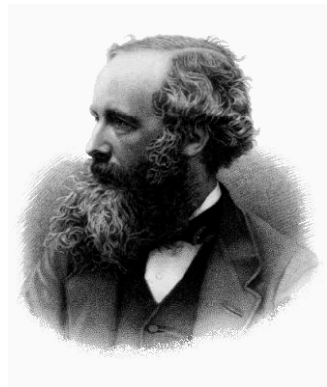


# Theoretical On-Demand Adiabatic Transfer of Light Between Adjacent Optical Cavities

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McMaster University

Some Other Dude, Ed Hinds  
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# The Setup



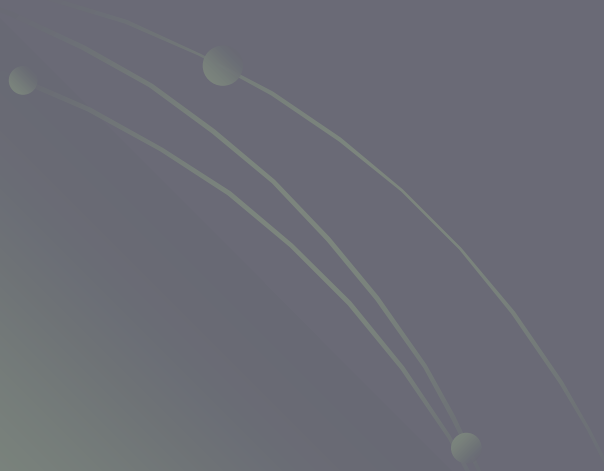
$$x = -L_1$$

$$x = 0$$

$$x = L_2$$

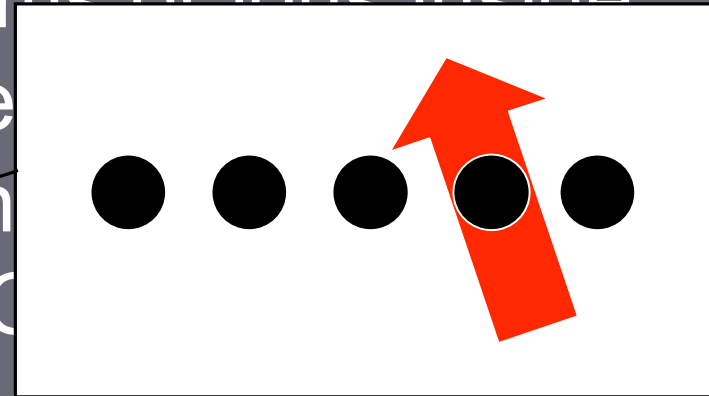
# Goal

- Coherently transfer light between two adjacent optical cavities connected by a slightly transmissive mirror



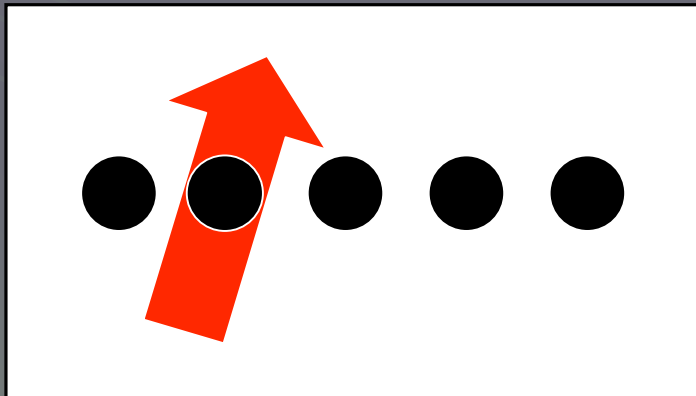
# ...why bother?

- A network of neutral atoms or ions inside optical cavities connected in a chain is one of the most promising systems for realizing QIC



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- One of the important tasks left to face is to transfer photons in a



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# Approach

-  Examine the modes of the double-cavity model with static boundary conditions for a classical light field
-  Look at the energy distribution for these modes
-  Adiabatically change the position of the transmissive central mirror in order to transfer light coherently from one cavity to the next

# Modes of Static Cavity

- We first solve Maxwell's equations, which in our case reduce to:

$$\frac{\partial^2 E(x, t)}{\partial x^2} - \mu_0 \epsilon_0 (1 + \alpha \delta(x)) \frac{\partial^2 E(x, t)}{\partial t^2} = 0$$

- The solutions are separable and we write them as:

$$E_n(x, t) = E_n(x) e^{-i\omega_n t}$$

- Using the continuity boundary condition:

$$E_n(x) = \begin{cases} A_n \sin(k_n(x + L_1)) & -L_1 \leq x \leq 0 \\ B_n \sin(k_n(x - L_2)) & 0 < x \leq L_2 \end{cases}$$

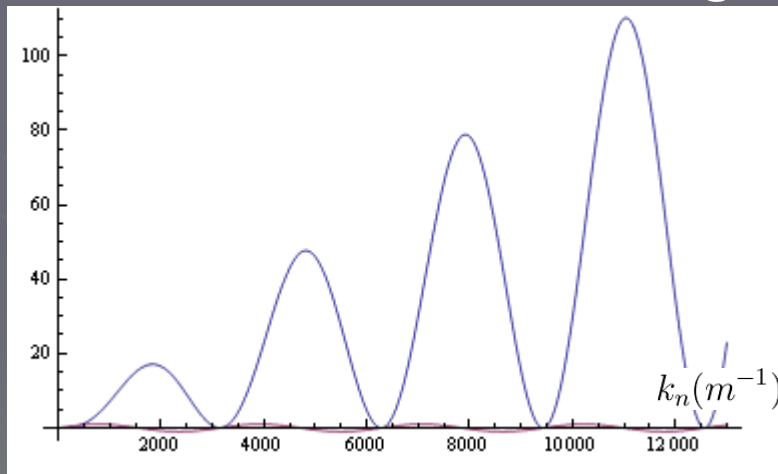
$$\frac{A_n}{B_n} = -\frac{\sin(k_n L_2)}{\sin(k_n L_1)}$$

# Modes of Static Cavity

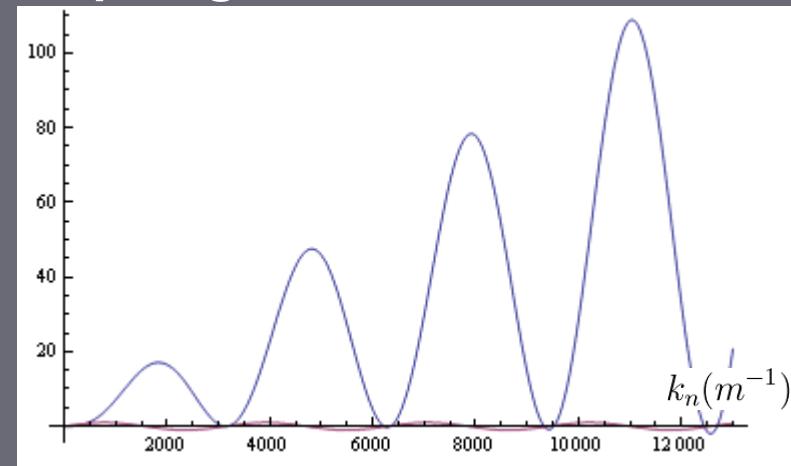
- Applying the boundary condition for the first derivative of the solution, we find:

$$\alpha k \sin(k_n L_1) \sin(k_n L_2) = \sin(k_n L)$$

## Finding $k_n$ By Graphing



Solving for  $k_n$  with  $L_1 = L_2 = 10^{-3}$  m,  
and  $\alpha = 10^{-2}$  m



Solving for  $k_n$  with  $L_1 = 0.99 \times 10^{-3}$  m  
 $L_2 = 1.01 \times 10^{-3}$  m, and  $\alpha = 10^{-2}$  m

# Modes of Static Cavity

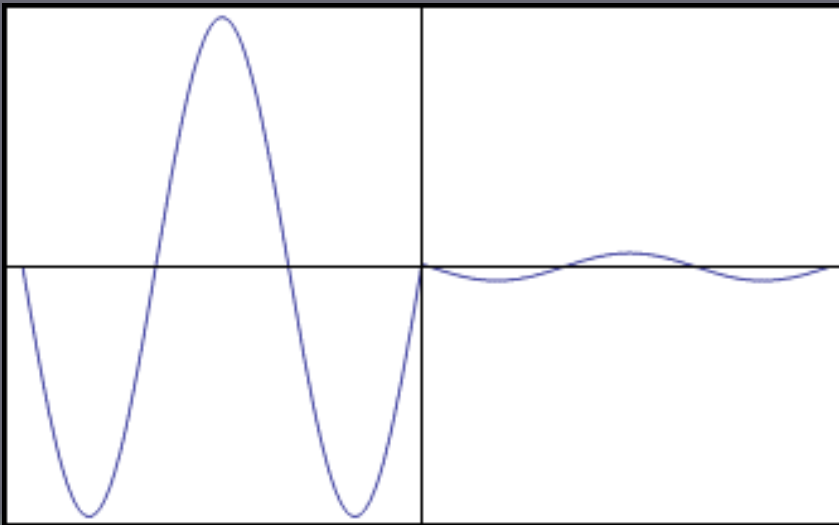
- Rearranging, and using the normalization condition, we find:

$$A_n^2 = \frac{1 - B_n^2 \left[ \frac{L_2}{2} - \frac{\sin(2k_n L_2)}{4k_n} + \alpha \sin^2(k_n L_2) \right]}{\frac{L_1}{2} - \frac{\sin(2k_n L_1)}{4k_n}}$$

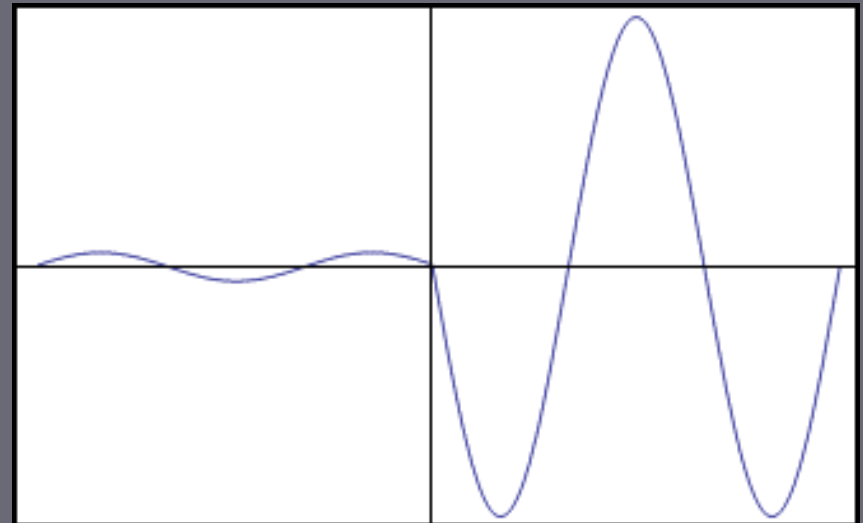
$$B_n^2 = \frac{1}{\frac{\sin^2(k_n L_2)}{\sin^2(k_n L_1)} \left[ \frac{L_1}{2} - \frac{\sin(2k_n L_1)}{4k_n} \right] + \left[ \frac{L_2}{2} - \frac{\sin(2k_n L_2)}{4k_n} + \alpha \sin^2(k_n L_2) \right]}$$

# Modes of Static Cavity

Double-Cavity Modes for  $\alpha = 10^{-2}$  m



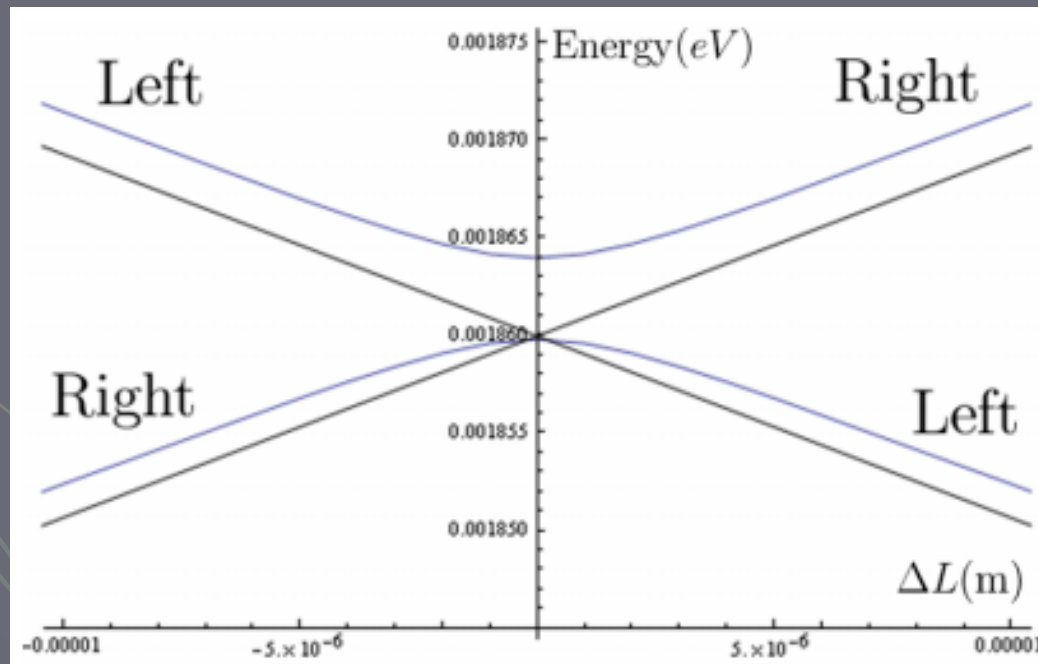
Double-cavity mode for  $k_n = 9531.15 \text{ m}^{-1}$ ,  
 $L_1 = 0.99 \times 10^{-4} \text{ m}$ ,  $L_2 = 1.01 \times 10^{-3} \text{ m}$



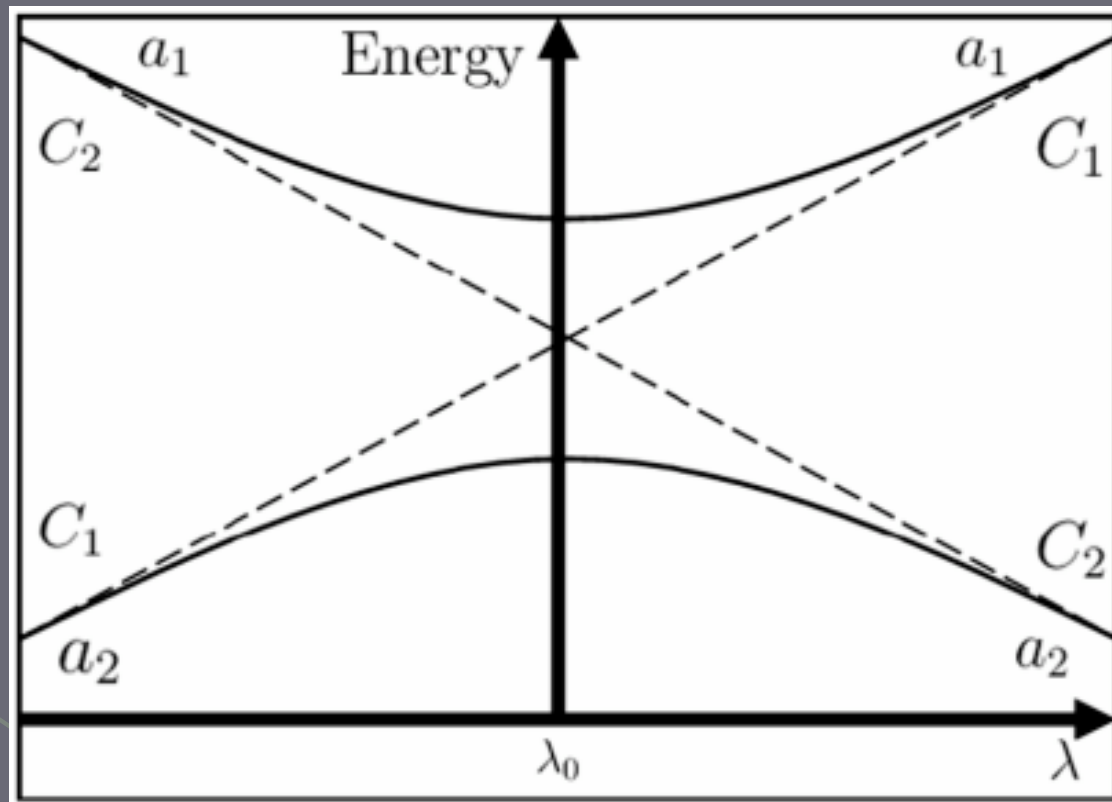
Double-cavity mode for  $k_n = 9341.46 \text{ m}^{-1}$ ,  
 $L_1 = 0.99 \times 10^{-4} \text{ m}$ ,  $L_2 = 1.01 \times 10^{-3} \text{ m}$

# Energy Distribution

- Choosing one energy “rung”, we fill each mode with one photon and calculate the corresponding energy

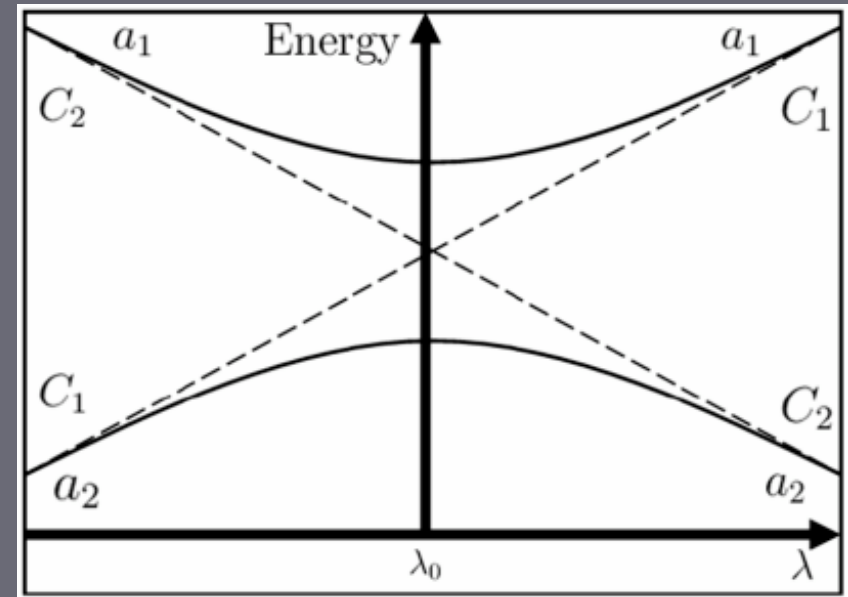
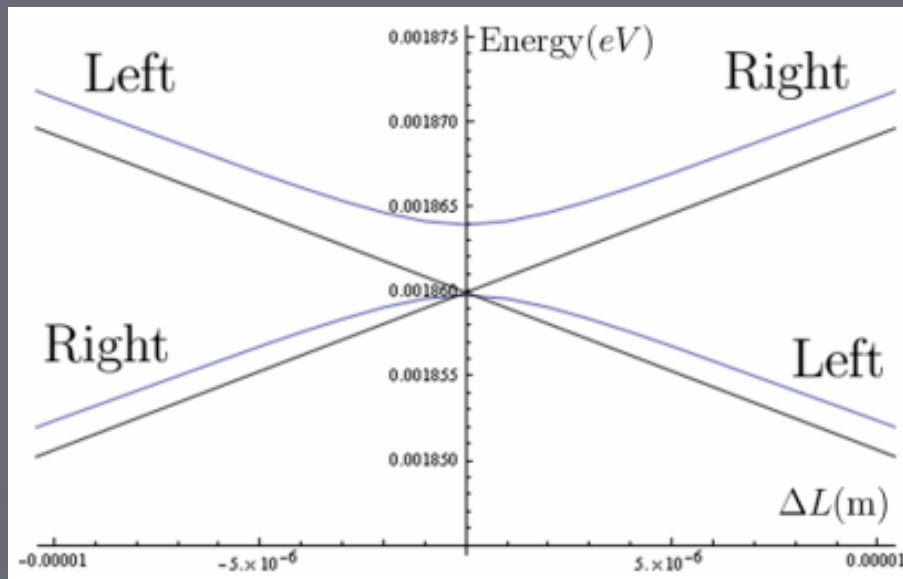


# Landau-Zener Problem



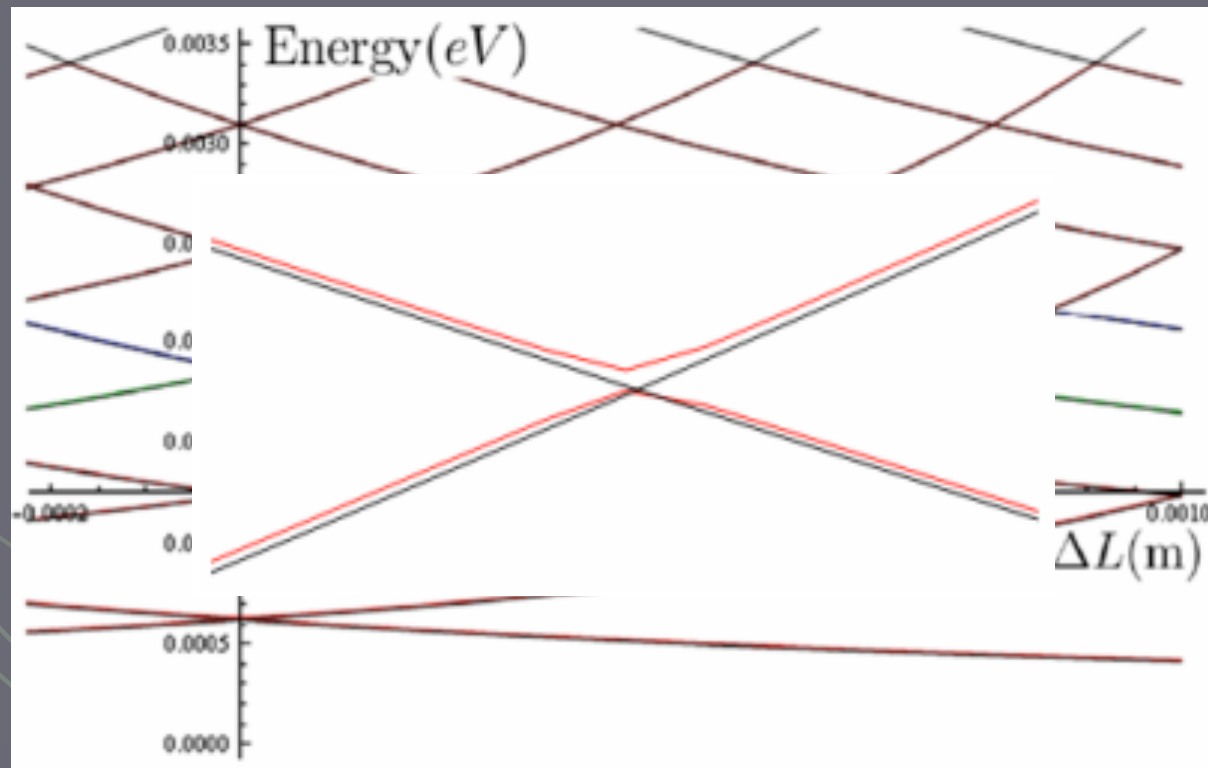
- Changing the parameter  $\lambda$  adiabatically, the adiabatic state changes which bare state it is associated with

# Light Field Transfer

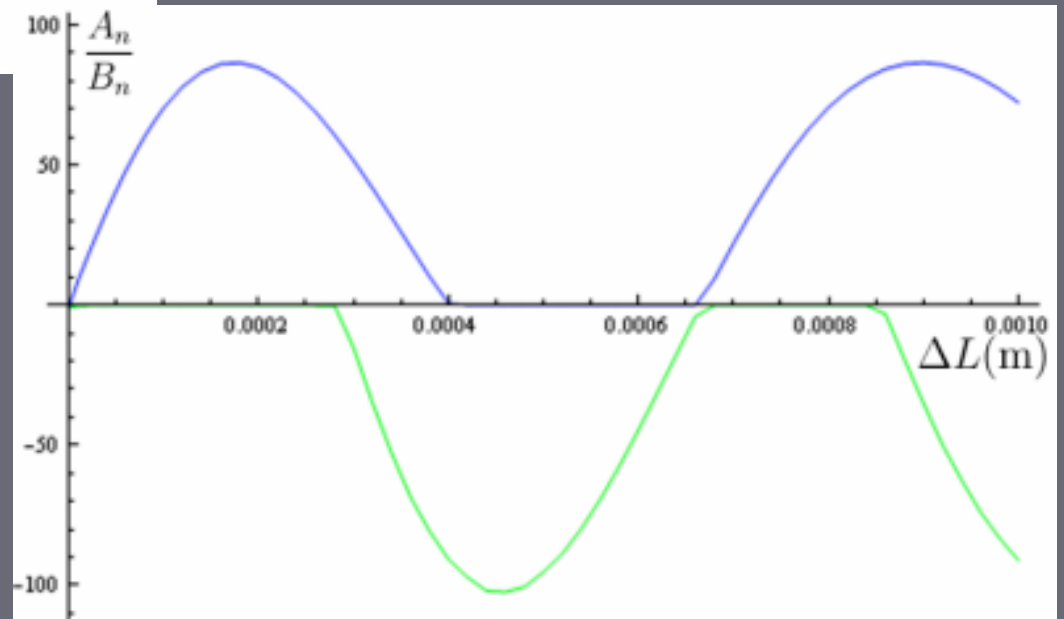
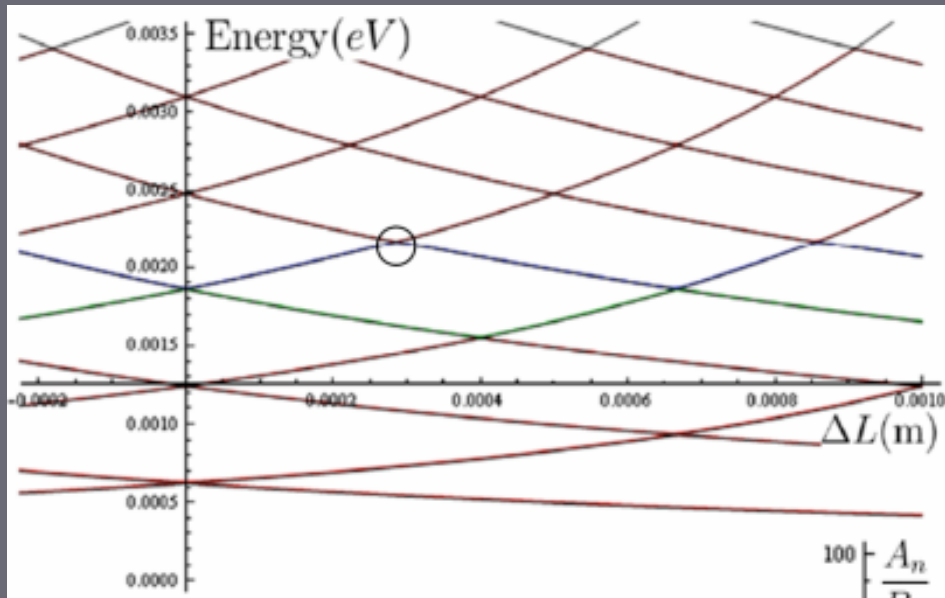


- Transfer the light field from one cavity to the next by adiabatically changing the position of the central mirror

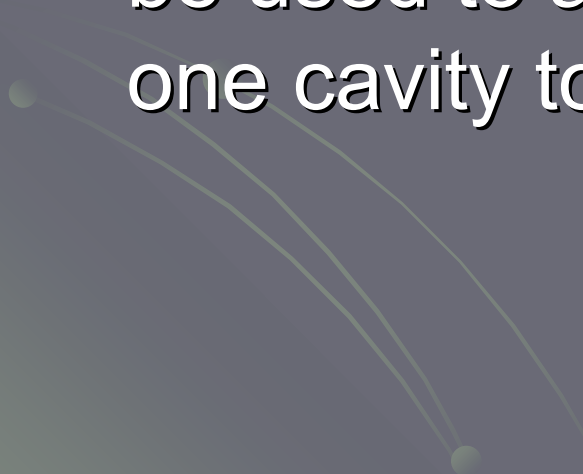
...and for other displacements?



# ...what about the amplitudes?



# Conclusions

- By taking  $x \neq 0$  for the shared mirror, we can localize the light modes into different cavities
  - We have identified level crossings that can be used to adiabatically transfer light from one cavity to the next
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# Future Direction

- Apply Landau-Zener formalism
- Apply more realistic values to our system
- Optimize the localization of the light modes
- Extend the double-cavity model to two optical cavities connected by an optical fibre
- Collaboration with the group of Ed Hinds?

# References

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- [3] “Quantum Information Science and Technology Roadmapping Project,” Los Alamos National Laboratory (2006)  
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The End

