

* Title credit goes to *Uncertain Principles on Science Blogs*.

Photons: Still Bosons*

Spectroscopic test of Bose-Einstein statistics for photons

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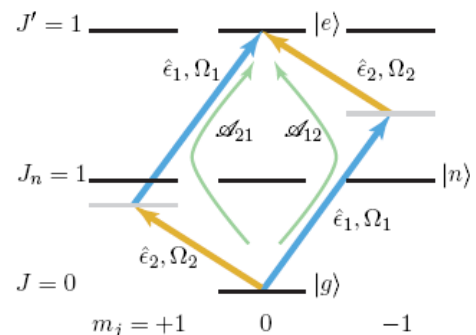
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Using Bose-Einstein-statistics-forbidden two-photon excitation in atomic barium, we have limited the rate of statistics-violating transitions, as a fraction ν of an equivalent statistics-allowed transition rate, to $\nu < 4.0 \times 10^{-11}$ at the 90% confidence level. This is an improvement of more than three orders of magnitude over the best previous result. Additionally, hyperfine-interaction enabling of the forbidden transition has been observed, to our knowledge, for the first time.

PACS numbers: 42.50.Xa, 82.50.Pt



September 1st, 2010 Group meeting – Rockson Chang

Spin statistics theorem

- Relates the spin of a particle to the statistics it obeys.
- Integer spin particles (0,1,2...) == Bosons
 - WF symmetric under particle exchange

$$\phi(x)\phi(y) = \phi(y)\phi(x).$$

-- Obey Bose-Einstein statistics

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Half-integer spin particles (1/2, 3/2, 5/2...) == Fermions

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Why??...

- Feynman's lectures on physics (Vol 3, Ch4 – Identical Particles)



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An explanation has been worked out by Pauli from complicated arguments of QFT and relativity... It appears to be one of the few places in physics where there is a rule which can be stated very simply, but for which no one has found a simple and easy explanation... This probably means that we do not have a complete understanding of the fundamental principle involved.



Spin statistics theorem

- A result from QFT
- Assumes:
 - Causality
 - Lorentz invariance in 3+1 spacetime dimensions
- Observing even a miniscule violation would rock the foundations of modern physics
- Quon algebra? $a_k a_l^\dagger - q a_l^\dagger a_k = \delta_{kl}$,
- Excitation in higher dimensions (string theory)?

Landau-Yang theorem

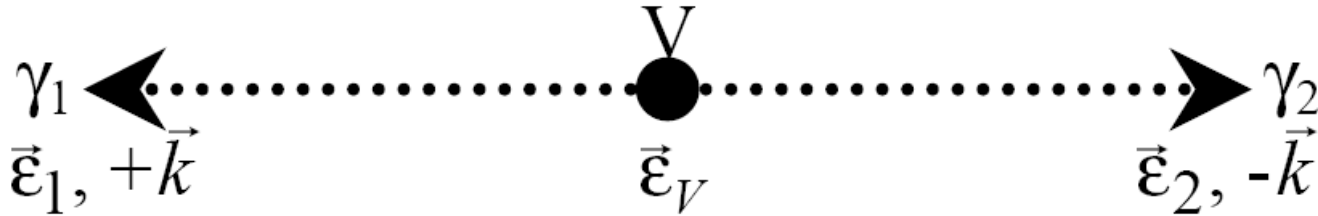


FIGURE 1. Kinematics of the decay $V \rightarrow \gamma\gamma$, in the rest frame of V .

- A neutral $J=1$ boson can't decay into two *degenerate* photons.
 - Resulting 2 photon state must have $J'=1$

$$\Psi = |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle$$

- Ψ ($J'=1$) anti-symmetric under particle exchange

Spectroscopic test of Bose-Einstein statistics for photons

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- Turn this into a table-top atomic physics experiment by considering the time-reversed scenario:

*Look for two photon
absorption in Barium-138*

$(I=0)$

$$J=0 + 2 \gamma \rightarrow J'=1$$

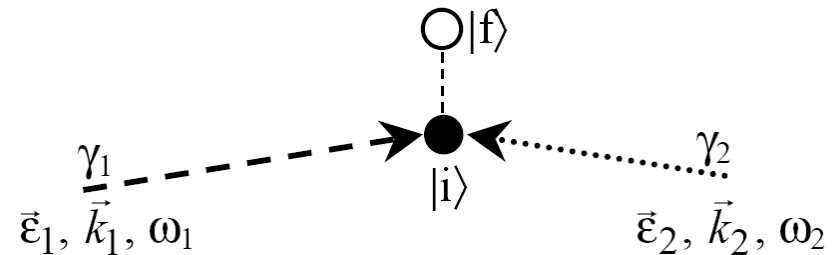
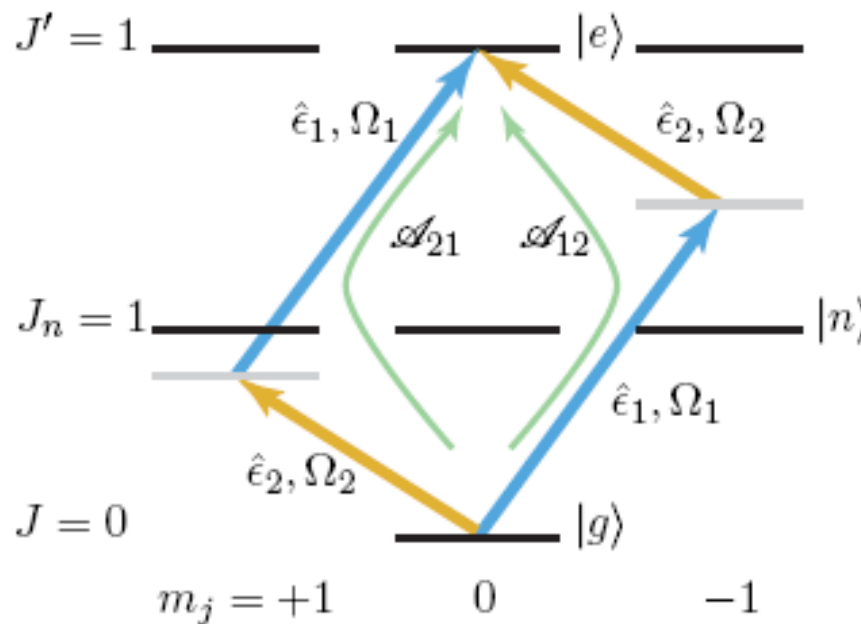


FIGURE 2. Kinematics of an atomic two-photon transition.



Orthogonal
circular
polarizations

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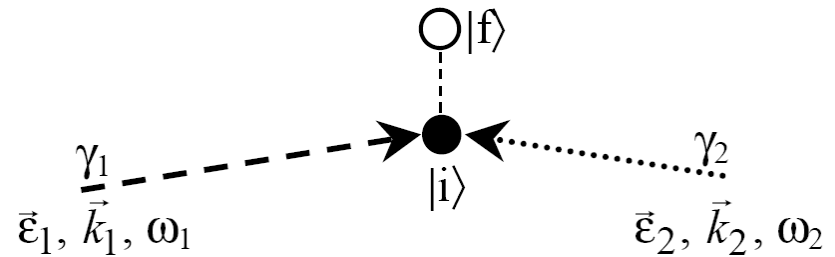
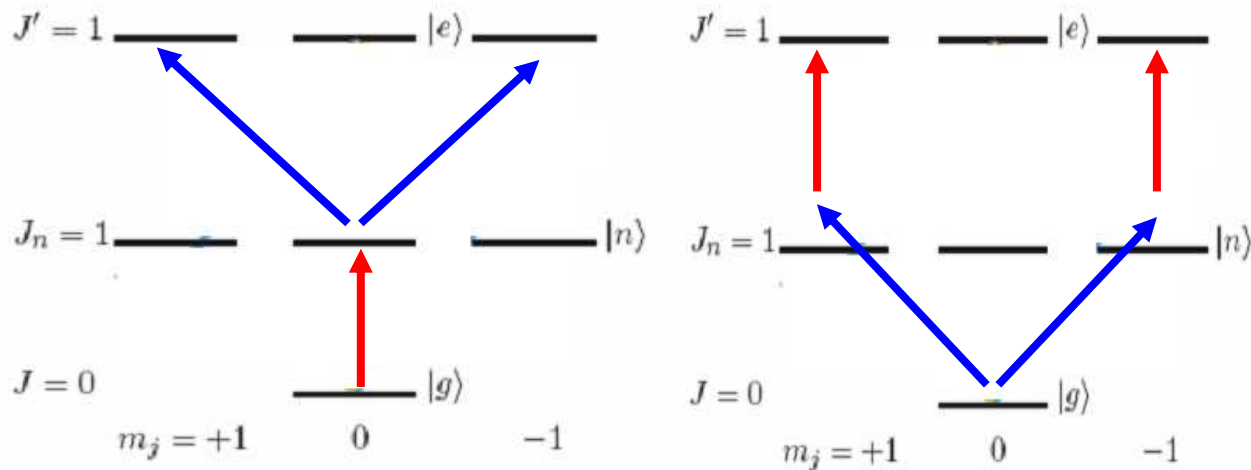


FIGURE 2. Kinematics of an atomic two-photon transition.



Orthogonal
linear
polarizations

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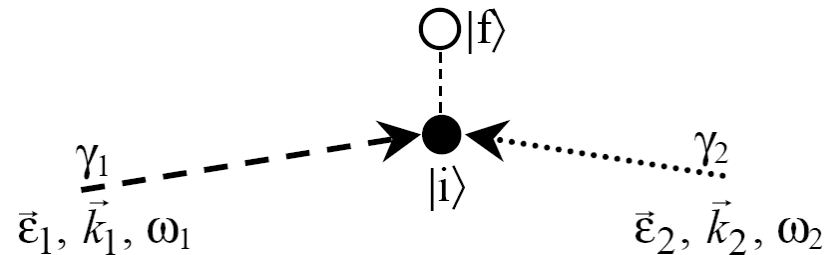
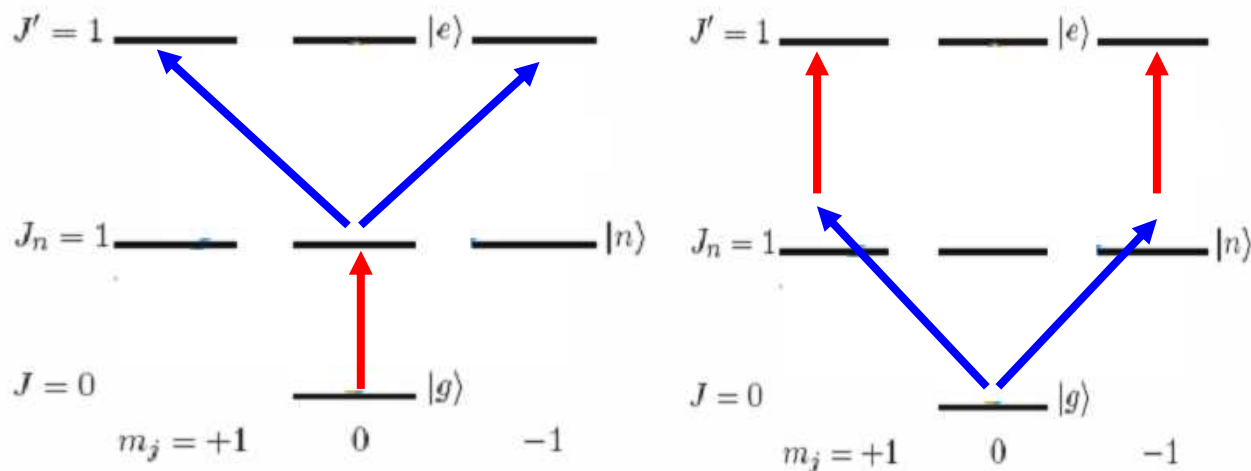


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Orthogonal
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Q1A: Pathways indistinguishable!
(what about photons?)

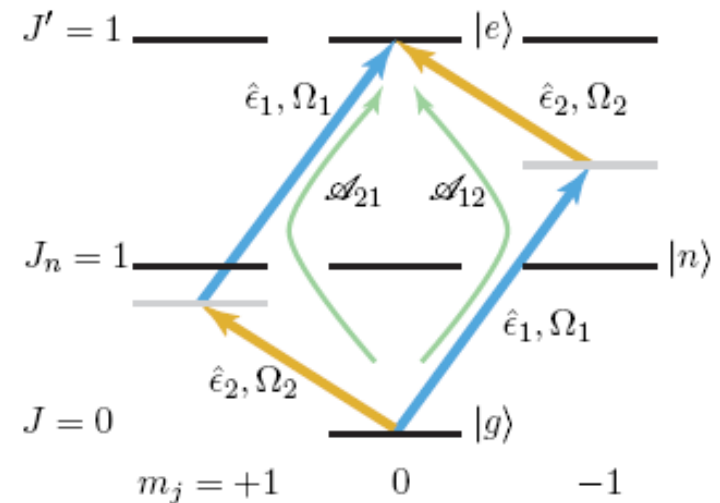
- Consider TPA rate

(2nd order time-dependent perturbation, Fermi's golden rule)

$$W_+ = \frac{2\pi}{\hbar^4} \times \left| \sum_n \mathcal{A}_{12}^{(n)} + \mathcal{A}_{21}^{(n)} \right|^2 \times \frac{1}{\pi} \frac{\Gamma/2}{(\Omega_1 + \Omega_2 - \omega_{eg})^2 + (\Gamma/2)^2} \frac{\bar{I}_1 \bar{I}_2}{4\epsilon_0^2 c^2},$$

$$\mathcal{A}_{jk}^{(n)} = \frac{\langle e | \hat{\epsilon}_k \cdot \mathcal{D} | n \rangle \langle n | \hat{\epsilon}_j \cdot \mathcal{D} | g \rangle}{\omega_{ng} - \Omega_j + i\Gamma_n/2}$$

- Generalize to allow for Bose-Einstein statistics violation
- Sum over intermediate states $|n\rangle$
- Amplitudes for each path interfere



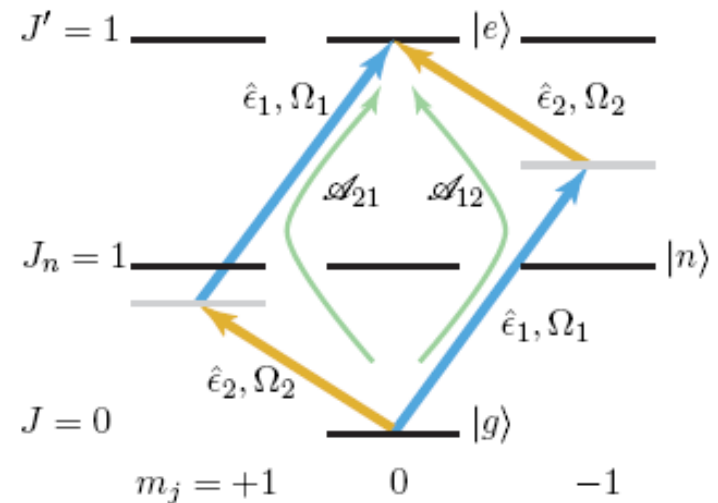
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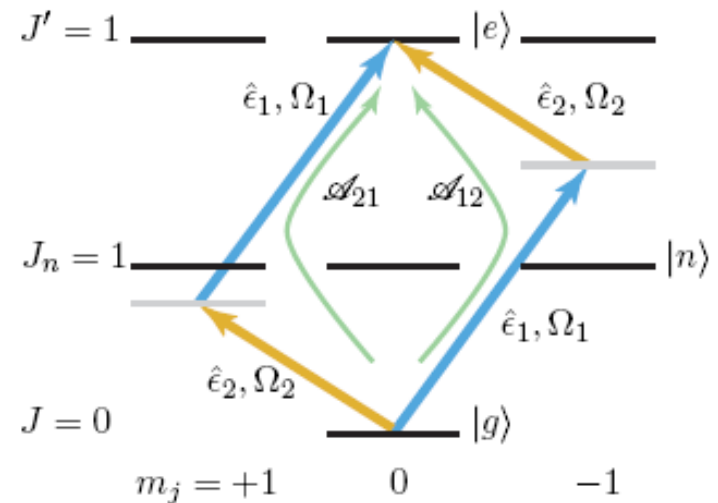
- Consider TPA rate

(2nd order time-dependent perturbation, Fermi's golden rule)

$$W_{\pm} = \frac{2\pi}{\hbar^4} \times \left| \sum_n \mathcal{A}_{12}^{(n)} \pm \mathcal{A}_{21}^{(n)} \right|^2 \times \frac{1}{\pi} \frac{\Gamma/2}{(\Omega_1 + \Omega_2 - \omega_{eg})^2 + (\Gamma/2)^2} \frac{\bar{I}_1 \bar{I}_2}{4\epsilon_0^2 c^2},$$

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- Generalize to allow for Bose-Einstein statistics violation
- Sum over intermediate states $|n\rangle$
- Amplitudes for each path interfere
 - Constructively, $\psi(2\gamma)$ symmetric
 - Destructively, $\psi(2\gamma)$ anti-symmetric (under *particle exchange*)



- Consider TPA rate

(2nd order time-dependent perturbation, Fermi's golden rule)

$$W_{\pm} = \frac{2\pi}{\hbar^4} \times \left| \sum_n \mathcal{A}_{12}^{(n)} \pm \mathcal{A}_{21}^{(n)} \right|^2 \times \frac{1}{\pi} \frac{\Gamma/2}{(\Omega_1 + \Omega_2 - \omega_{eg})^2 + (\Gamma/2)^2} \frac{\bar{I}_1 \bar{I}_2}{4\epsilon_0^2 c^2},$$

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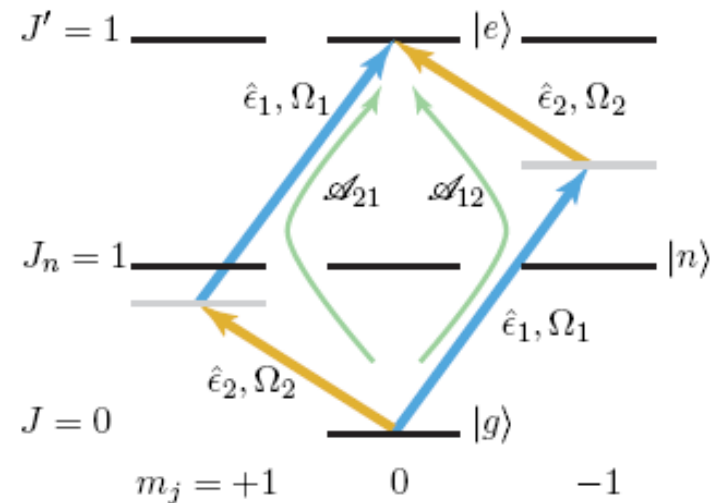
Q1B:

What determines the relative phase between paths? Is this equation valid for all 3 cases of photon states?

3 cases: distinguishable

indistinguishable symm / antisymm

Or am I thinking about this wrong?



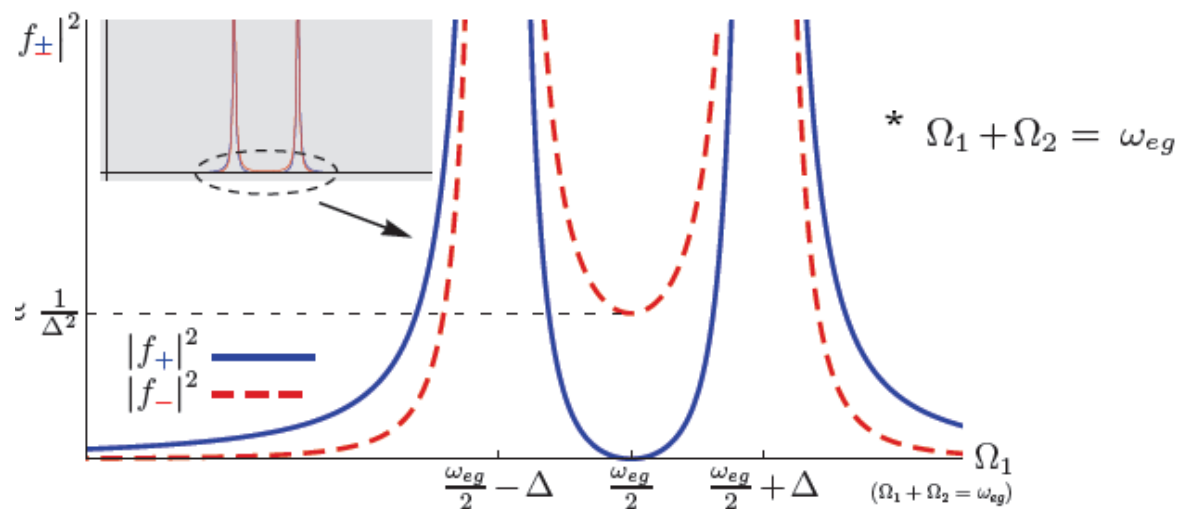
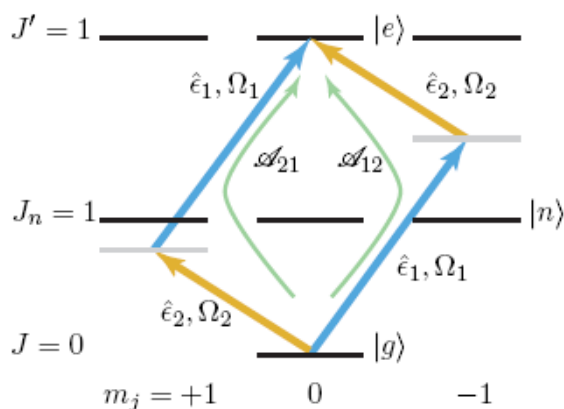
- Sum over intermediate levels, apply Wigner-Eckart theorem extracting the rank-1 irreducible component of the transition operator, and assume orthogonal light polarizations

$$W_{\pm} = |f_{\pm}|^2 \frac{\Gamma/2}{(\Omega_1 + \Omega_2 - \omega_{eg})^2 + (\Gamma/2)^2} \frac{\mathcal{D}_{en}^2 \mathcal{D}_{ng}^2 \bar{I}_1 \bar{I}_2}{3\epsilon_0^2 c^2 \hbar^4},$$

$$f_{+} = \frac{(\Omega_1 - \Omega_2)/2}{(\omega_{ng} - \Omega_1 + i\Gamma_n/2)(\omega_{ng} - \Omega_2 + i\Gamma_n/2)},$$

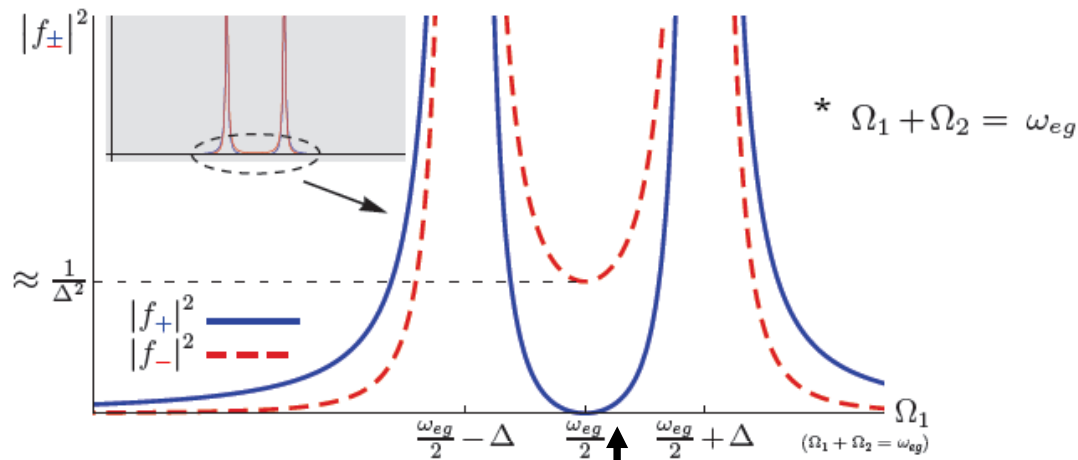
$$f_{-} = \frac{\omega_{ng} - (\Omega_1 + \Omega_2)/2 + i\Gamma_n/2}{(\omega_{ng} - \Omega_1 + i\Gamma_n/2)(\omega_{ng} - \Omega_2 + i\Gamma_n/2)},$$

Q2: What symmetry is leading to this selection rule for the symmetric state?



- Measured fluorescence signal:

$$\mathcal{S}(\delta) = \gamma \{W_+(\delta) + \nu W_-(\delta)\}$$



Assume ν small, Taylor expand about $\delta=0$

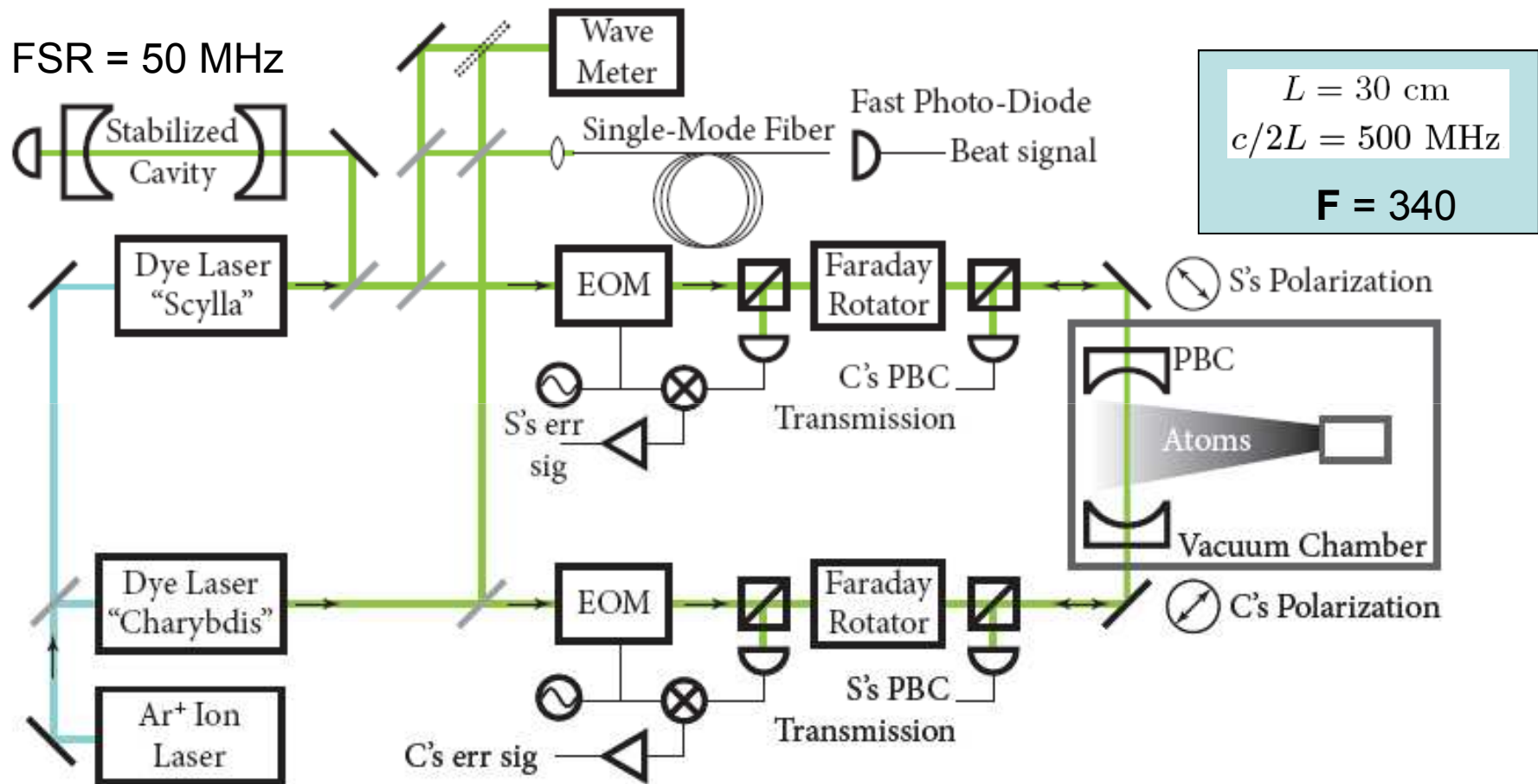
$$W_-(\delta) = \frac{2 \mathcal{D}_{en}^2 \mathcal{D}_{ng}^2 \bar{I}_1 \bar{I}_2}{\Gamma \hbar^4 \Delta^2 \epsilon_0^2 c^2} \left(1 - \frac{\delta^2}{\Delta^2}\right)^{-2}$$

$$W_+(\delta) = \frac{\delta^2}{\Delta^2} W_-(\delta).$$

$$\mathcal{S}_{\text{cal}} = \mathcal{S}(\delta_{\text{cal}}) = \gamma W_-(0) \frac{\delta_{\text{cal}}^2}{\Delta^2} + \mathcal{O}\left(\frac{\delta_{\text{cal}}^4}{\Delta^4}\right)$$

$$\mathcal{S}_{\text{lim}} = \mathcal{S}(0) = \gamma \nu W_-(0)$$

$$\nu = \frac{\mathcal{S}_{\text{lim}}}{\mathcal{S}_{\text{cal}}} \frac{\delta_{\text{cal}}^2}{\Delta^2}.$$



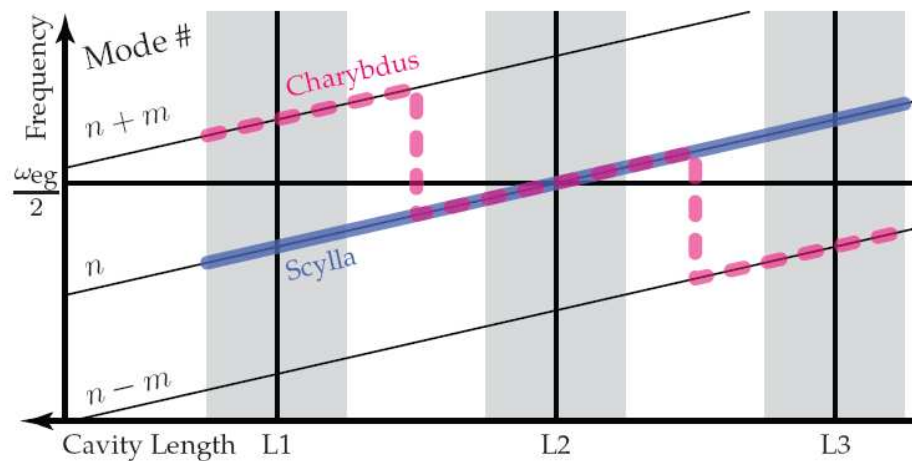
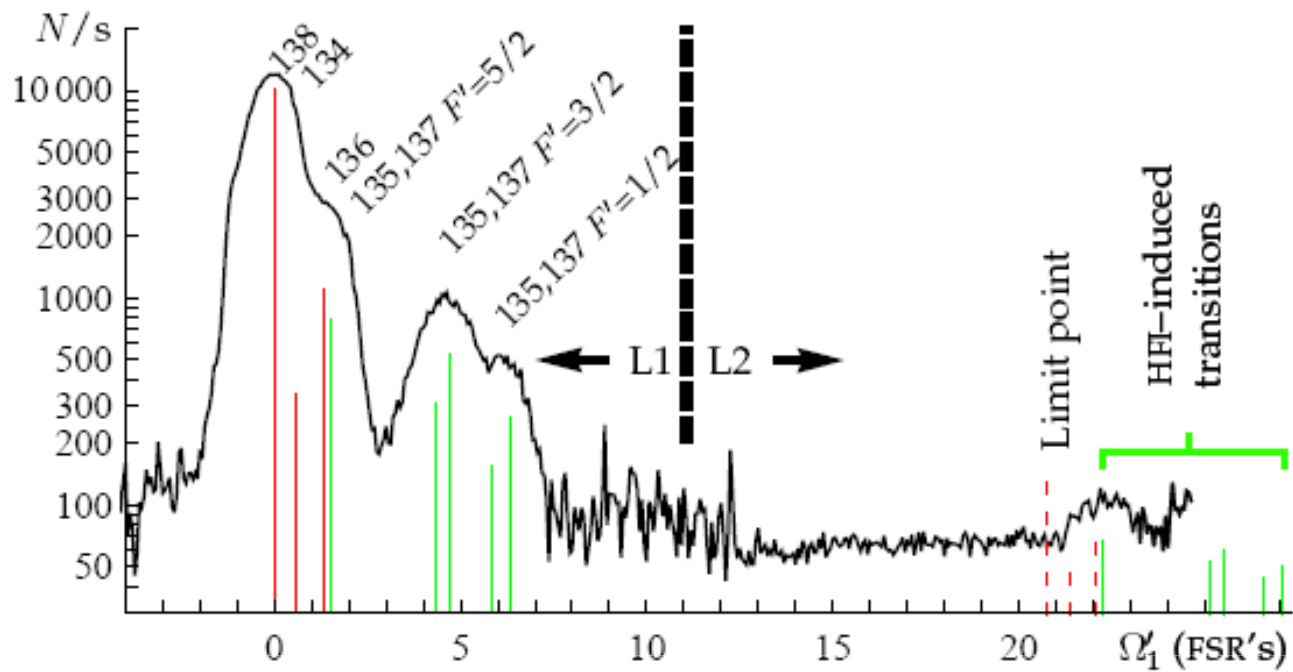
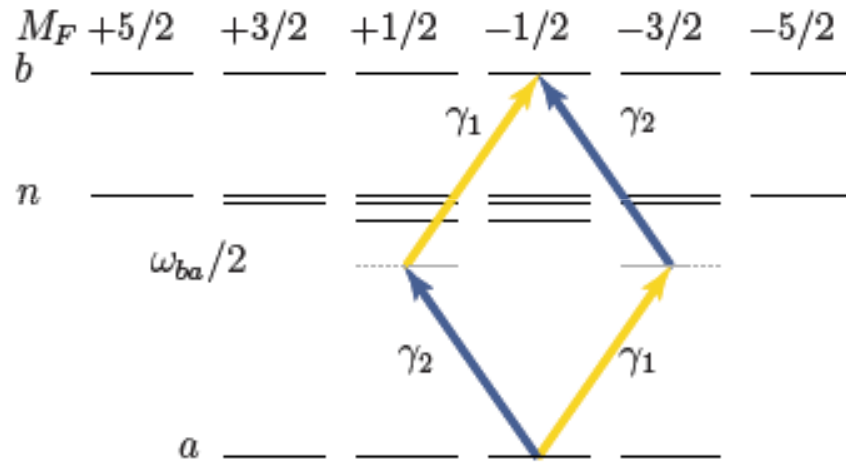
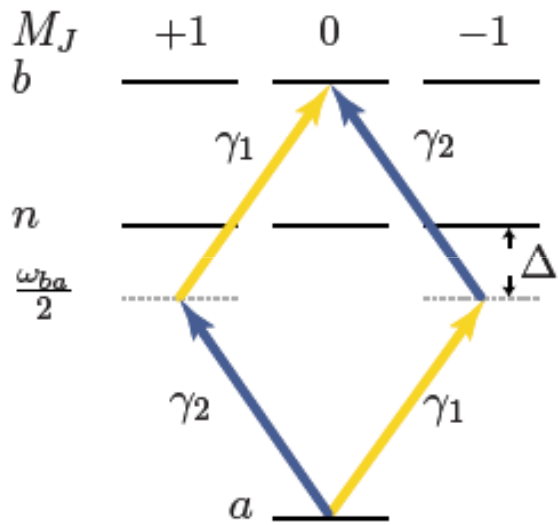


FIG. 5. (color online). Laser tuning path during run. The lasers are tuned in concert with the PBC, either separated by $m=4$ longitudinal cavity modes, or in the same mode.

HFI-induced transitions



$$\mathcal{S}(\delta) = \gamma \{W_+(\delta) + \nu W_-(\delta)\}$$

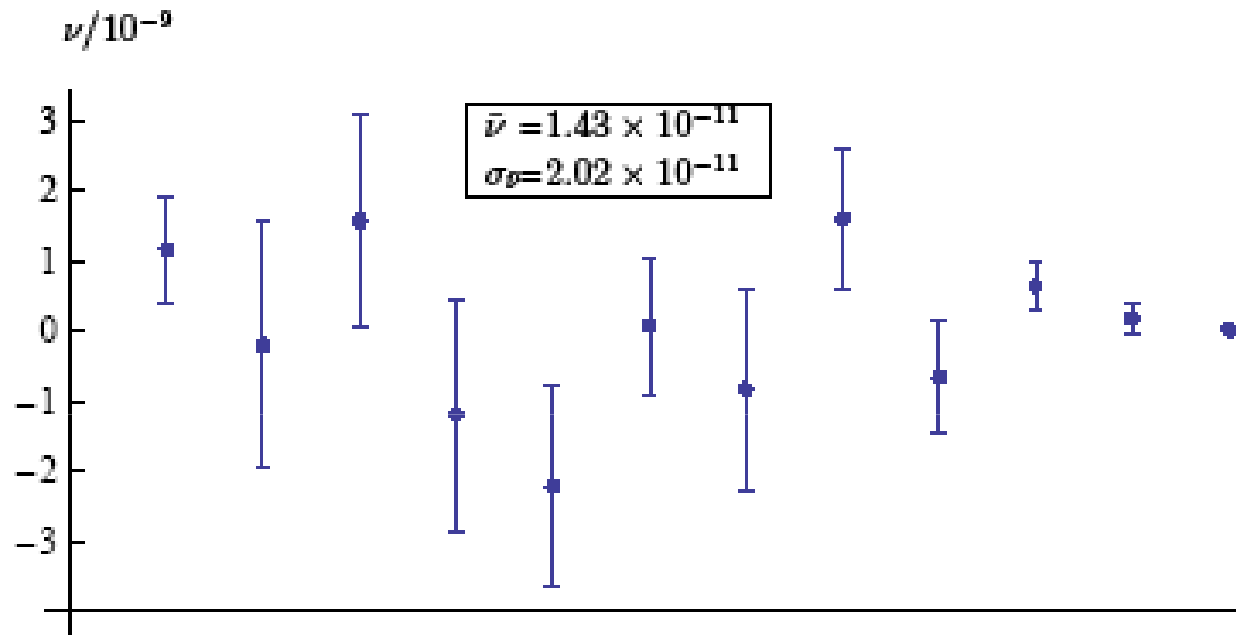


FIG. 7. (color online). ν from twelve runs. The error bars decrease over the lifetime of the experiment as improvements in excitation and detection efficiency were achieved.

Systematic errors

- False positives
 - Ba-137 (11%) and Ba-135 (6.6%), $I \neq 0$
 - Hyperfine splitting breaks symmetry between paths leading to imperfect destructive interference. $\approx 10^{-9}$
 - Zeeman splitting due to b/g magnetic fields $\approx 10^{-11}$ @ 10 G
 - Non-zero spectral width of lasers (3 MHz) $\approx 10^{-12}$
 - Line broadening due to transverse atomic motion (13 MHz) $\approx 10^{-15}$
 - Light shifts $\approx 10^{-12}$
 - E1M2, E2M1 transitions $\approx 10^{-19}$

Conclusion

- Photons: still bosons*.

* to 1 part in 10^{11}

(I think...)