

diversion - prompted by question:

(27.1) ①

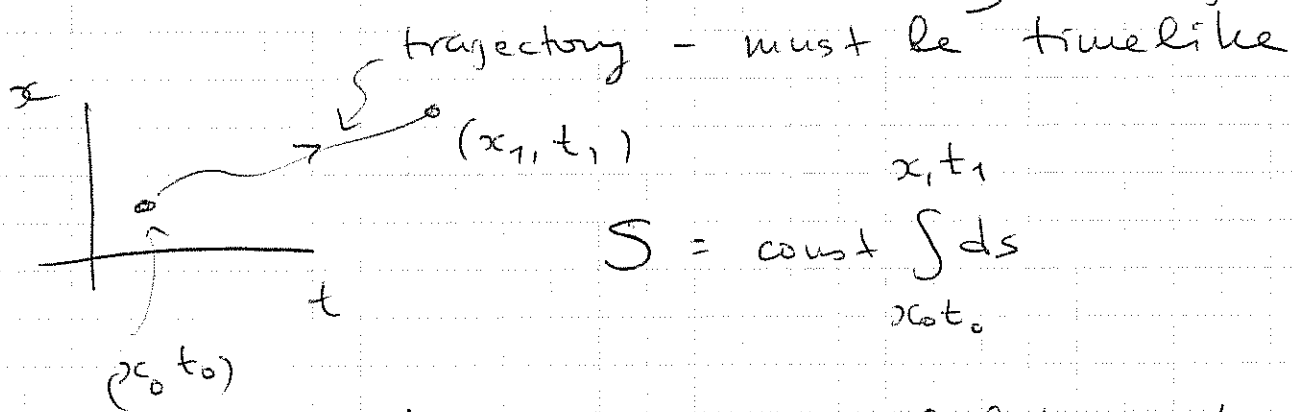
for a relativistic particle:

Galileo's principle \rightarrow

\rightarrow Lorentz invariance principle

$S(\text{trajectory } x_0 t_0 \rightarrow x_1 t_1)$ should be same in any Lorentz frame

(i.e. invariant under Lorentz transf's)



ds = spacetime interval between two neighboring points

$$\Rightarrow ds^2 = -c^2 dt^2 + d\vec{x}^2$$

$$ds = \sqrt{-ds^2} = \sqrt{c^2 dt^2 - d\vec{x}^2} =$$

\uparrow remember, timelike

$$= c dt \sqrt{1 - \frac{d\vec{x}^2}{c^2 dt^2}}$$

$$= c dt \sqrt{1 - \frac{v^2}{c^2}}$$

5. $S(\bar{x}_0 t_0 \rightarrow \bar{x}_1 t_1) = \text{const } c \int_{t_0}^{t_1} dt \sqrt{1 - \frac{\vec{v}^2}{c^2}}$

$$\vec{v} = \frac{d\vec{x}(t)}{dt}$$

$$\vec{x}(t_0) = \bar{x}_0$$

$$\vec{x}(t_1) = \bar{x}_1$$

const = ?

take nonrelativistic limit:

$$\sqrt{1 - \frac{\vec{v}^2}{c^2}} \approx 1 - \frac{1}{2} \frac{\vec{v}^2}{c^2} + \dots$$

($\sqrt{1-x} \approx 1 - \frac{1}{2}x + \dots$)

$$S(\bar{x}_0 t_0 \rightarrow \bar{x}_1 t_1) = \text{const} \cdot c \int_{t_0}^{t_1} dt \sqrt{1 - \frac{\vec{v}^2}{c^2}} \approx$$

$$\approx \underbrace{\text{const } c \int_{t_0}^{t_1} dt}_{\text{indep. of } \vec{x}(t)} - \frac{\text{const}}{2} c \int_{t_0}^{t_1} \frac{\vec{v}^2}{c^2} dt + \dots$$

-drop

$$\approx - \frac{\text{const}}{2} \frac{1}{c} \int_{t_0}^{t_1} dt \vec{v}^2 = \int_{t_0}^{t_1} \frac{m \vec{v}^2}{2} dt$$

$\Rightarrow \text{const} = -mc$ \neq $S = -mc \int ds$
↑
relativistic

29.3

3

$$S = -mc^2 \int_{t_0}^{t_1} dt \sqrt{1 - \frac{\vec{v}^2}{c^2}}$$

$L(\vec{v}^2)$

$\left\{ \begin{array}{l} \text{homogeneity +} \\ \text{isotropy} \end{array} \right.$
 but NOT Galilean relativity!!

Equation of motion:

$$\frac{d}{dt} \frac{\partial L}{\partial \vec{v}} = 0 \Rightarrow \frac{d}{dt} \left(\frac{\vec{v}}{\sqrt{1 - \frac{\vec{v}^2}{c^2}}} \right) = 0$$

\uparrow
 (relativistic acceleration = 0)

For a charged particle:

$$S = -mc^2 \int dt \sqrt{1 - \frac{\vec{v}^2}{c^2}} + q \int dt \left(\varphi(\vec{x}, t) - \frac{d\vec{x}}{dt} \cdot \vec{A}(\vec{x}, t) \right)$$

$$= -mc^2 \int dt \sqrt{1 - \frac{\vec{v}^2}{c^2}} + q \int \left(dt \varphi(\vec{x}, t) - d\vec{x} \cdot \vec{A}(\vec{x}, t) \right)$$

$(dt, d\vec{x})$: 4-vector

(φ, \vec{A}) : 4-vector

$dt \cdot \varphi - d\vec{x} \cdot \vec{A} = \text{Lorentz invariant}$

hw1#1:
 nonrelativistic
 limit
 power!