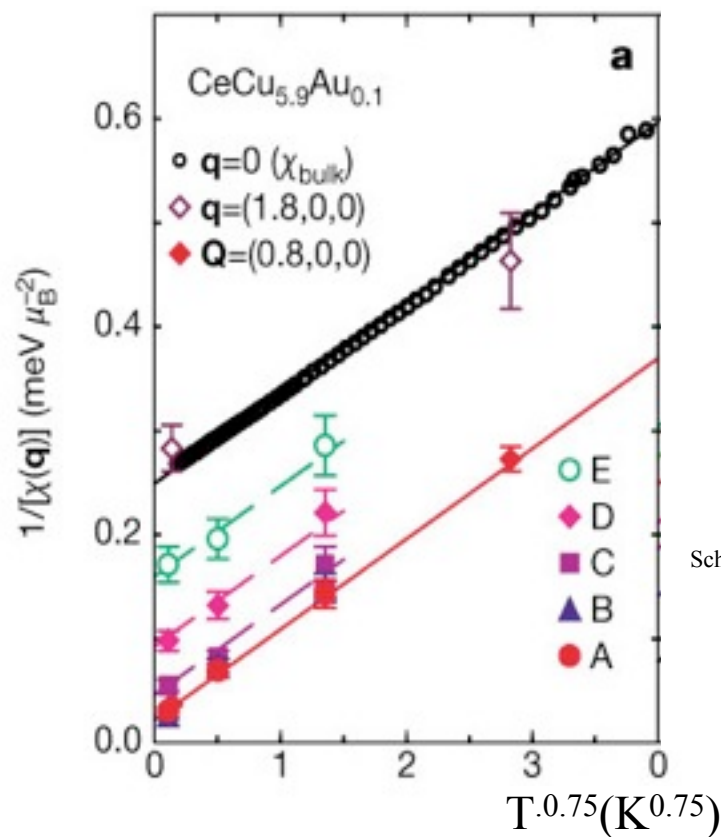
Fluctuations critical in time.



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Locality of critical fluctuations

$$\chi^{-1} = \chi_0^{-1} + AT^\alpha$$



Schroeder et al, Nature 407,351(2000).

New Ideas

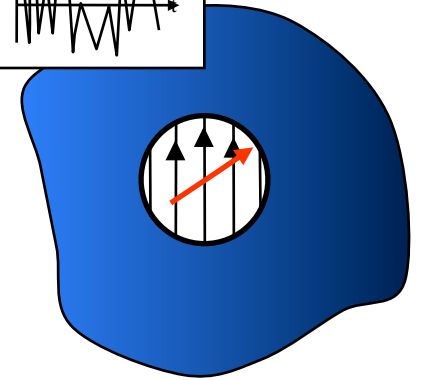
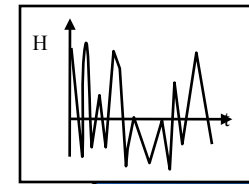
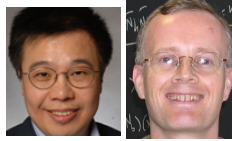
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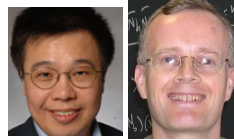
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New Ideas

Si, Ingersent

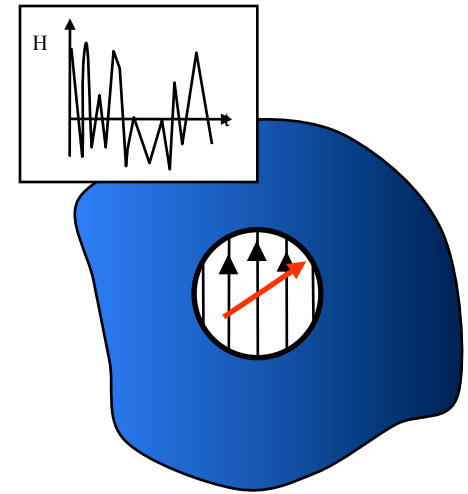


- **Local quantum criticality**

(Si, Ingersent, Smith, Rabello, Nature 2001):

Spin is the critical mode,
Fluctuations critical in time.

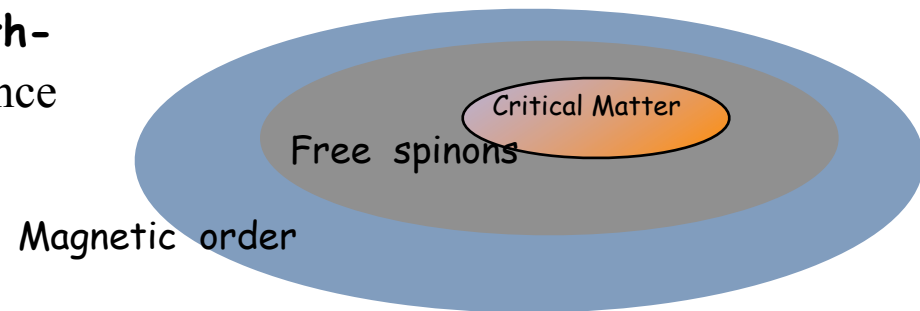
Requires a two dimensional spin fluid



- **Deconfined Criticality: Two diverging length-scales.** (Hermele et al 2004; Senthil et al, Science 2004).

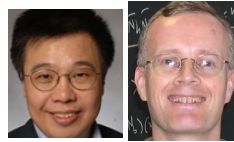


Senthil Sachdev Vishwanath



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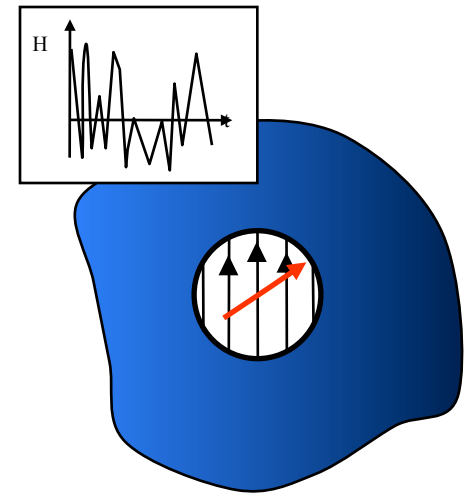


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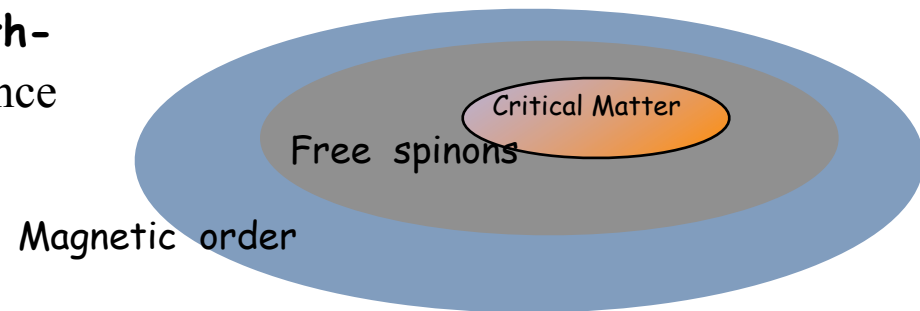
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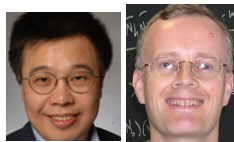
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New Ideas

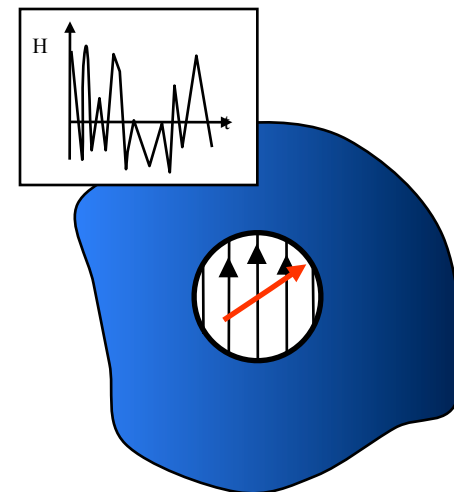
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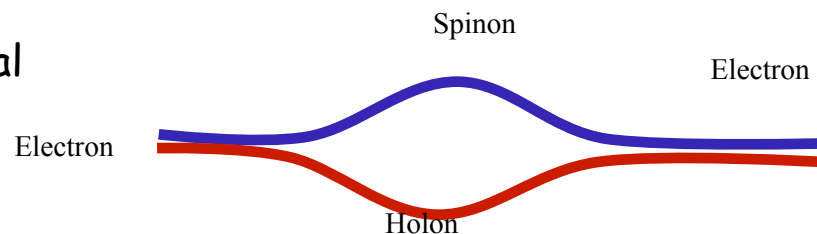
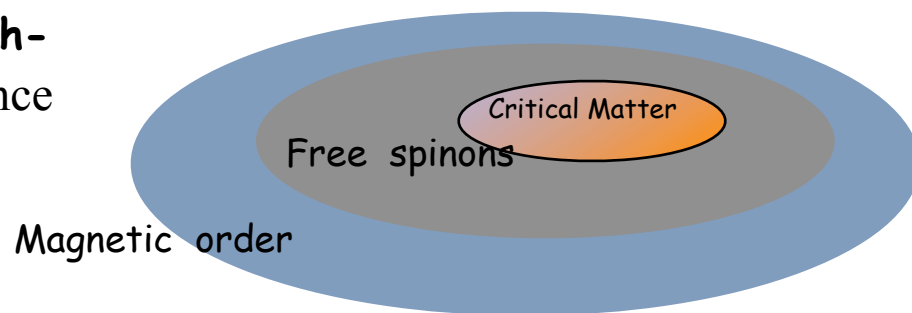


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Pepin



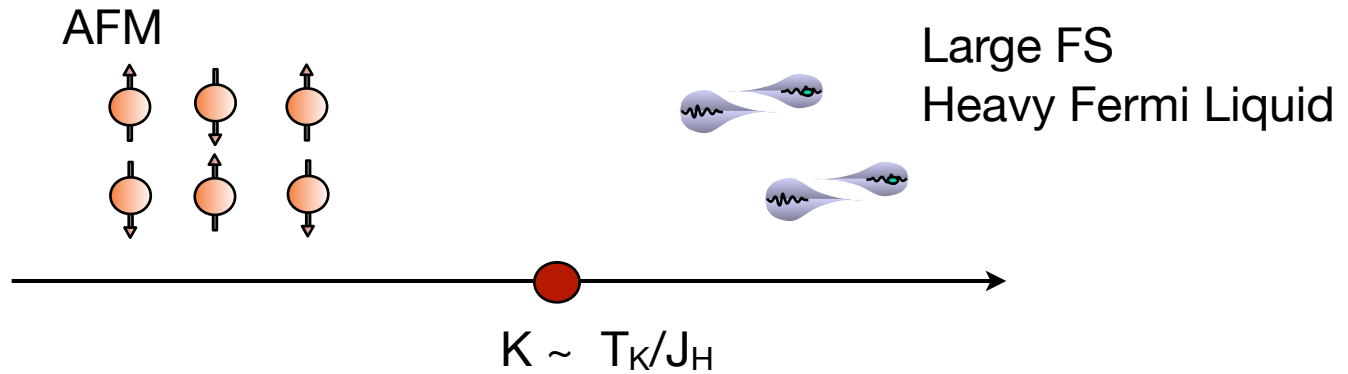
Q-Frustration

Kondo meets frustration

- Frustration and Kondo have different effects.

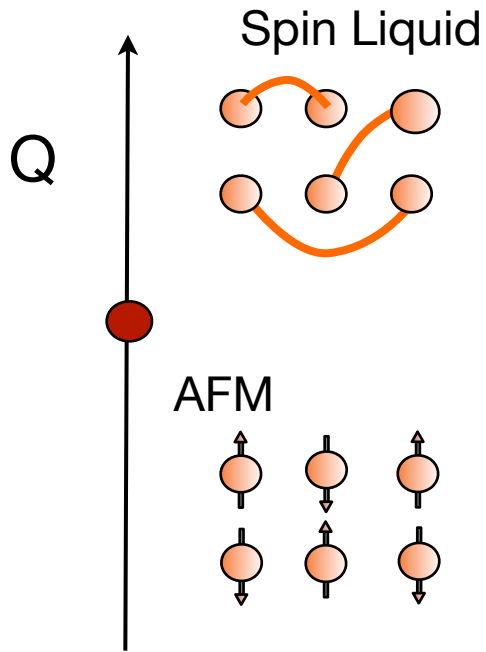
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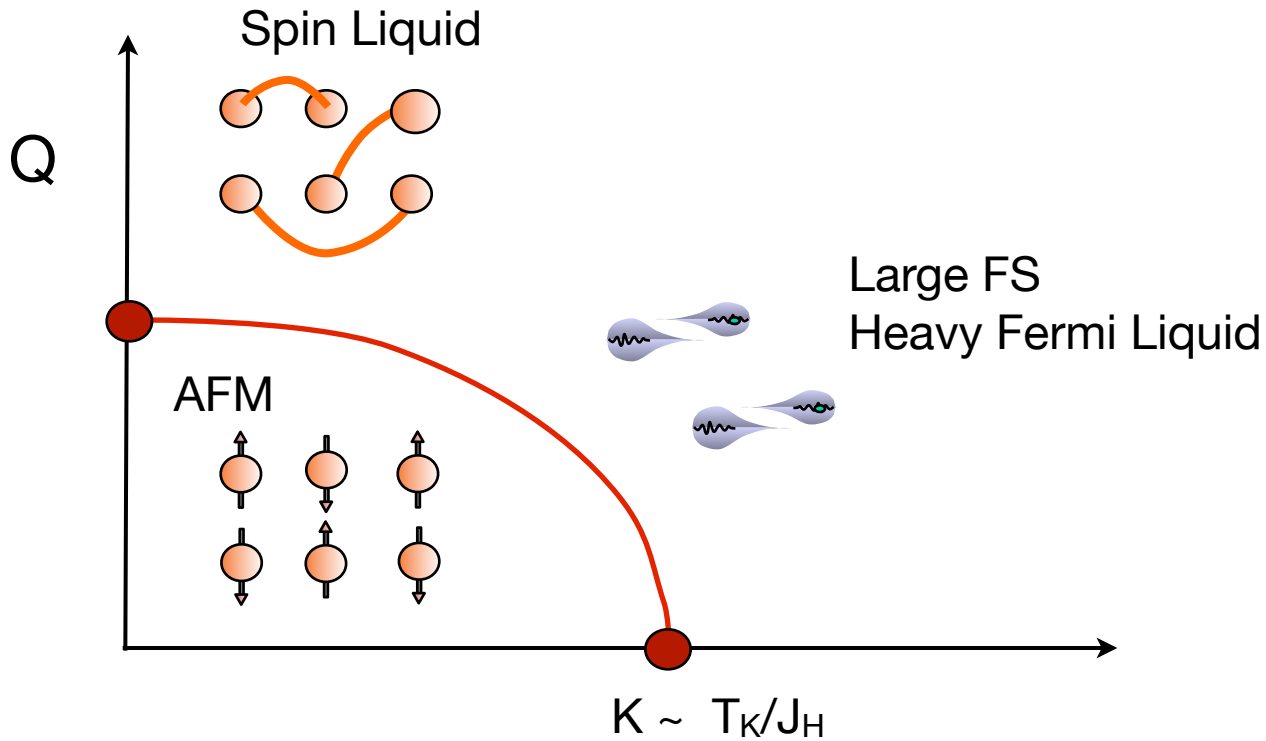
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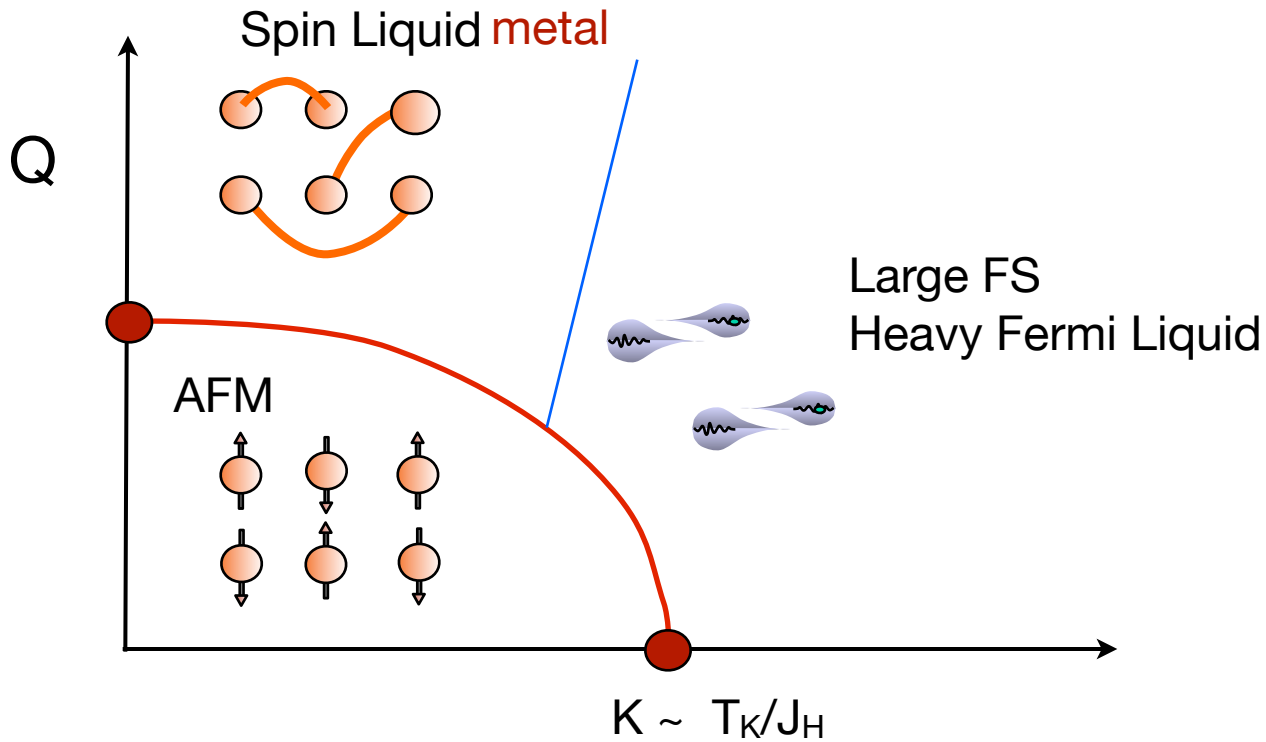
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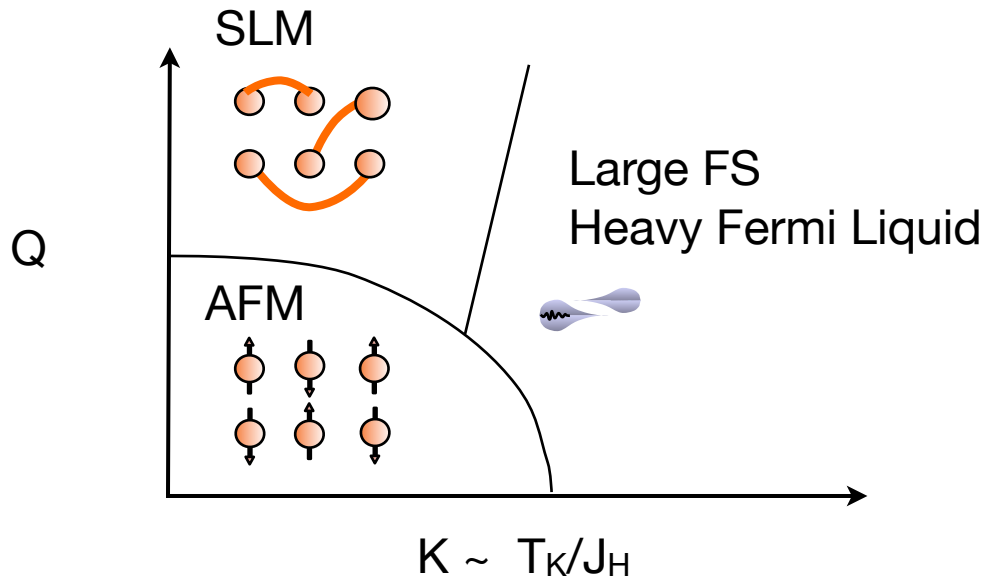


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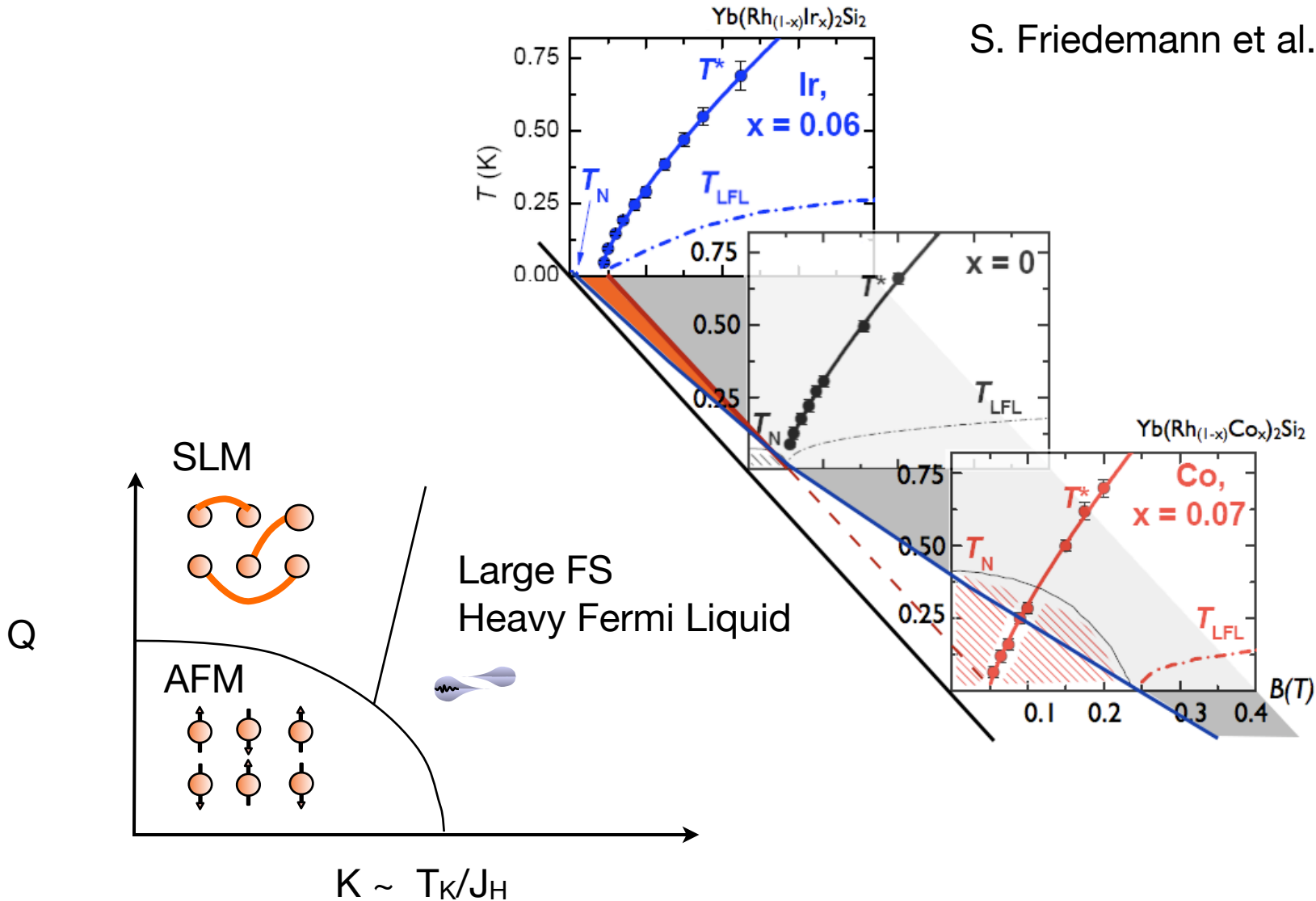


Experimental Support I



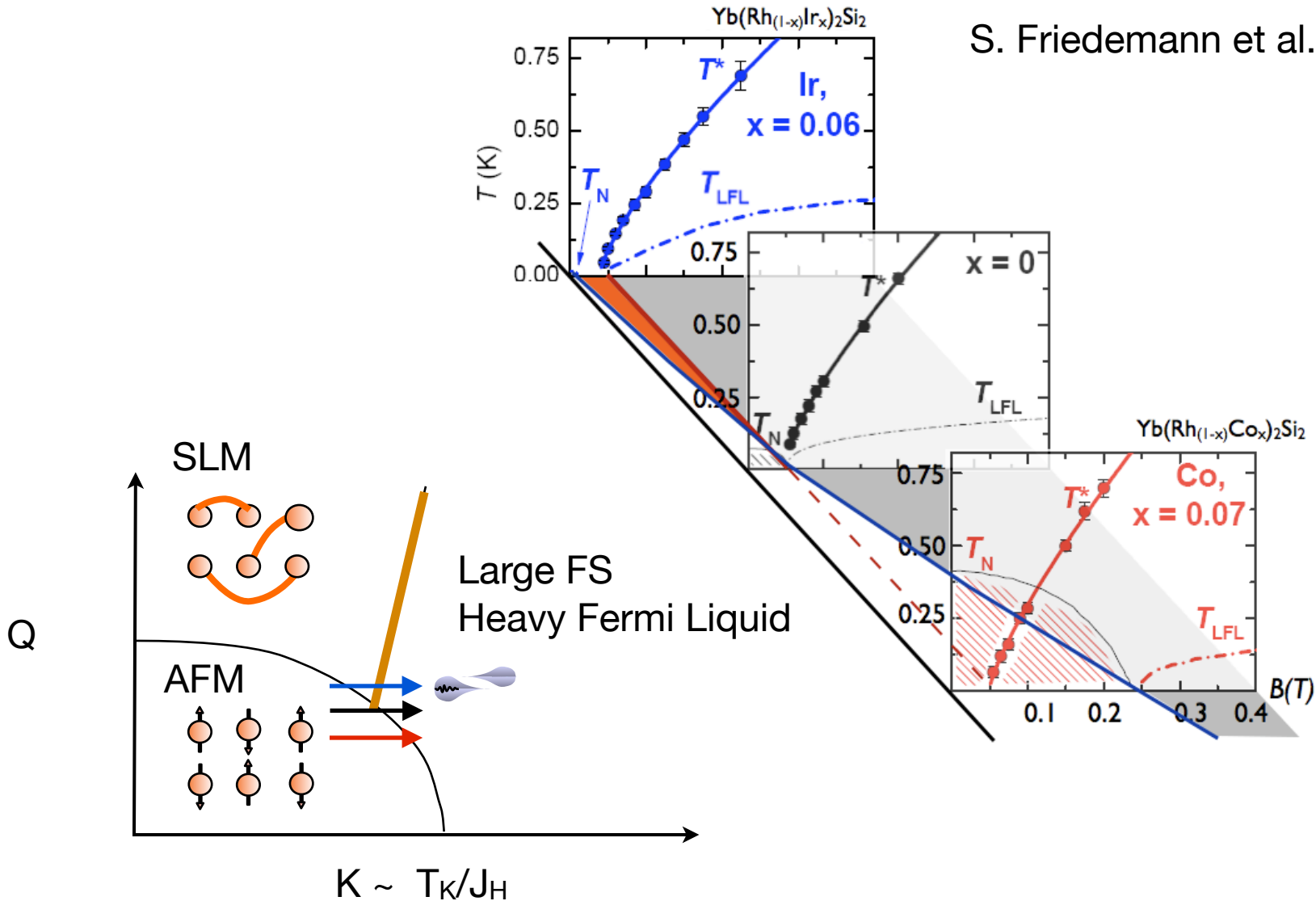
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S. Friedemann et al. (2009).



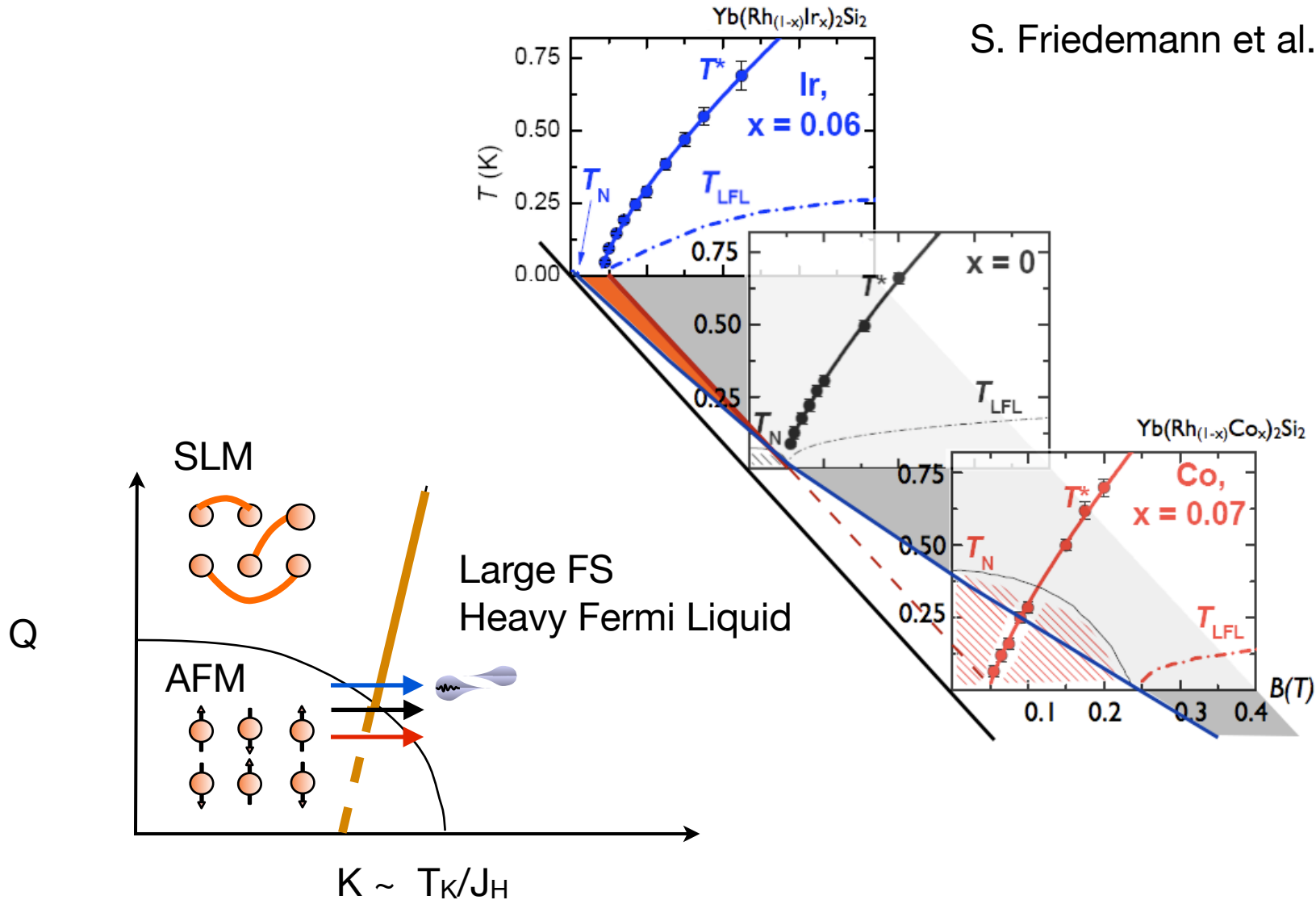
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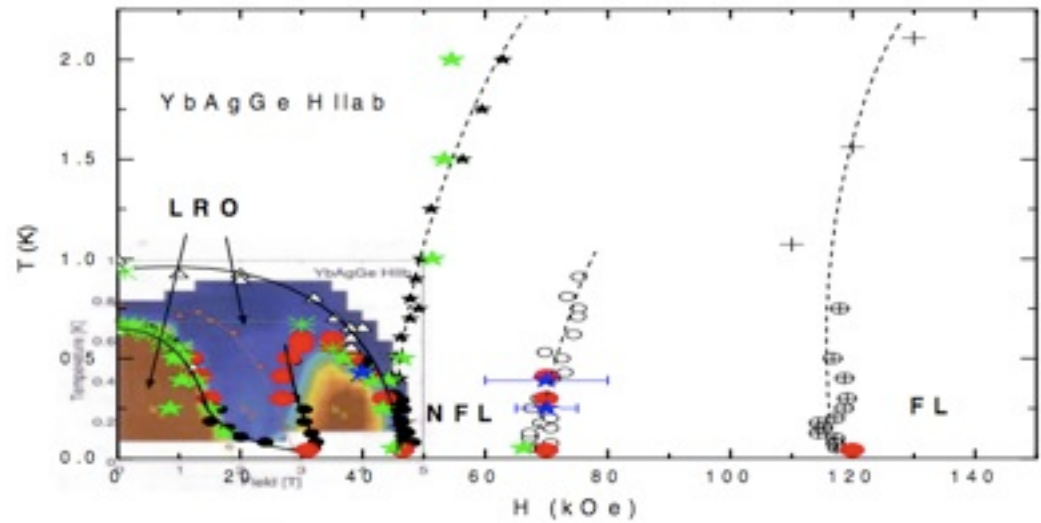


Experimental Support I

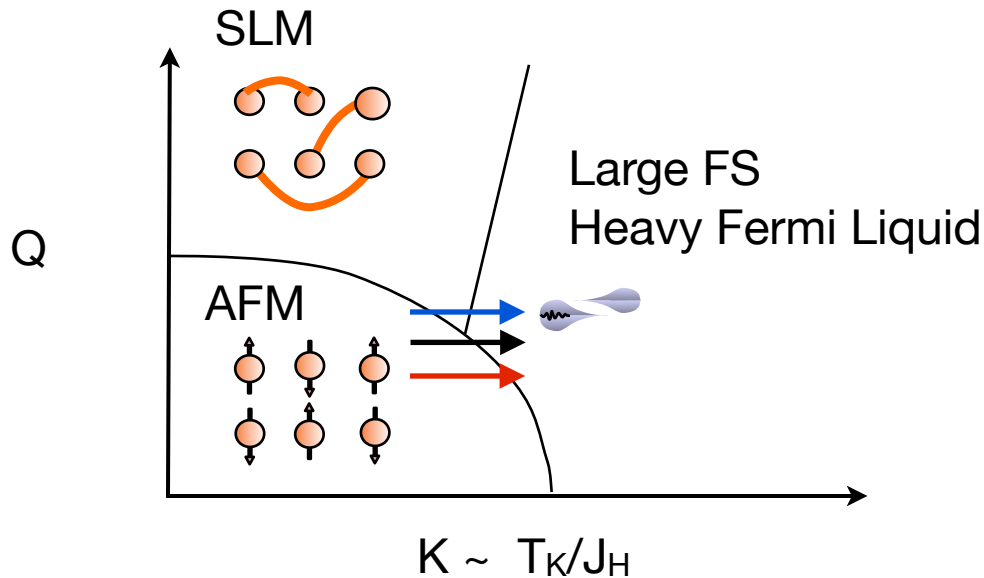
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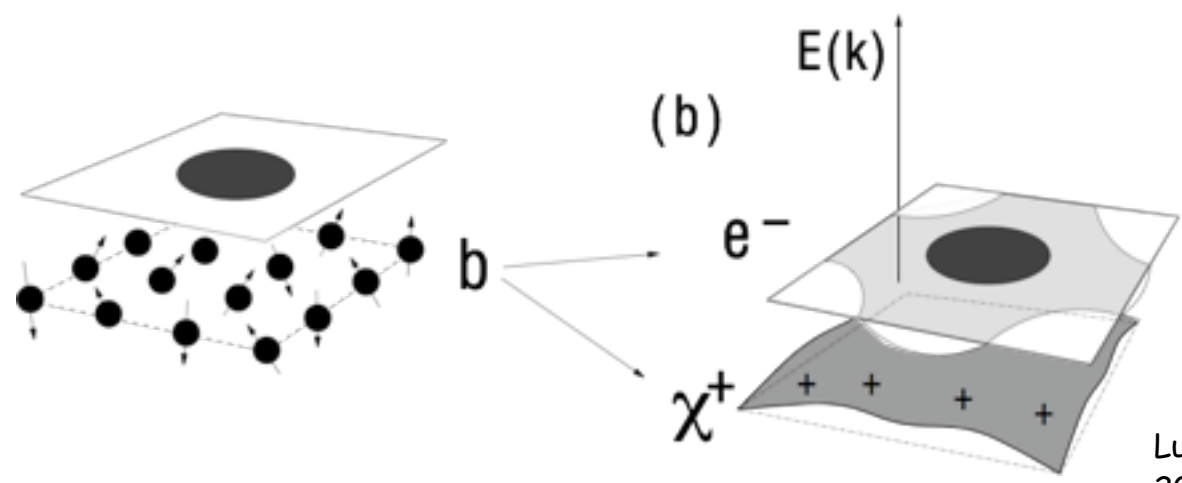


Experimental Support II



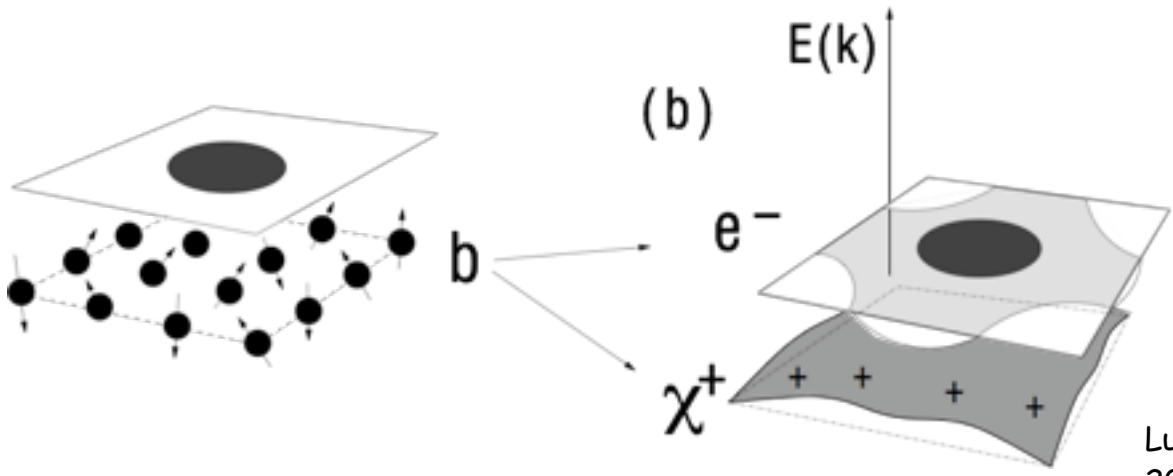
Canfield et al (unpublished)





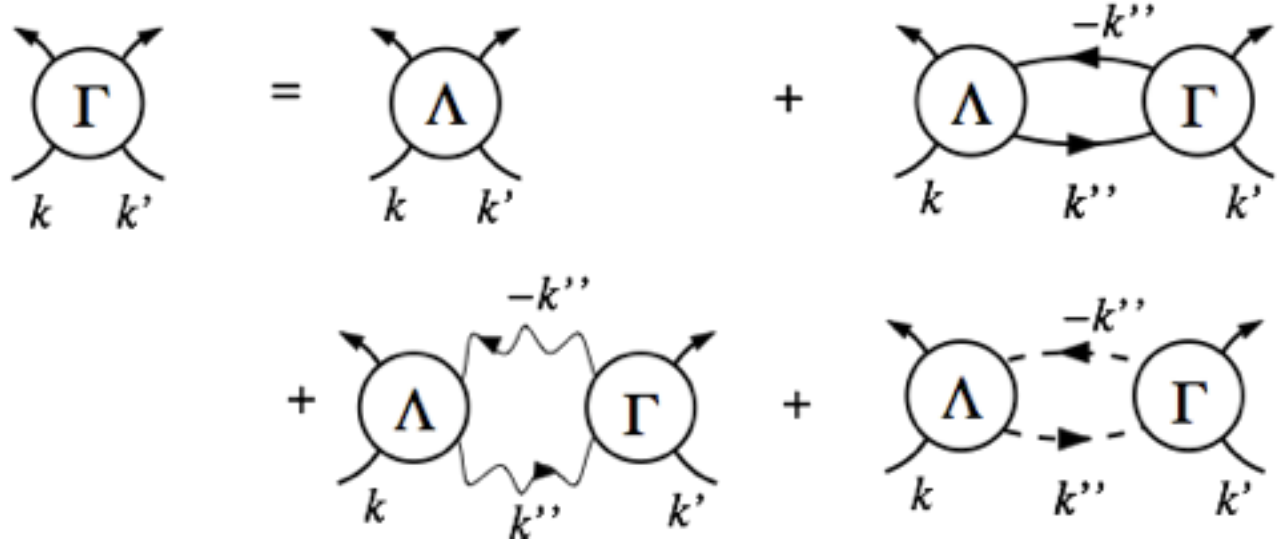
$$\frac{n_e}{K} = N \frac{v_{FS}}{(2\pi)^D} - \underbrace{\frac{v_\chi}{(2\pi)^D}}_{=1}$$

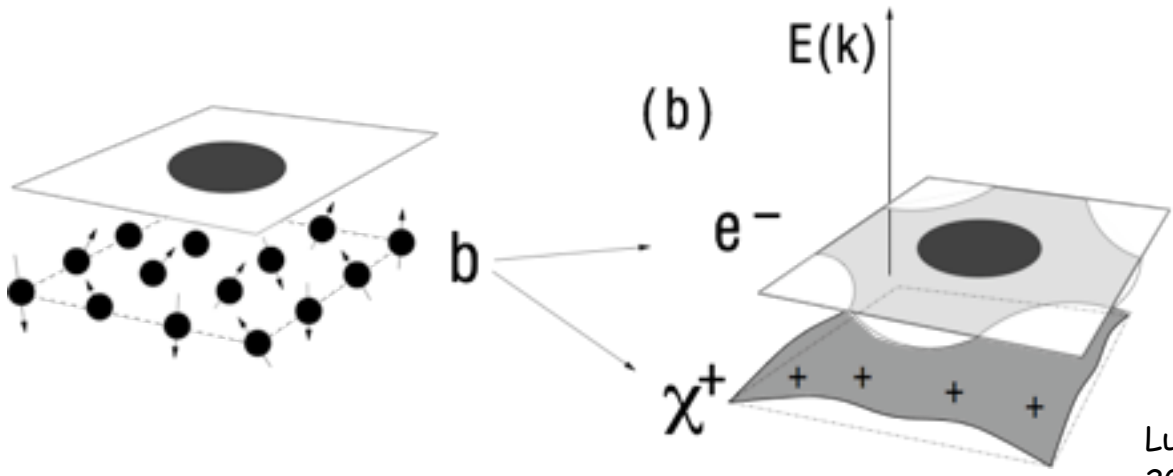
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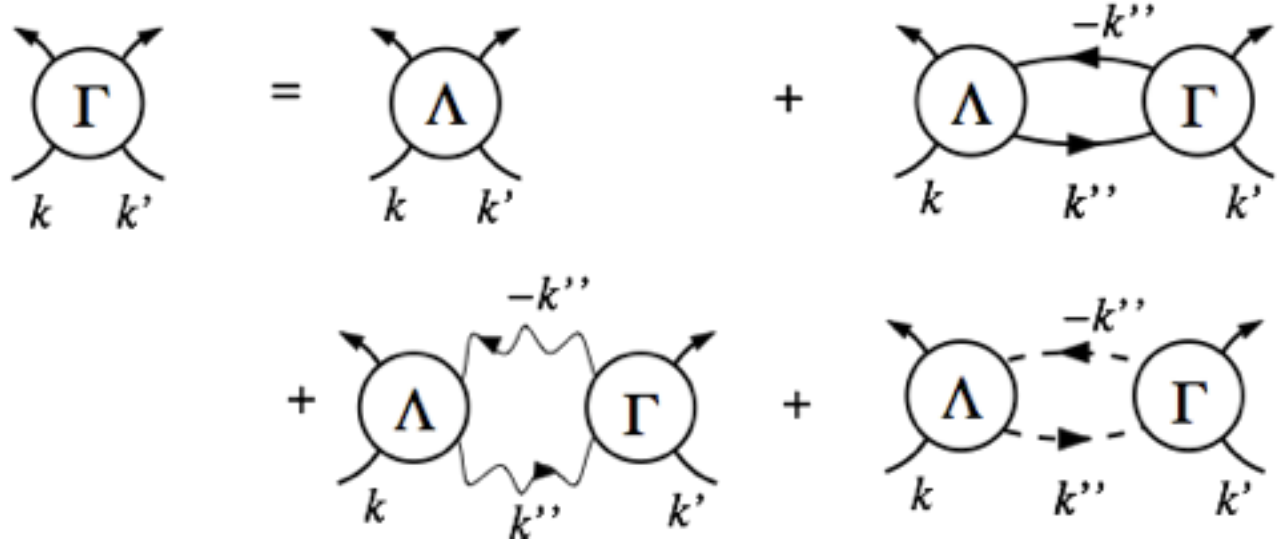
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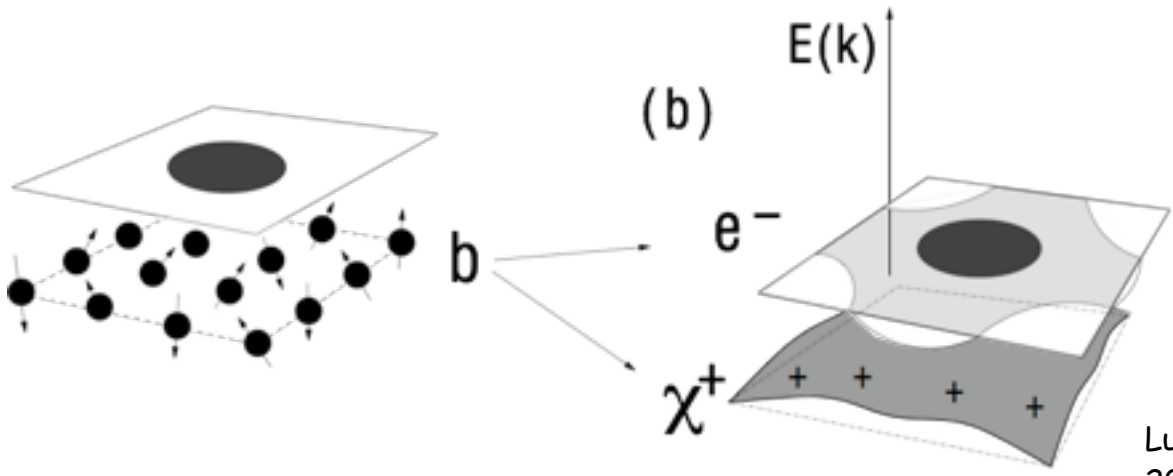


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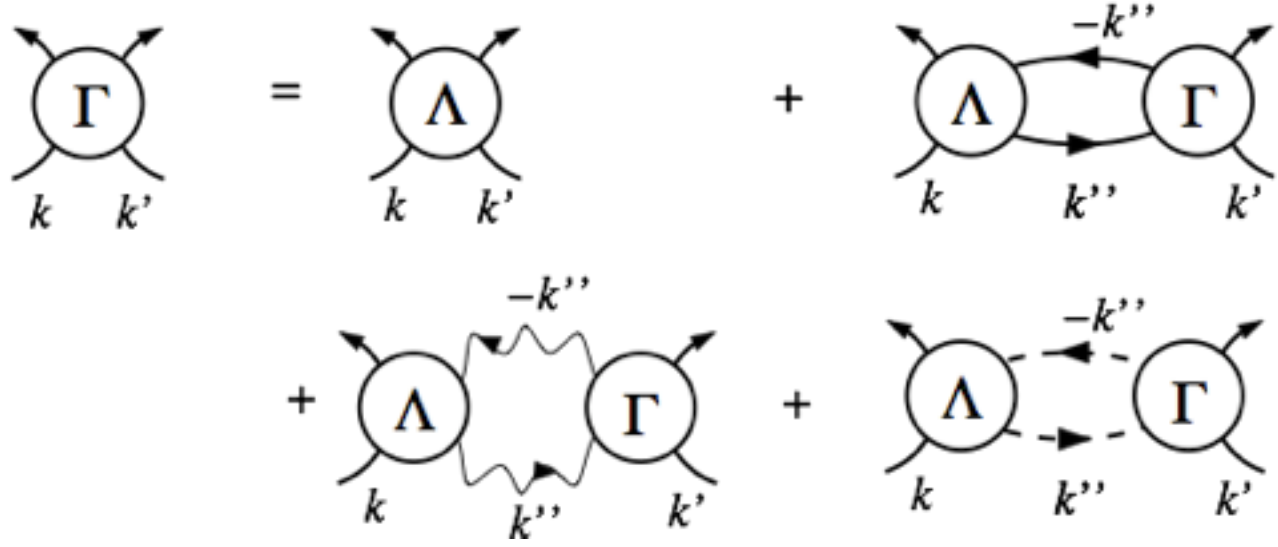


Virtual spinons.



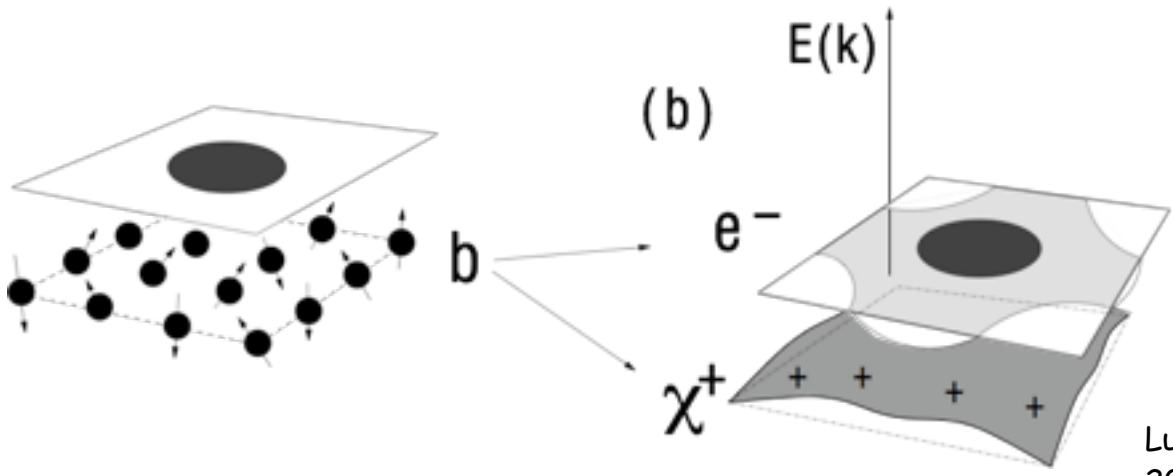
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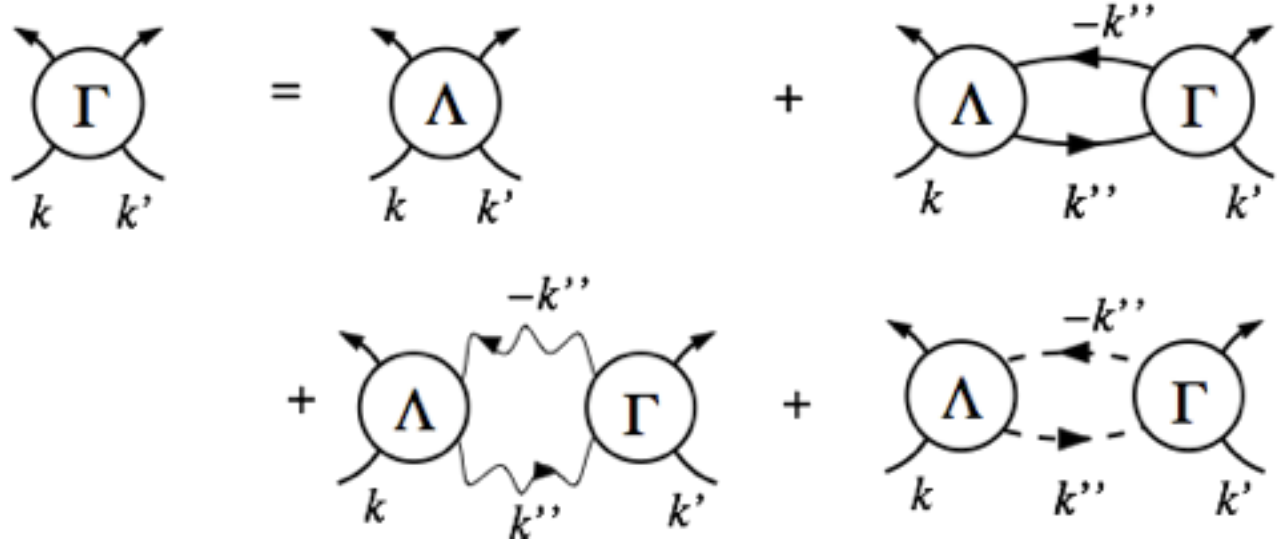
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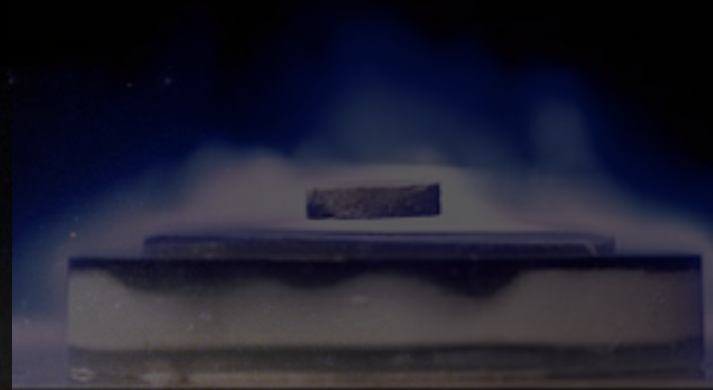
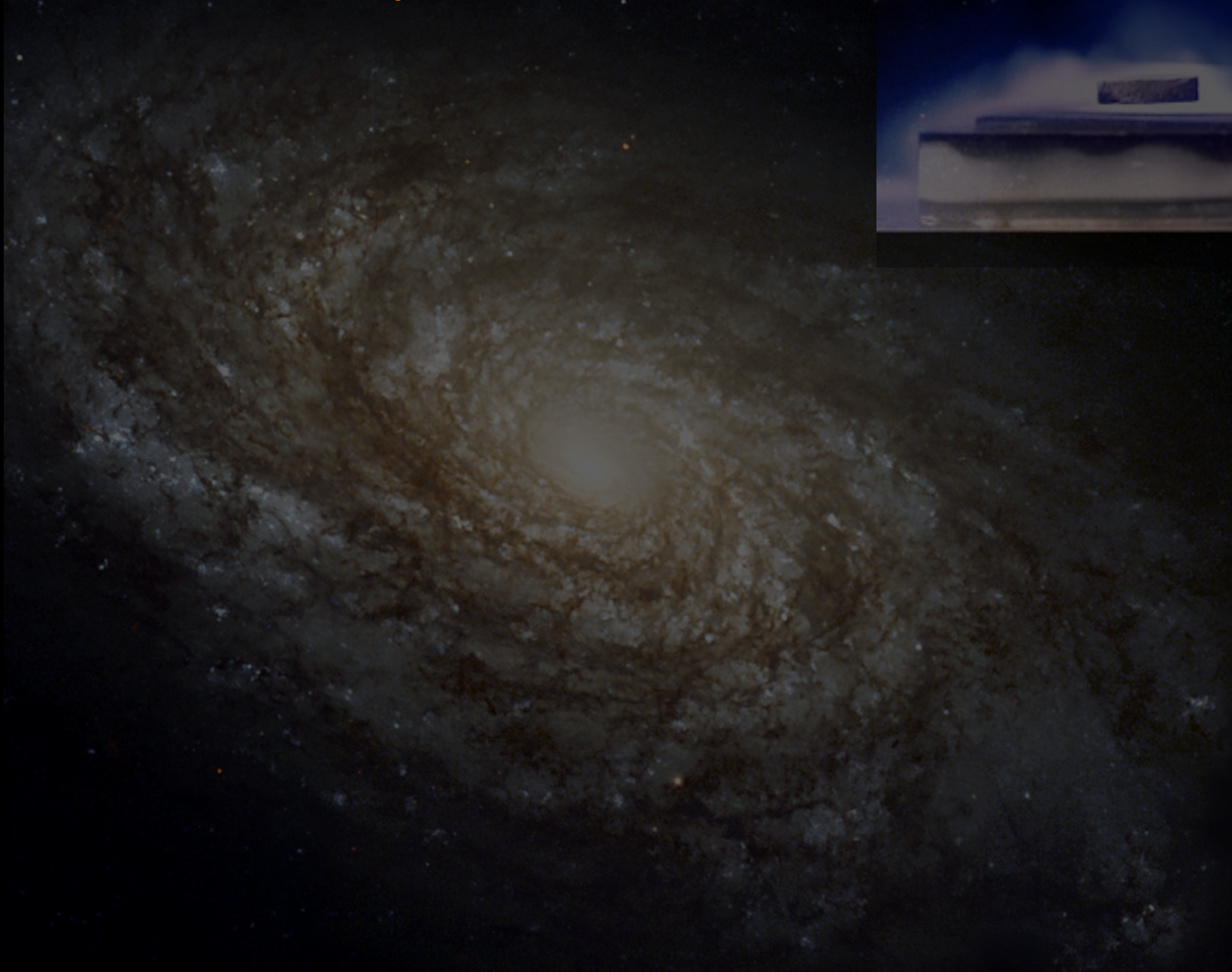


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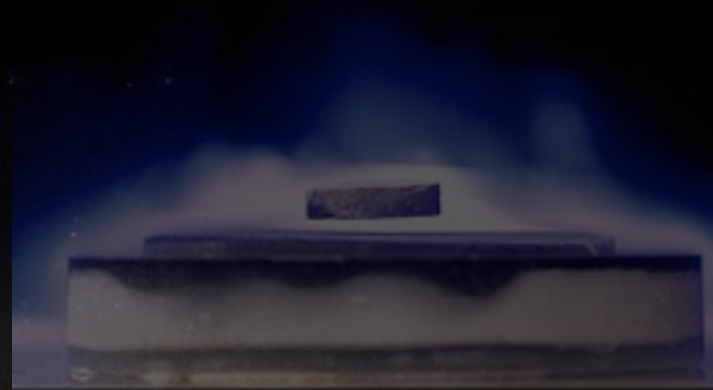
Fermi Liquid scattering parameters determined primarily by excitation of low-lying spinon and holon states..

Towards a new Conjunction of ideas.

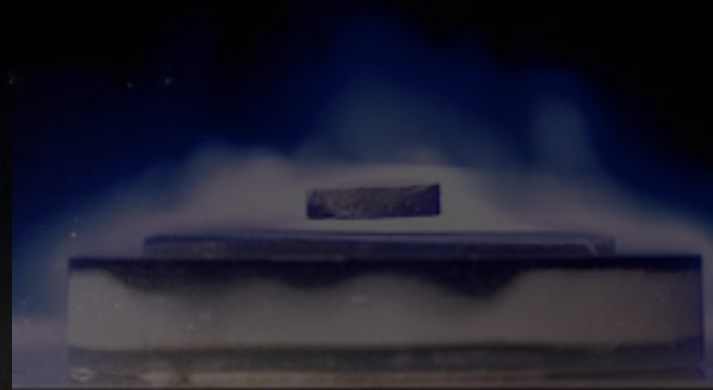


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“Non-Fermi liquids from holography” (large N /AdsCFT).

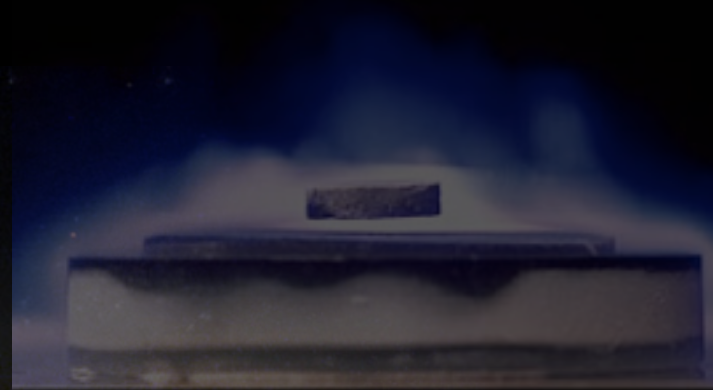


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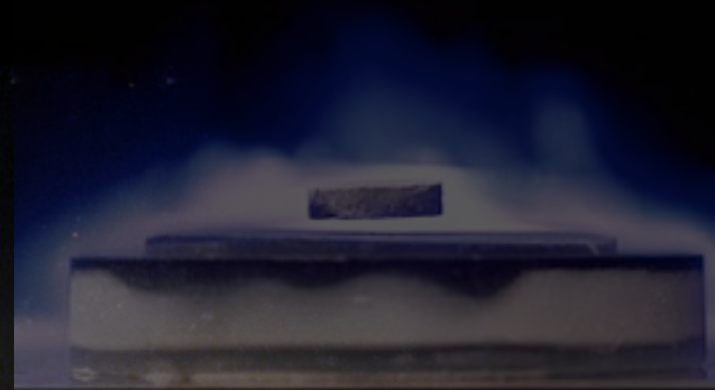
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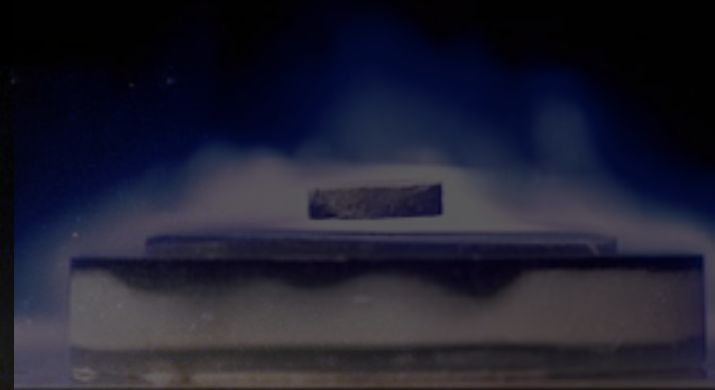
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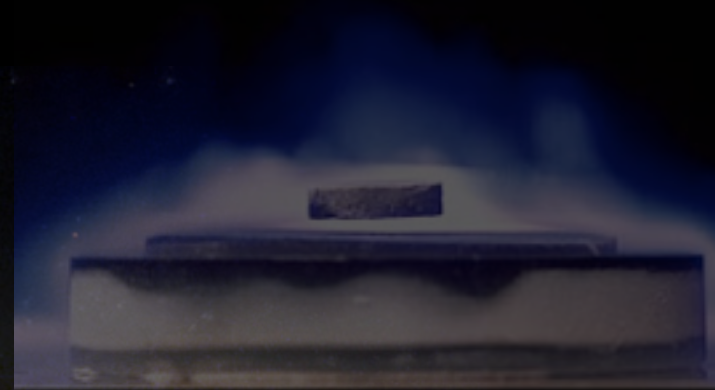
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Collaborators



Flint



Nev.



Dzero

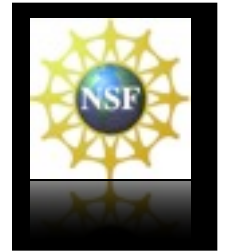


Rech

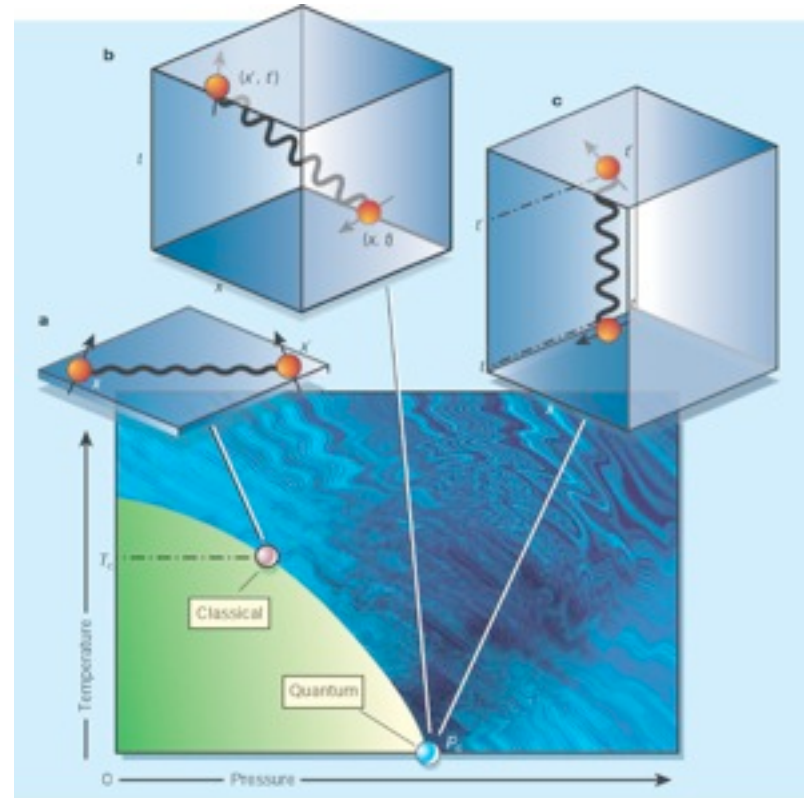


Lebanon

Rebecca Flint	Rutgers
Andriy Nevidomskyy	Rutgers
Maxim Dzero	Columbia/Rutgers
Jerome Rech,	ANL/Munich.
Eran Lebanon,	Israel.
Indranil Paul,	CNRS, Grenoble
Lucia Palova	Rutgers
Premi Chandra	Rutgers
Gergely Zarand	Budapest
Olivier Parcollet	SpHT Paris.
Andy Schofield	Birmingham
Qimiao Si	Rice, Houston
Catherine Pepin	SpHT Paris.
Almut Schroeder	Kent State
Gabriel Aeppli	LCN
Hilbert v. Lohneysen	Karlsruhe

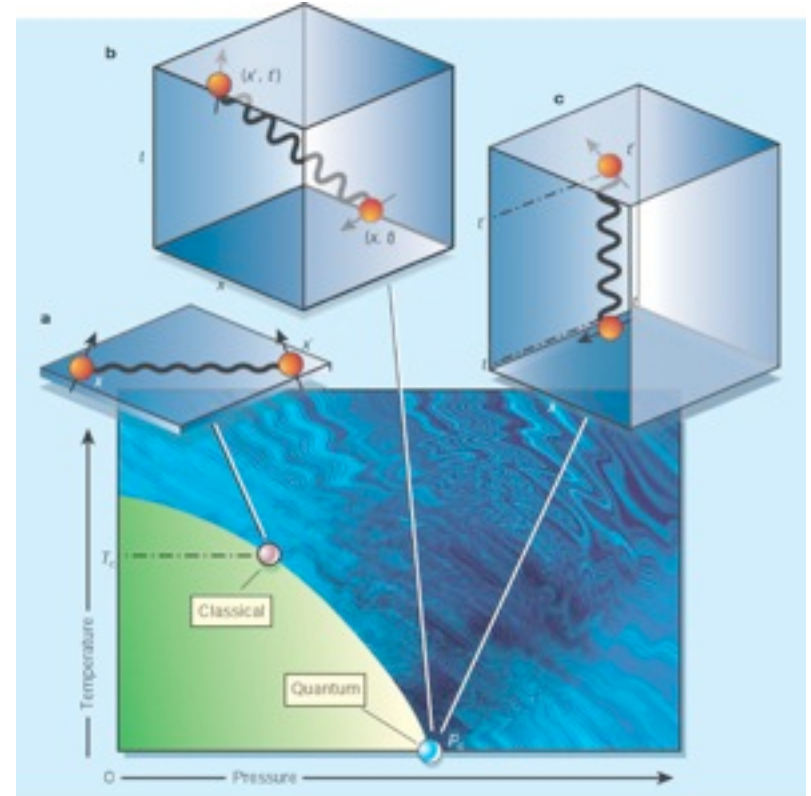


Conclusions



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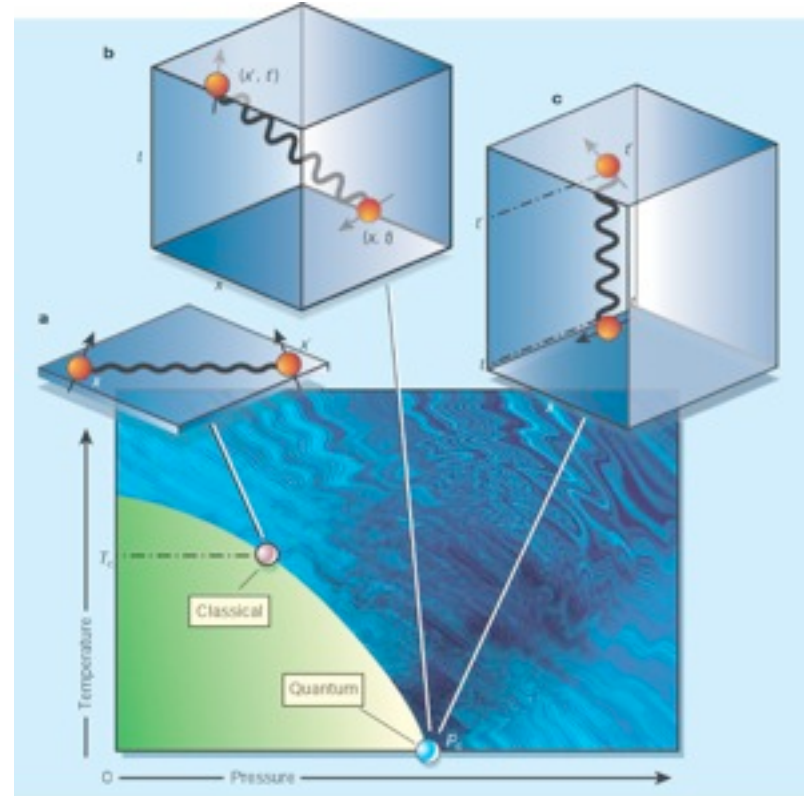
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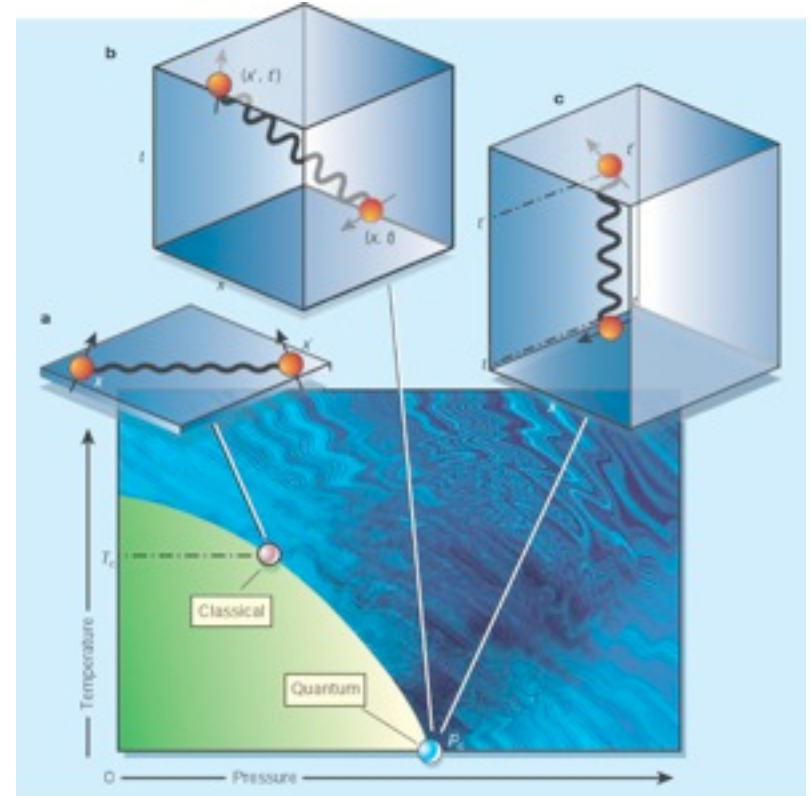


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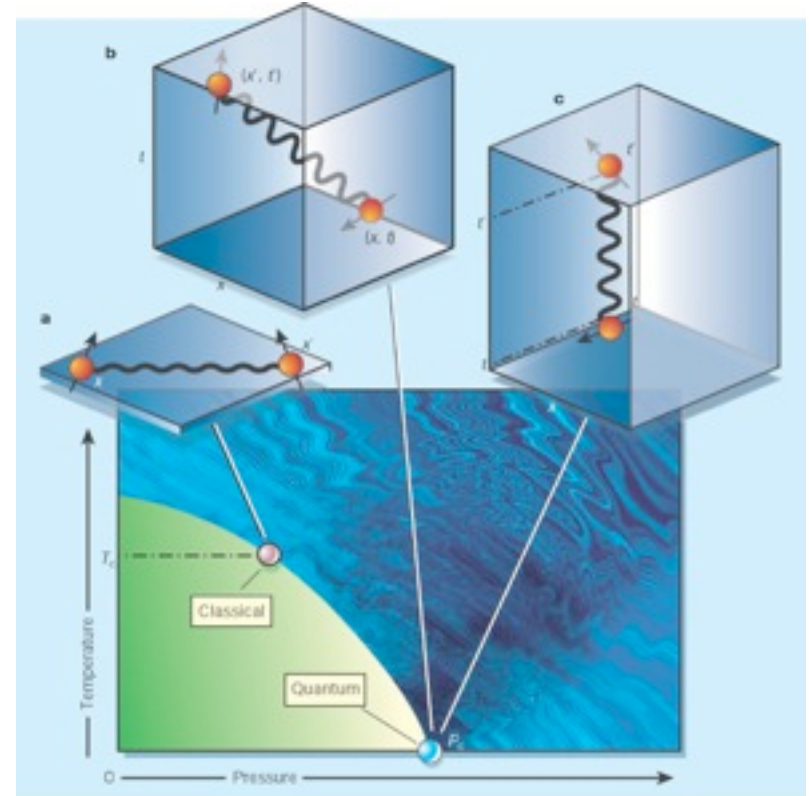
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New phases develop in order to avoid the
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