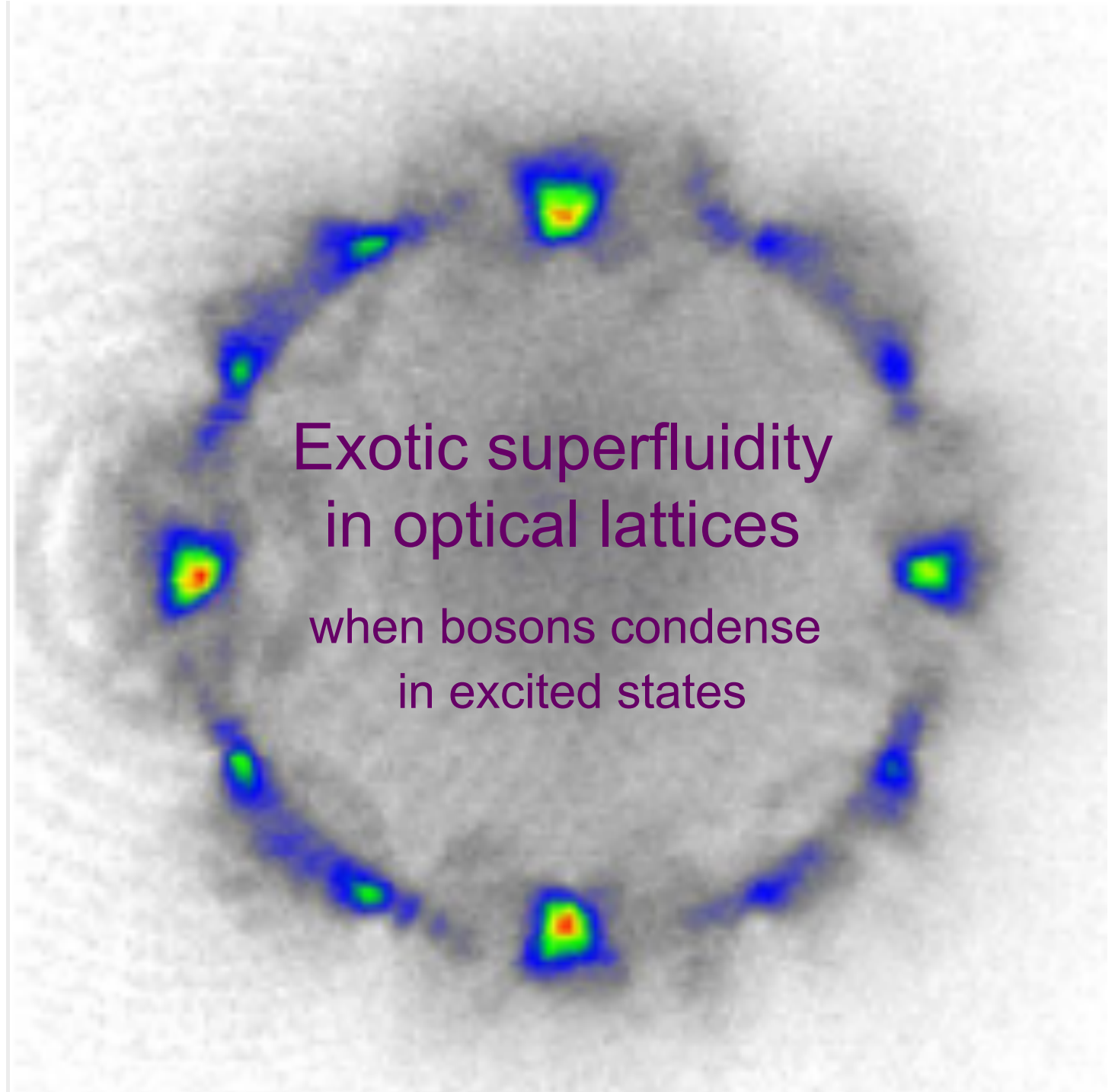




Universität Hamburg

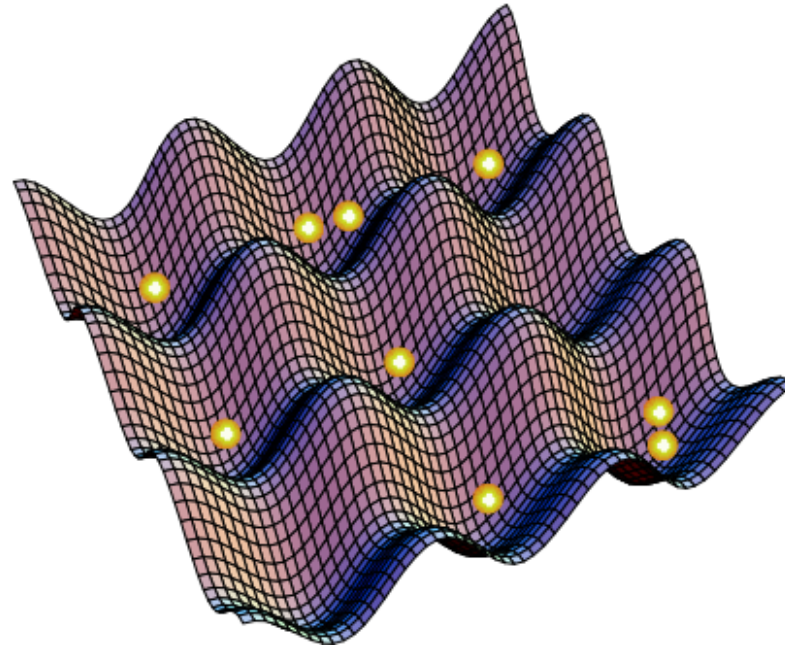


Andreas Hemmerich

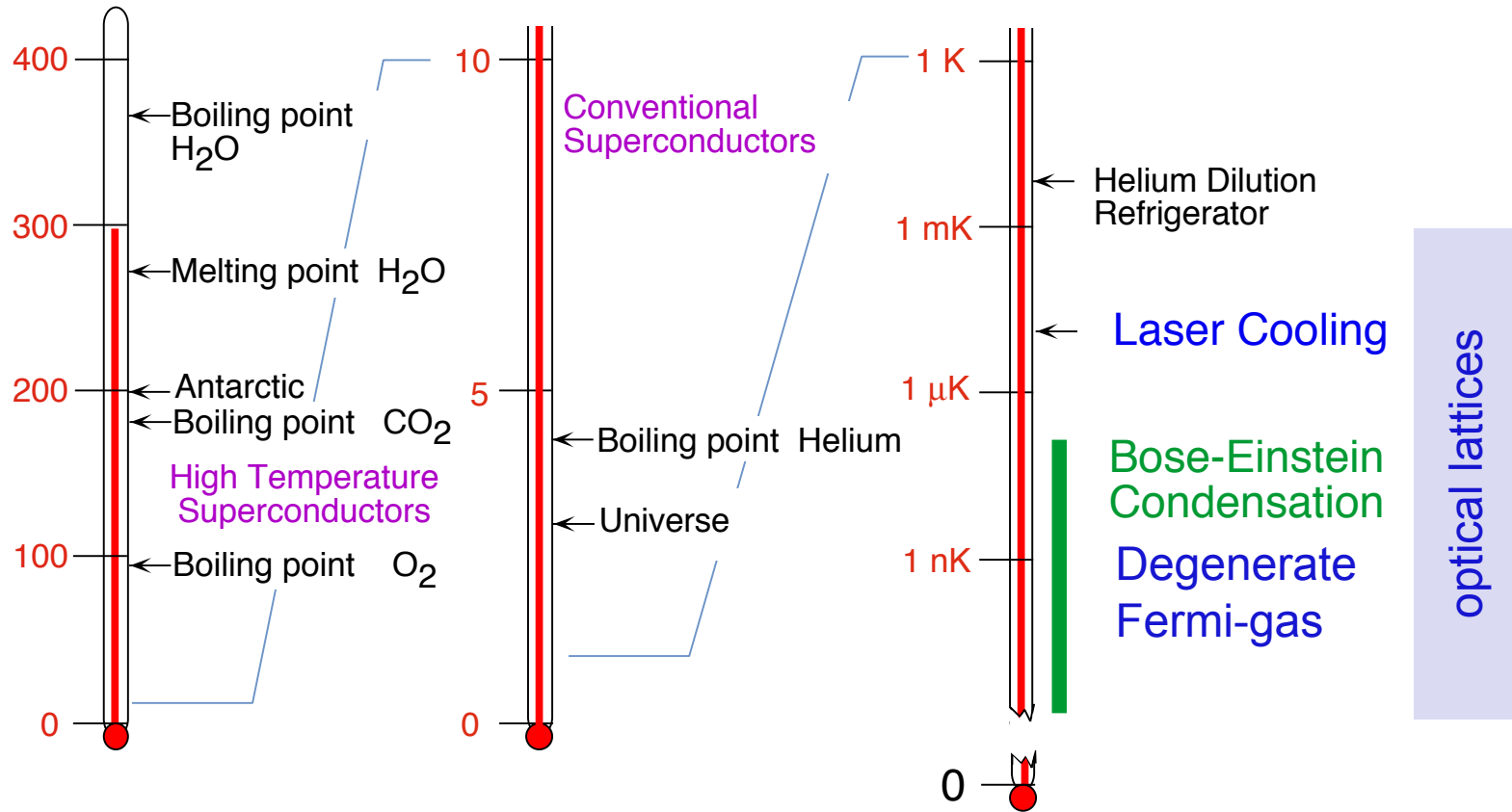


Optical lattice =

Ultracold atoms in a lattice made of light



# How cold is cold?



# Optical lattices, why bother ?

Study many-body phenomena at ultralow temperatures in synthetic lattices with

## extreme purity

no impurities or disorder unless designed  
bottom-up design of lattice systems

## extreme control

quantum statistic of particles: fermions, bosons, mixtures

interaction: attractive, repulsive, short range, long range

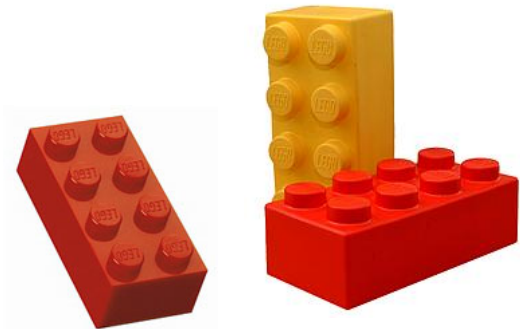
dimensionality: 1D, 2D, 3D

lattice geometry: any crystallographic class & quasiperiodic & disordered

unit cell: monopartite, multipartite

band structure: custom-tailored topology, orbital physics in higher bands

artificial gauge potentials: abelian and non-abelian





# Light shift potentials ...

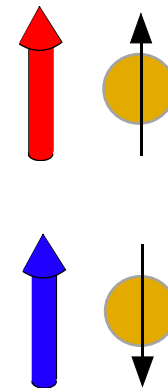
1968 Vladilen Lethokov: lattices of neutral atoms confined in the interference patterns of superimposed laser beams

sub-wavelength confinement → eliminate Doppler-effect

well depths only a few  $100 \mu\text{K}$  → need extremely efficient cooling

red detuned light → atoms are trapped in antinodes

blue detuned light → atoms are trapped in nodes

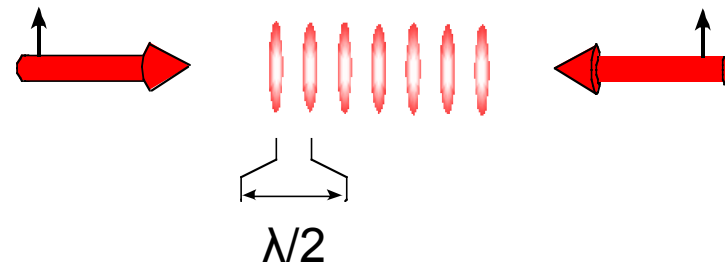


E-field      Induced  
dipole

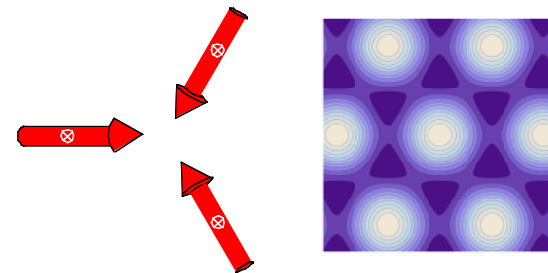
A. Ashkin, Phys. Rev. Lett. 40, 729, (1978)

G.P. Gordon & A. Ashkin, Phys. Rev. A 21, 1606, (1980)

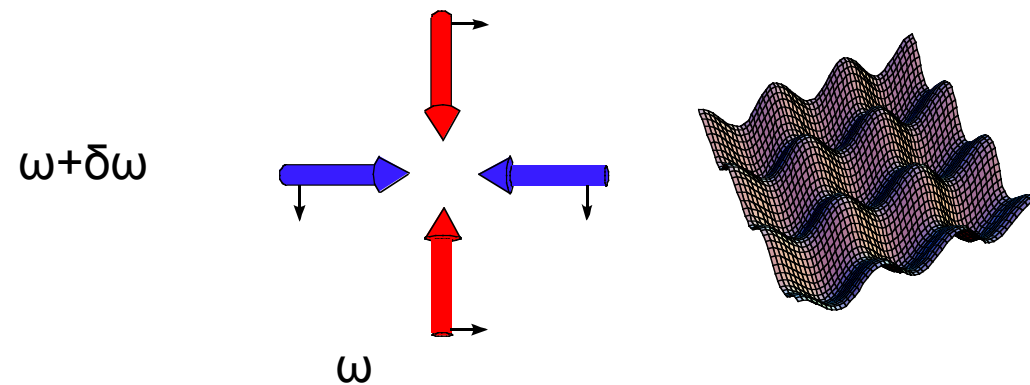
2 Beams  $\rightarrow$  1D lattice of disks



3 Beams  $\rightarrow$  2D triangular or hexagonal (cf. Graphene) lattice



4 Beams  $\rightarrow$  2D square lattice of tubes



Also: 3D, multipartite, quasi-periodic, disorder potentials, ...

# Lethokovs vision becomes reality ...

1980 - 1990: Development of laser cooling, Nobelprize 1997

1992: First 1D optical lattice of laser-cooled atoms



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29 JUNE 1992

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## **Dynamics and Spatial Order of Cold Cesium Atoms in a Periodic Optical Potential**

P. Verkerk, B. Lounis, C. Salomon, and C. Cohen-Tannoudji

*Collège de France et Laboratoire de Spectroscopie Hertzienne de l'Ecole Normale Supérieure, 24 rue Lhomond,  
F-75231 Paris, CEDEX 05, France*

J.-Y. Courtois and G. Grynberg

*Laboratoire de Spectroscopie Hertzienne de l'Ecole Normale Supérieure, Université Pierre et Marie Curie,  
F-75252 Paris, CEDEX 05, France*

(Received 23 March 1992)

VOLUME 69, NUMBER 1

PHYSICAL REVIEW LETTERS

6 JULY 1992

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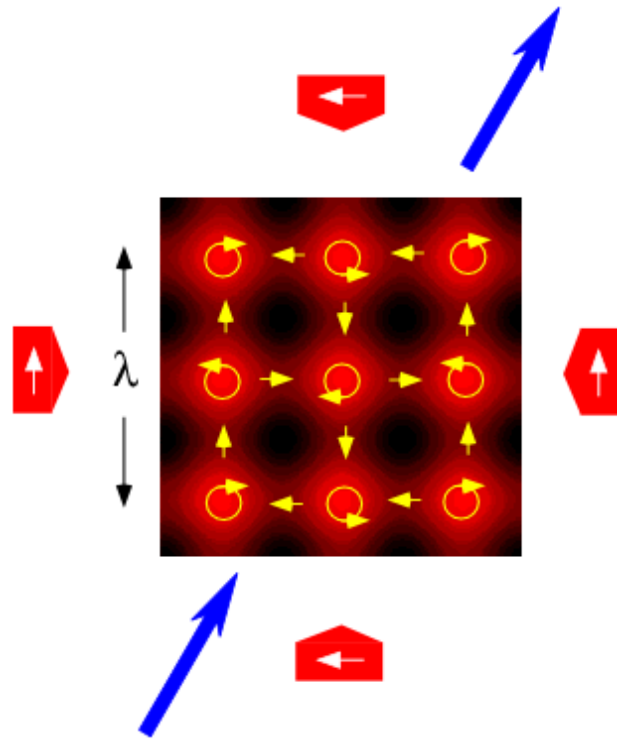
## **Observation of Quantized Motion of Rb Atoms in an Optical Field**

P. S. Jessen,<sup>(a)</sup> C. Gerz, P. D. Lett, W. D. Phillips, S. L. Rolston, R. J. C. Spreeuw, and C. I. Westbrook

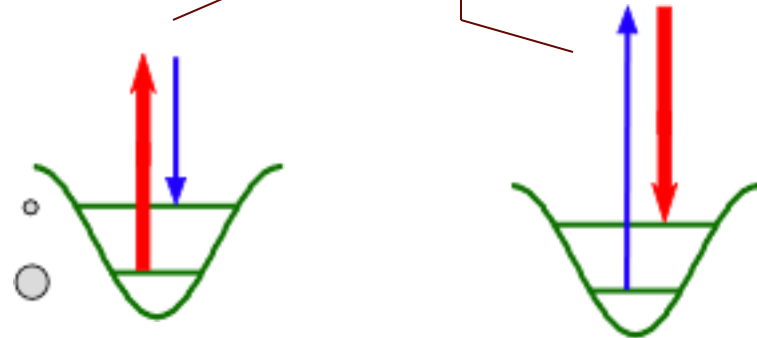
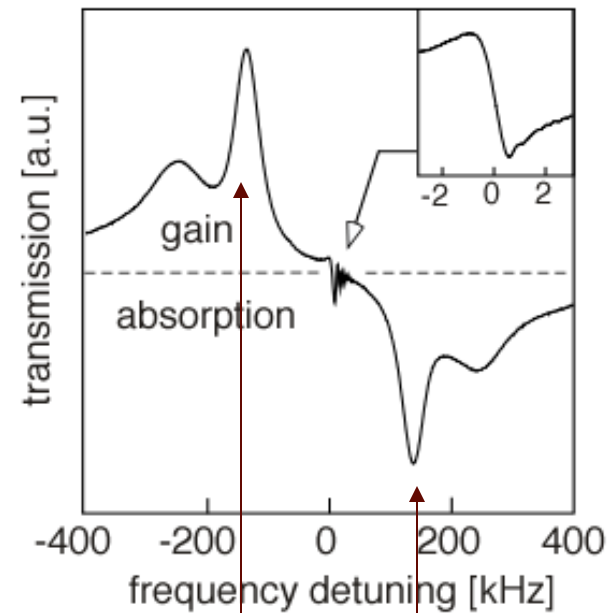
*National Institute of Standards and Technology, U.S. Department of Commerce, Technology Administration,  
PHYS A167, Gaithersburg, Maryland 20899*

(Received 6 May 1992)

# 1993: First 2D and 3D optical lattices with laser-cooled atoms



Occupation  $< 5\%$   $\rightarrow$  interaction negligible  
Deep wells  $\rightarrow$  tunneling negligible  
Array of isolated quantum oscillators  
with few vibrational levels occupied  
Temperature = few  $\mu\text{K}$



A. Hemmerich & T. Hänsch, PRL 70, 410 (1993)

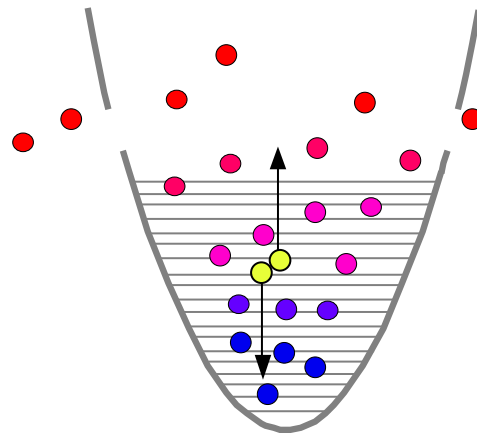
# Cooling beyond the limits of laser cooling

Cooling by evaporation:



H. Hess, Phys. Rev. B, 34, 3476 (1986).

trap potential



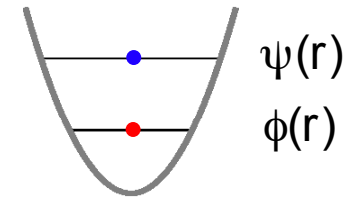
Below critical temperature: quantum degeneracy  
bosons or fermions ?

# Bosons & Fermions

In quantum mechanics, particles are described by wave functions

Distinguishable particles

$$\chi(r_1, r_2) = \phi(r_1)\psi(r_2)$$



No individualism for quantum particles of same sort

Indistinguishable particles

$$\chi(r_1, r_2) = \phi(r_1)\psi(r_2) \pm \psi(r_1)\phi(r_2) \quad \Rightarrow \quad \chi(r_1, r_2) = \pm \chi(r_2, r_1)$$

+ **Bosons**  $\phi = \psi \Rightarrow \chi(r_1, r_2) \propto \phi(r_1)\phi(r_2)$

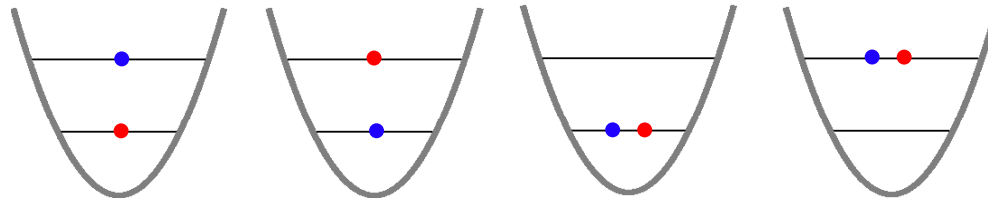
- **Fermions**  $\phi = \psi \Rightarrow \chi(r_1, r_2) = 0$

Pauli principle:

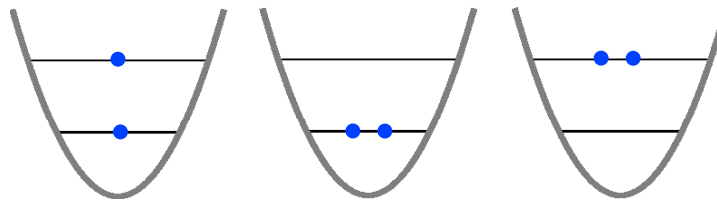
Fermions cannot share the same single particle state

# Fermions avoid each other but bosons like to be together

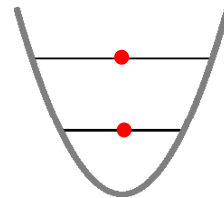
Distinguishable particles



Bosons

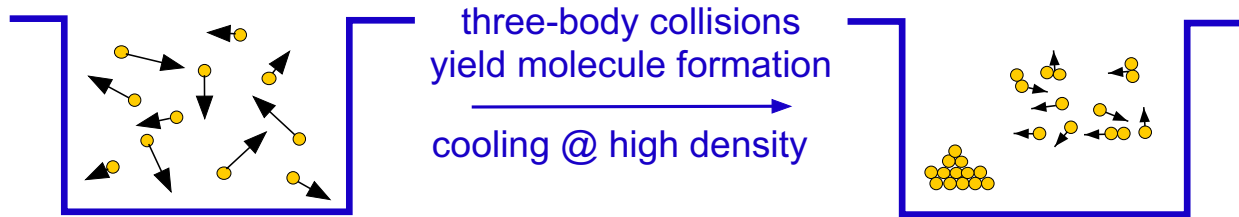


Fermions

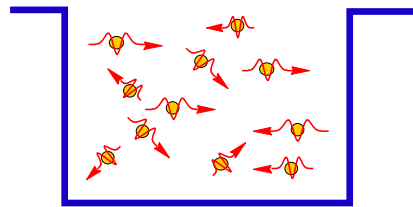




# Bosons condense in the ground state



↓ cooling @ low density



$$k_B T = \frac{1}{2} m \bar{V}^2$$

$$m \bar{V} = \hbar \frac{2\pi}{\Lambda}$$

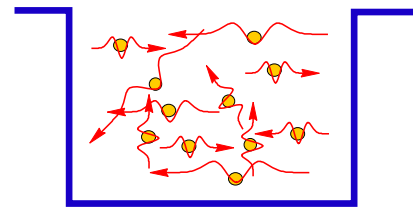
⇒

Thermal de Broglie Wavelength:

$$\Lambda = \left( \frac{\hbar^2}{2m k_B T} \right)^{1/2}$$



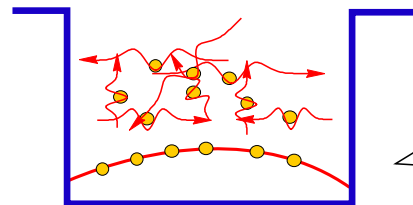
Luis Victor de Broglie



Critical Temperature:

de Broglie Wavelength  $\Lambda \approx$  mean particle separation  $d \approx 1/\rho^{1/3}$

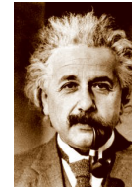
$$T = \frac{\hbar^2 \rho^{2/3}}{2m k_B}$$



← Bose-Einstein-Condensate



Satyendra Nath Bose



Albert Einstein

1995 Bose-Einstein condensation of ultracold dilute gases  
Nobelprize 2001



■ **REPORTS** ■ SCIENCE • VOL. 269 • 14 JULY 1995

**Observation of Bose-Einstein Condensation  
in a Dilute Atomic Vapor**

M. H. Anderson, J. R. Ensher, M. R. Matthews, C. E. Wieman,\*  
E. A. Cornell

PHYSICAL REVIEW  
LETTERS

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VOLUME 75

27 NOVEMBER 1995

NUMBER 22

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**Bose-Einstein Condensation in a Gas of Sodium Atoms**

K. B. Davis, M.-O. Mewes, M. R. Andrews, N. J. van Druten, D. S. Durfee, D. M. Kurn, and W. Ketterle  
*Department of Physics and Research Laboratory of Electronics, Massachusetts Institute of Technology,  
Cambridge, Massachusetts 02139*  
(Received 17 October 1995)

1999 quantum degenerate Fermi-gases: D. Jin 's group, JILA

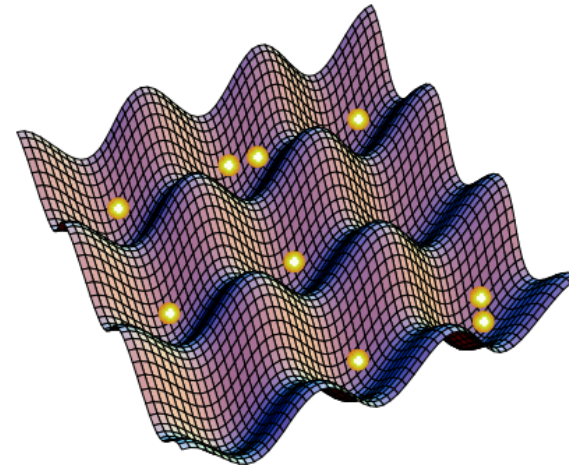
# Quantum gases in optical lattices

load quantum degenerate gas (Bose, Fermi, Mixture)  
into periodic light shift potentials

- very low temperatures ( $< 100$  nK)
- multiply occupied lattice sites

## relevant physical mechanisms

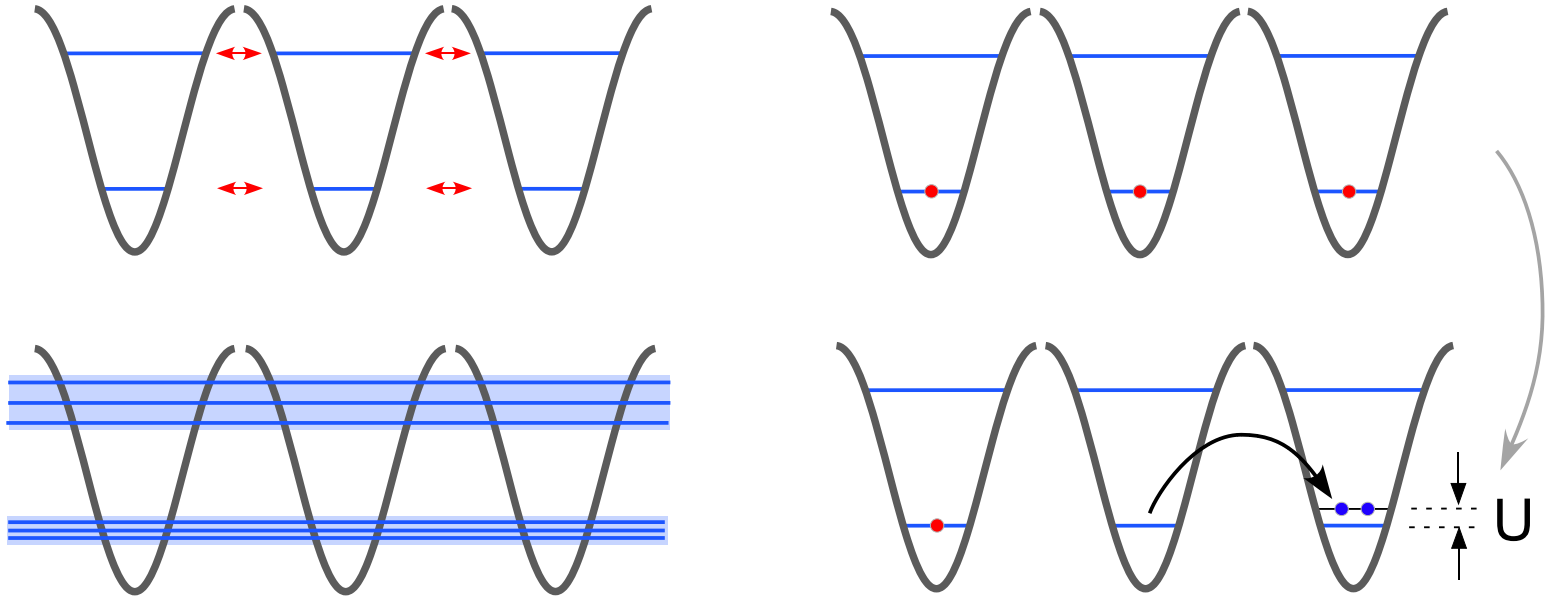
- quantum statistics
- tunneling
- collisional (on-site) interactions



construction kit for synthesising strongly correlated many body systems  
in perfectly controlled environments

- study quantum phase transitions
- study non-equilibrium many-body dynamics

# Competition of tunneling and (repulsive, on-site) collisions



Tunneling yields allowed energy bands separated by forbidden gaps  
Atoms are delocalized across the entire lattice

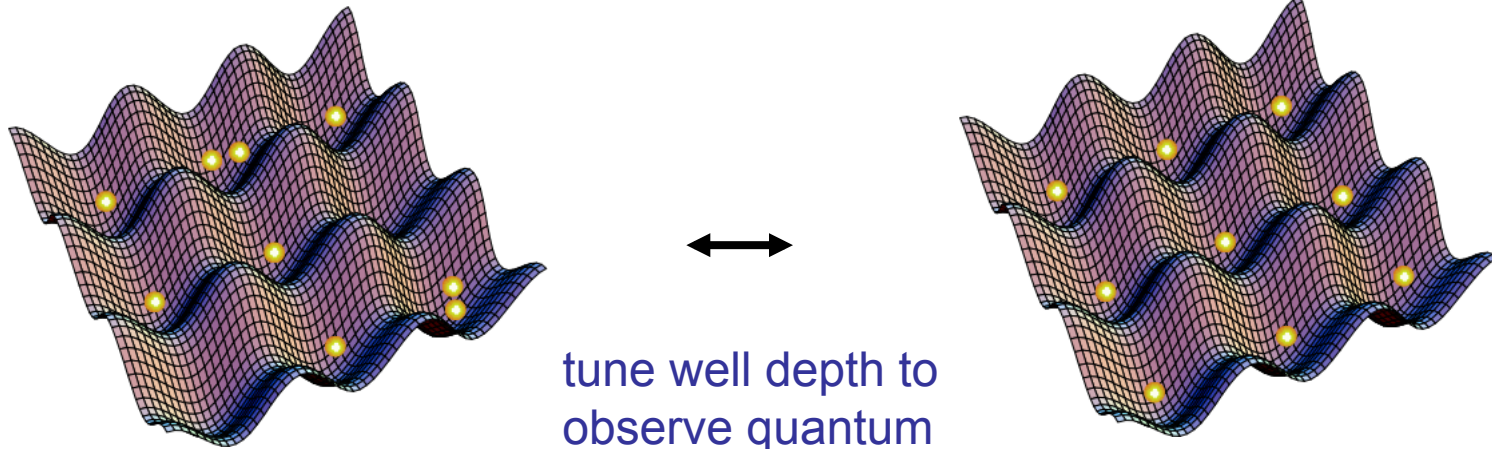
Double occupancy costs Interaction energy  $U$

# Optical Lattice with Bosons

in the ground state of the lowest band

Emulation of bosonic Hubbard-model

M. Fisher et al., Phys. Rev. B 40, 546 (1989)



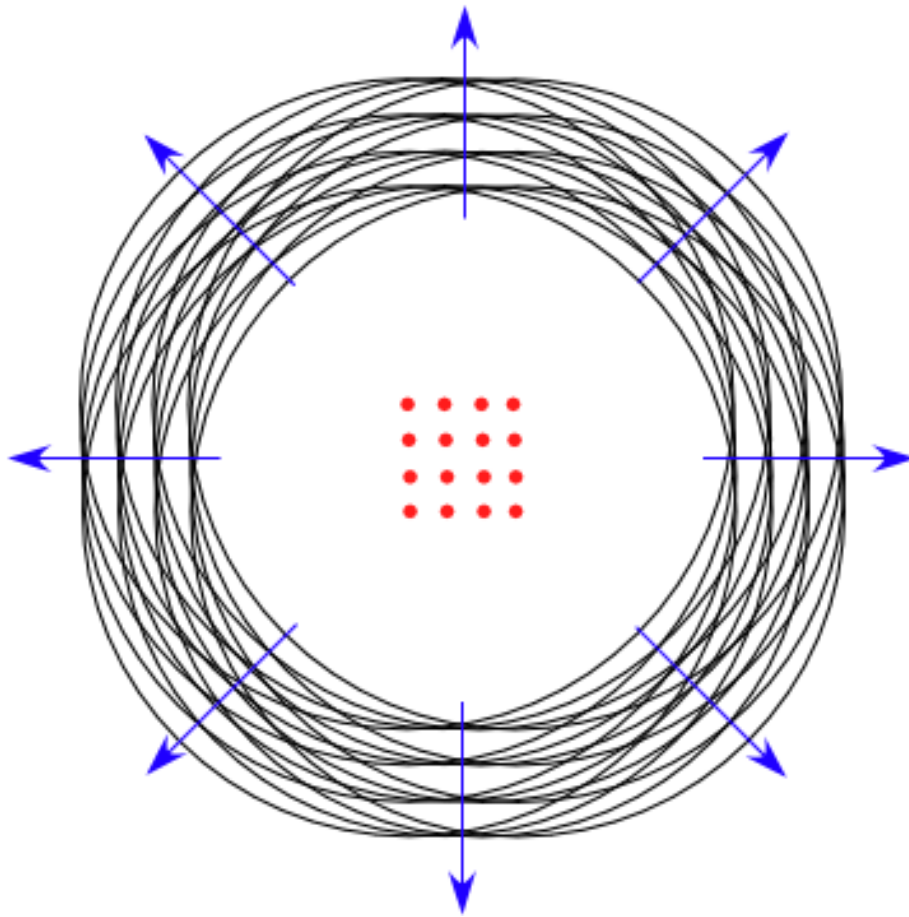
## superfluid

Shallow wells  
Tunneling dominates  
Poissonian particle number fluctuations  
Defined relative phases  
→ coherence

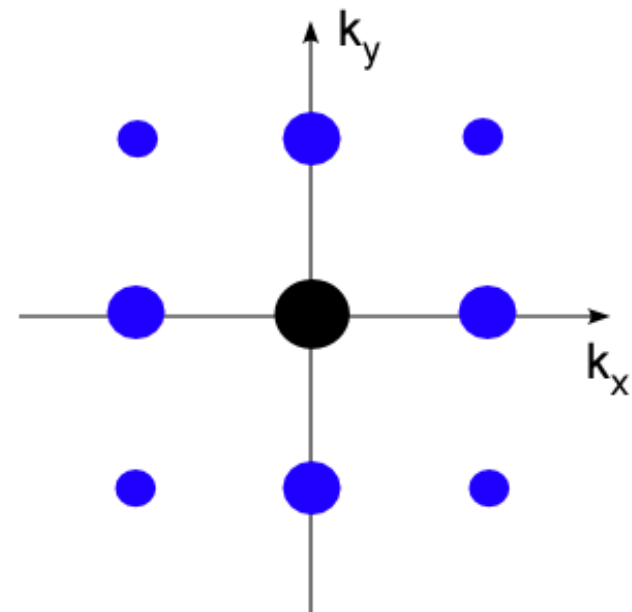
## Mott insulator

Deep wells  
Collisional repulsion dominates  
Fixed local particle number  
Relative phases undetermined  
→ no coherence

# Localized Bragg maxima indicate coherence



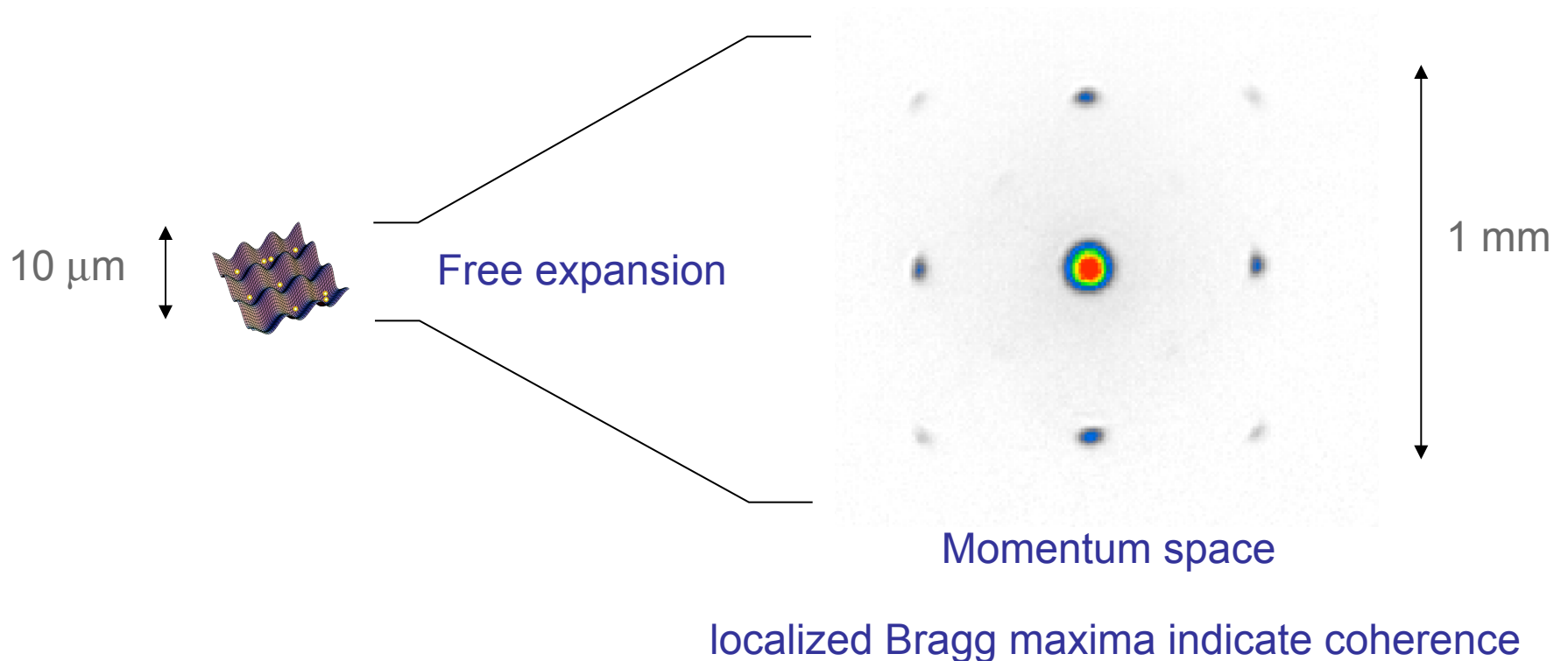
Configuration space



Momentum space

# Observation of momentum spectra

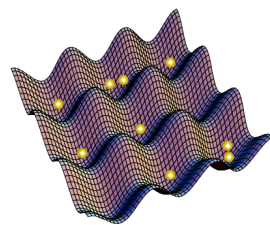
- lattice potential **rapidly** ( $1 \mu\text{s}$ ) switched off
- free expansion (typically 30 ms)
- absorption image



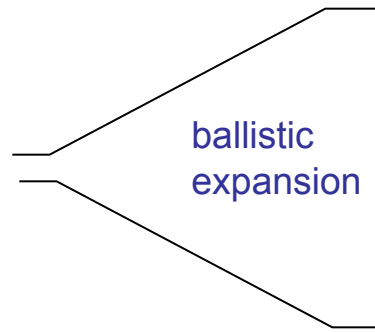


# Observation of quantum phase transition

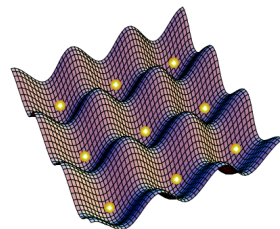
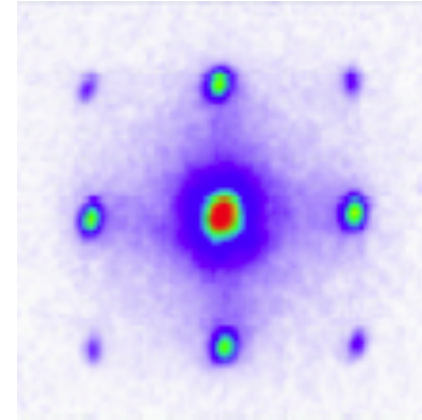
Image momentum space via ballistic expansion experiment



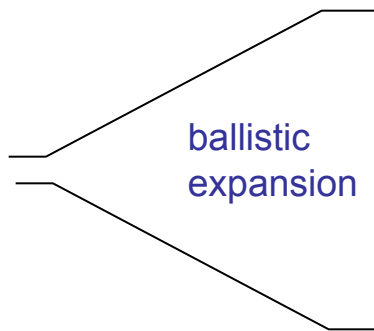
superfluid



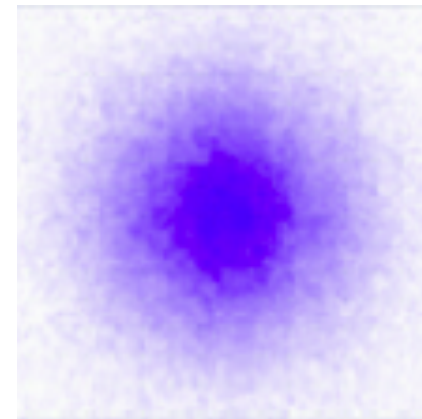
ballistic  
expansion



Mott-Insulator



ballistic  
expansion

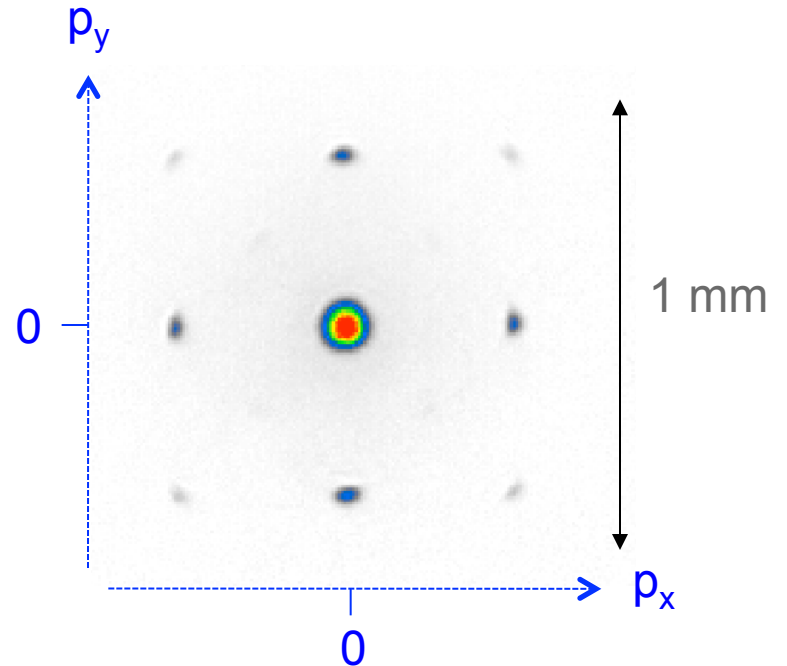
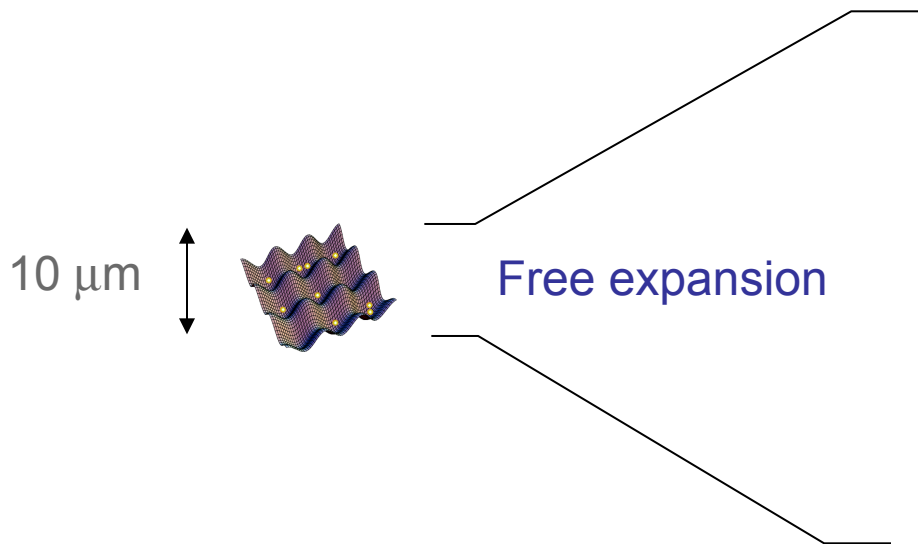
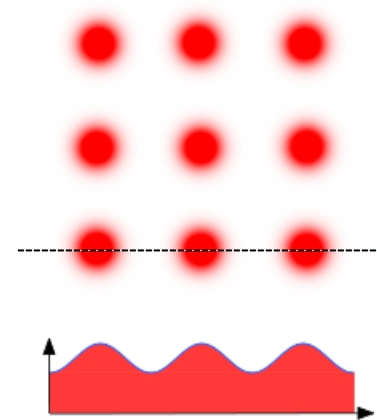


# Standard optical lattice with bosons

In standard bosonic optical lattices atoms are condensed in the ground state of lowest band

- local s-orbital at each site  
ground state wave function positive definite  
(cf. Feynmans no node theorem for bosons)

Lowest order Bragg peak at zero momentum



# More advanced lattices with bosons?

Can we get closer to scenarios of electronic condensed matter?

magnetism & superconductivity in rare earth or transition metal compounds, Kondo-systems, quantum Hall systems, . . .

- tailored artificial gauge fields for ground state lattices, dynamical lattices
- orbital degrees of freedom in optical lattices excited to higher bands

Incomplete list of pioneering theory groups  
(alphabetic order)

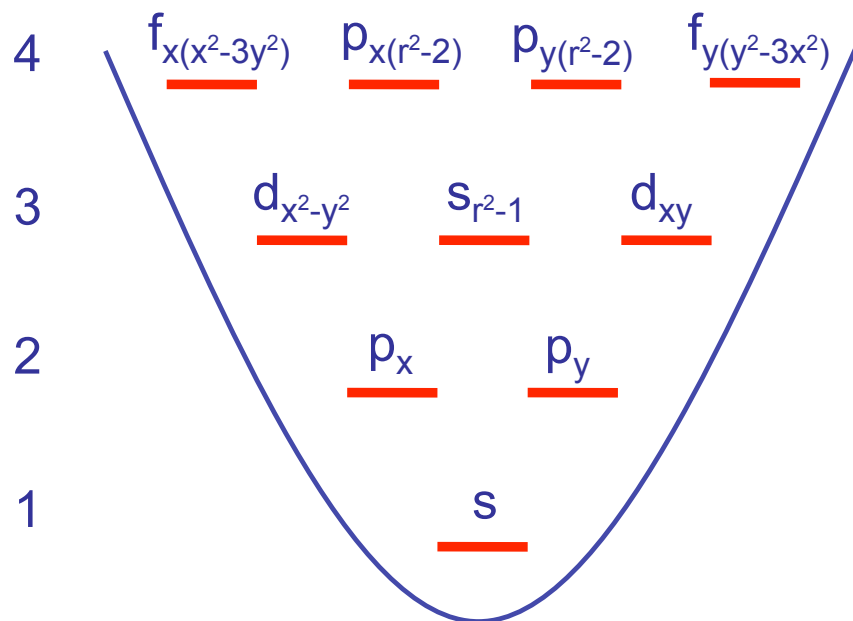
S. Girvin, J. Larson, J.-P. Martikainen, M. Lewenstein,  
W. V. Liu, S. Das Sarma, C. Wu, ...

# Orbital Optical Lattices

New physics in higher bands

→ anisotropic orbits

→ freedom of orientation



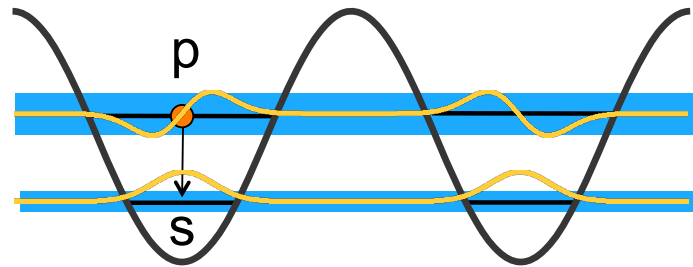
harmonic oscillator orbits (two dimensions)

# Challenges

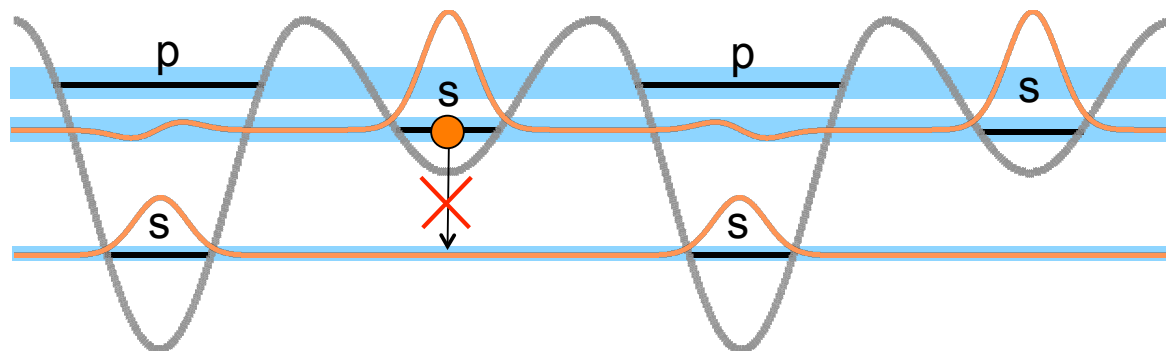
- exciting higher bands with moderate entropy increase
- providing cross-dimensional tunneling
- control of band relaxation

# Control of band relaxation

simple lattice → band life time limited by collisions

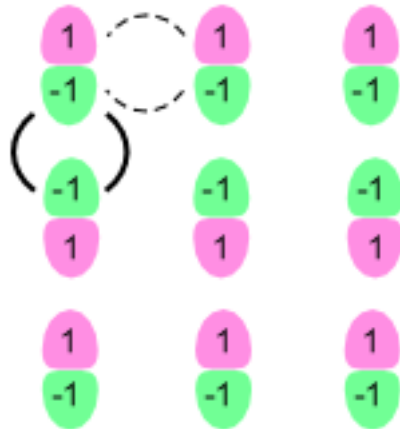


bipartite lattice → band life times up to several 100 ms

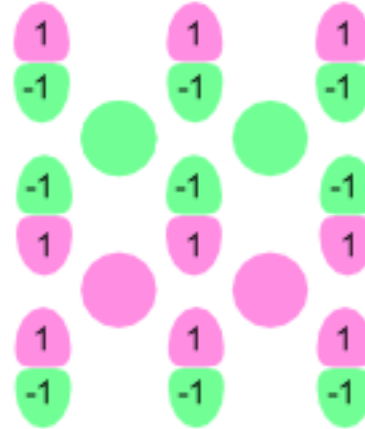


# Providing cross-dimensional coherence

monopartite lattice



bipartite lattice

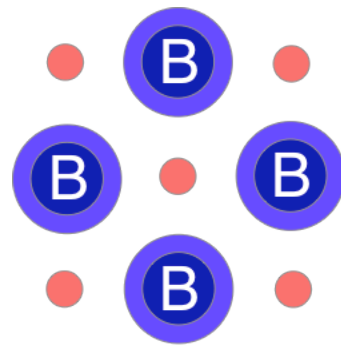


A. Isacsson and S. Girvin, Phys. Rev. A 72, 053604 (2005).  
W. Vincent Liu and C. Wu, Phys. Rev. A 74, 013607 (2006).

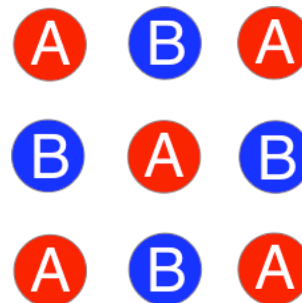


# Chequerboard Lattice

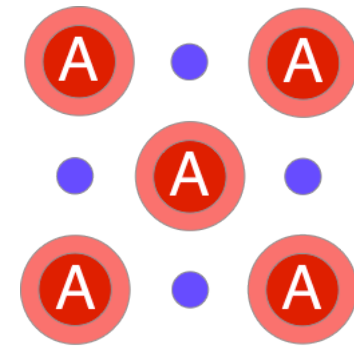
Energy difference  $\Delta V$  between A- and B-sites can be tuned



$$\Delta V < 0$$

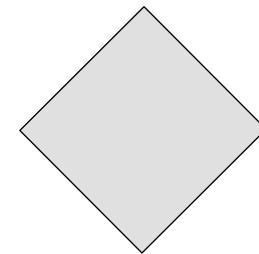
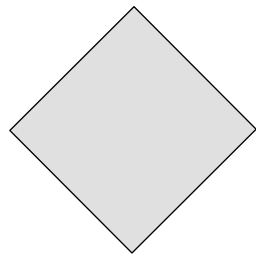


$$\Delta V = 0$$

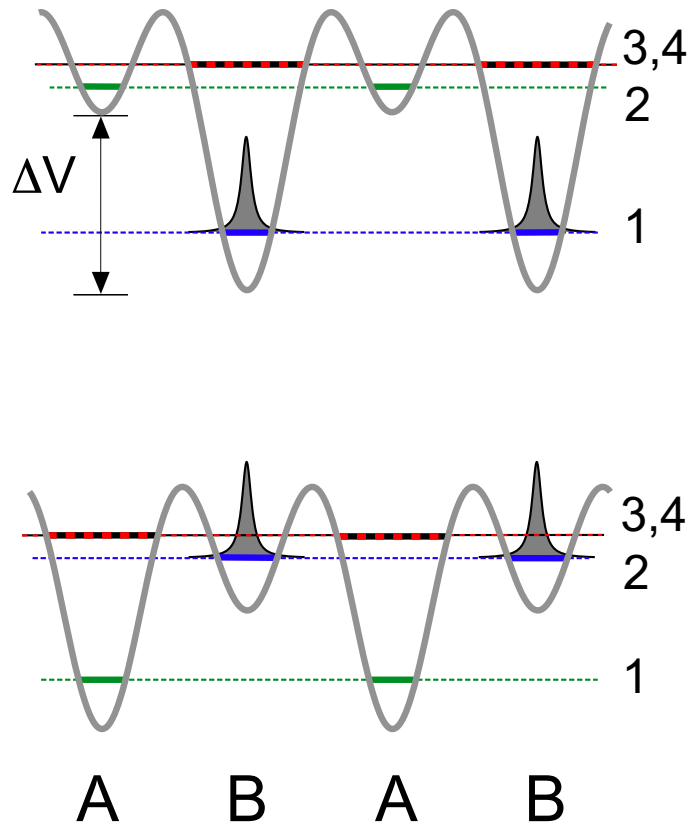


$$\Delta V > 0$$

Wigner-Seitz  
unit cells:



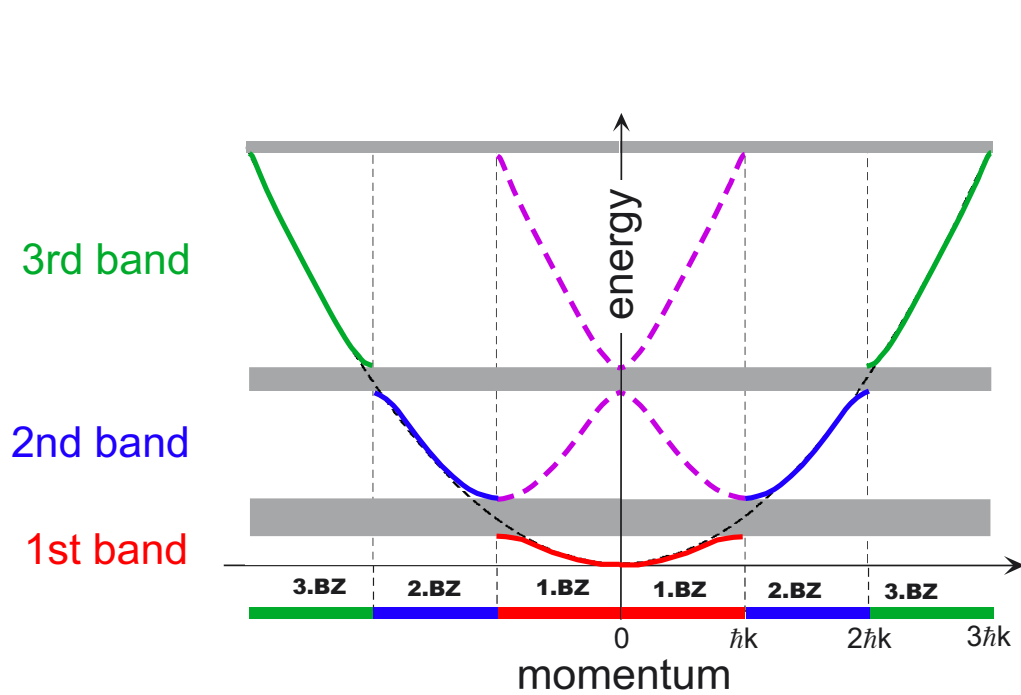
# Populating higher bands



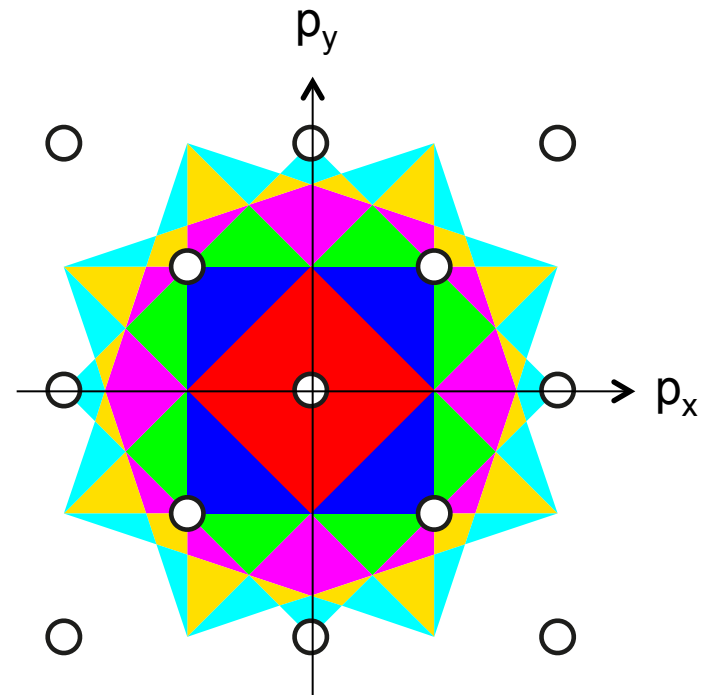
$\ll$  tunneling time

$\gg$  on-site oscillation period

# Observing atoms in higher bands



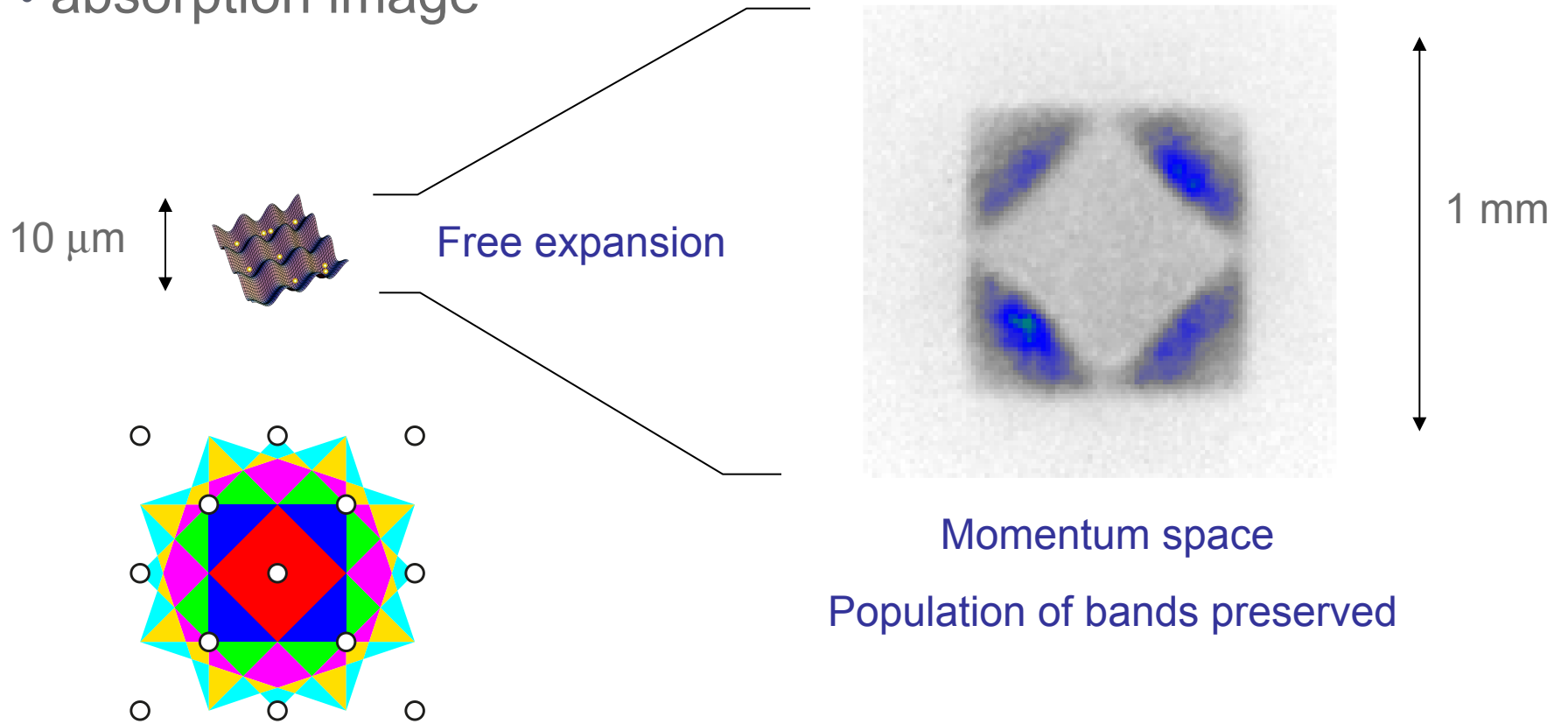
1 dimension



2 dimensions

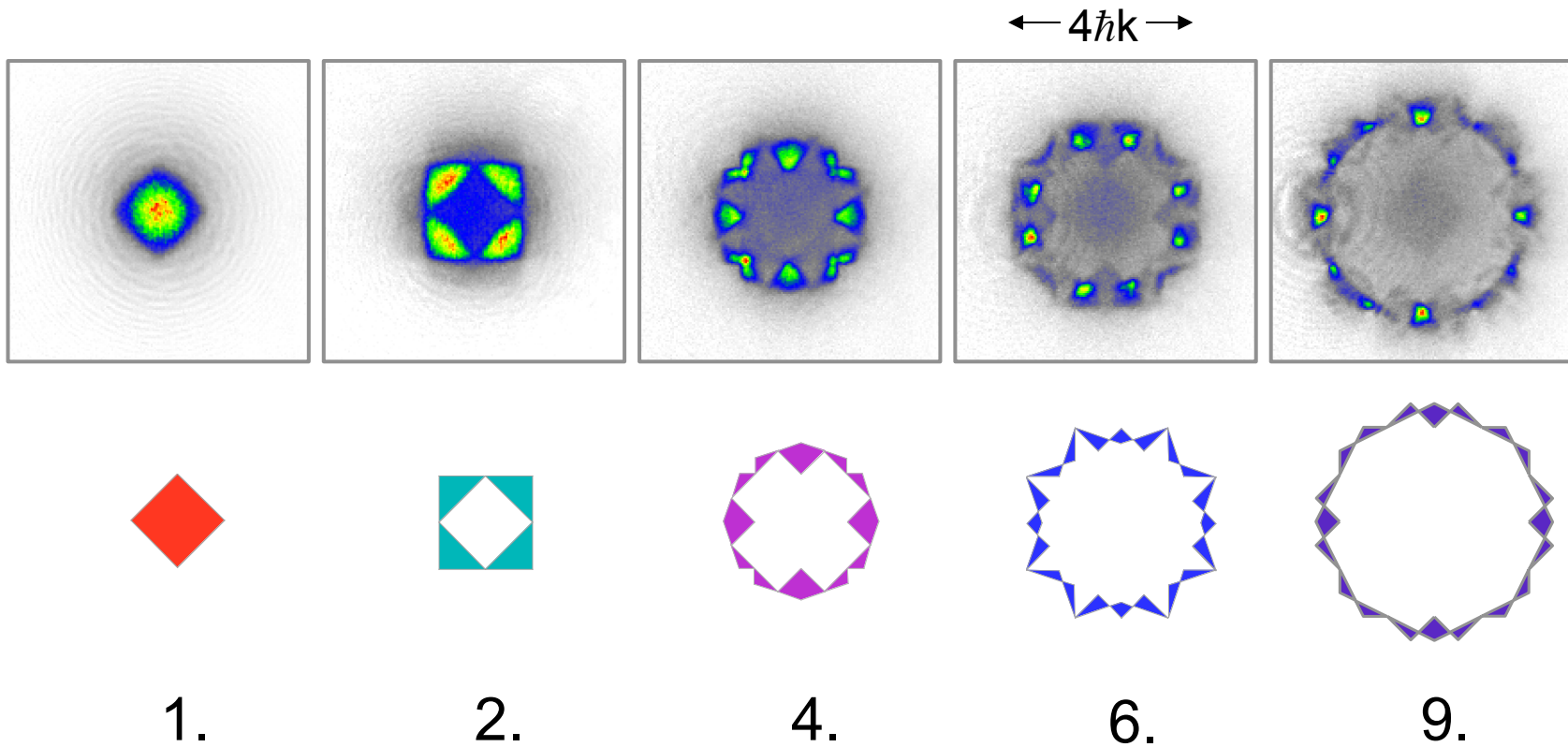
# Band mapping

- lattice potential **slowly** ( $> 400 \mu\text{s}$ ) switched off
- free expansion for 30 ms
- absorption image



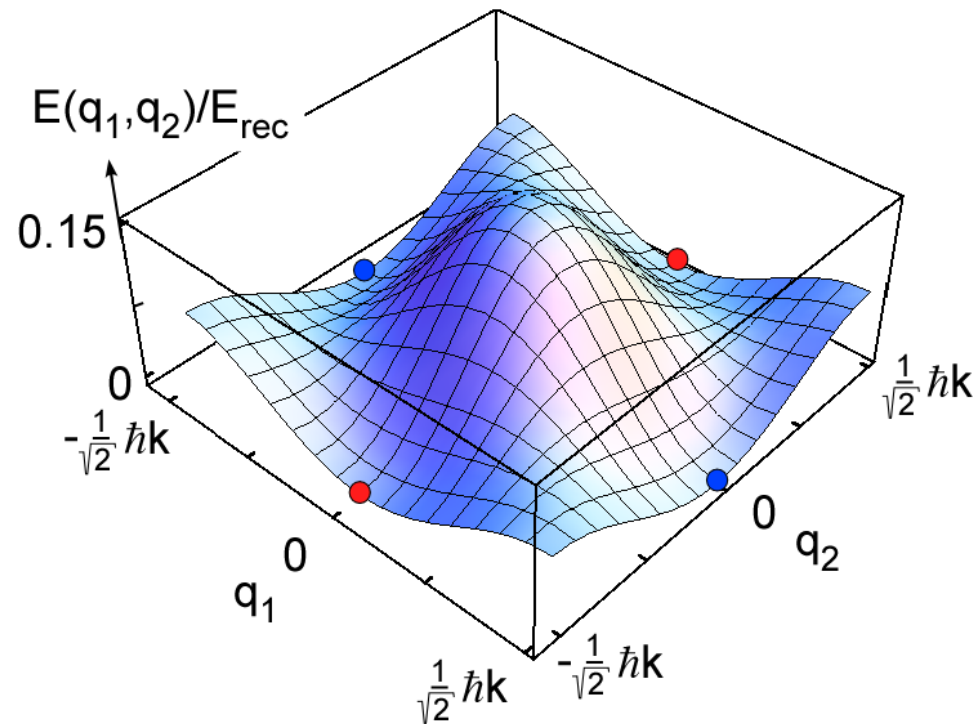
# Addressing different bands

G. Wirth, M. Ölschläger, A. Hemmerich, Nature Physics 7, 147 (2011)



2nd band  $(2P_x \pm i 2P_y)$

2nd band  $\rightarrow$  two inequivalent minima



Lattice with perfect  $90^\circ$  rotation symmetry:  
energy difference between band minima  $\Delta E = 0$ .  
Uniaxial distortion:  $\Delta E \neq 0 \rightarrow$  degeneracy lifted

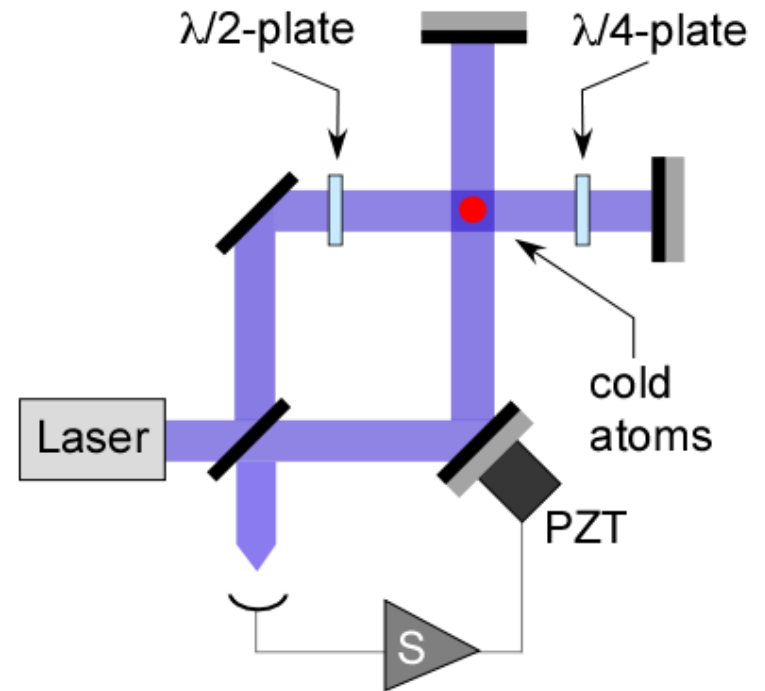
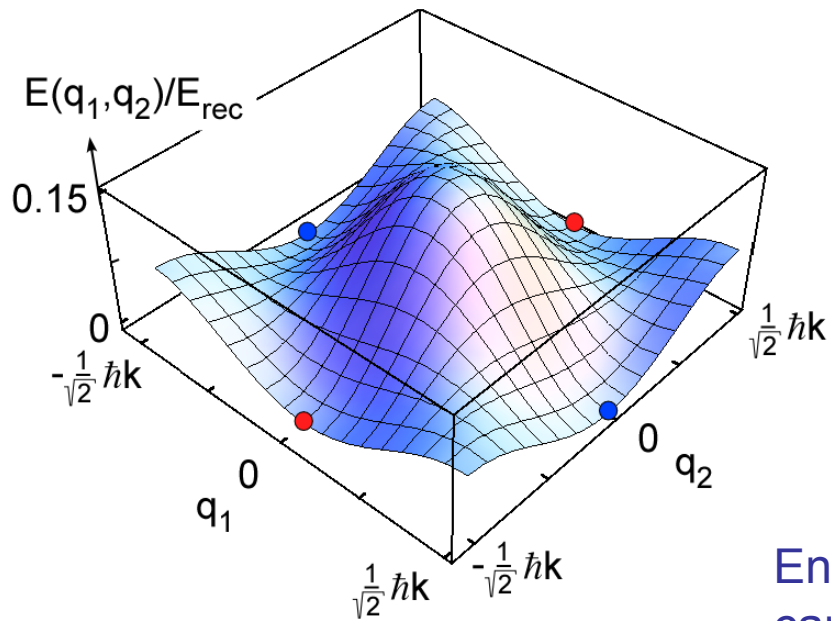


# Uniaxial distortion can be controlled in experiment

$$V_{real}(x,y) = -\frac{V_0}{4} \left| \eta \left[ e^{ikx} + \varepsilon_x e^{-ikx} \right] + e^{i\theta} \left[ e^{iky} + \varepsilon_y e^{-iky} \right] \right|^2$$

$\eta$  unequal intensities coupled to x and y directions

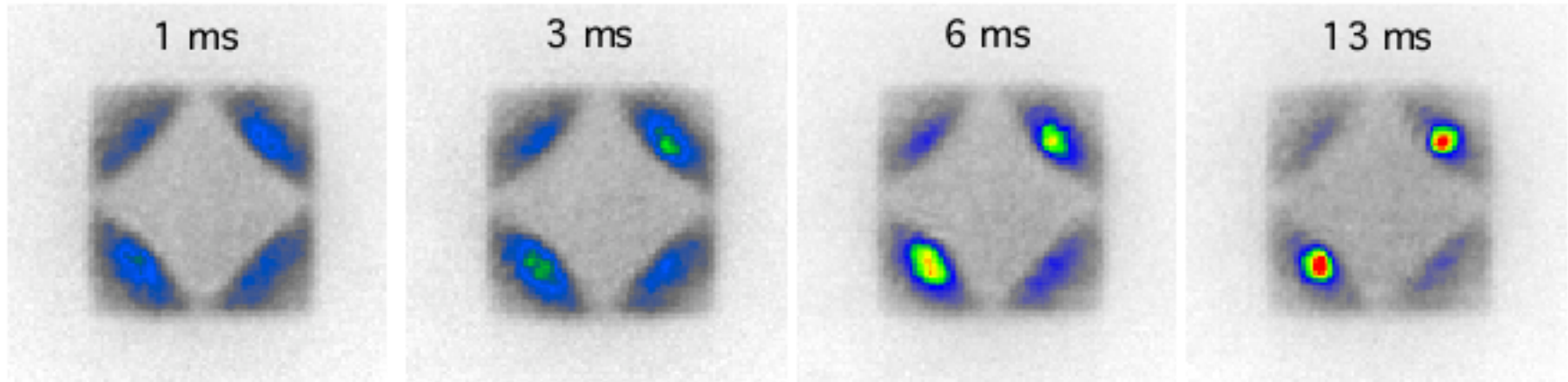
$\varepsilon_x, \varepsilon_y$  imperfect reflection of beams coupled to x and y directions



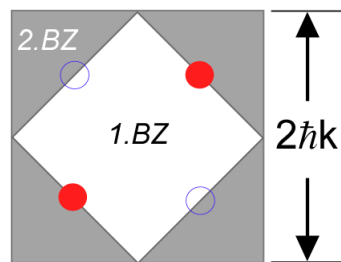
Energy difference  $\Delta E$  between P-band minima can be tuned via  $\varepsilon_y$

# Time evolution of band population: $\Delta E \neq 0$

Condensation at lowest band minimum

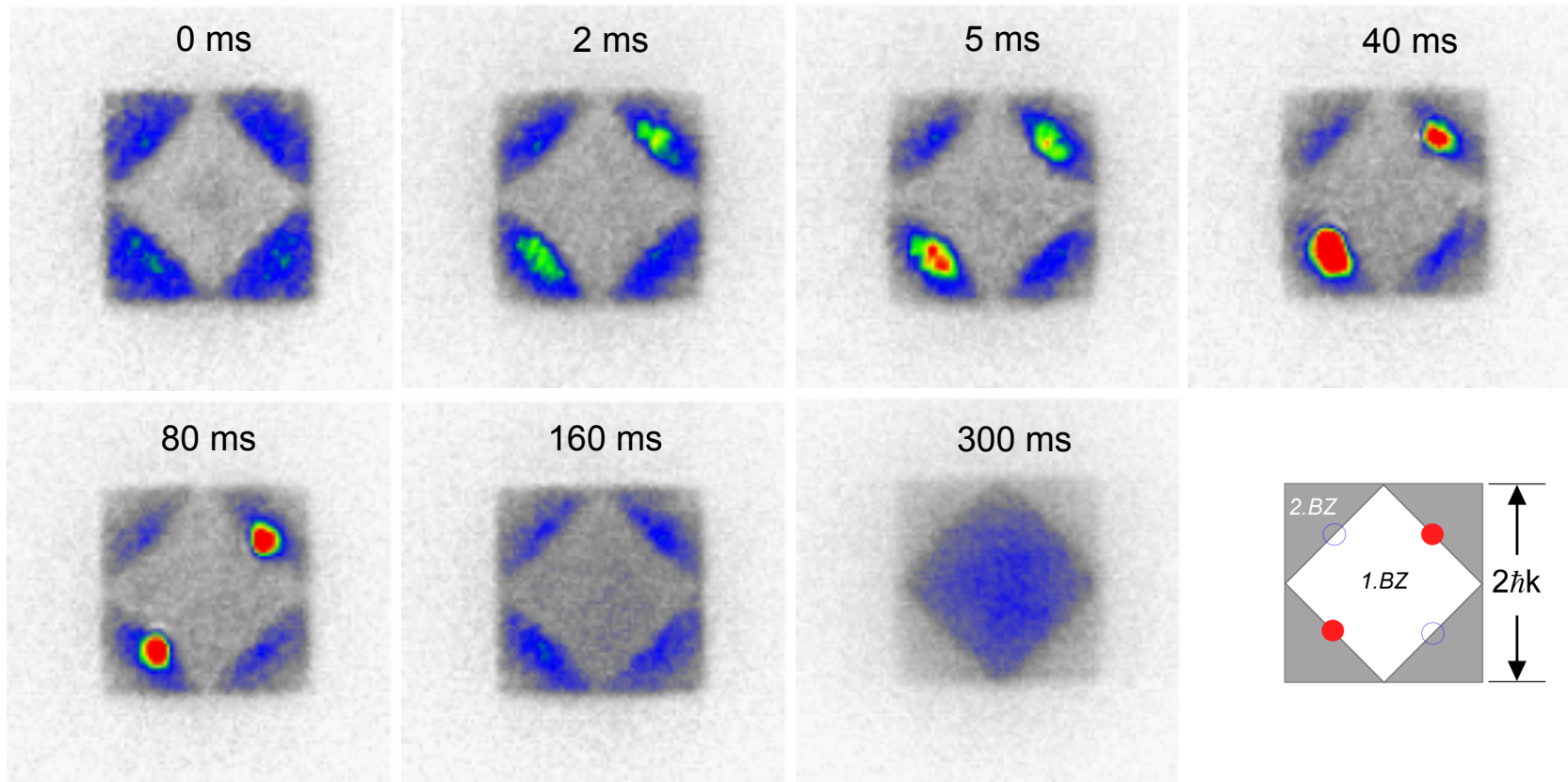


G. Wirth, M. Ölschläger, A. Hemmerich, Nature Physics 7, 147 (2011)

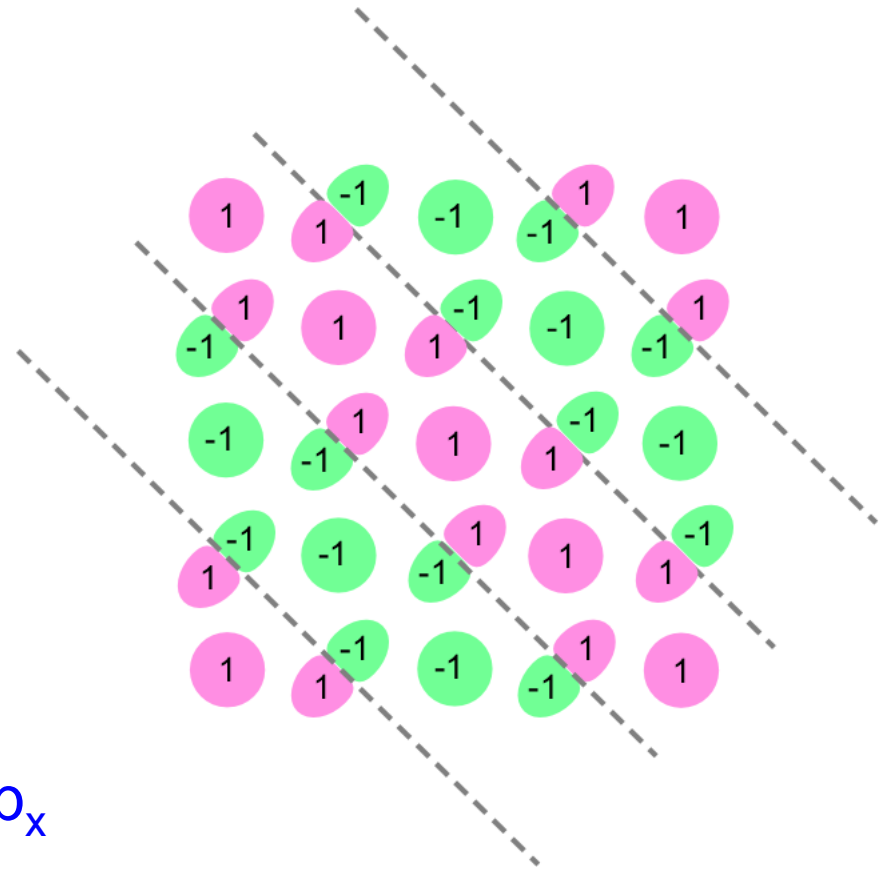
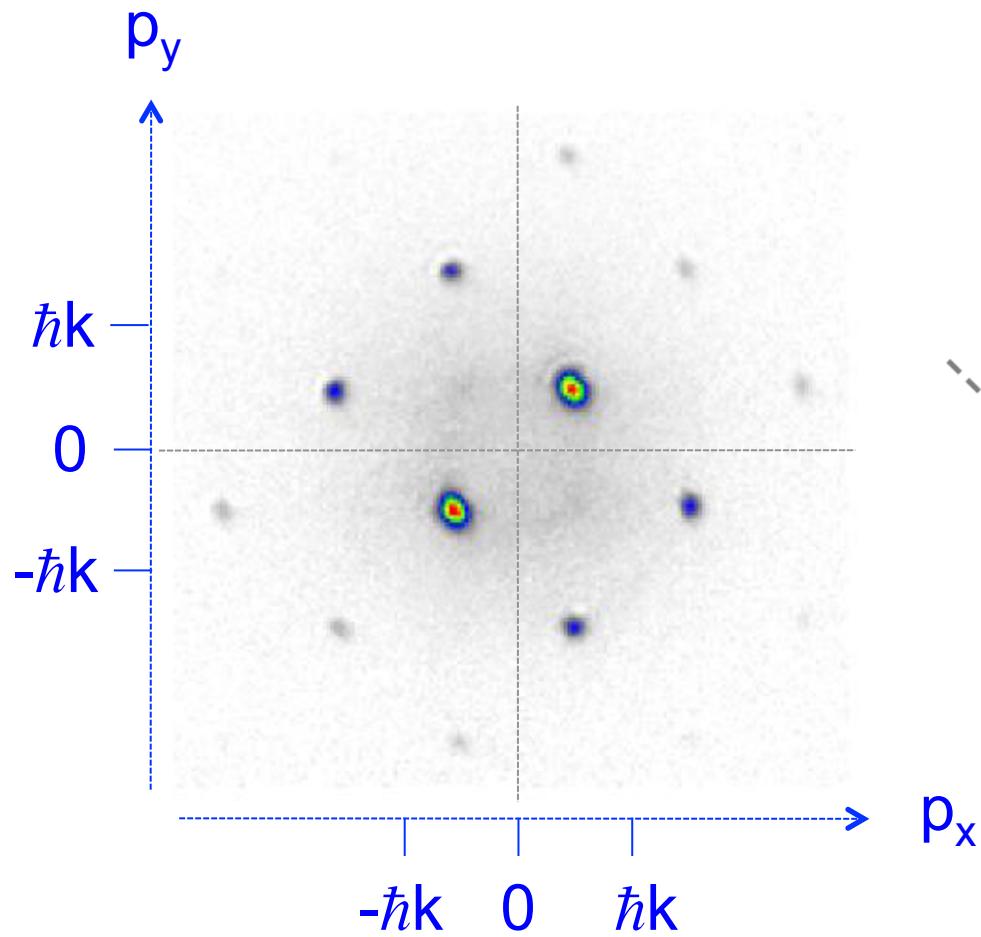


- After 1 ms condensation at finite momenta on the edge between 1st and 2nd Brillouin zone
- Decay after 100 ms
- Distortion selects condensation points

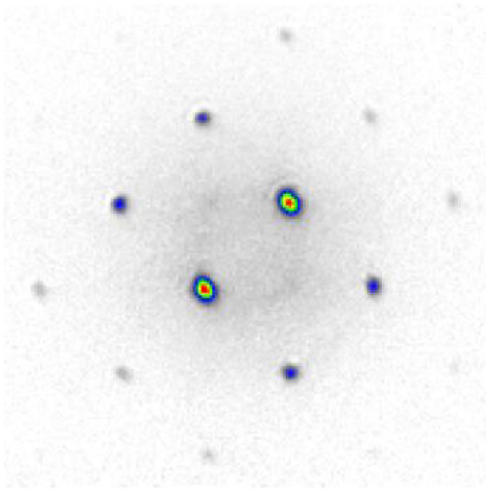
# Band relaxation sets in after 100 ms



# Finite momentum superfluidity

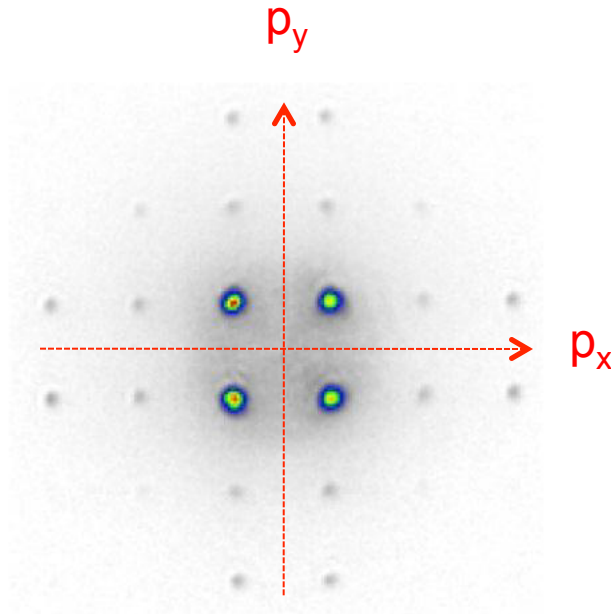
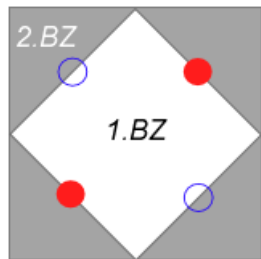


# Tuning the distortion $\Delta E$

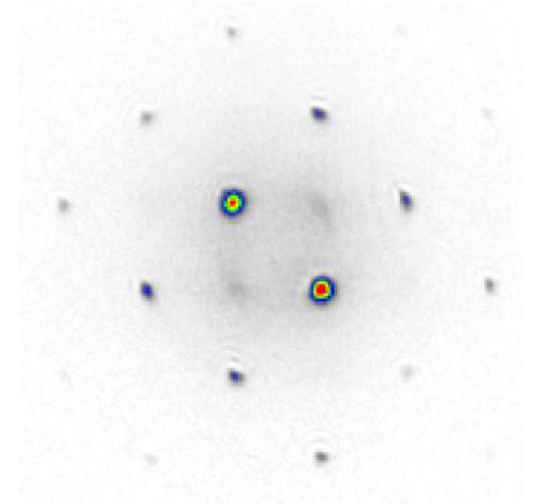
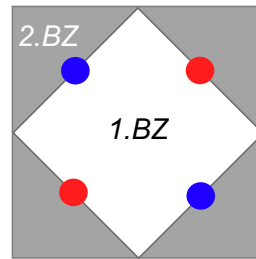


$$\Delta E \approx 0.005 E_{\text{rec}} > 0$$

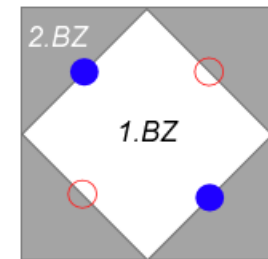
$$\rightarrow \Delta T \approx 0.5 \text{ nK}$$



$$\Delta E \approx 0$$



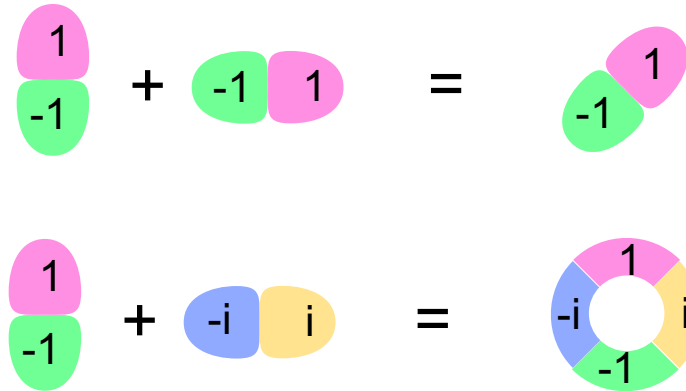
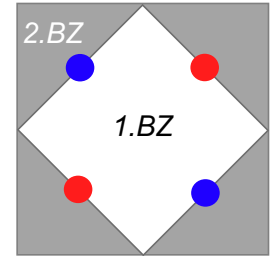
$$\Delta E < 0$$



For  $\Delta E = 0$

superposition of condensates at both band minima

$90^\circ$  relative phase chosen in order to minimize interaction of atoms in  $p_x$ - and  $p_y$ -orbitals

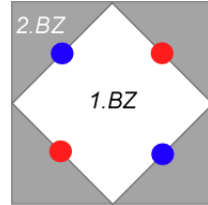


minimizes interaction energy

A. Isacsson and S. Girvin, Phys. Rev. A 72, 053604 (2005).  
W. Vincent Liu and C. Wu, Phys. Rev. A 74, 013607 (2006).

# A chiral wave function arises

Atoms condense at both minima



$$E_{\bullet} = E_{\bullet}$$

Deep wells support  $P_x$ ,  $P_y$ -orbitals

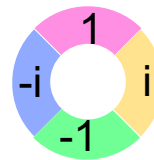


Shallow wells support local S-orbitals

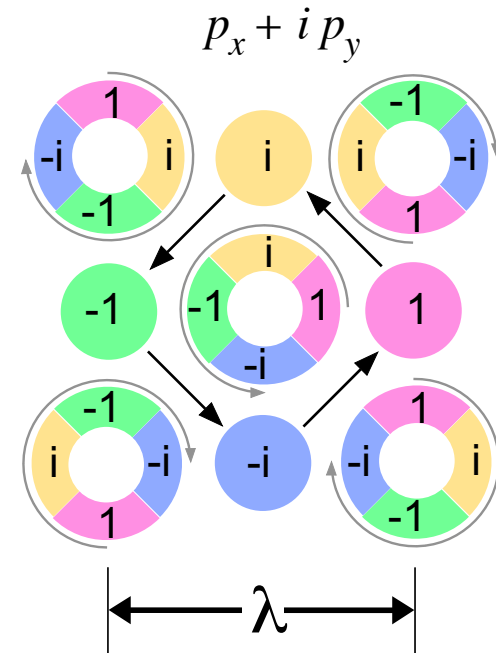


$P_x$ ,  $P_y$  degenerate with respect to kinetic energy

$P_x \pm i P_y$  minimize local interaction energy



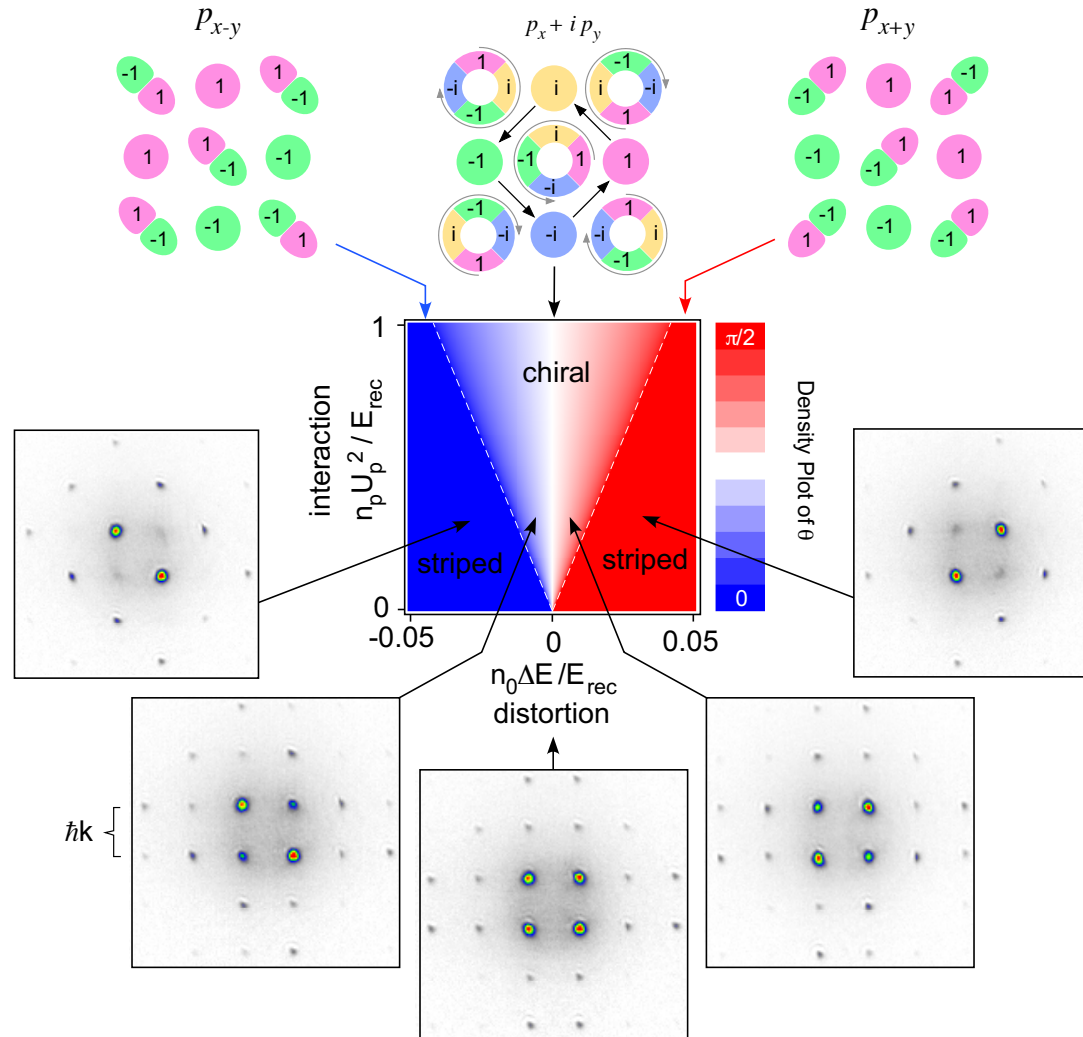
Tunneling maximized by matching local phases at tunneling junctions



→ broken time reversal symmetry

→ interaction induced vortical flow, artificial magnetic field

# Observed quantum phases confirmed by Hubbard model



M. Ölschläger, T. Kock, G. Wirth, A. Ewerbeck, C. Morais Smith, A. Hemmerich  
 New Journal of Physics 15, 083041 (2013)



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Ultracold Alkaline Earth Atoms

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