

A Universal Correction to the Inflationary Vacuum

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Introduction and motivation

Quantum Gravity in de Sitter space?

- Even if $\Lambda \ll M_p^2$, interesting physics because of dS temperature $T_{GH} = \frac{H}{2\pi}$
 - Eternal inflation, holography, complementarity, ...
 - Cosmological constant problem?
- Moreover: observationally relevant
 - Inflation and the primordial power spectrum: $P_{\mathcal{R}}(k) \propto T_{GH}^2$
 - Current phase of acceleration

Cosmology: best bet to observationally probe string scale physics

Inflation and string scale physics

Inflation quite generally suggests the following possibilities

- CMB as a probe of the inflationary vacuum structure ($\frac{H}{M_s} > 10^{-6}$)
- CMB as a probe of effective dynamics during inflation ($\frac{H}{M_s} > 10^{-3}$)
- CMB as a probe of primordial non-gaussiannities ($c_s < 1$)
- Cosmic string production in string inflation models (?)

Not entirely hopeless if $H \sim 10^{14}$ GeV. Future prospect: the precision of the primordial scalar density perturbation might reach the 10^{-6} level [Spergel]

Worthwhile to explore every option

Probing the inflationary vacuum

Inflation is an excellent probe of the small scale vacuum structure

- Theory and observation agree: state in $\langle \text{in} | \delta\phi^2 | \text{in} \rangle$ is very close to the BD vacuum
- Pragmatic point of view: small departures cannot be entirely excluded, but are subject to strong consistency constraints (backreaction)
- The proposals so far can lead to H/Λ corrections, but in general do not have to, and are typically either somewhat fine-tuned [BEFT] or somewhat ad-hoc [NPH]
- The proposals do lead to rather distinctive signatures (modulations)
- In general dynamical (quantum) corrections [EFT] give H^2/Λ^2 effects

Greene, Schalm, Shiu, JPvdS, [astro-ph/0503458](#)

Schalm, Shiu, JPvdS, [hep-th/0412288](#)

Kaloper, Kleban, Lawrence, Shenker, Susskind, [hep-th/0209231](#)

Easter, Greene, Kinney, Shiu, [hep-th/0204129](#), Danielsson, [hep-th/0203198](#)

Martin, Brandenberger, [hep-th/0005209](#)

Inflationary vacuum corrections

Can we come up with calculable model-independent corrections to the inflationary vacuum?

- Proof of principle
- Obvious question: could it be observable?
- Interesting underlying physics?

One argument in favor of the BD state is that it gives rise to a thermal spectrum from the point of view of a free falling observer

Turning this argument around: do we really expect the particle spectrum to be exactly thermal?

De Sitter tunneling

Temperature in dS can be described explicitly in terms of tunneling through the dS horizon

- Energy conservation: spectrum cannot be precisely thermal
 - Energy of emitted quanta cannot exceed Narai bound
 - De Sitter horizon 'relocates' after emitting quanta
- Effect of backreaction can be incorporated for (dominant) s-waves

$$\Gamma_k = N \exp \left(-2 \operatorname{Im} \int dt L(\omega_k) \right) = N \exp \Delta S(\omega_k) \approx N \exp (-\beta \omega_k)$$

with $L(\omega_k) = p\dot{r} - \omega_k$ and $\beta = 1/T_{GH}$

Kraus, Wilczek, gr-qc/9408003

Parikh, Wilczek, hep-th/9907001

Parikh, hep-th/0204107

De Sitter Vacua

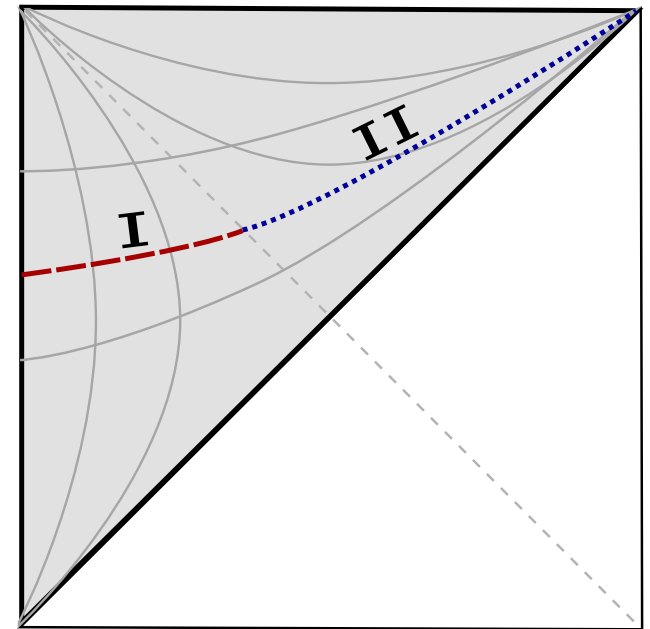
We will be assuming Painlevé coordinates covering dS

$$ds^2 = - (1 - H^2 r^2) dt^2 - 2Hr dr dt + dr^2 + r^2 d\Omega^2$$

with $t_{\text{Painleve}} = t_{\text{planar}} = t_{\text{static}} + \frac{1}{2H} \ln(1 - H^2 r^2)$ and $r_{\text{Painleve}} = r_{\text{static}} = \rho_{\text{planar}} e^{Ht}$

These coordinates are useful for several reasons

- Cover the full planar patch
- Between static and planar
 - Planar sections of constant t
 - Static sections of constant r
- Stationary



De Sitter Vacua

Introduce $|\text{out}\rangle = |I\rangle \otimes |II\rangle$, $|\text{BD}\rangle$ and $|\text{BD}'\rangle$, then with $b_k |\text{out}\rangle = 0$ the probability for particle creation can be written as

$$\left| \frac{\langle \text{out} | b_k b_{-k} | \text{BD} \rangle}{\langle \text{out} | \text{BD} \rangle} \right|^2 = e^{-\beta \omega_k}, \quad \left| \frac{\langle \text{out} | b_k b_{-k} | \text{BD}' \rangle}{\langle \text{out} | \text{BD}' \rangle} \right|^2 = e^{\Delta S(\omega_k)}$$

- The state $|I\rangle$ should be thought of as the static vacuum
- Consequently $|II\rangle$ corresponds to a 2nd static vacuum
- Annihilation operators each act on different sides

Exercise: determine the Bogolyubov rotation between the $|\text{BD}\rangle$ and $|\text{BD}'\rangle$ state

$$a'_k = \alpha_k^* a_k - \beta_k^* a_{-k}^\dagger \quad \text{with} \quad a_k |\text{BD}\rangle = a'_k |\text{BD}\rangle = 0$$

- Transformations are diagonal because all states are defined in the same patch
- The Bogolyubov transformations wrt the $|\text{out}\rangle$ state are known

Relating the states

$$\tilde{\gamma}_k \equiv -\frac{\tilde{\beta}_k^*}{\tilde{\alpha}_k} = \frac{\langle \text{out} | b_k b_{-k} | \text{BD} \rangle}{\langle \text{out} | \text{BD} \rangle}, \quad \tilde{\gamma}'_k \equiv -\frac{\tilde{\beta}'_k^*}{\tilde{\alpha}'_k} = \frac{\langle \text{out} | b_k b_{-k} | \text{BD}' \rangle}{\langle \text{out} | \text{BD}' \rangle}$$

After some basic linear manipulations one obtains

$$\alpha_k^* = \tilde{\alpha}'_k \tilde{\alpha}_k^* - \tilde{\beta}'_k^* \tilde{\beta}_k, \quad \beta_k^* = \tilde{\alpha}'_k \tilde{\beta}_k^* - \tilde{\beta}'_k^* \tilde{\alpha}_k$$

Defining $\gamma_k = -\frac{\beta_k^*}{\alpha_k}$ then gives $\gamma_k = \frac{\tilde{\gamma}_k - \tilde{\gamma}'_k}{1 - \tilde{\gamma}_k \tilde{\gamma}'_k}$

- Nontrivial phases are absent
- Consequence of tunneling description and consistent with α -state calculations

Since $|\tilde{\gamma}_k|^2 = e^{-\beta\omega_k}$ and $|\tilde{\gamma}'_k|^2 = e^{\Delta S(\omega_k)}$ we find, after expanding to leading order,

$$|\gamma_k| \approx \frac{e^{\pi \frac{\omega_k}{H}}}{e^{2\pi \frac{\omega_k}{H}} - 1} \frac{\omega_k H}{8M_p^2}$$

Primordial spectrum

Our claim: the state $|BD'\rangle$ is the ‘true’ inflationary vacuum

For small corrections we then immediately find

$$P(k) \approx P_{BD} \left[1 + \frac{1}{4e^\pi} \left(\frac{H(k)}{M_p} \right)^2 \right]$$

- Evaluated at time when a mode crosses the horizon, i.e when $\omega_k = H$
- Absence of k -dependent phases, so no modulation
- Tiny effect, already constrained at the $H^2/M_p^2 \lesssim 10^{-10}$ level
- Indistinguishable from a (specific) dimension 6 operator correction

Final thoughts

Some remarks and questions

- Obviously too small an effect to be observable
- Nevertheless formally interesting
 - Model independent, i.e. universal
 - Backreaction effect
 - Explicitly calculable
- Effective field theory interpretation?
- Breaks de Sitter $SO(1, 4)$ isometries (spontaneously)
 - Consequences for de Sitter stability?

**Finding models that can make testable predictions
remains an important challenge**