

# What drives the annual cycle in tropical tropopause temperatures?

Peter Haynes

University of Cambridge  
Department of Applied Mathematics and Theoretical Physics (DAMTP)

Cambridge Centre for Climate Science (CCfCS)

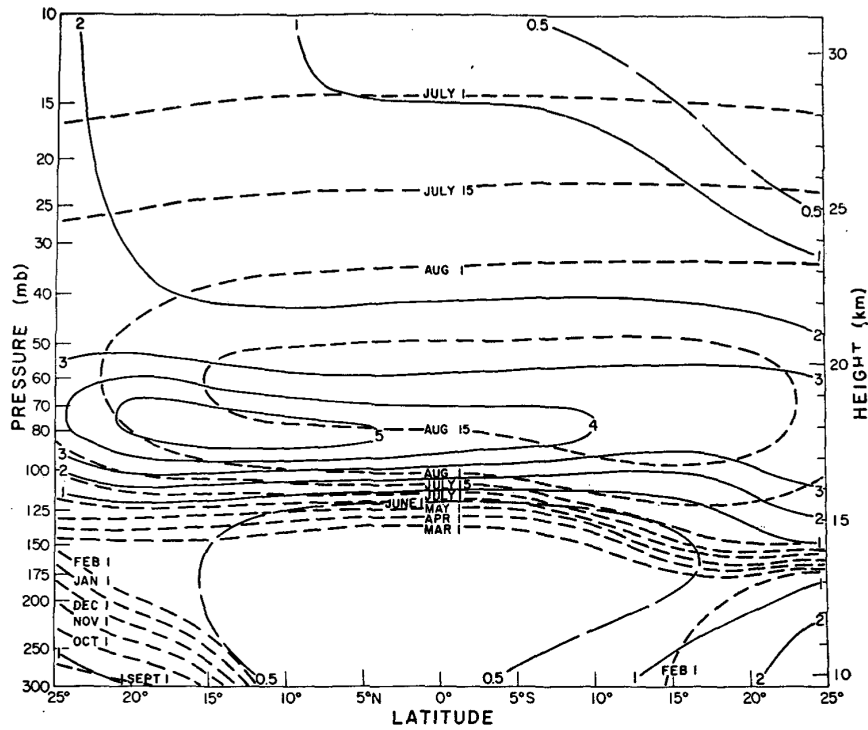
Université Fédérale Toulouse Midi-Pyrénées

(PHH acknowledges collaborations with Stephan Fueglistaler, Alison Ming, Peter Hitchcock, Amanda Maycock and others)

**DAMTP**

**CCfCS**

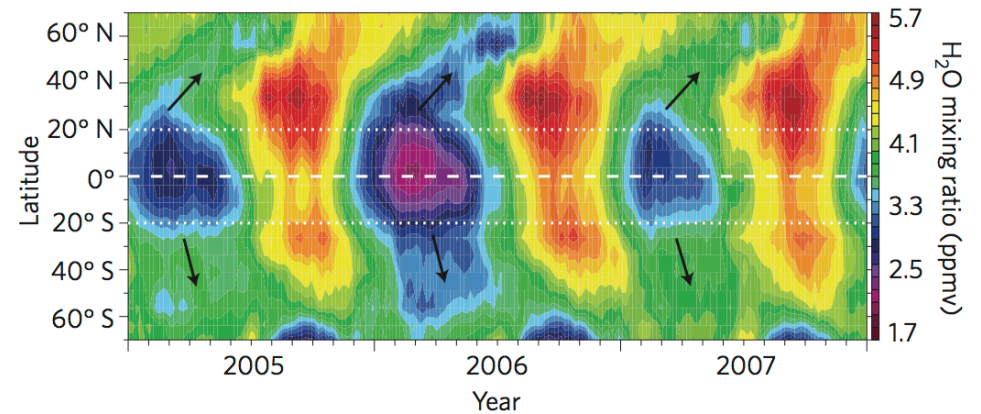




Reed and Vlcek (1969)

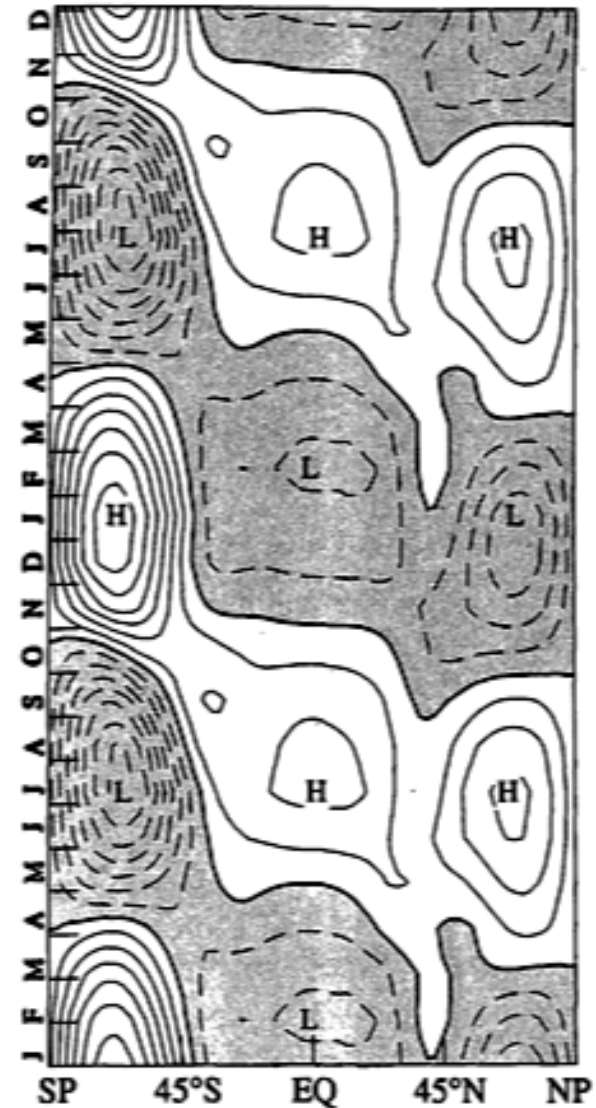
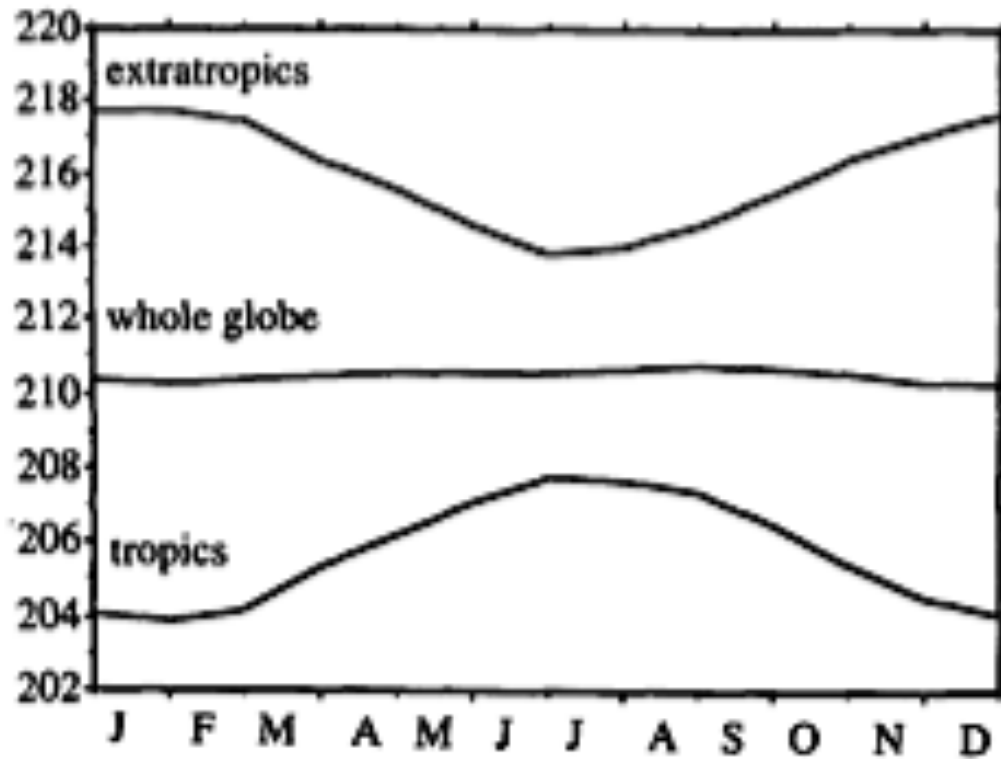
MLS annual cycle in water vapour

Randel and Jensen (2013)



# Yulaeva et al (1994)

Annual cycle in 100mb temperatures  
(MSU channel 4)



# Transformed Eulerian-mean equations

(Eliassen-Palm flux divergence appears as complete eddy forcing of mean flow)

$$\bar{u}_t + \frac{1}{a \cos \phi} \bar{v}^* \partial_\phi (\bar{u} \cos \phi) + \bar{w}^* \partial_z \bar{u} - f \bar{v}^* = \frac{1}{\rho_0 a \cos \phi} \nabla \cdot \vec{F}, \quad (1)$$

$$f \partial_z \bar{u} + \frac{R}{H} \frac{1}{a} \partial_\phi \bar{T} = 0, \quad (2)$$

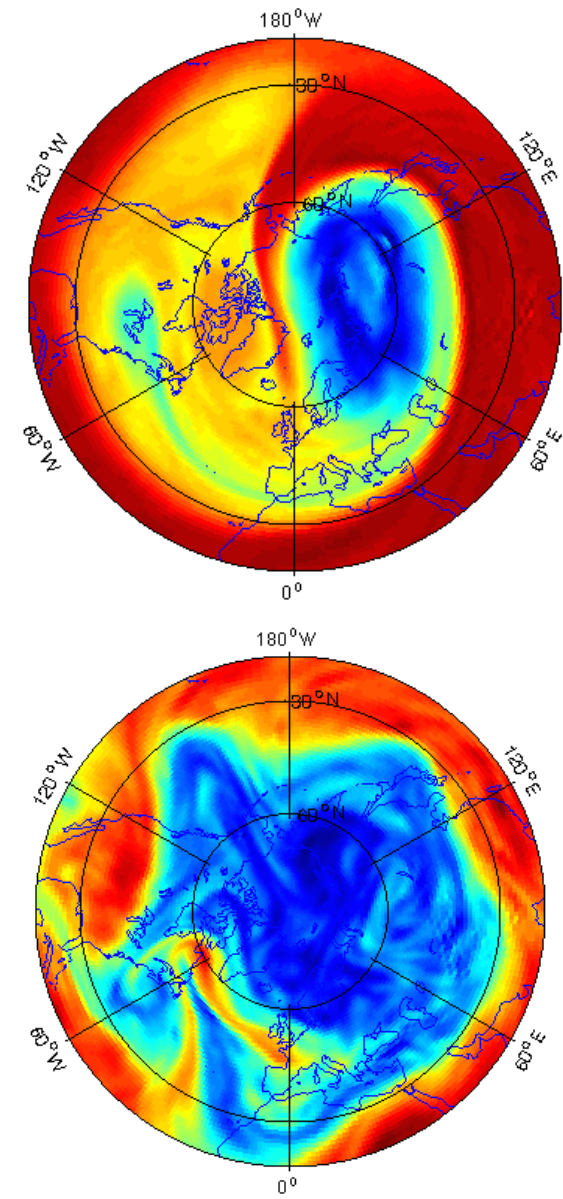
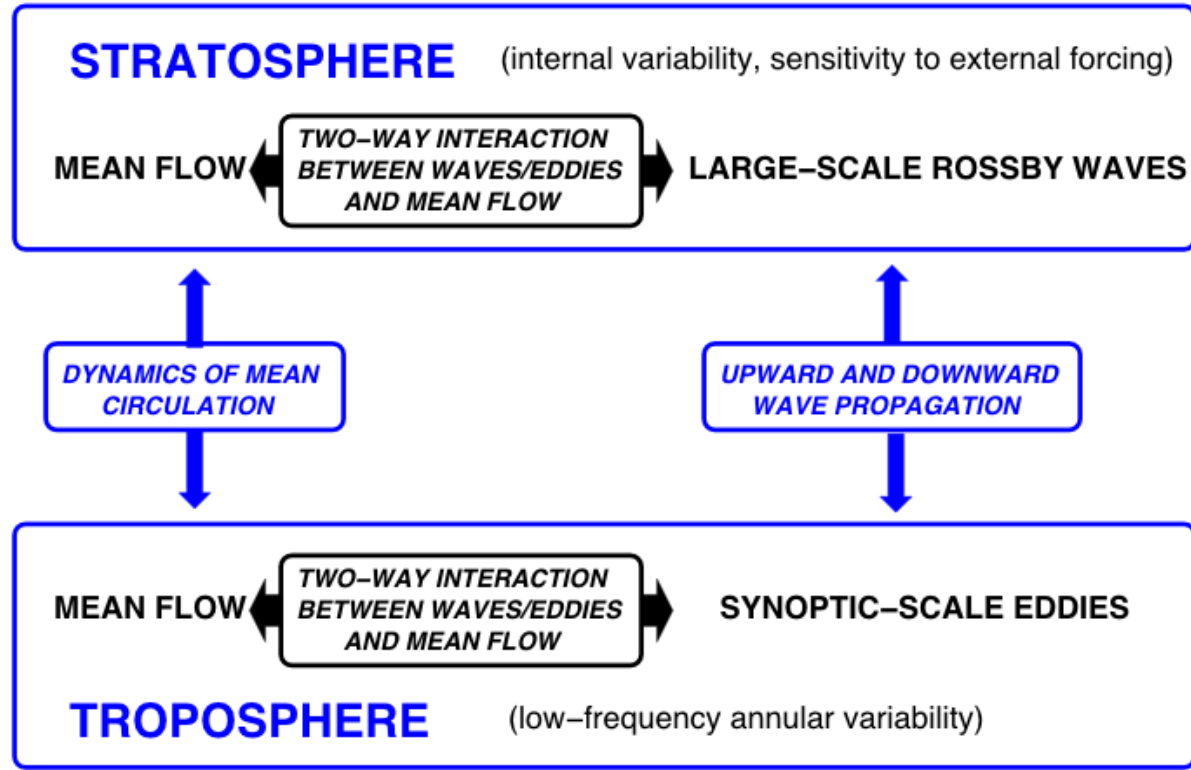
$$\frac{1}{a \cos \phi} \partial_\phi (\bar{v}^* \cos \phi) + \frac{1}{\rho_0} \partial_z (\rho_0 \bar{w}^*) = 0, \quad (3)$$

$$\bar{T}_t + \frac{1}{a} \bar{v}^* \partial_\phi \bar{T} + \bar{S} \bar{w}^* = \bar{Q} - \alpha \bar{T}. \quad (4)$$

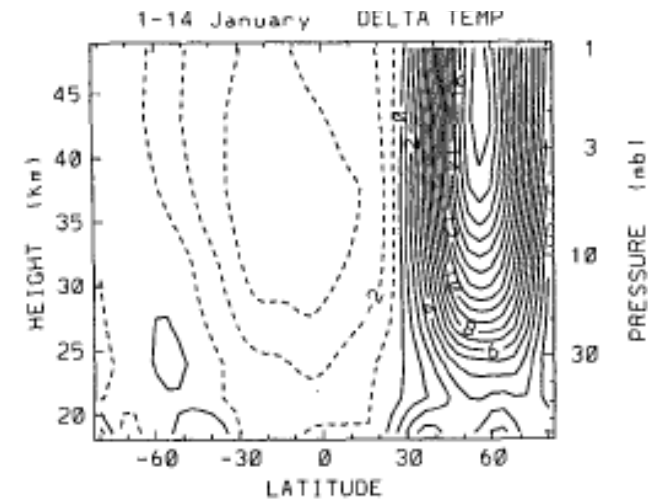
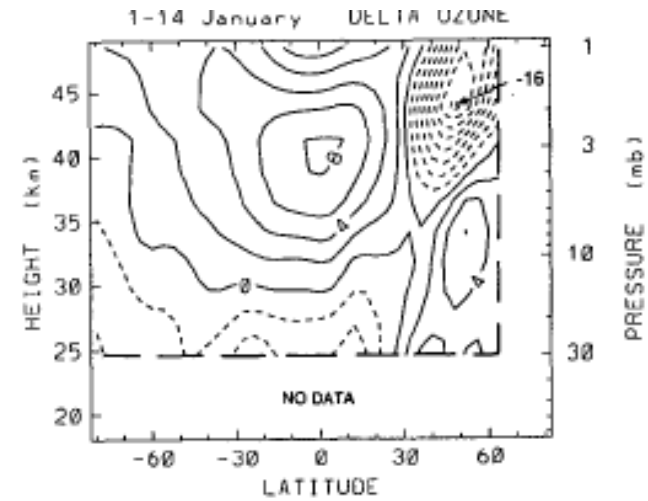
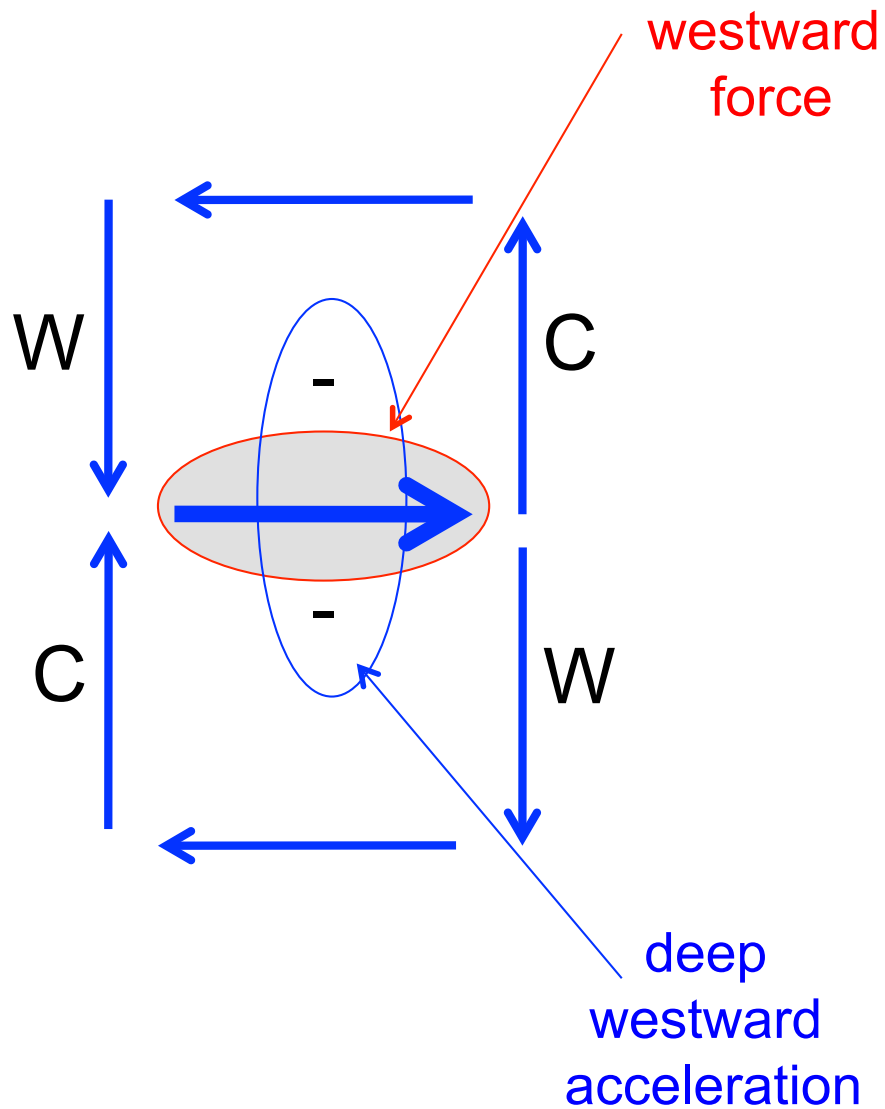
simple representation of radiative transfer

Forcing gives rise to response in both velocity and meridional circulation – proportion of each depends on shape of forcing distribution and on timescale (relative to radiative timescale)





# Schematic response to localised force



Stratospheric wave event (Randel 1993)



# Effect of radiative damping

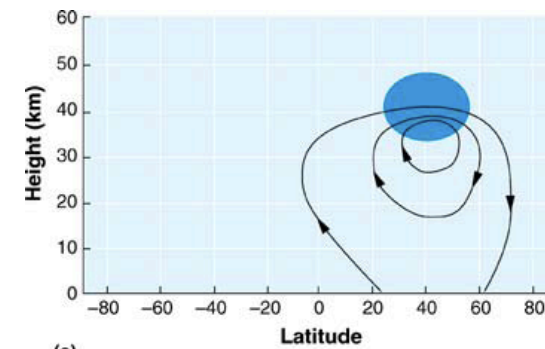
Response to localised forcing with frequency  $\omega$

$\omega \gg \alpha$   
short time scales

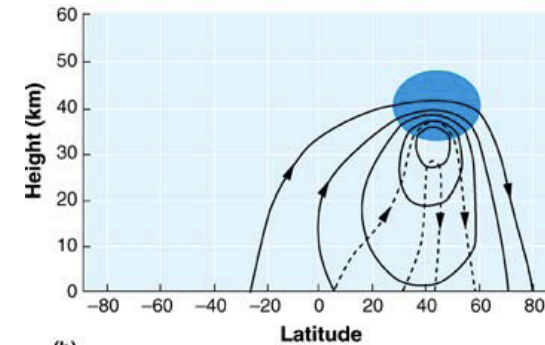
$\omega \sim \alpha$

$\omega \ll \alpha$   
long time scales

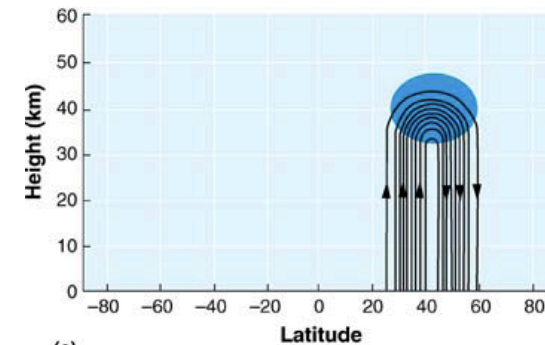
Holton et al (1995)



(a)



(b)



(c)



## Steady state limit

$$\overline{w^*}(\phi, z) = \frac{1}{a\rho_0(z) \cos \phi} \frac{\partial}{\partial \phi} \left[ \int_z^\infty \frac{\nabla \cdot \mathbf{F}}{2\Omega a \sin \phi} \Big|_{\phi=\text{const.}} dz' \right]$$

(Haynes et al, 1991)

Steady upwelling at any location  $(\phi, z)$  is determined by wave force above that location – ‘downward control principle’. Steady circulation *requires* wave force.

“Low latitude singularity”.

Does it make any sense to consider  $\nabla \cdot \mathbf{F}$  as specified?

Time averaged BDC is controlled by time average wave force.

Time averaged BDC upwelling in tropics is controlled by time average wave force in subtropics



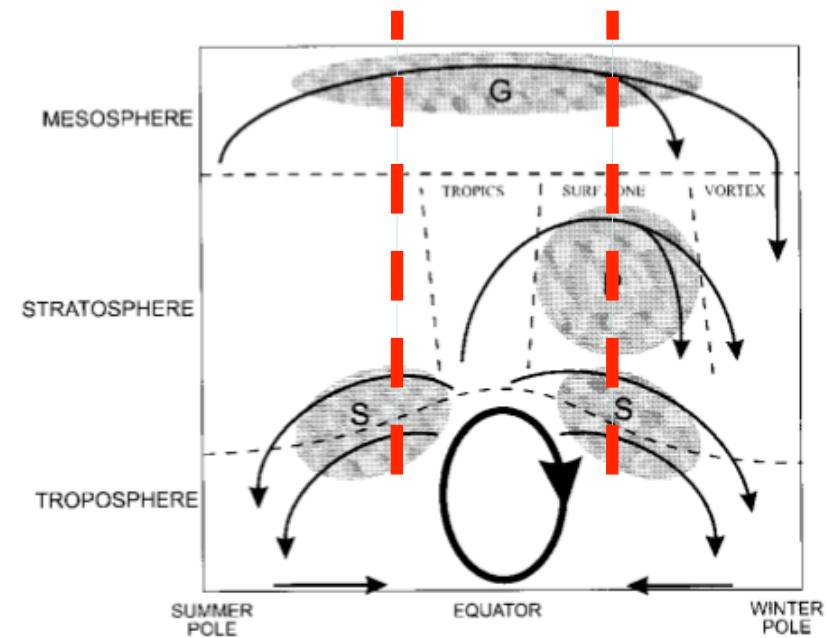


# Annual cycle in wave-driven circulation

Rosenlof and Holton (1993)

$$\bar{w}^*(\phi, z) = \frac{1}{a\rho_0(z) \cos \phi} \frac{\partial}{\partial \phi} \left[ \int_z^\infty \frac{\nabla \cdot \mathbf{F}}{2\Omega a \sin \phi} \Big|_{\phi=\text{const.}} dz' \right]$$

	UKMO	Holton [1990]
DJF	114.0	93.3
MAM	76.4	56.1
JJA	55.8	47.2
SON	70.3	62.4
Mean	79.1	64.8

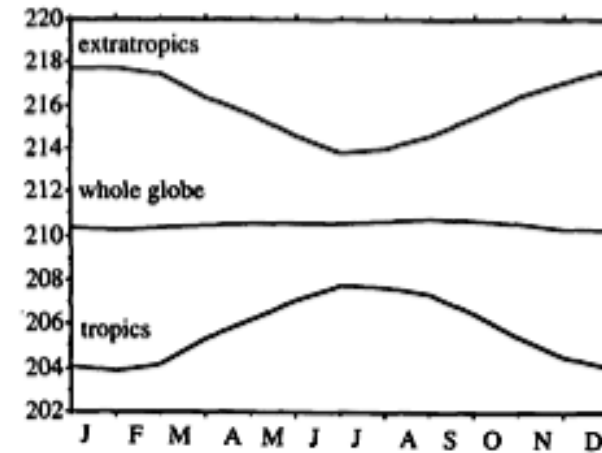


Plumb (2002)



## Yulaeva et al (1994)

$$\frac{\partial \bar{T}}{\partial t} + \underbrace{\bar{w}^* \bar{S}}_{-\bar{Q}_{\text{dyn}}} = -\alpha(\bar{T} - T_0) + \underbrace{\bar{Q}_{\text{rad}}}_{\text{radiation}}$$



Global Average

$$\frac{\partial \langle \bar{T} \rangle_{\text{global}}}{\partial t} + \langle \bar{w}^* \bar{S} \rangle_{\text{global}} = -\alpha(\langle \bar{T} \rangle_{\text{global}} - \langle T_0 \rangle_{\text{global}})$$

Driven by annual variation in ozone heating in tropical stratosphere?

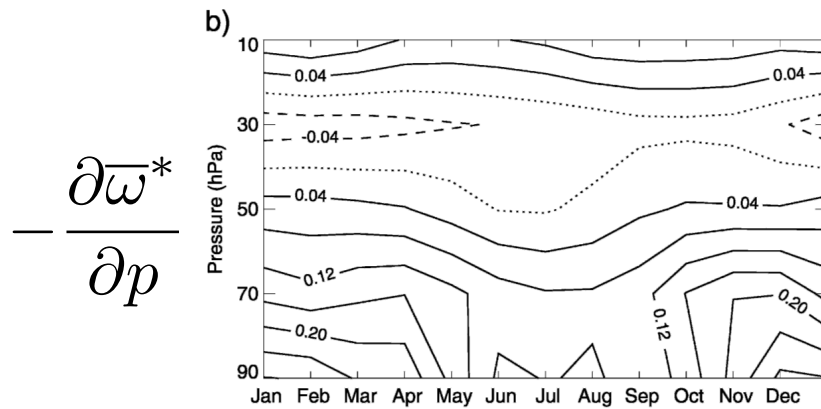
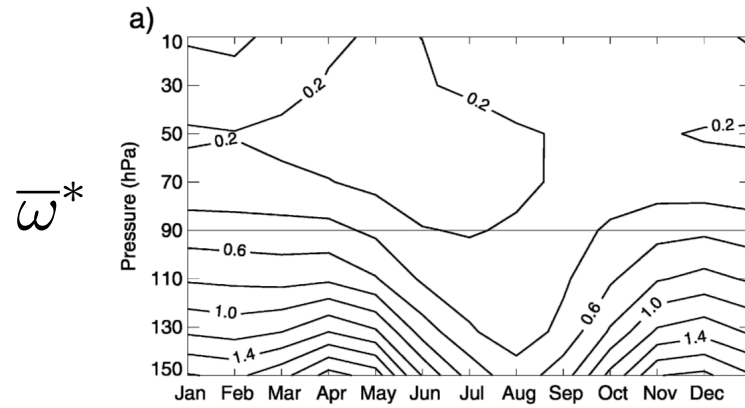
Tropics - Extratropics

$$\frac{\partial \langle \bar{T} \rangle_{T-E}}{\partial t} + \langle \bar{w}^* \bar{S} \rangle_{T-E} = -\alpha(\langle \bar{T} \rangle_{T-E} - \langle T_0 \rangle_{T-E})$$

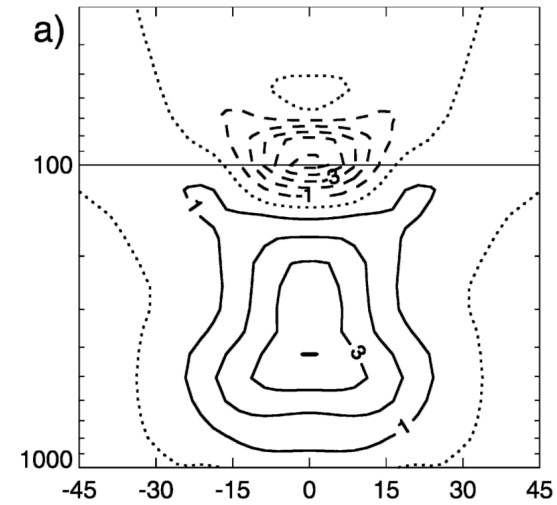
Dynamical response to annual variation in wave fluxes in extratropical stratosphere?



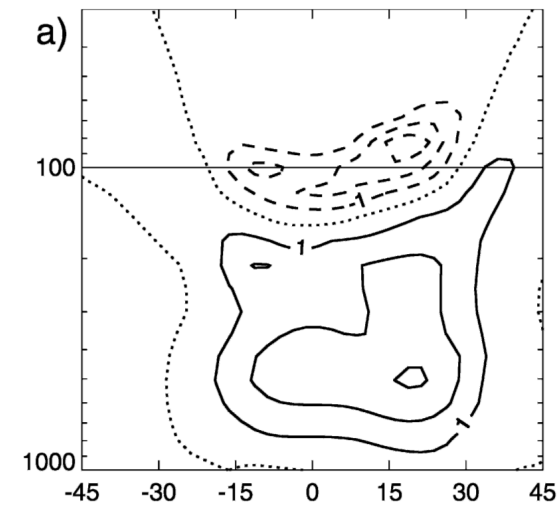
# Role of low-latitude waves in $\nabla \cdot \mathbf{F}$ ?



Kerr-Munslow and Norton (2006)



Equatorial  
localised  
heating



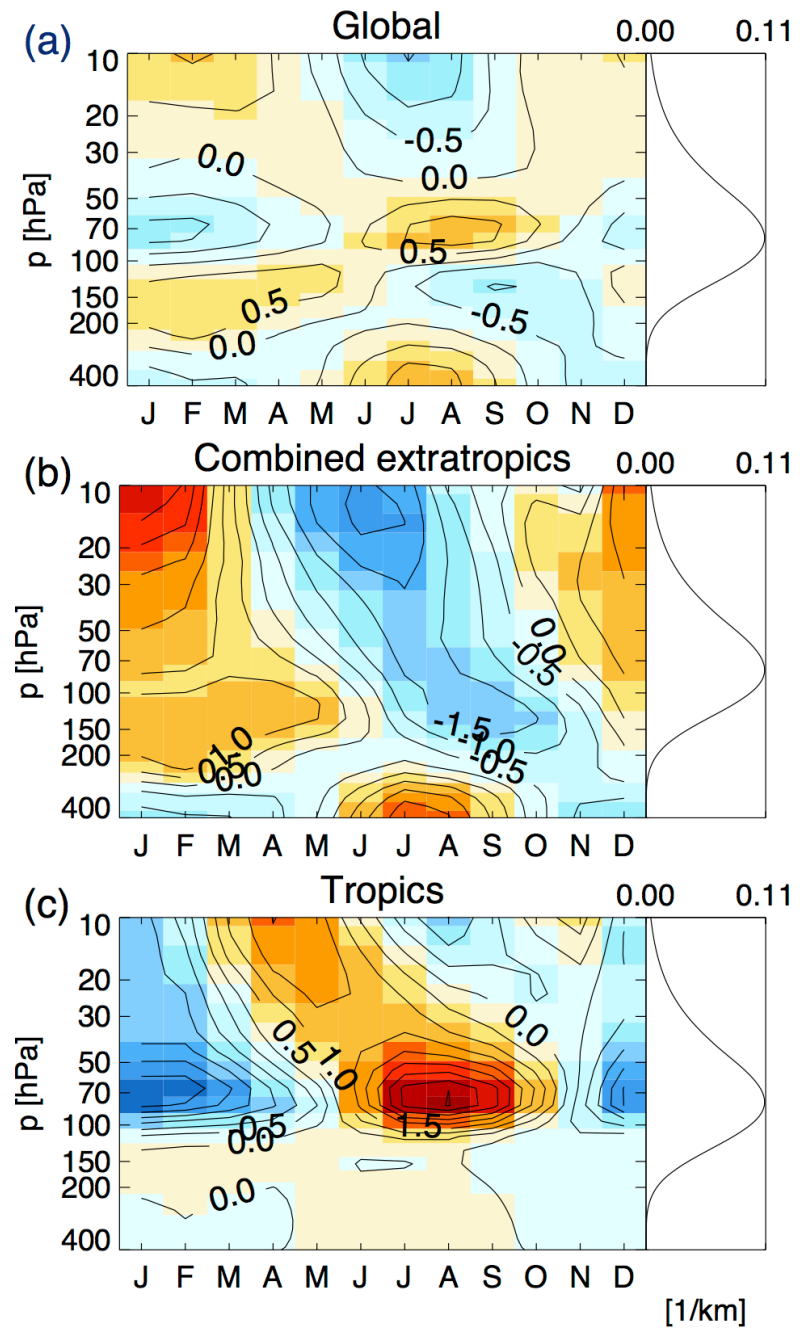
Off-equatorial  
localised  
heating

Norton (2006)

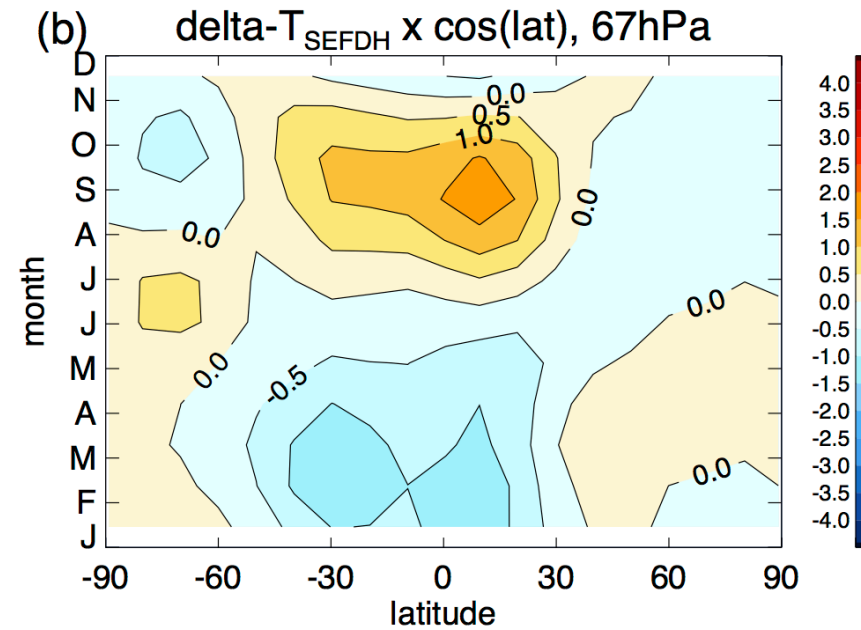
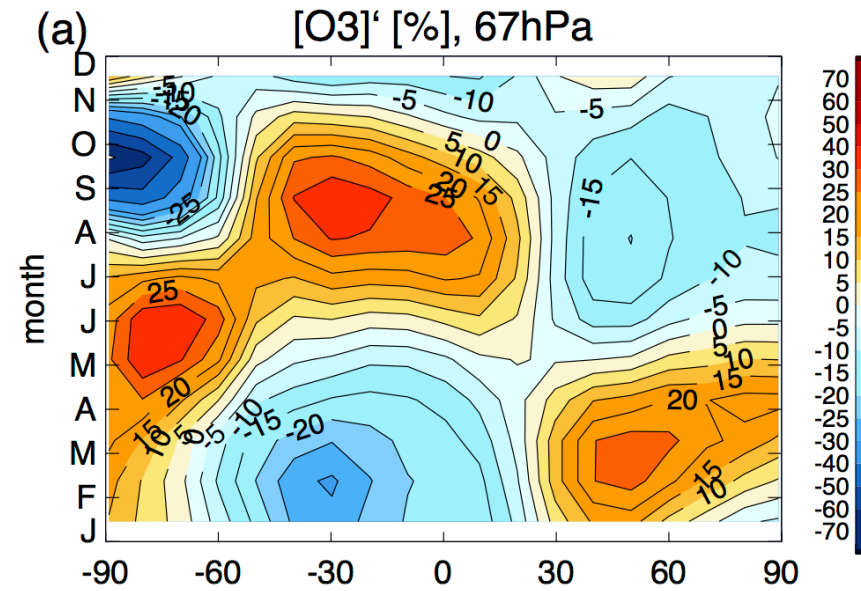




# Fueglistaler et al (2011)



# Ozone-driven temperature annual cycle

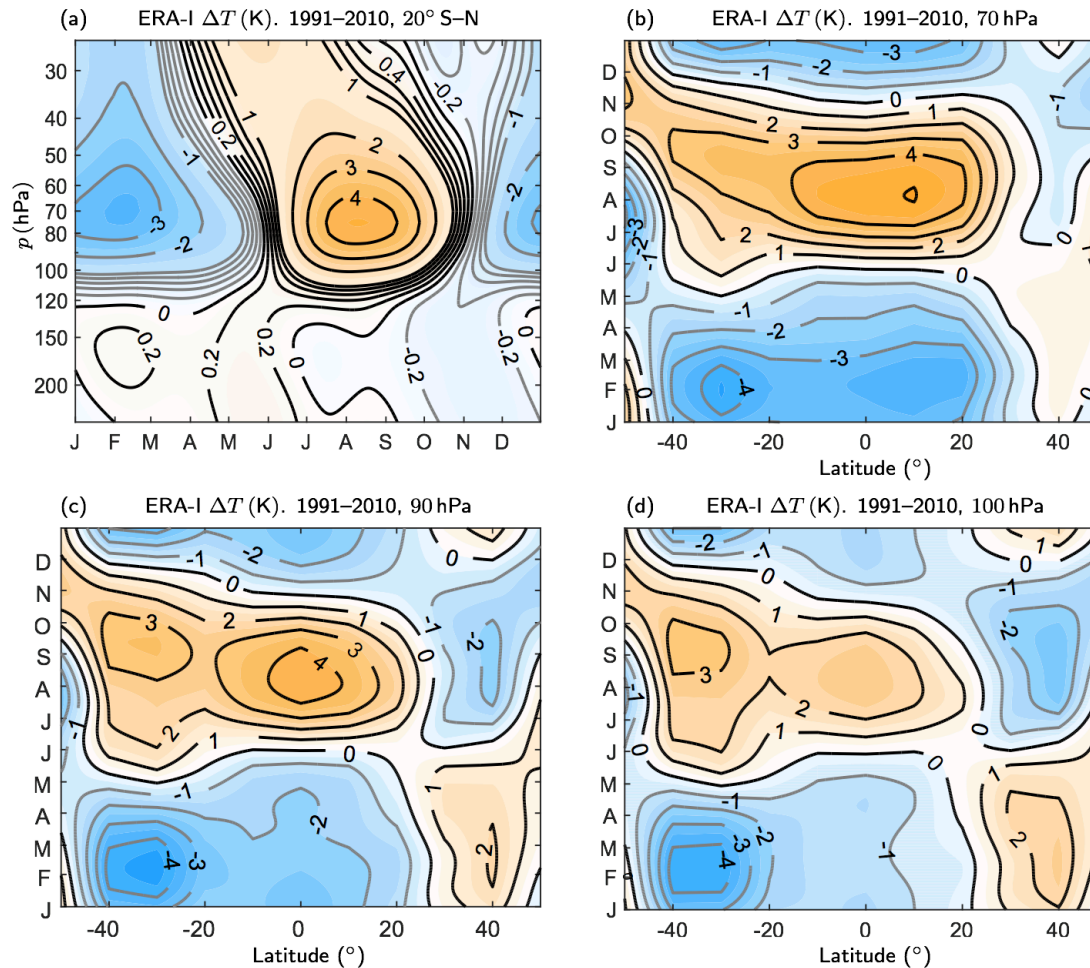


Fueglistaler et al (2011)



# Re-examination of driving of annual cycle

Ming, Maycock,  
Hitchcock and H  
(2017)



amplitude  
largest at  
70hPa

90-100 hPa  
relevant for  
dehydration

What physical processes determine amplitude and structure (vertical and horizontal)?



# SEFDH calculation

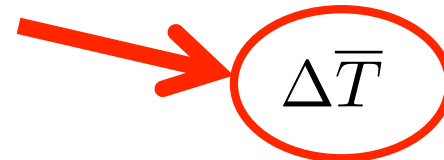
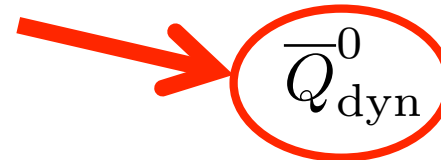
$$\frac{\partial \bar{T}}{\partial t} + \bar{w}^* \bar{S} = Q_{\text{rad}}(\bar{T}, \bar{\chi}_{\text{O}_3}, \bar{\chi}_{\text{H}_2\text{O}})$$

specified *t*-varying  $\bar{T}^0 \quad \bar{\chi}_{\text{O}_3}^0 \quad \bar{\chi}_{\text{H}_2\text{O}}^0$

$$\frac{\partial \bar{T}^0}{\partial t} = Q_{\text{rad}}(\bar{T}^0, \bar{\chi}_{\text{O}_3}^0, \bar{\chi}_{\text{H}_2\text{O}}^0) + \bar{Q}_{\text{dyn}}^0$$

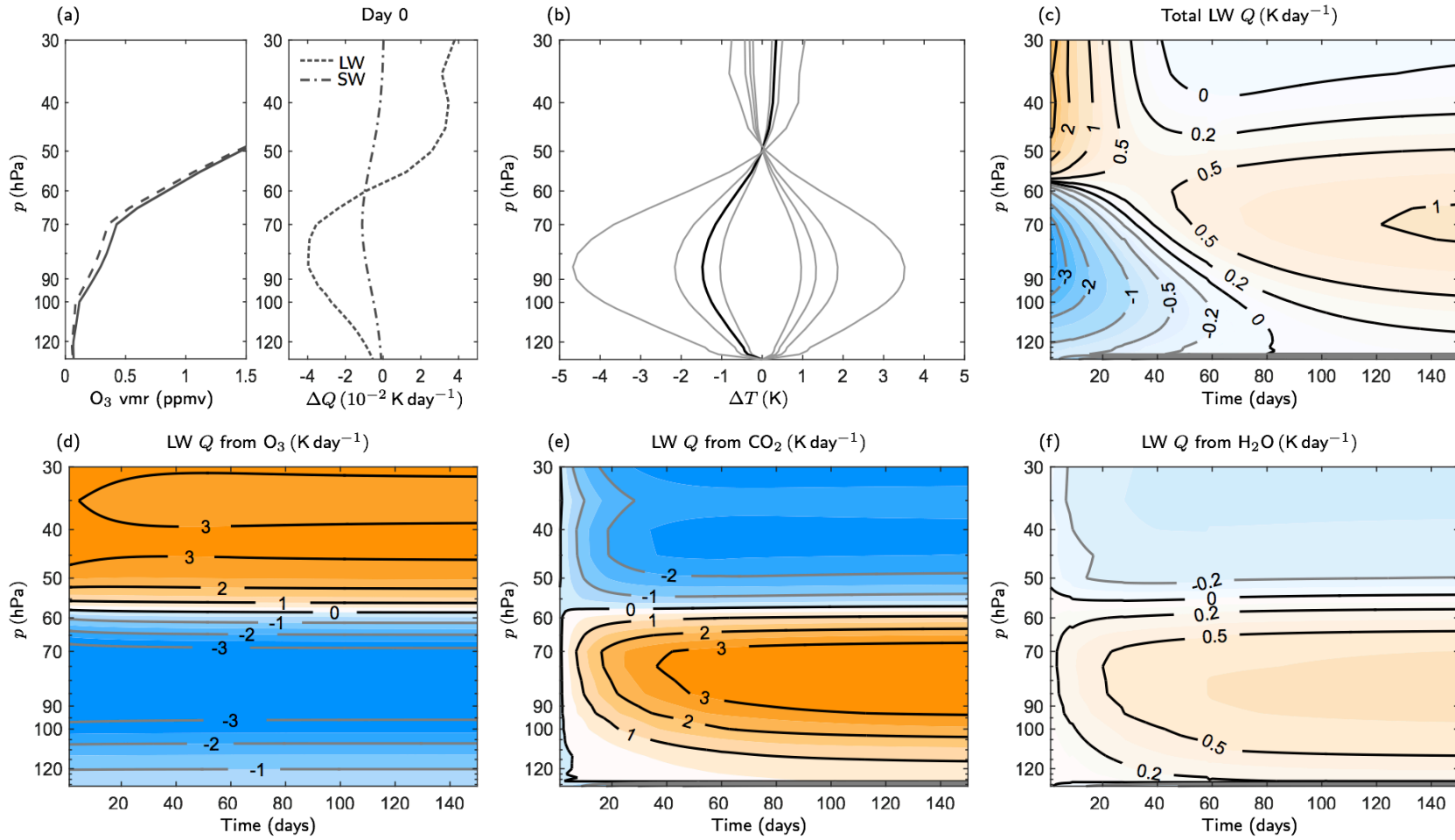
specified *t*-varying  $\Delta \bar{\chi}_{\text{O}_3}^0 \quad \Delta \bar{\chi}_{\text{H}_2\text{O}}^0$

$$\frac{\partial(\bar{T}^0 + \Delta \bar{T})}{\partial t} = Q_{\text{rad}}(\bar{T}^0 + \Delta \bar{T}, \bar{\chi}_{\text{O}_3}^0 + \Delta \bar{\chi}_{\text{O}_3}^0, \bar{\chi}_{\text{H}_2\text{O}}^0 + \Delta \bar{\chi}_{\text{H}_2\text{O}}^0) + \bar{Q}_{\text{dyn}}^0$$

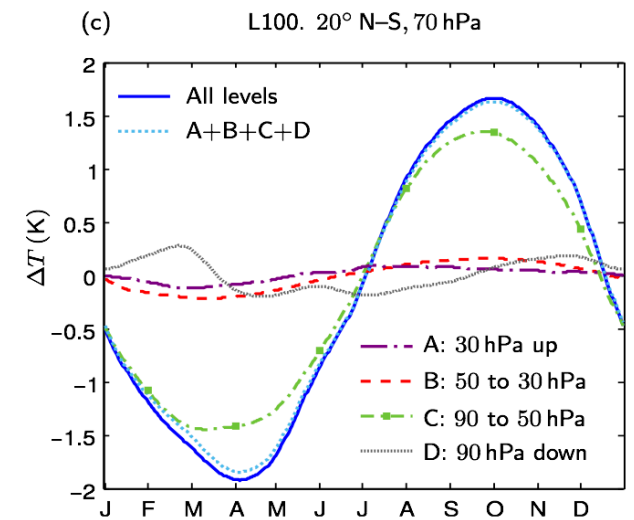
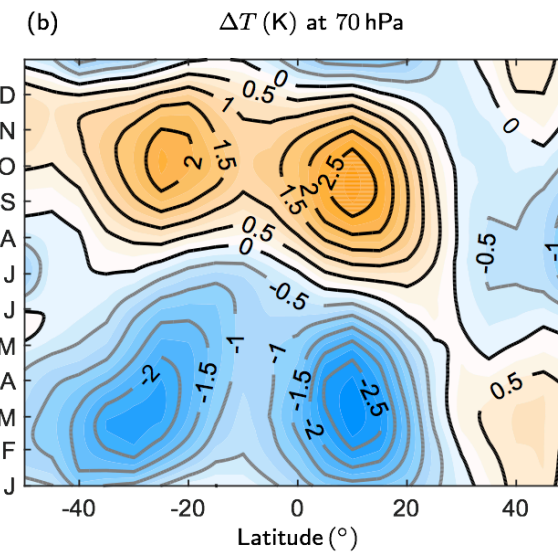
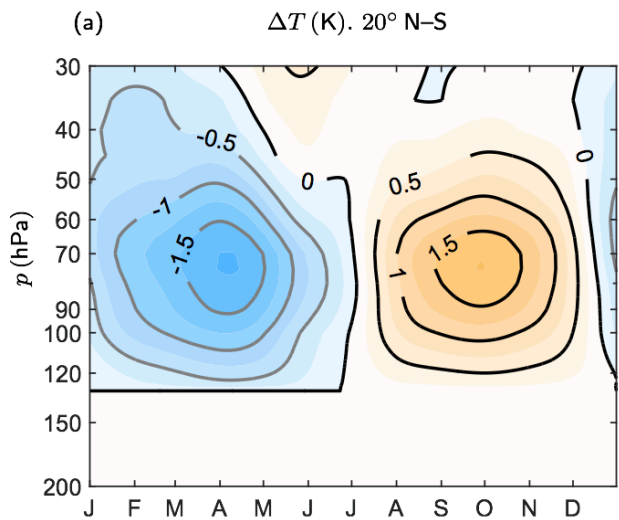
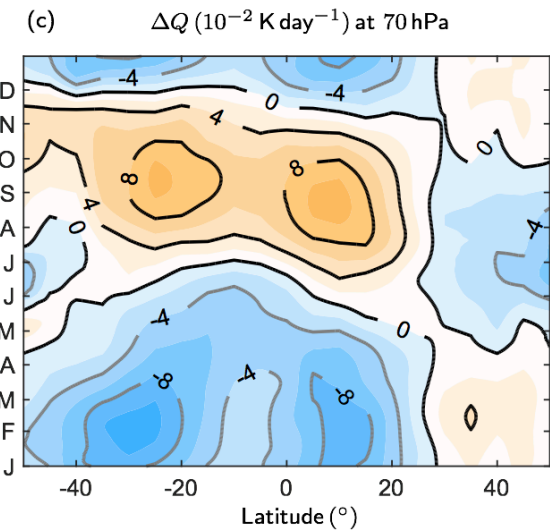
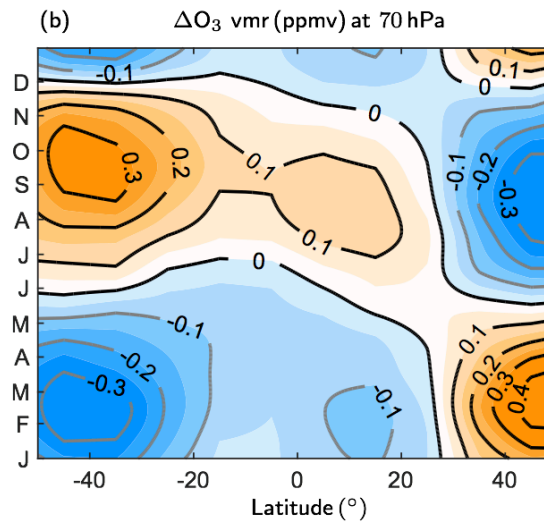
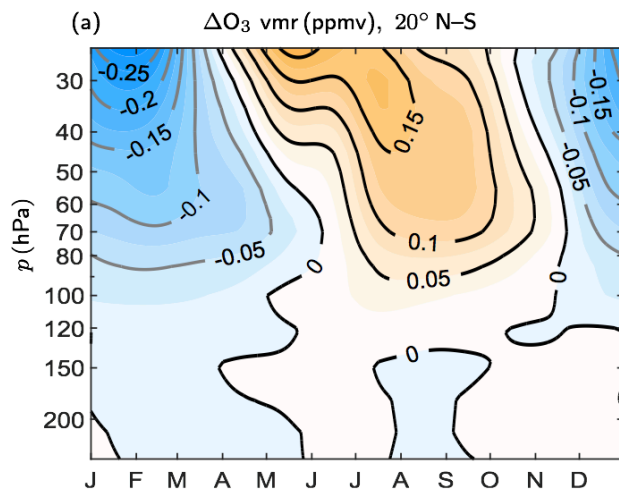




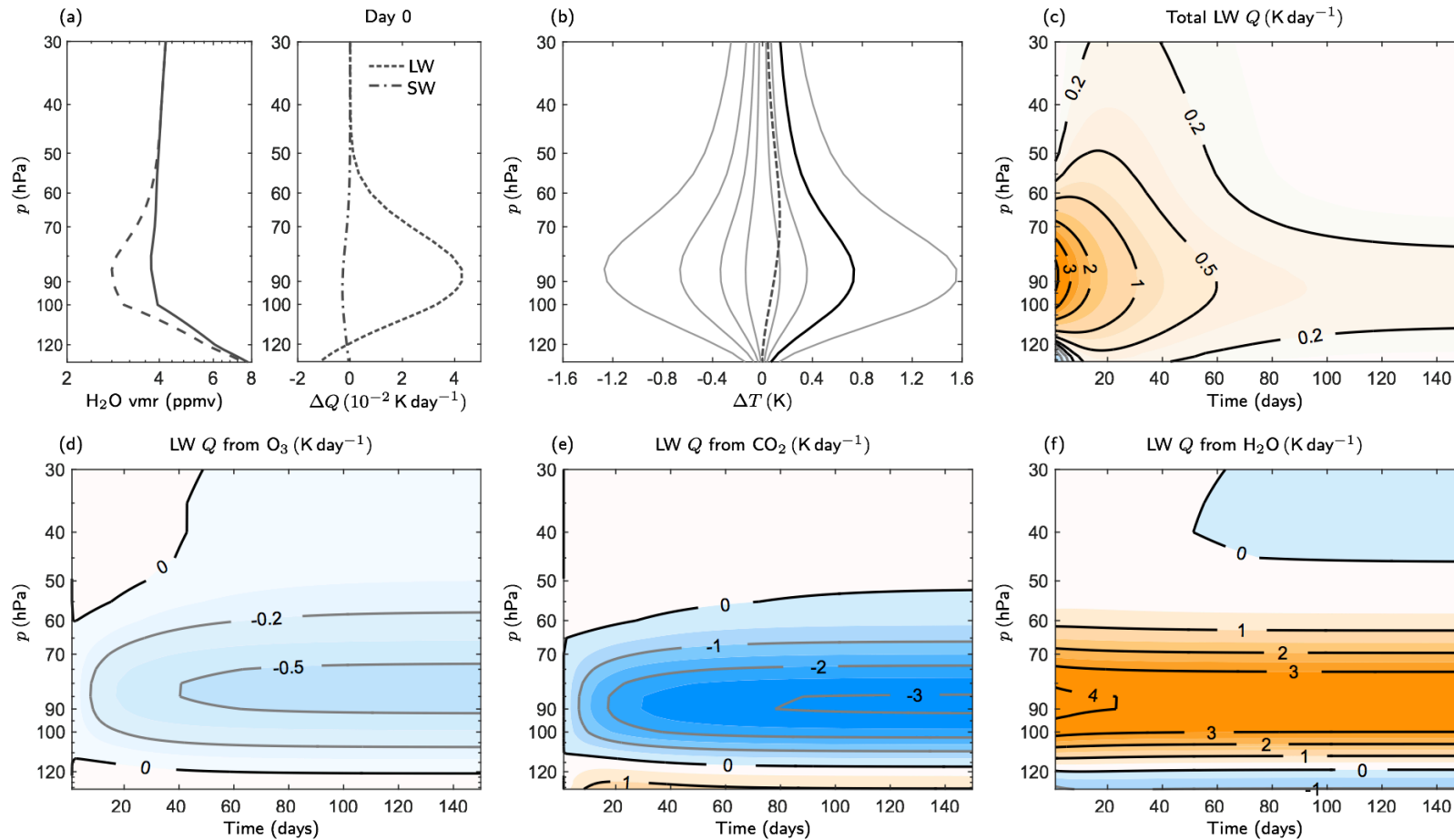
# Effect of imposed $\Delta \bar{\chi}_{\text{O}_3}^0$



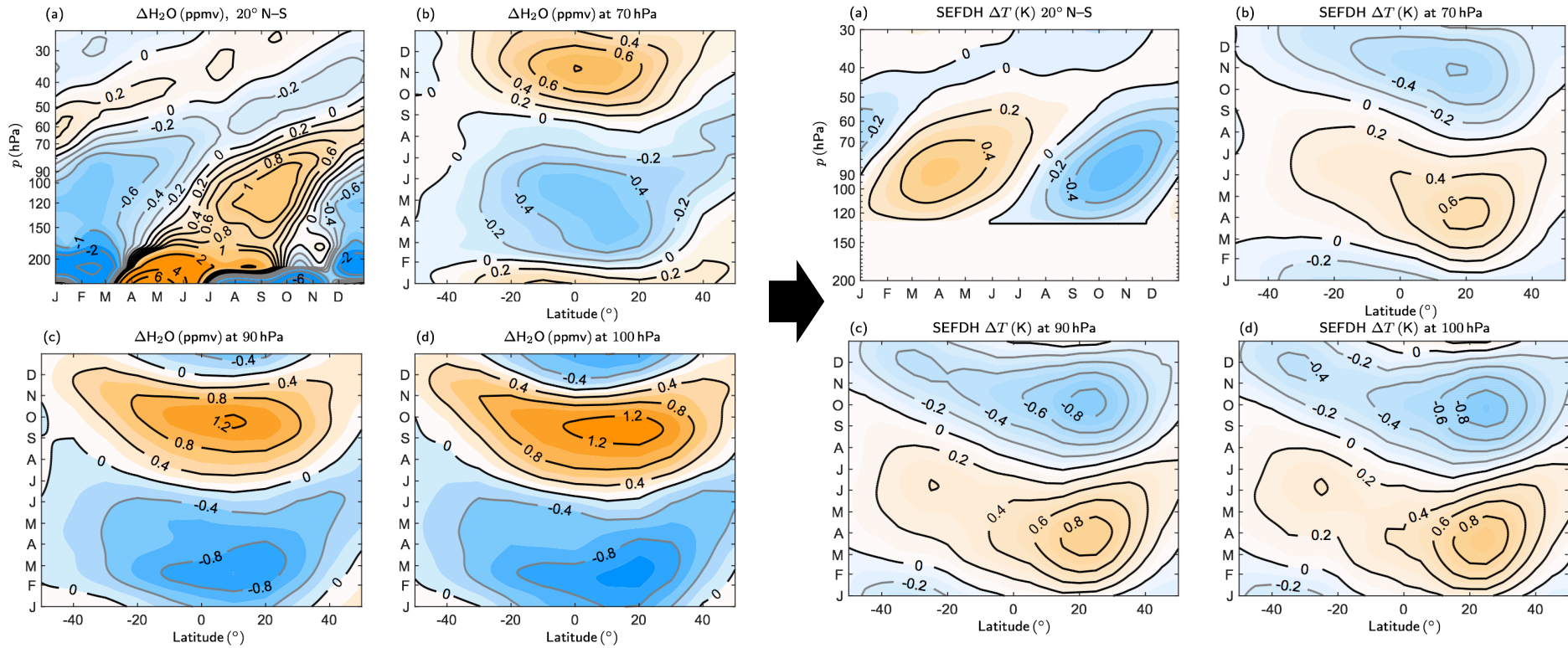
# Ozone



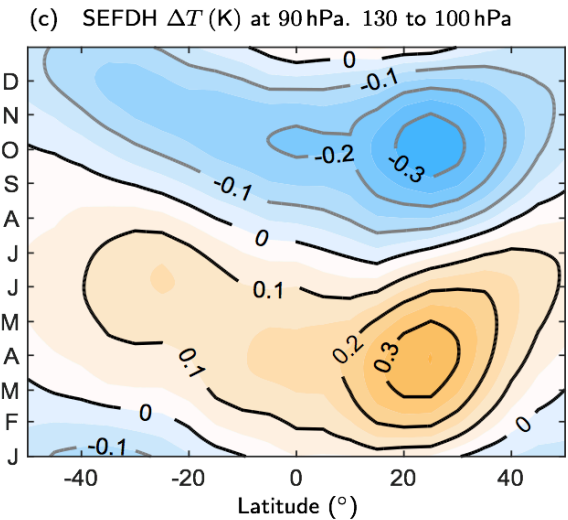
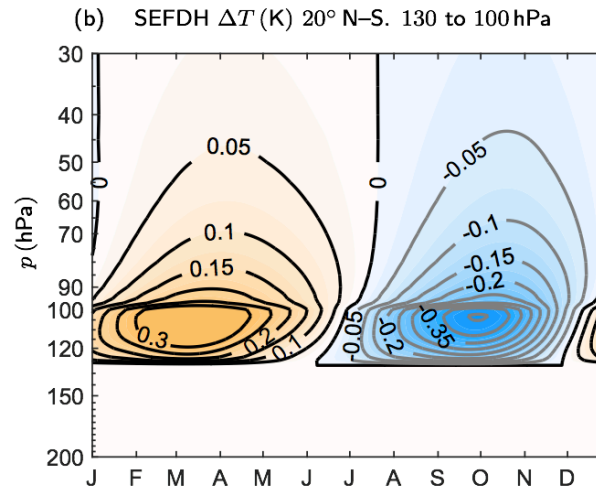
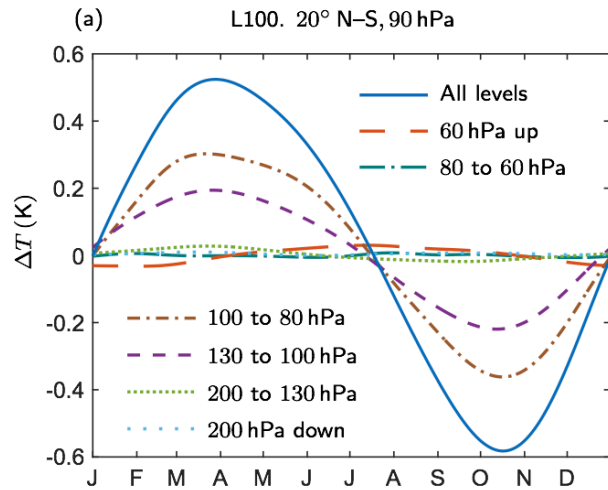
# Effect of imposed $\Delta\bar{\chi}_{\text{H}_2\text{O}}^0$



# Water vapour



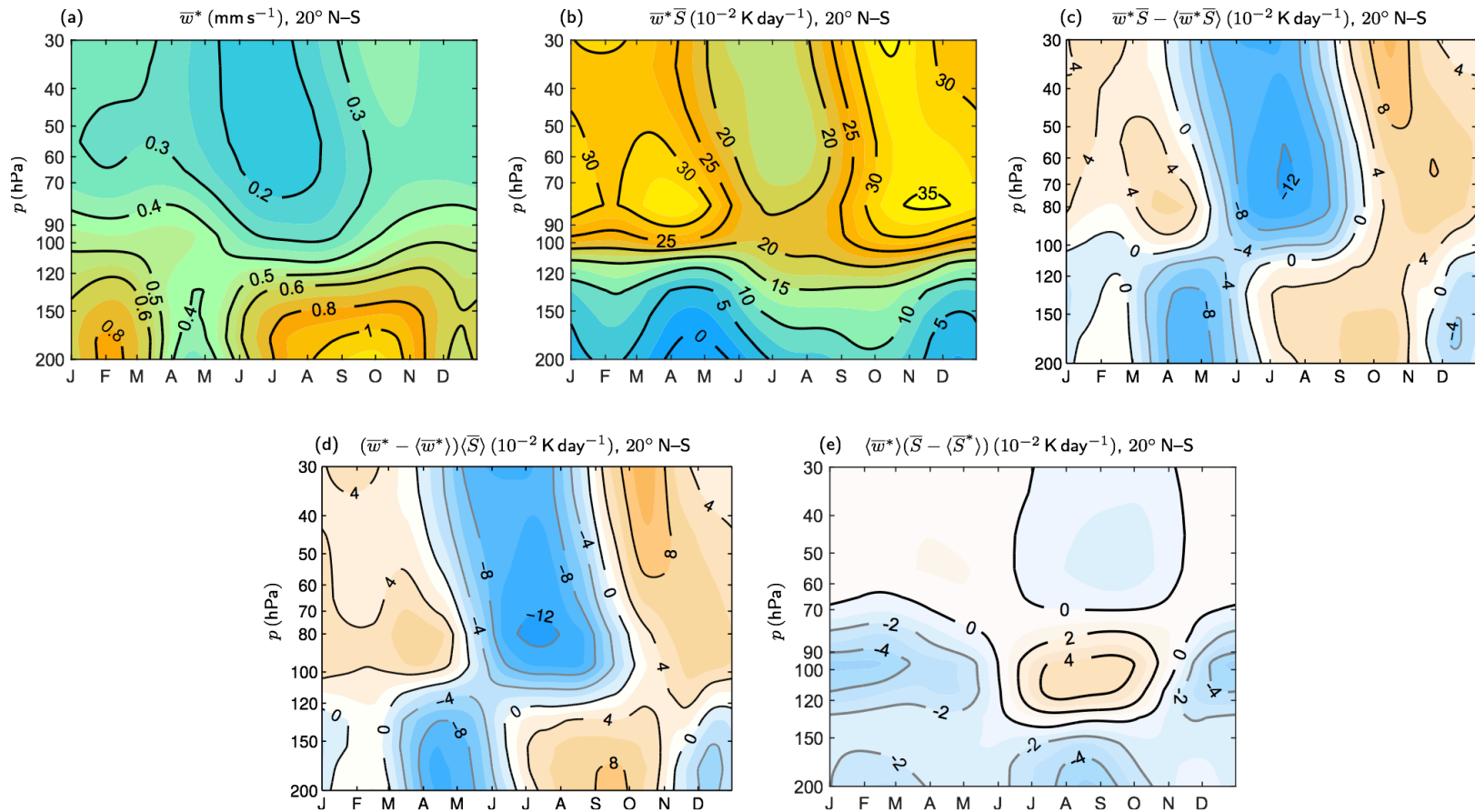
# Water vapour – non-local effects in vertical



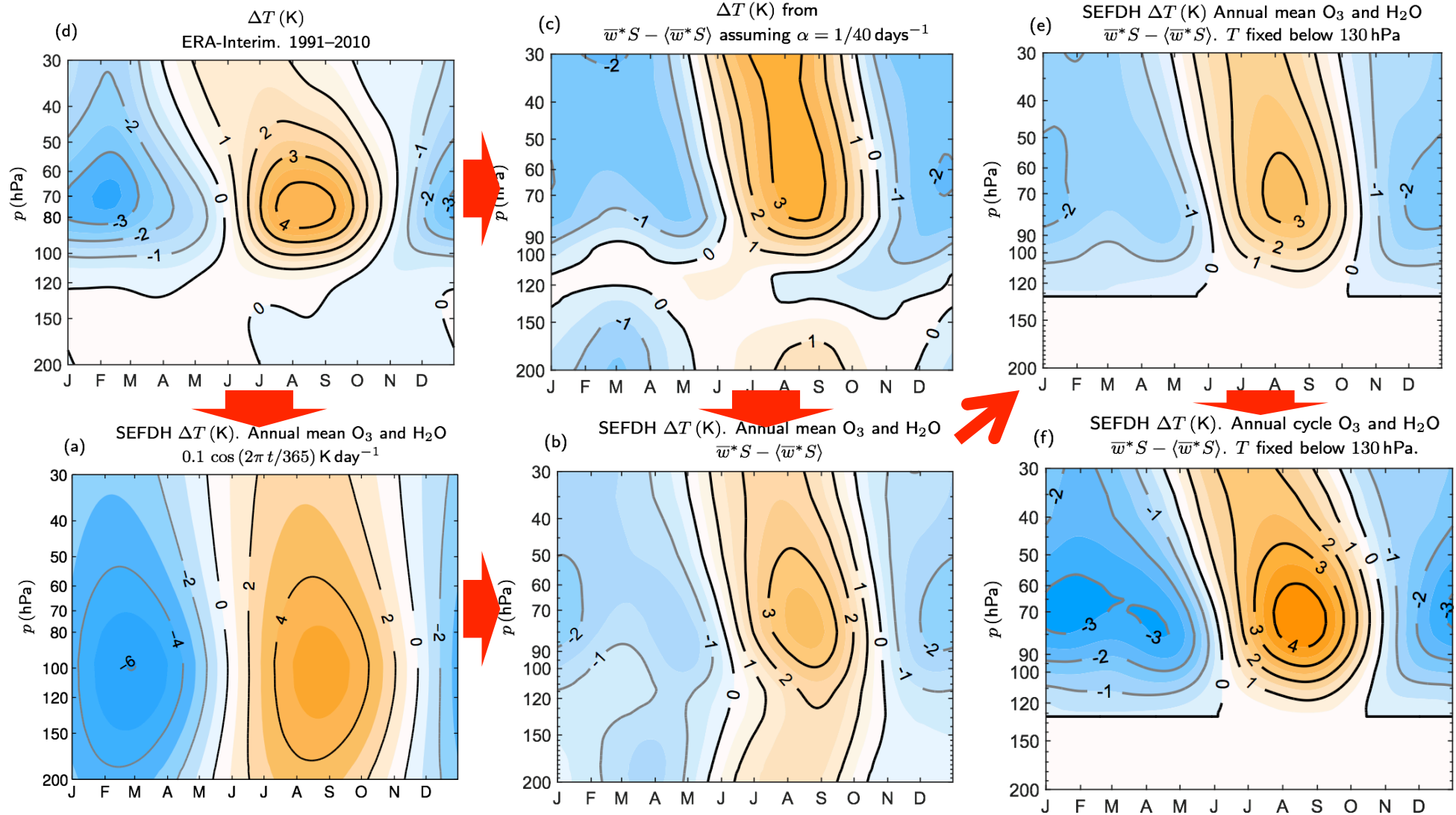
# Temperature cycle driven by dynamical heating

$$\frac{\partial(\bar{T}^0 + \Delta\bar{T})}{\partial t} = Q_{\text{rad}}(\bar{T}^0 + \Delta\bar{T}, \bar{\chi}_{\text{O}_3}^0, \bar{\chi}_{\text{H}_2\text{O}}^0) + \bar{Q}_{\text{dyn}}^0 + \Delta\bar{Q}_{\text{dyn}}$$

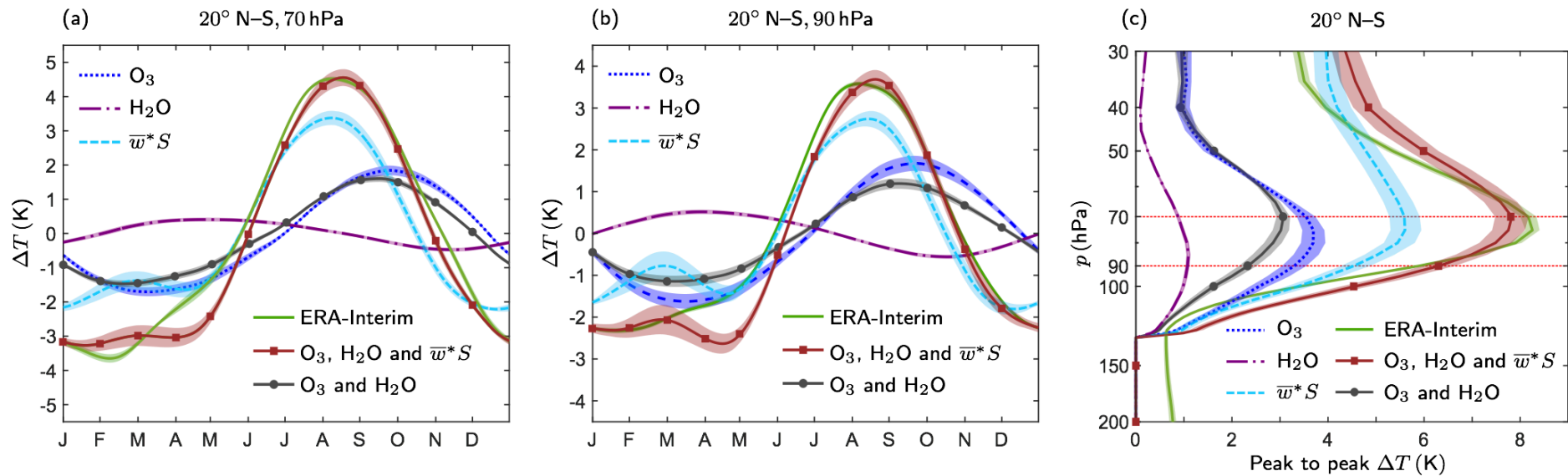
$\Delta\bar{Q}_{\text{dyn}}?$



# $\Delta \bar{T}$ driven by dynamical heating

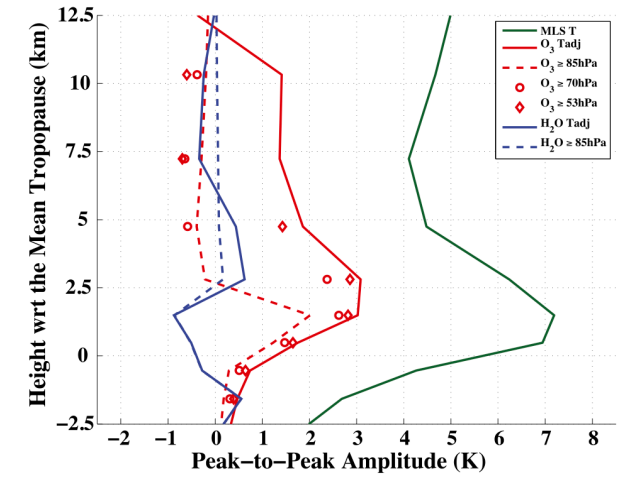
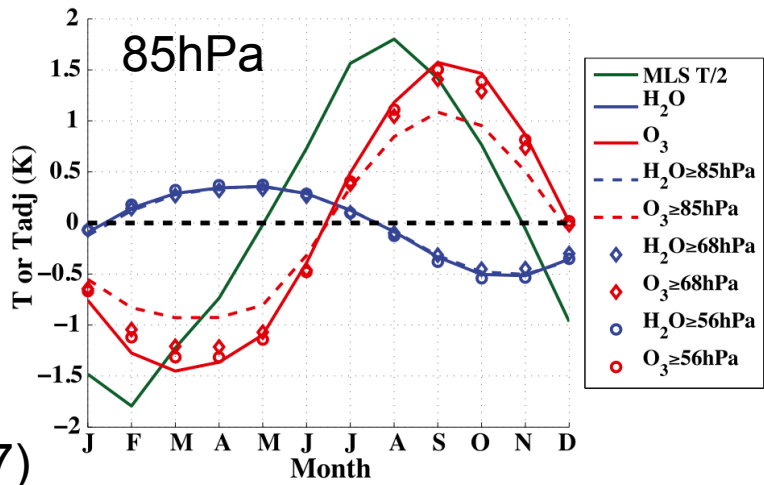
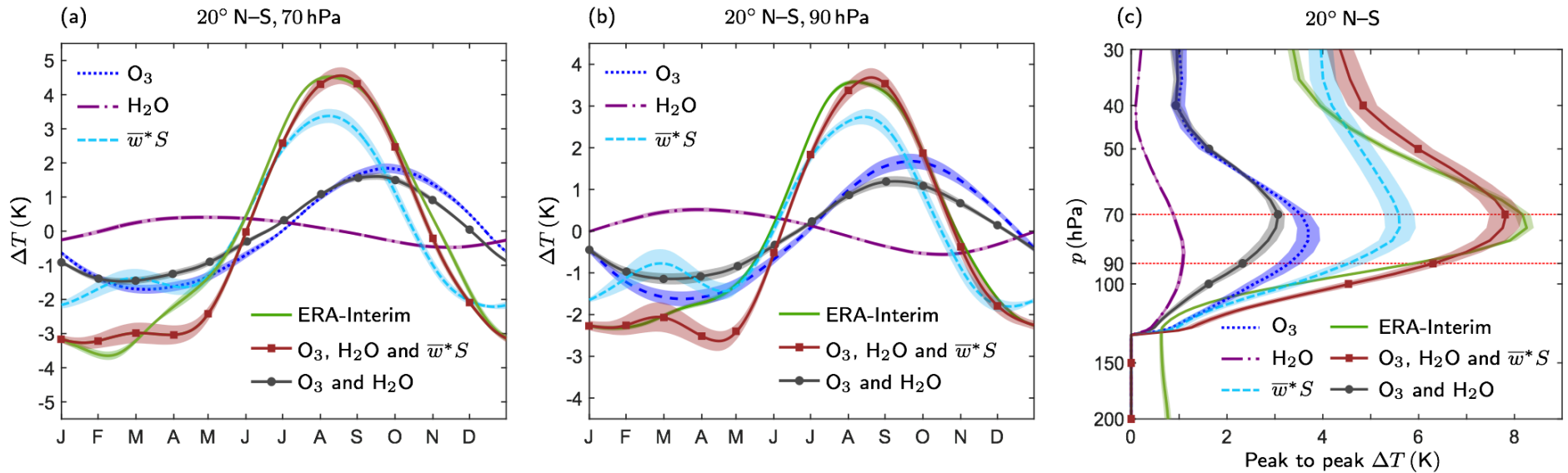


$$\frac{\partial \Delta \bar{T}}{\partial t} = \Delta Q_{\text{rad}}[\Delta T] + \Delta Q_{\text{rad}}[\Delta \bar{\chi}_{\text{O}_3}^0] + \Delta Q_{\text{rad}}[\Delta \bar{\chi}_{\text{H}_2\text{O}}^0] + \Delta \bar{Q}_{\text{dyn}}$$





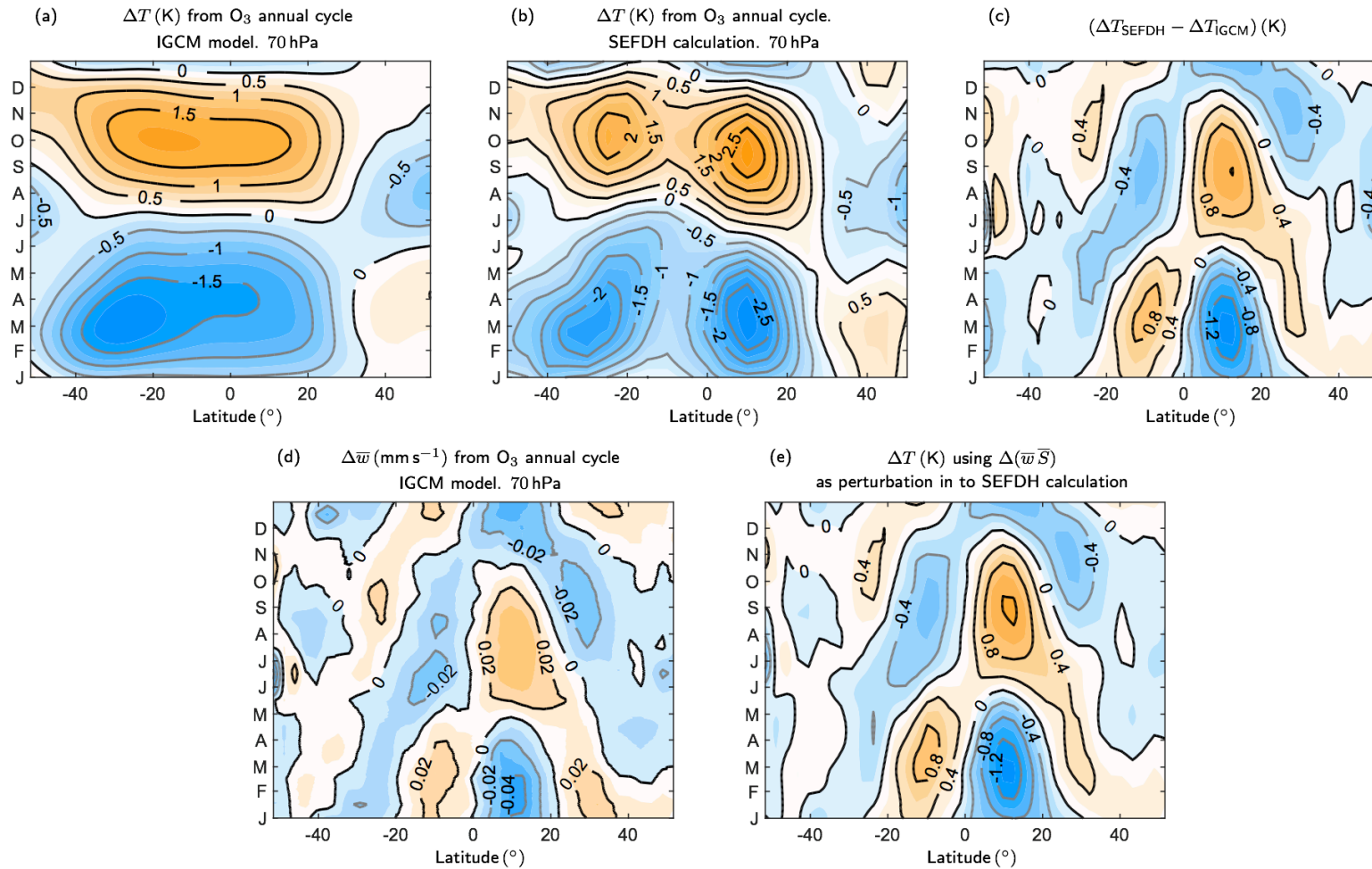
# Ming et al (2017)



Gilford and Solomon (2017)

# Relax SEFDH assumption

## Radiation code + zonally symmetric dynamics



$$\partial_t \bar{u} + \frac{1}{a \cos \phi} \bar{v}^* \partial_\phi (\bar{u} \cos \phi) + \bar{w}^* \partial_z \bar{u} - f \bar{v}^* = \frac{1}{\rho_0 a \cos \phi} \nabla \cdot \vec{F}, \quad (1)$$

$$f \partial_z \bar{u} + \frac{R}{H} \frac{1}{a} \partial_\phi \bar{T} = 0, \quad (2)$$

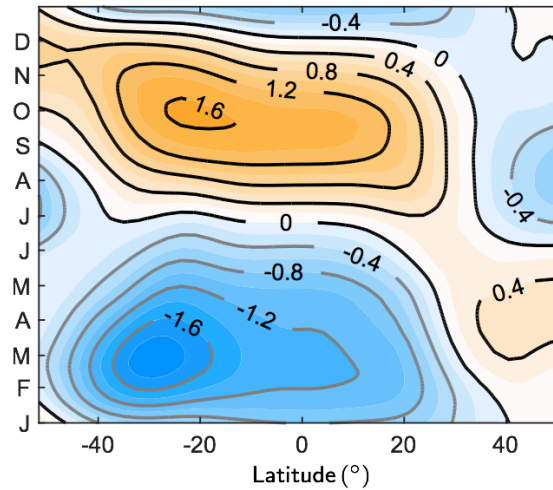
$$\frac{1}{a \cos \phi} \partial_\phi (\bar{v}^* \cos \phi) + \frac{1}{\rho_0} \partial_z (\rho_0 \bar{w}^*) = 0, \quad (3)$$

$$\partial_t \bar{T} + \frac{1}{a} \bar{v}^* \partial_\phi \bar{T} + \bar{S} \bar{w}^* = \Delta \bar{Q}_{\text{rad}}[\bar{T}, \Delta \bar{\chi}_{\text{O}_3}, \Delta \bar{\chi}_{\text{H}_2\text{O}}]. \quad (4)$$

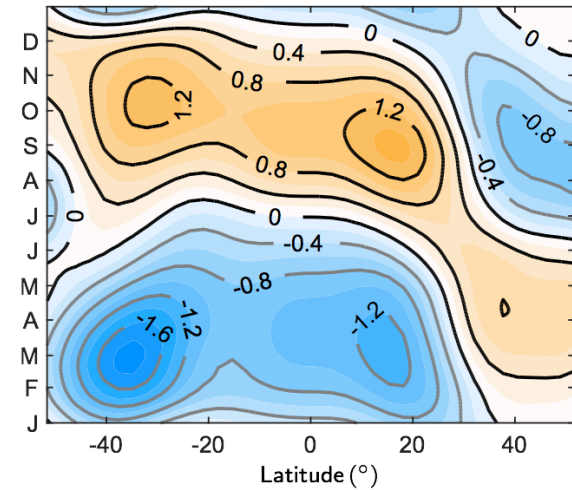


# Effect on tropical average 20N-20S

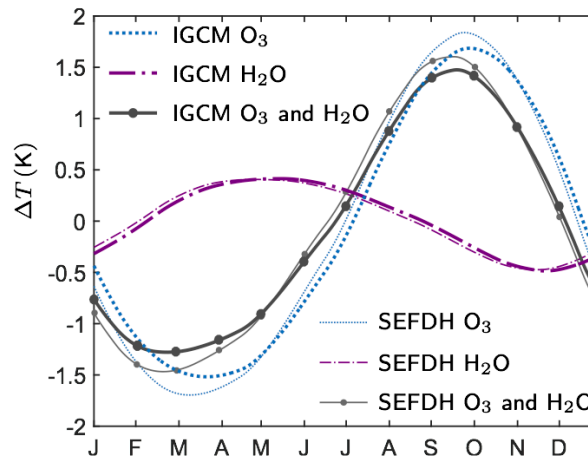
(a)  $\Delta T$  (K) from O<sub>3</sub> and H<sub>2</sub>O annual cycle  
IGCM model. 70 hPa



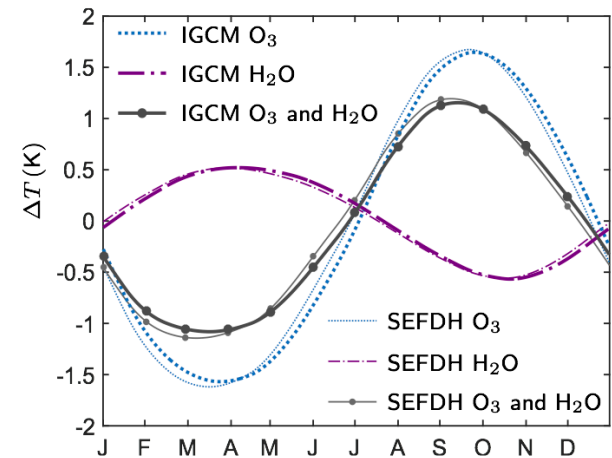
(b)  $\Delta T$  (K) from O<sub>3</sub> and H<sub>2</sub>O annual cycle  
IGCM model. 90 hPa



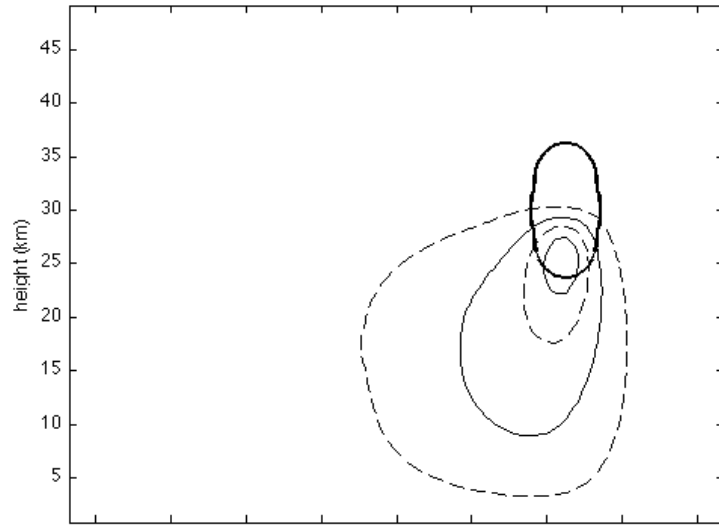
(c)  $\Delta T$  (K). IGCM model and SEFDH  
20° N-S, 70 hPa



(d)  $\Delta T$  (K). IGCM model and SEFDH  
20° N-S, 90 hPa

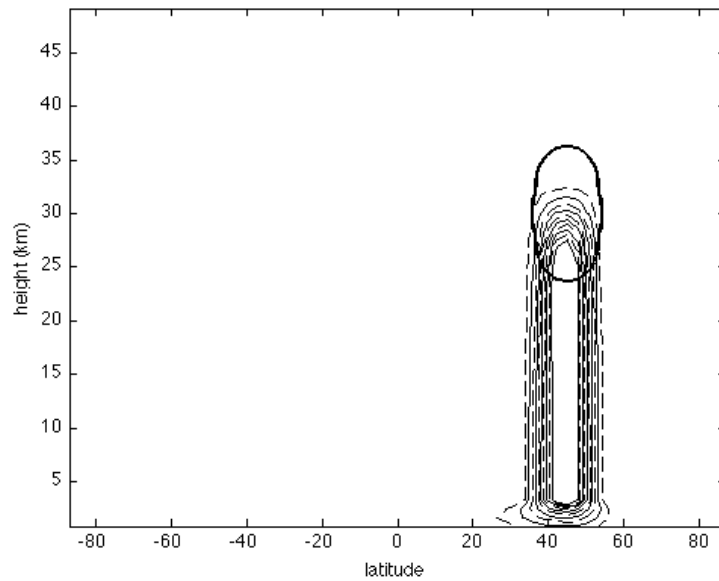
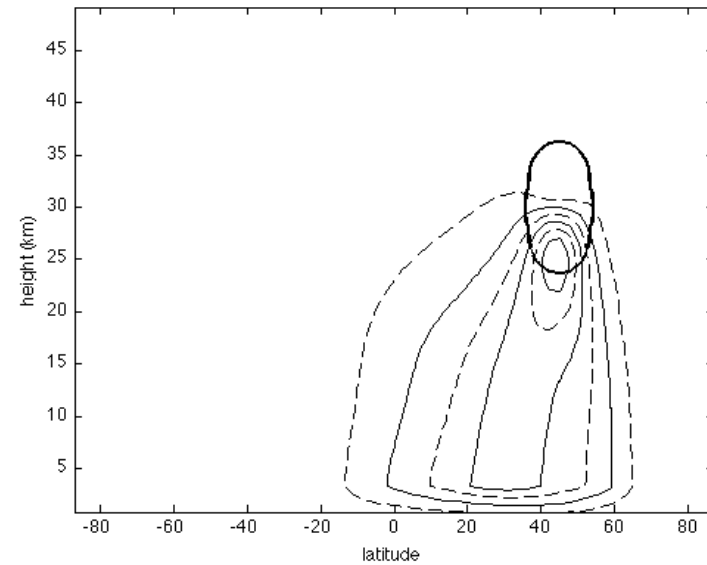


# Meridional circulation response on different timescales



rapid variation

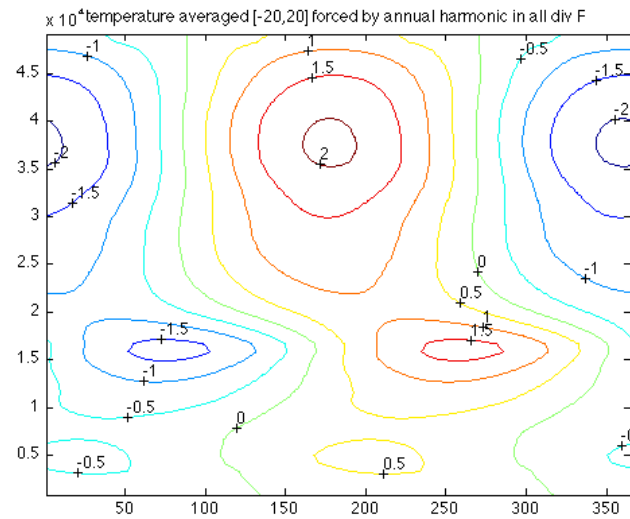
annual variation



slow variation

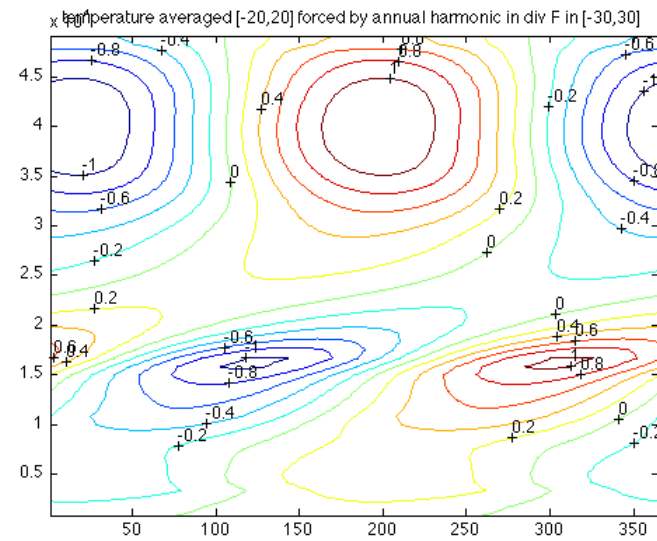
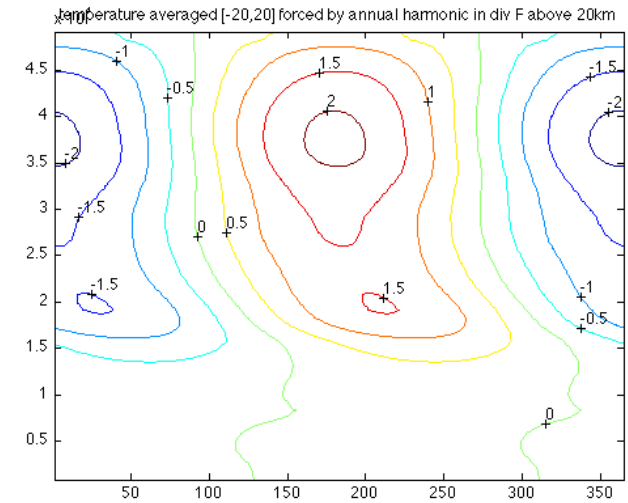


# $\Delta \bar{T}$ driven by $\nabla \cdot \mathbf{F}$ from different regions



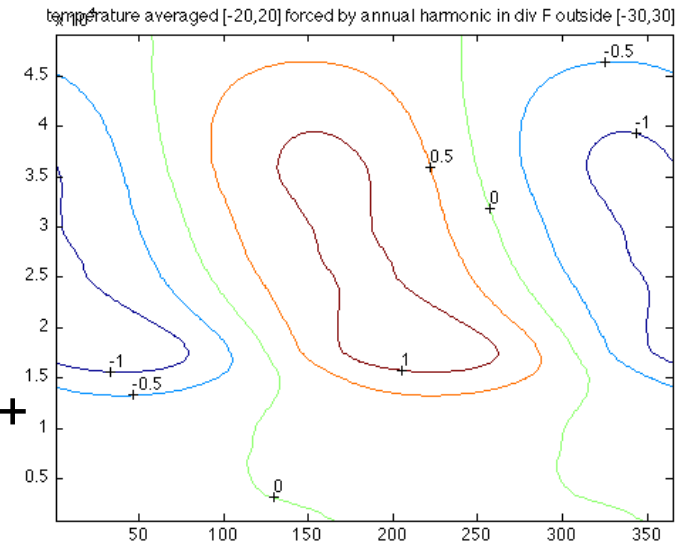
all

$z > 20\text{km}$



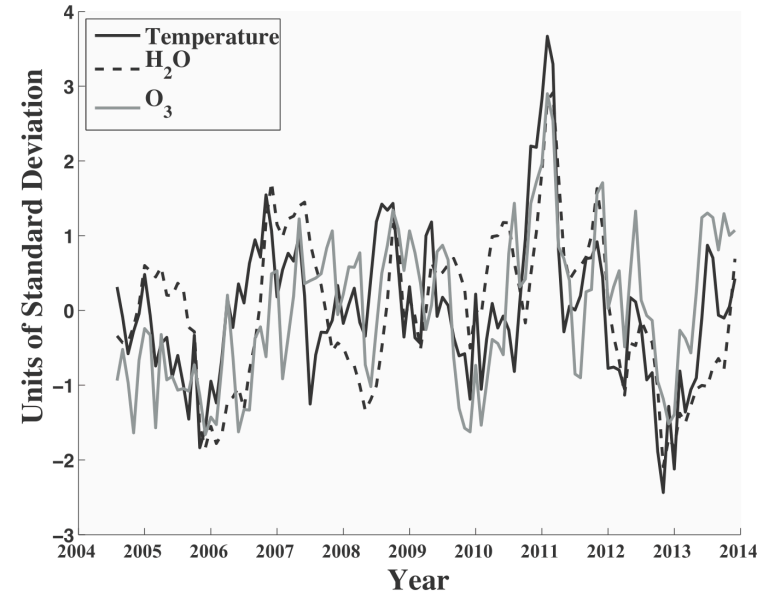
[-30 ,30]

[-90 , -30] +  
[30,90]

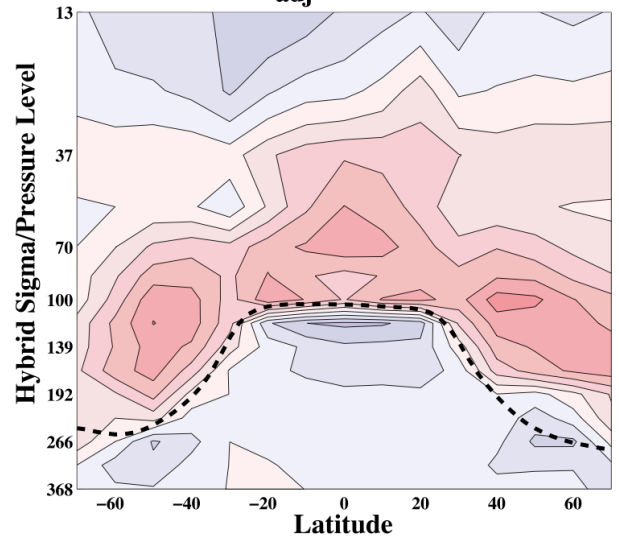


# 2011 'water vapour drop'

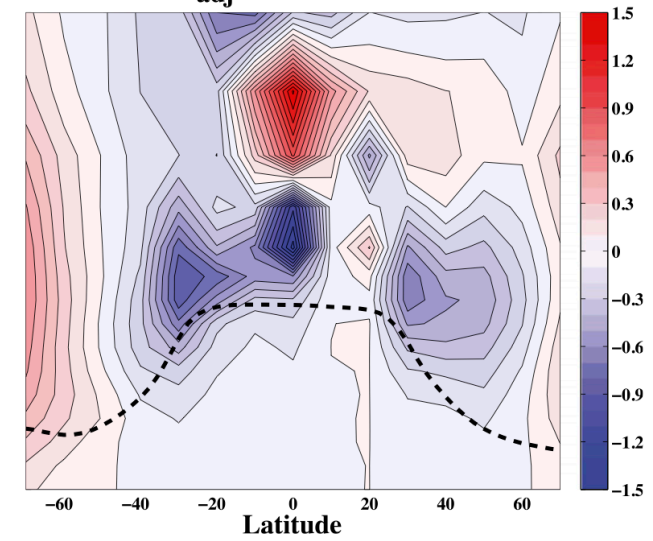
Gilford et al (2016)



a) Water Vapor  $T_{adj}$  (K), during 2011 Abrupt Drop



b) Ozone  $T_{adj}$  (K), during 2011 Abrupt Drop



# Summary

1. Annual cycle in tropical tropopause temperatures is driven both by dynamical effects (response to wave force) and by radiative effects (response to variations in ozone and water vapour)
2. Dynamical effects – non-locality of response to wave force needs to be taken into account. Wave forces due to waves propagating from extratropics and from tropics are likely to be important.
3. Radiative effects – SEFDH calculations give useful insight into tropical average, but not into latitudinal structure.
4. Similar considerations apply to longer term variations, e.g. interannual variations.
5. Calculations presented are for *given* change in wave force or *given* change in chemical species – but in reality these are determined by feedbacks and must be understood as part of the response as a whole.

