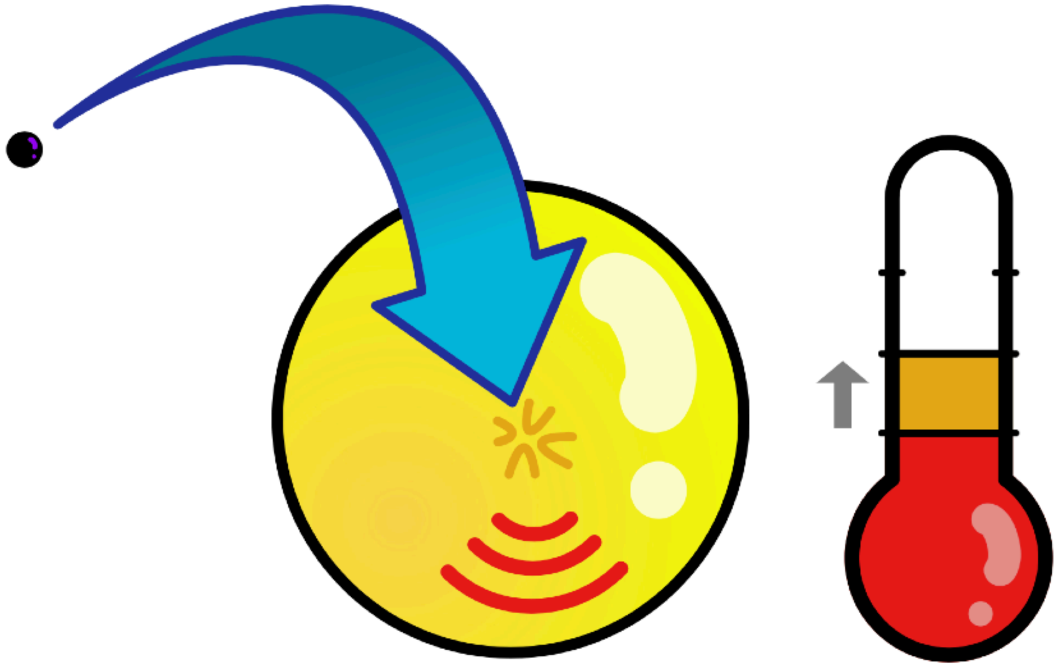
Nucleons, Electrons, and Pasta:

Discovering Dark Matter by Reheating the Neutron Star Soup

Nirmal Raj

TRIUMF

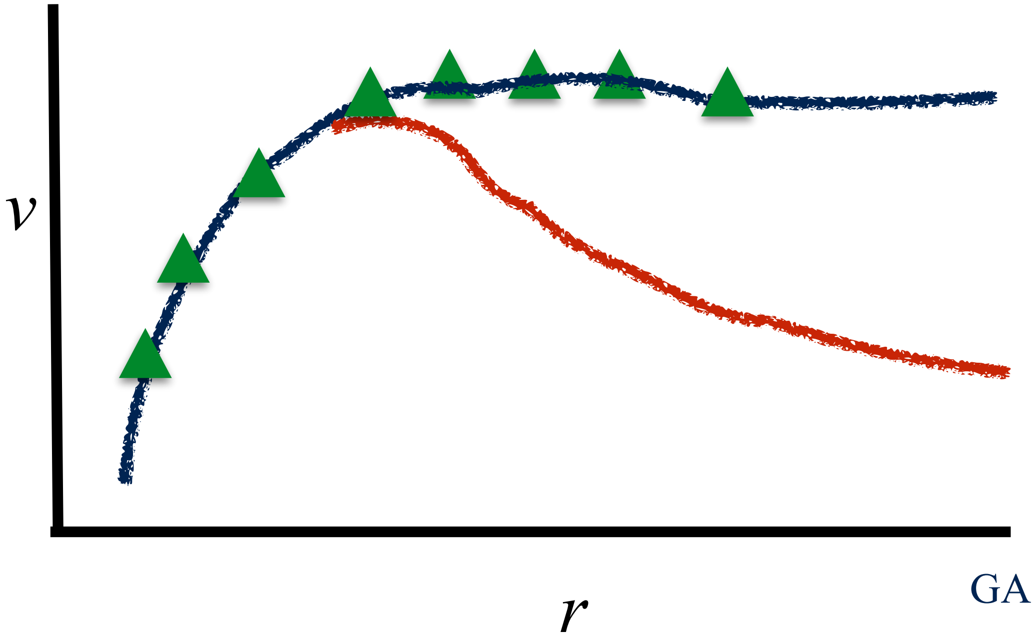


with

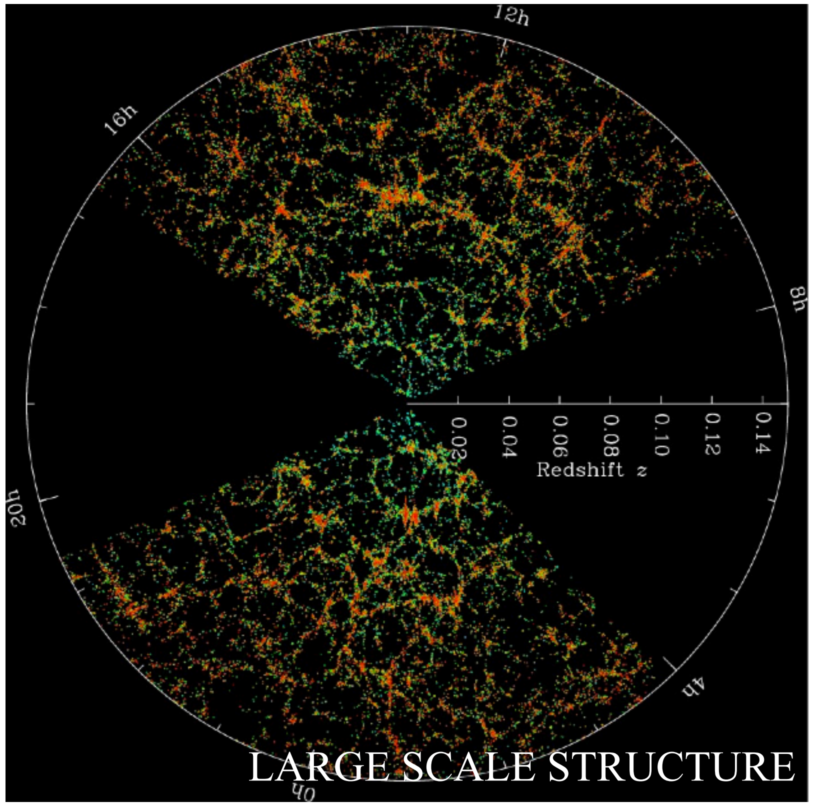
Javier Acevedo | Masha Baryakhtar | Joe Bramante | Aniket Joglekar
Rebecca Leane | Shirley Li | Tim Linden | Flip Tanedo | Hai-Bo Yu

28 Jan 2020, U Toronto

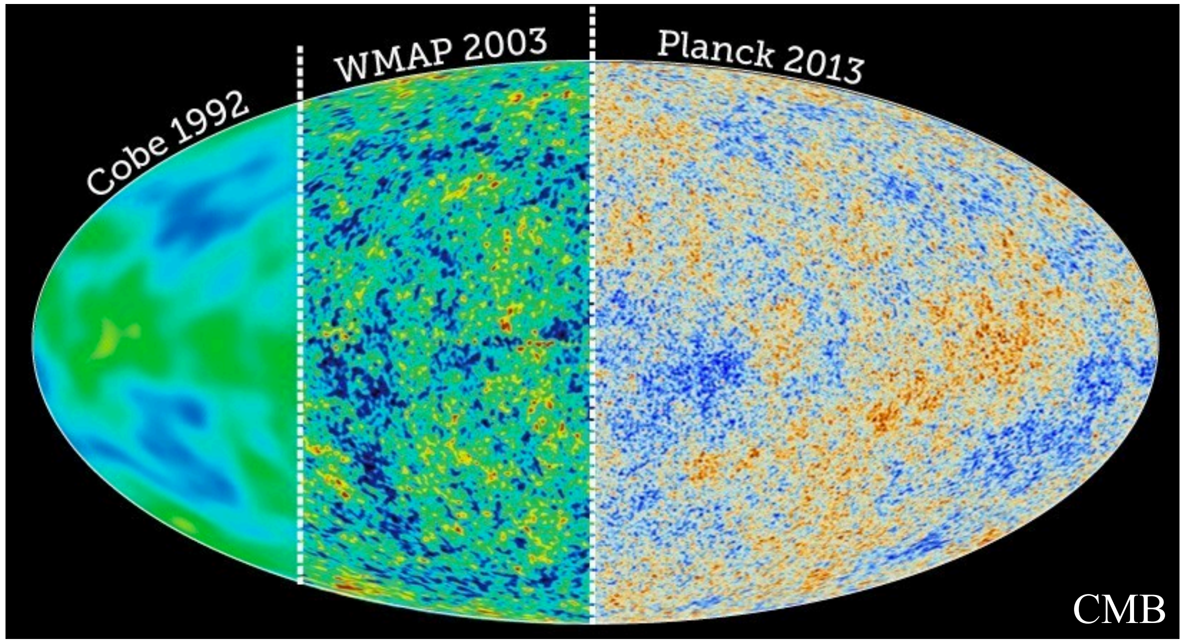
Dark reality



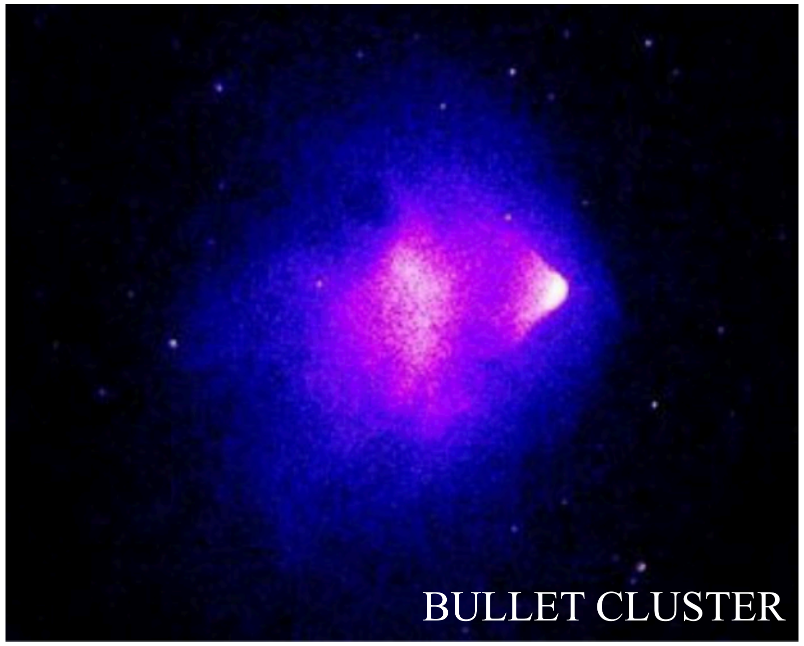
GALACTIC ROTATION



LARGE SCALE STRUCTURE



CMB

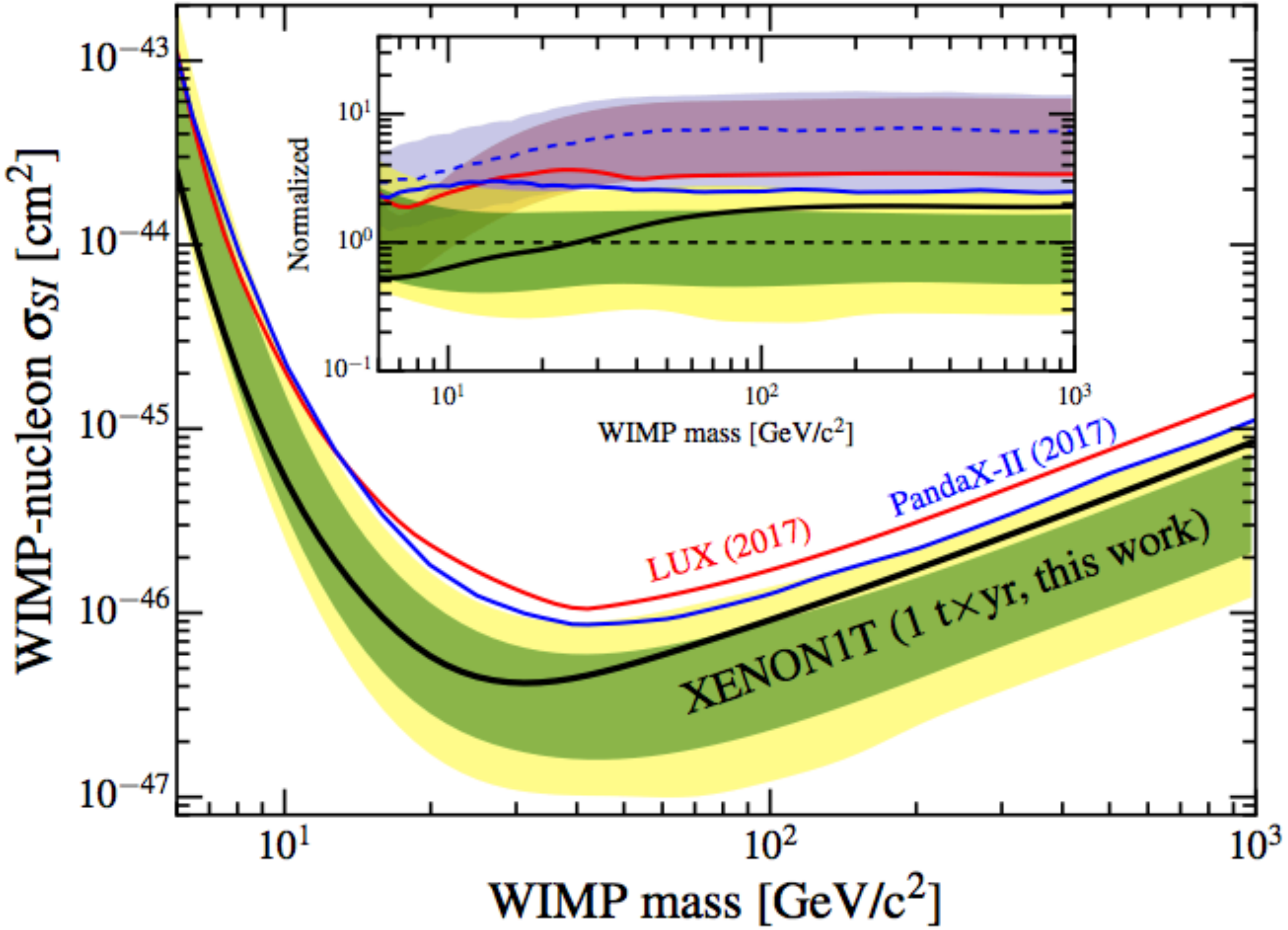


BULLET CLUSTER

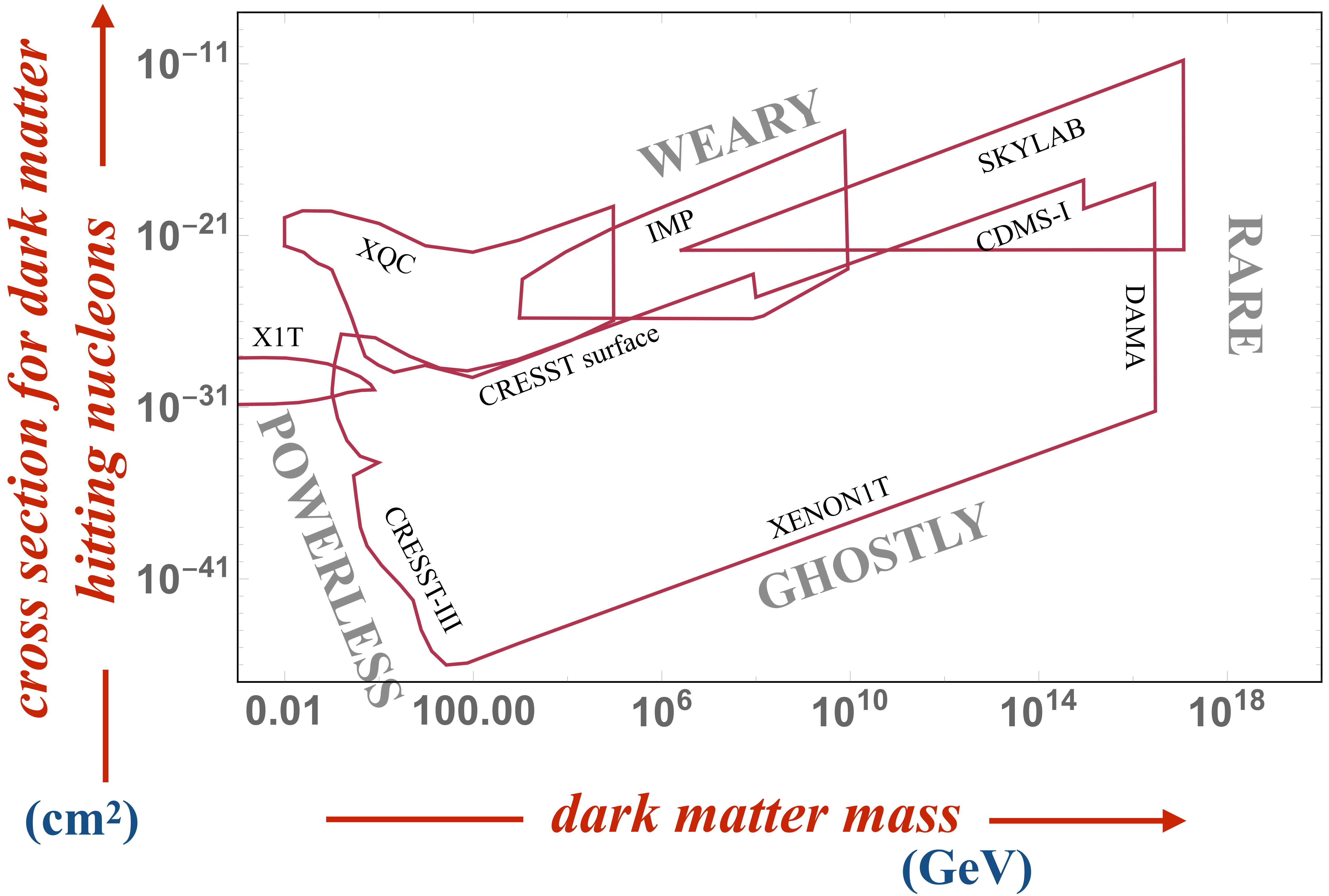
Direct searches: status

After 1000 kg-year exposure:

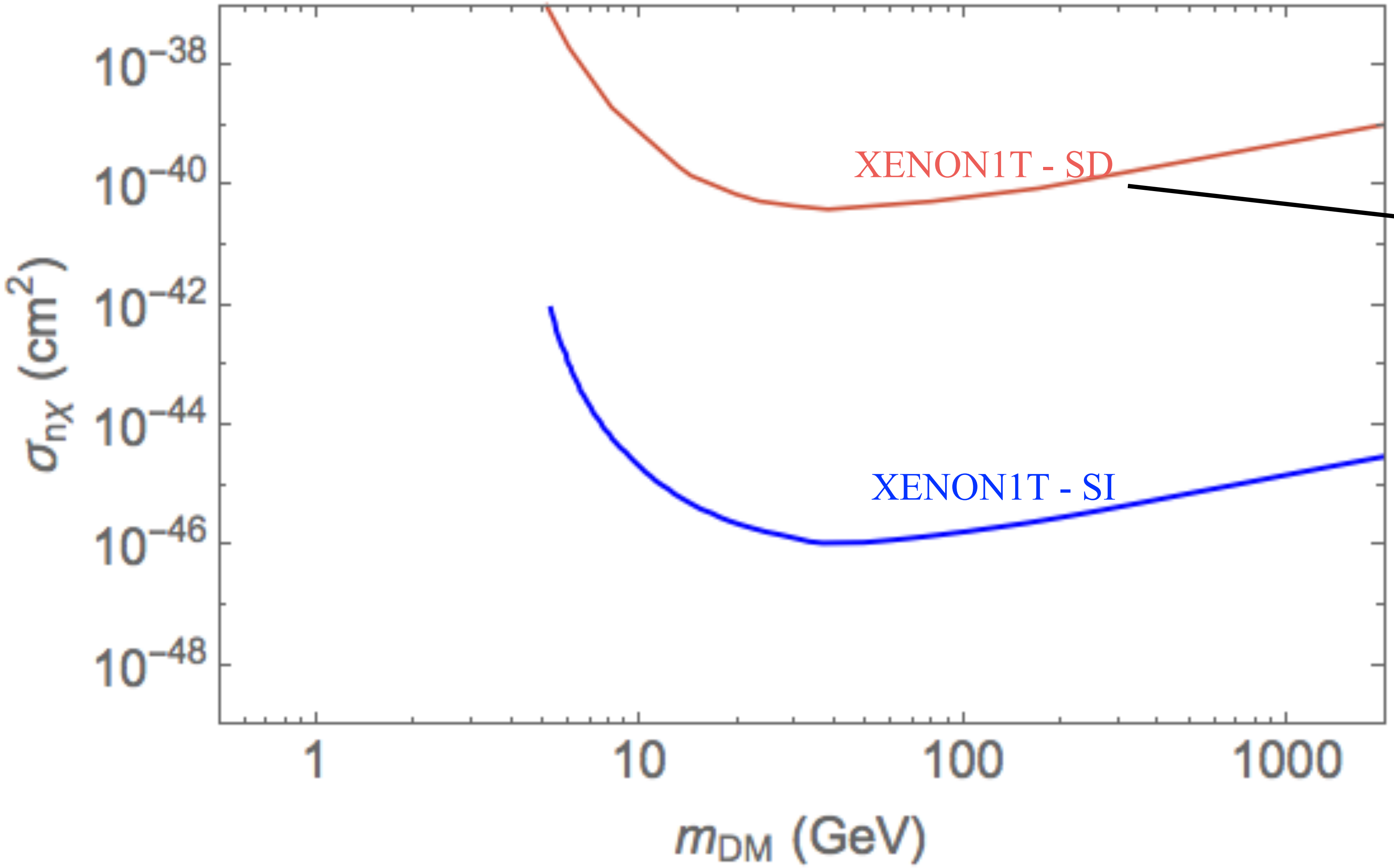
1805.12562



Challenges of direct searches



More challenges: spin dependence

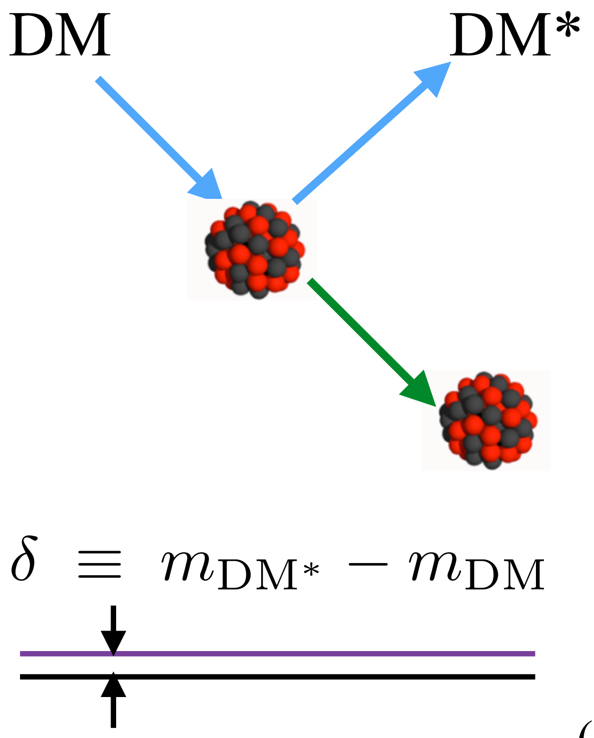


No nuclear coherence
if scattering spin-dependent:

$$\sigma_{\text{nucleon}} = \cancel{\left(\frac{1}{A}\right)^2} \left(\frac{\mu_{n\chi}}{\mu_{N\chi}}\right)^2 \sigma_{\text{Nucleus-DM}}$$

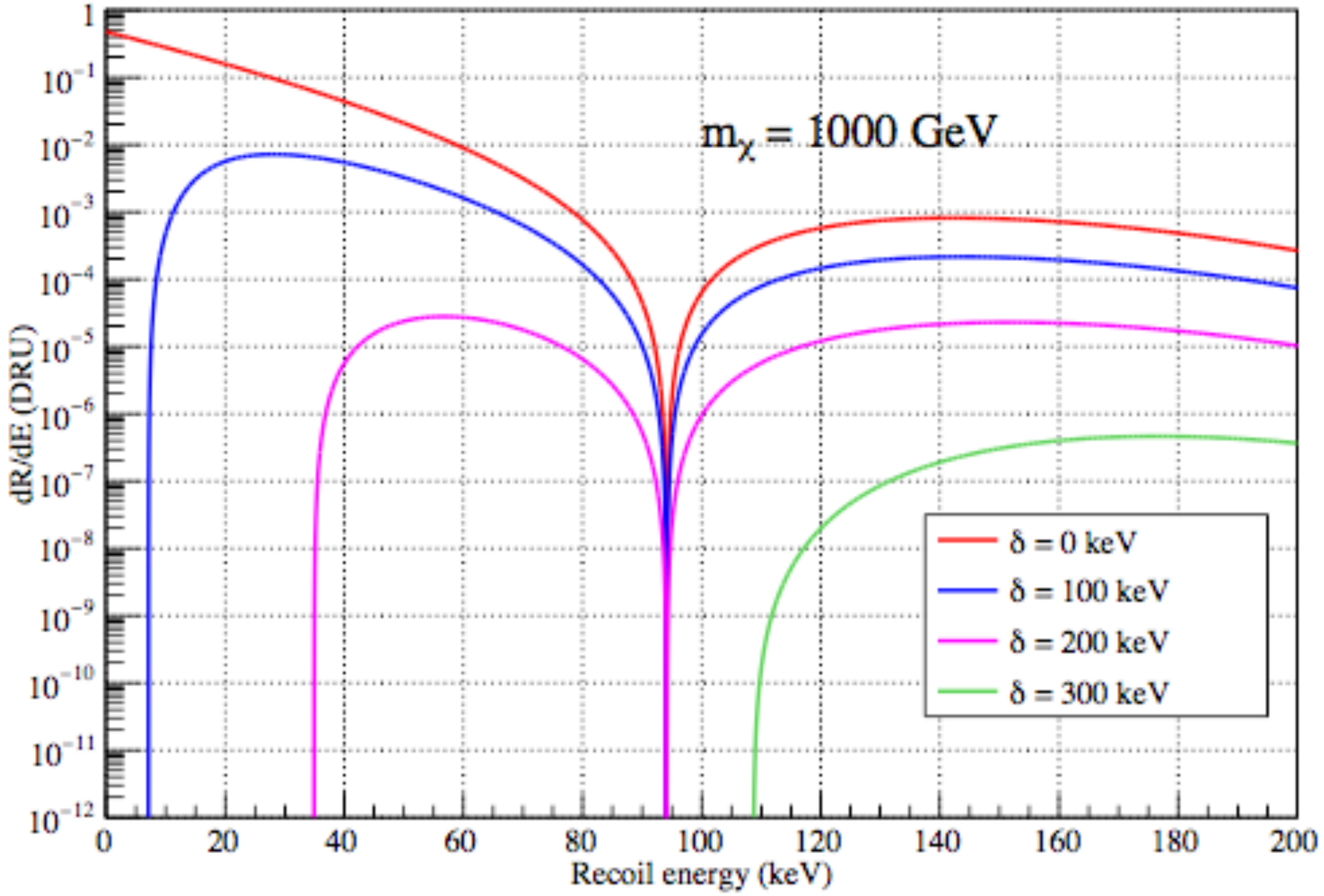
E.g. (Majorana) DM coupling to axial quark current $\bar{q}\gamma_\mu\gamma_5 q$

More challenges: inelasticity



$$\delta \equiv m_{\text{DM}^*} - m_{\text{DM}}$$

(Image: G. Kribs)



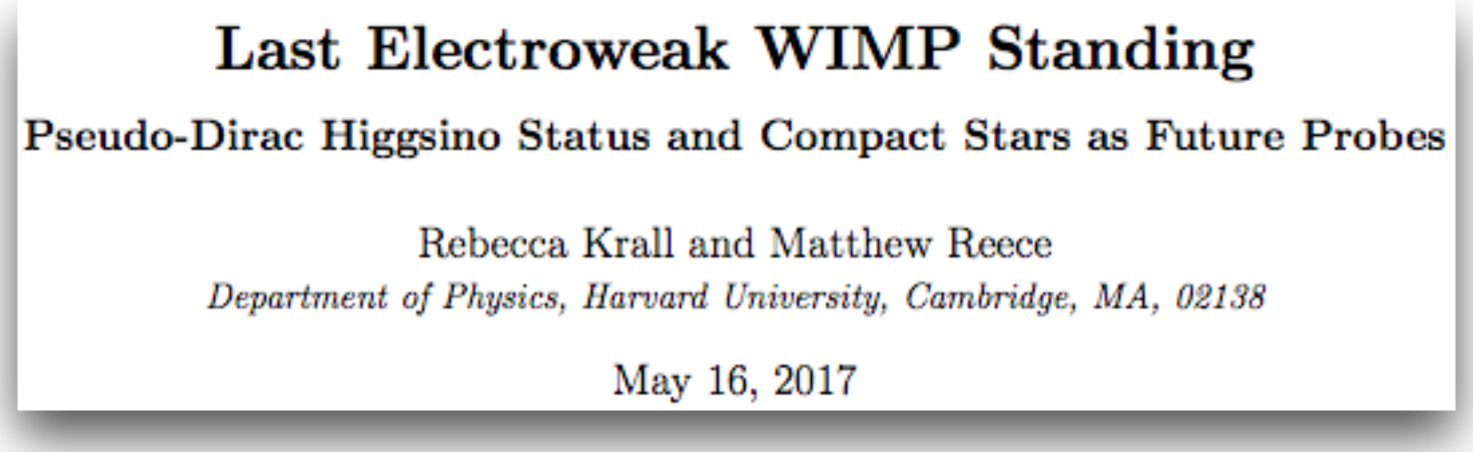
PandaX-II, 1708.05825

If scattering <u>inelastic</u>,
no recoil when

$$\delta > 2\mu_{N\chi} v_{\text{DM}}^2 = \mathcal{O}(100\text{keV})$$

Tucker-Smith, Weiner 0101138, 0402065,
Barello, Chang, Newby 1409.0536

More challenges: inelasticity



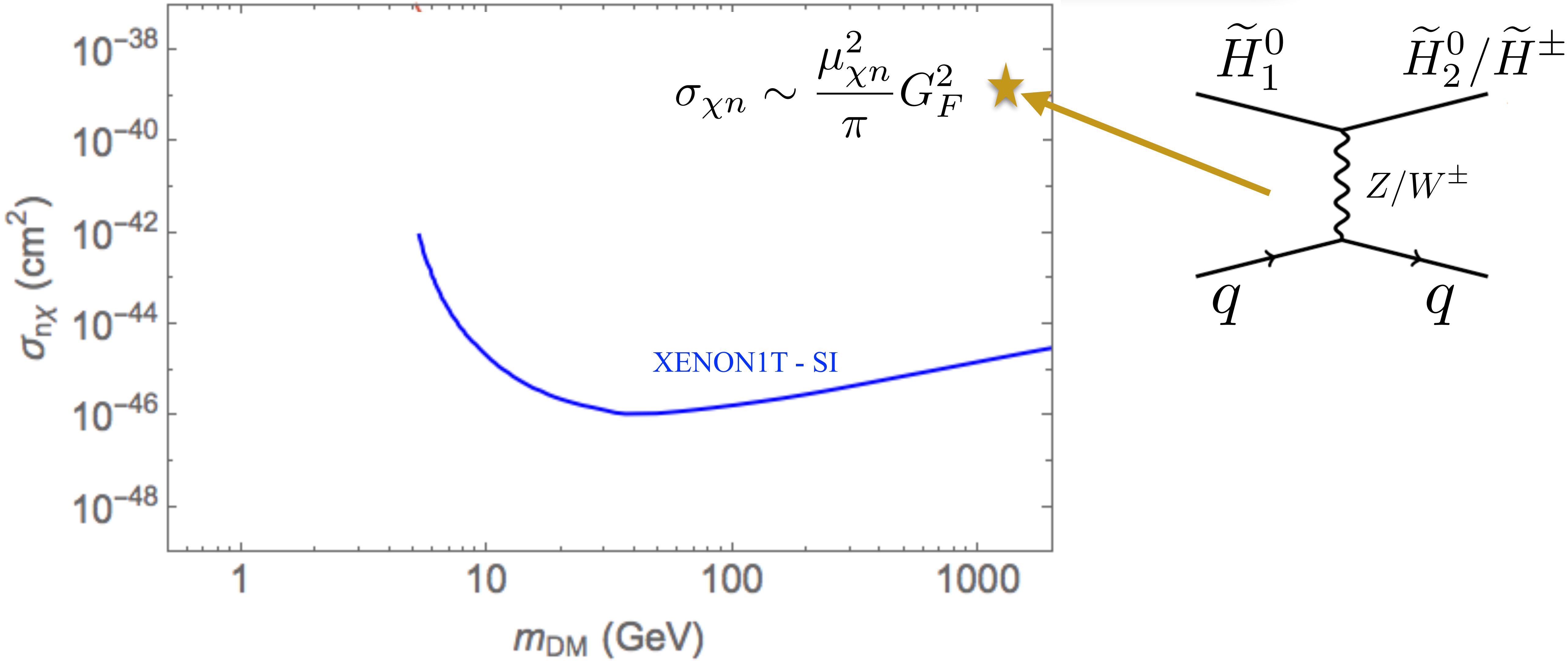
Last Electroweak WIMP Standing

Pseudo-Dirac Higgsino Status and Compact Stars as Future Probes

Rebecca Krall and Matthew Reece

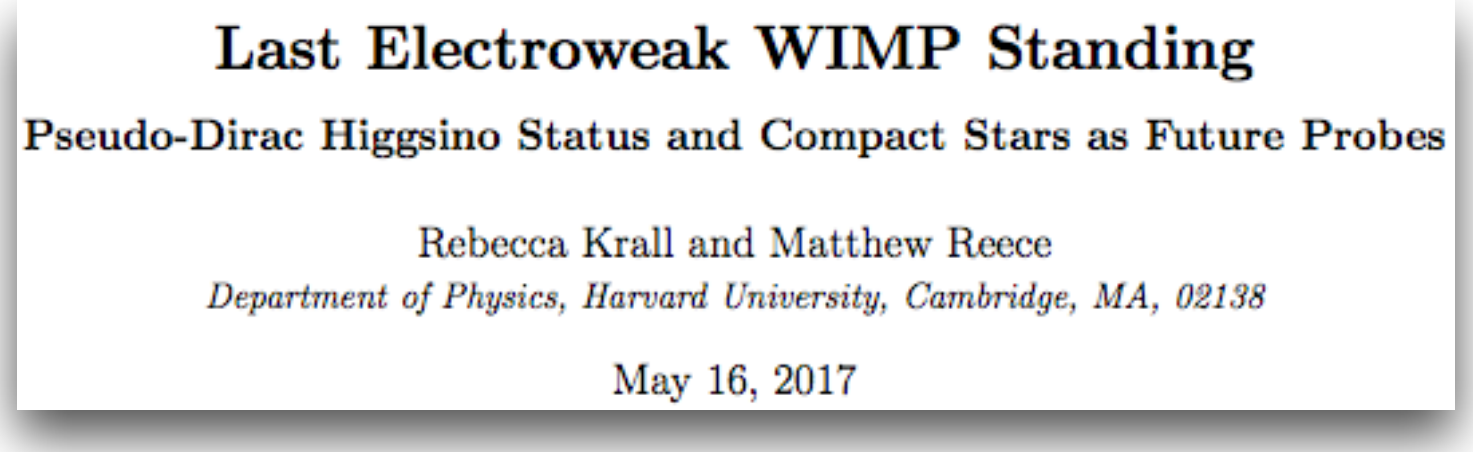
Department of Physics, Harvard University, Cambridge, MA, 02138

May 16, 2017



$$\sigma_{\chi n} \sim \frac{\mu_{\chi n}^2}{\pi} G_F^2$$

More challenges: inelasticity



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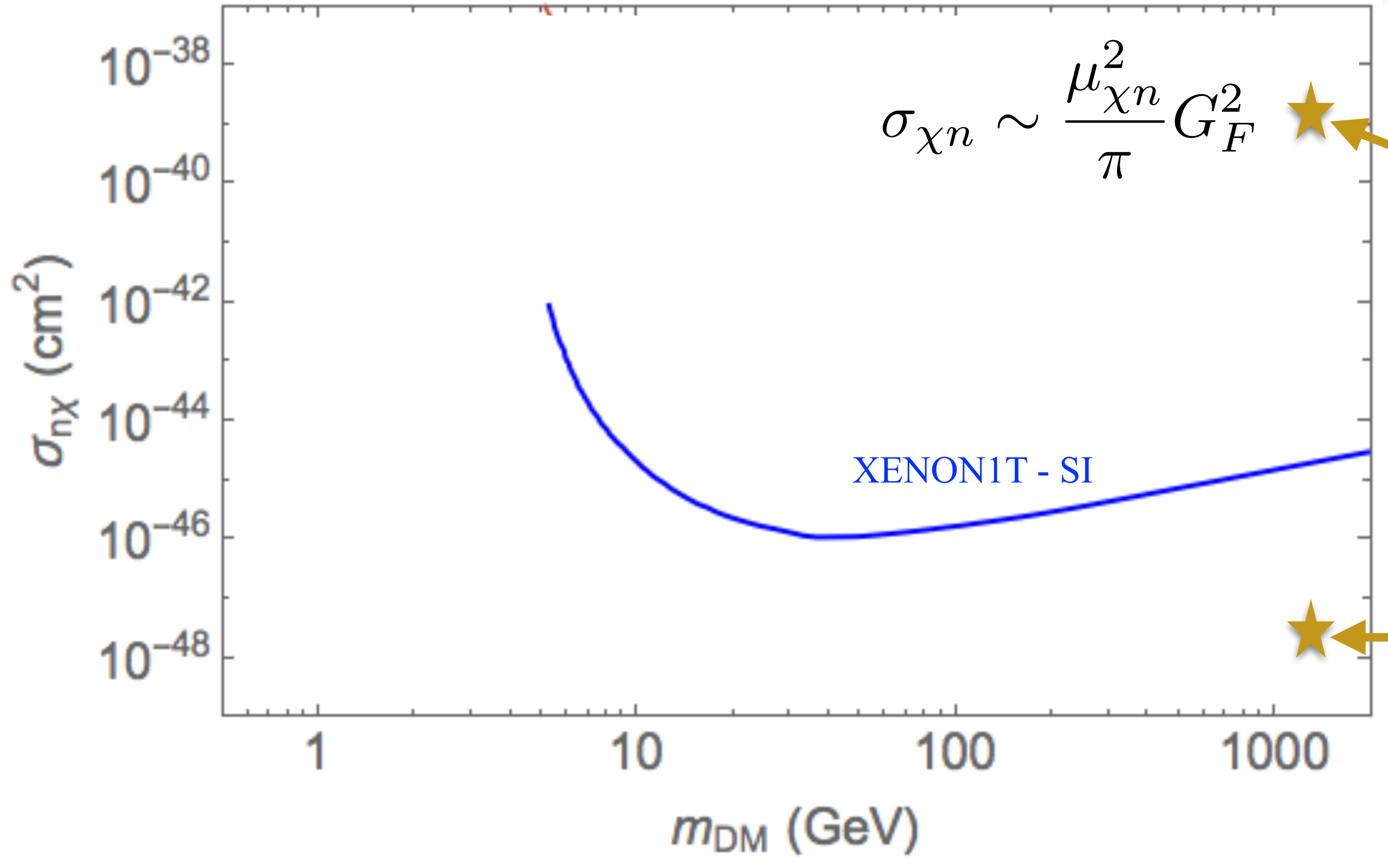
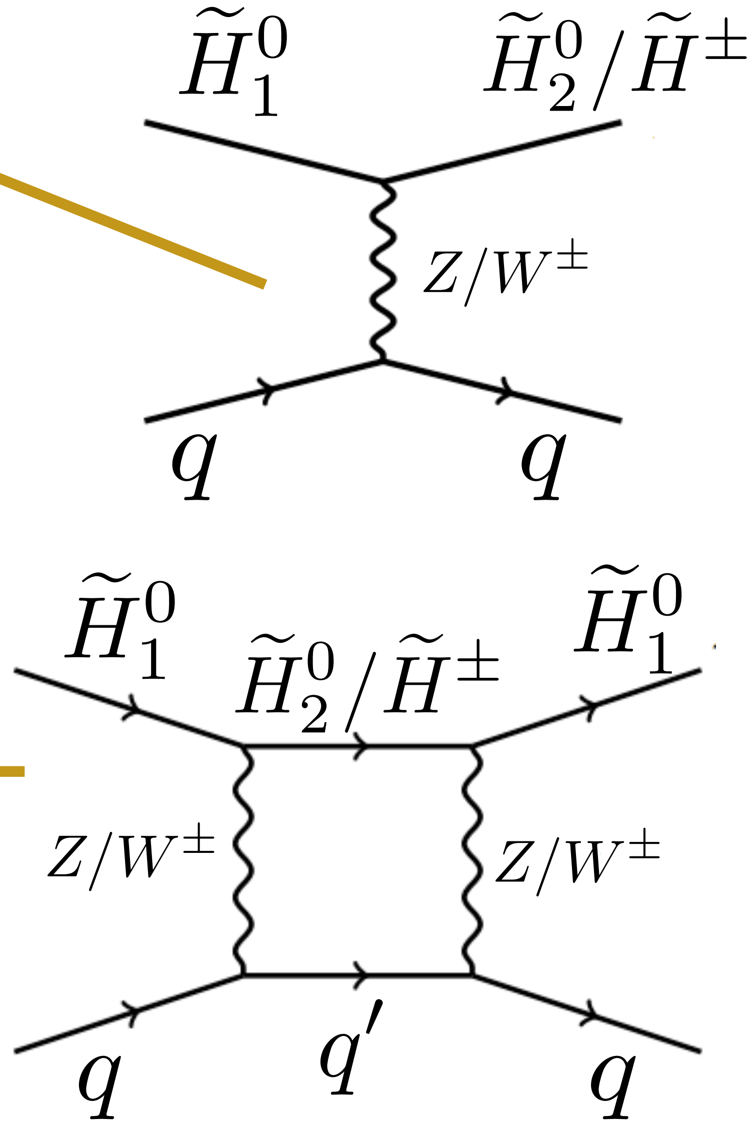


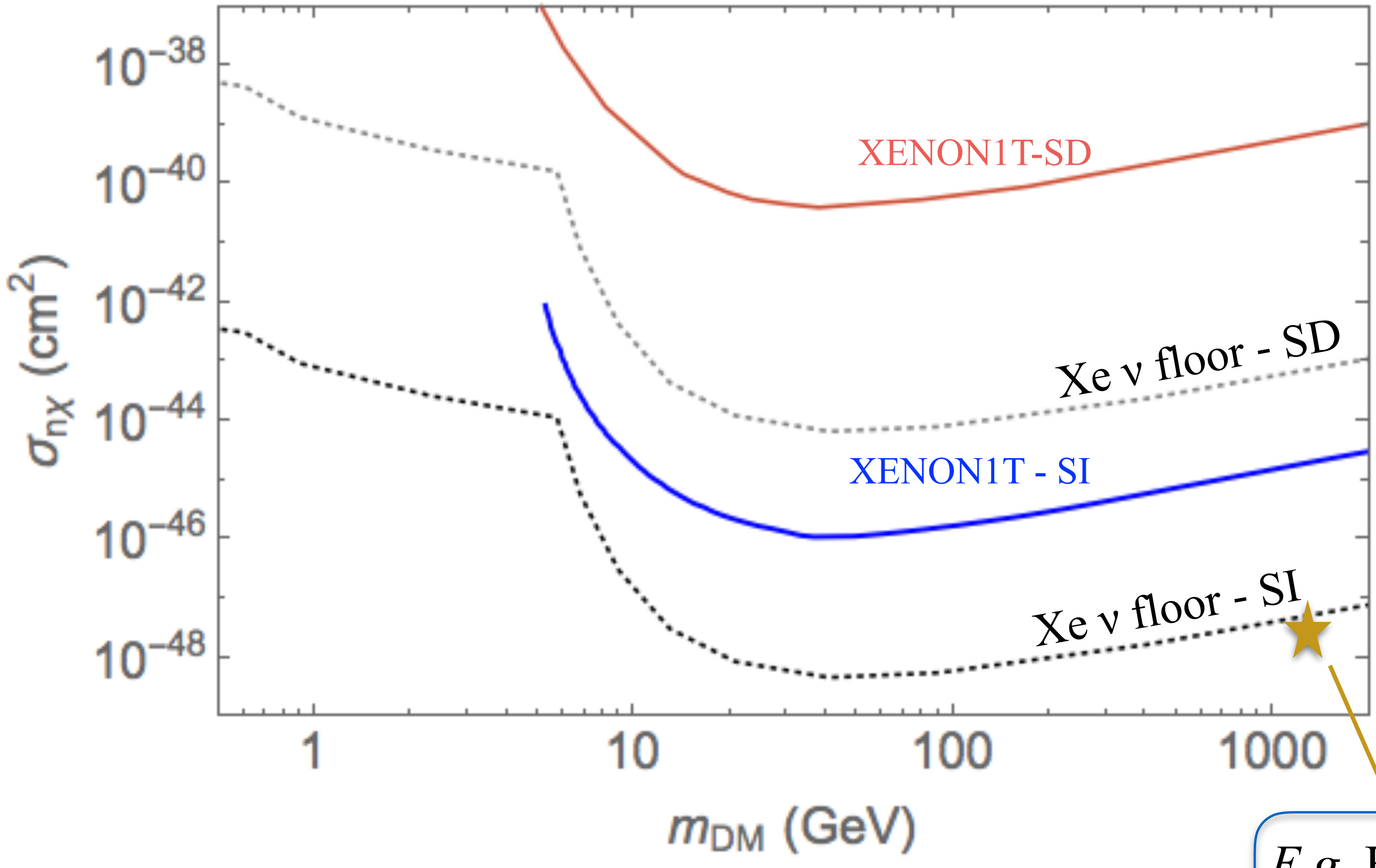
$$\sigma_{\chi n} \sim \frac{\mu_{\chi n}^2}{\pi} G_F^2$$



Thermal Higgsino — elusive!

Hill & Solon 1408.8290

More challenges: irreducible backgrounds



Atmospheric +
diffuse supernovae +
solar
neutrino background

E.g. Higgsino elastic scattering (loops)

Hill & Solon 1408.8290

Challenges: summary

(1) Low mass ☒
(2) High mass ☒
(3) Strongly interacting ☒
(4) Spin-dependent ☒
(5) Inelastic ☒
(6) Neutrino floors ☒

Crucial frontiers — beyond which dark matter could be.

(Dark) Kinetic Heating

heating rate = cooling rate

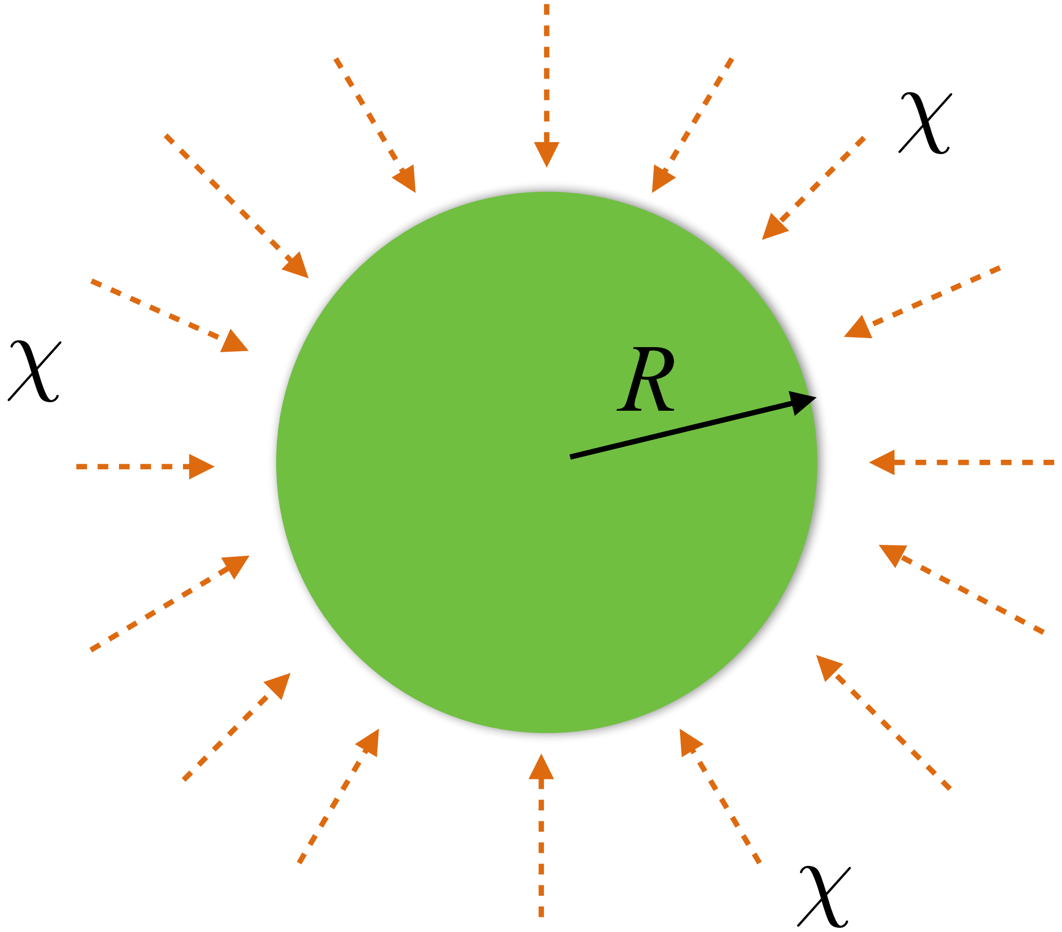
heating rate = cooling rate

$\propto \mathrm{KE}$

$\propto \dfrac{dN}{dt}$

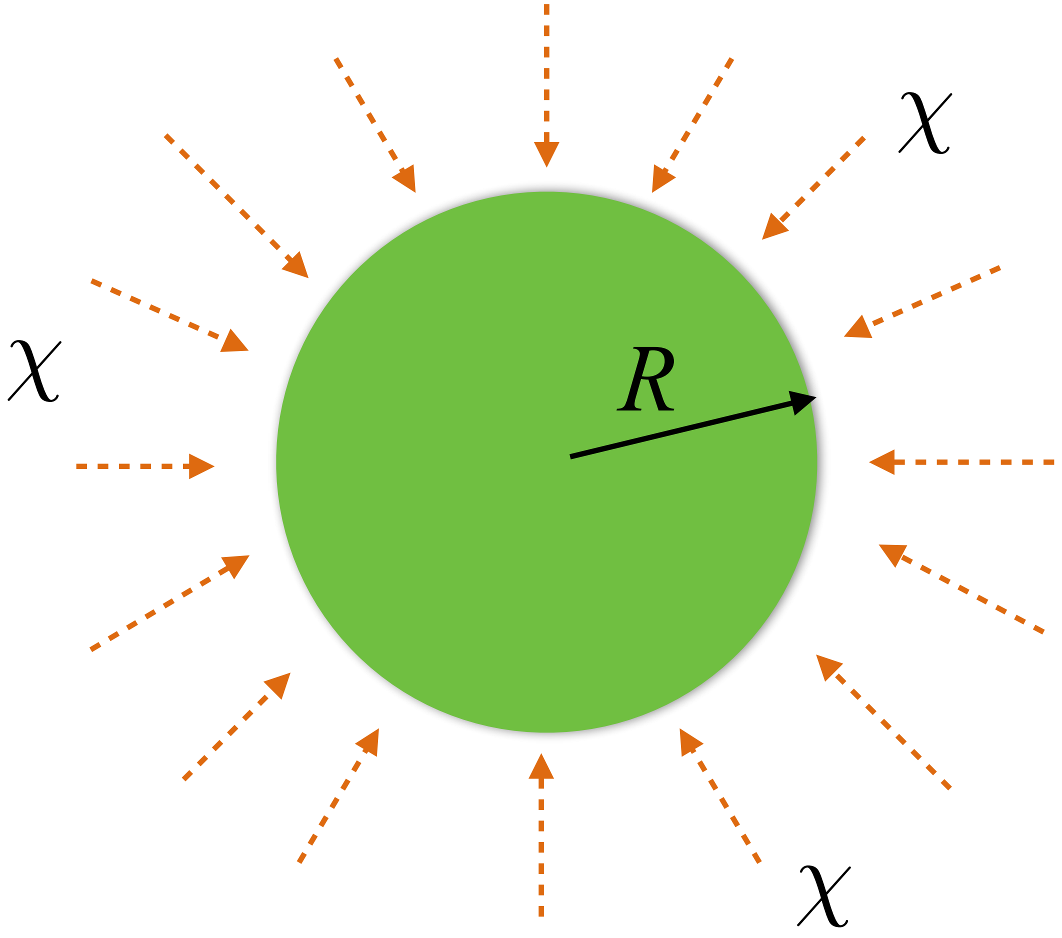
$\propto T^4$

Dark fire



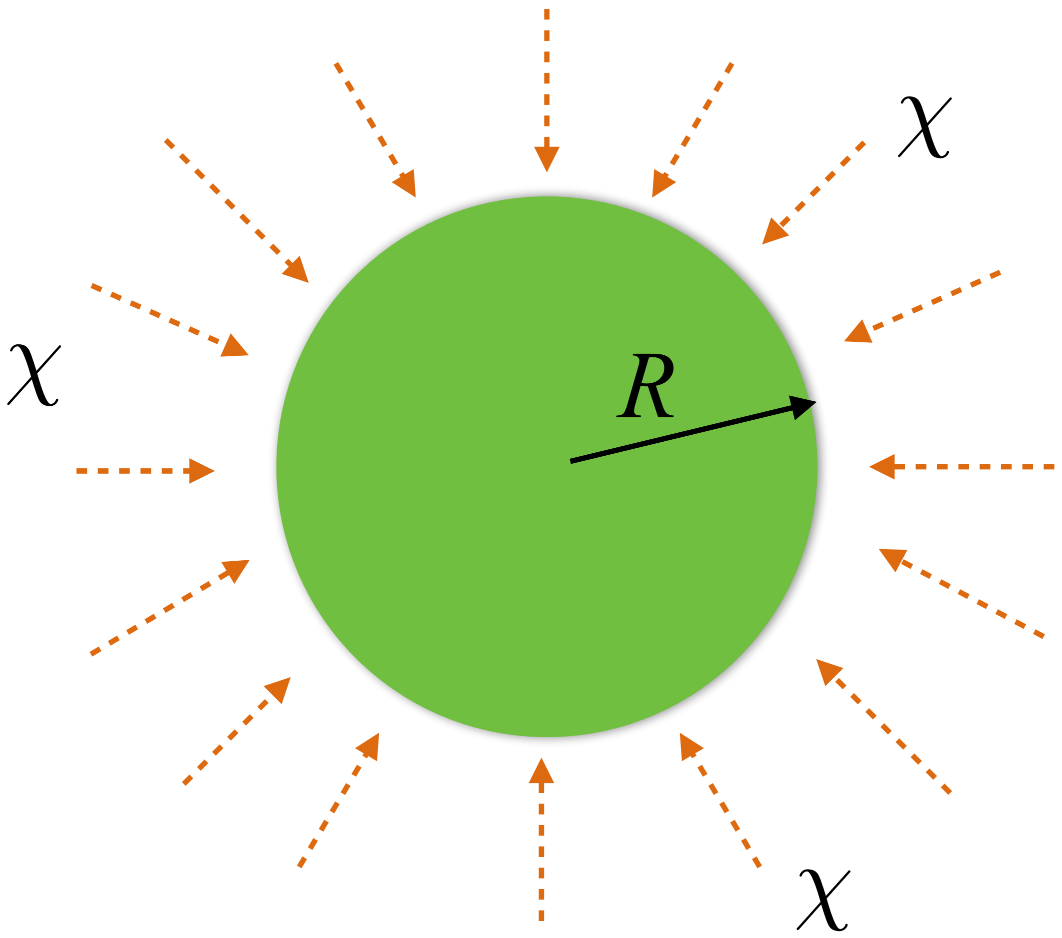
Dark fire

$$\text{heating rate} = \text{KE}_{\text{DM}} \times \frac{dN_{\text{DM}}}{dt}$$



Dark fire

$$\text{heating rate} = \text{KE}_{\text{DM}} \times \frac{dN_{\text{DM}}}{dt}$$



How hot can dark matter keep my soup?

Dark fire: soup temperature

$$\text{heating rate} = \mathrm{KE}_{\mathrm{DM}} \times \frac{dN_{\mathrm{DM}}}{dt}$$

one hit per transit:

$$\frac{R}{\text{mean free path}} = \sigma n R = 1$$

$$\sigma_{\text{threshold}} = \boxed{10^{-29}\ \mathrm{cm}^2}$$

Dark fire: soup temperature

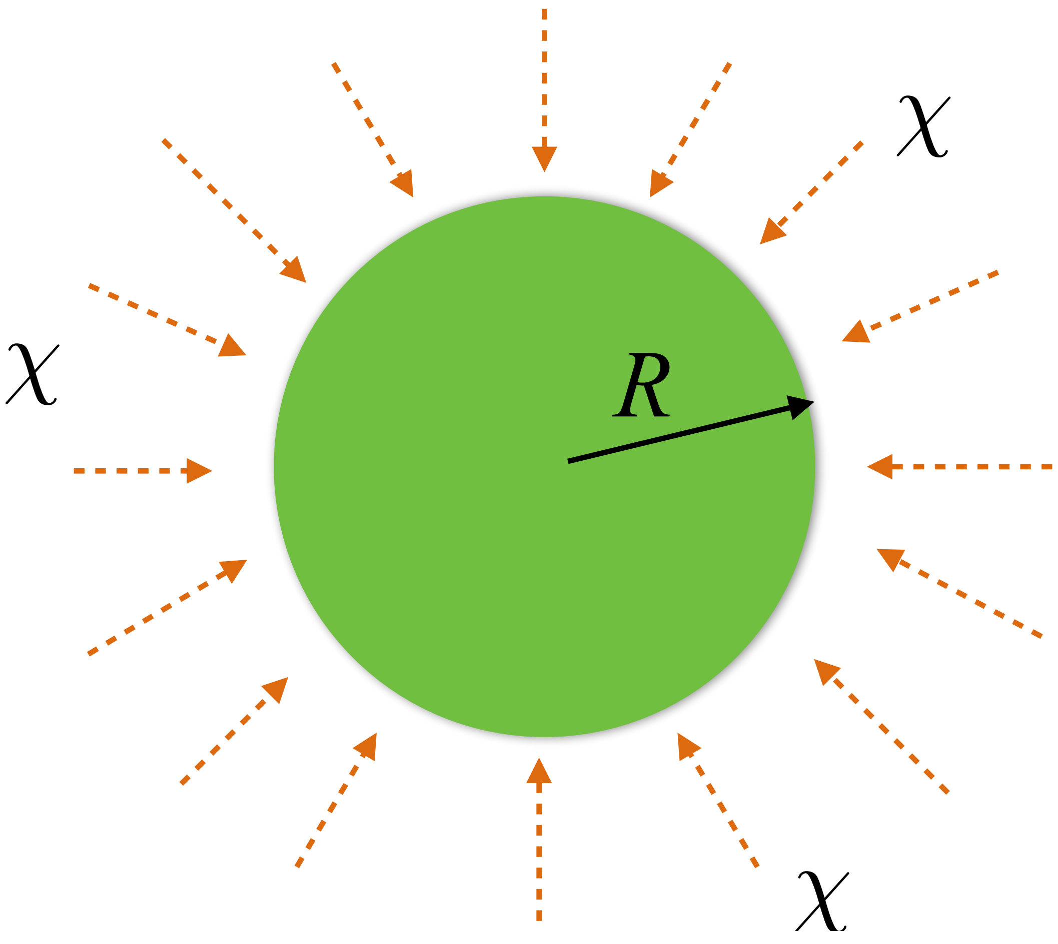
heating rate = $\mathrm{KE}_{\mathrm{DM}} \times \frac{dN_{\mathrm{DM}}}{dt}$

mass drops out!

$$\frac{1}{2} m_{\mathrm{DM}} (300\ \mathrm{km/s})^2$$

$$\frac{0.3\ \mathrm{GeV/cm^3}}{m_{\mathrm{DM}}} \mathrm{A}_{\mathrm{soup}} (300\ \mathrm{km/s})$$

Dark fire: soup temperature



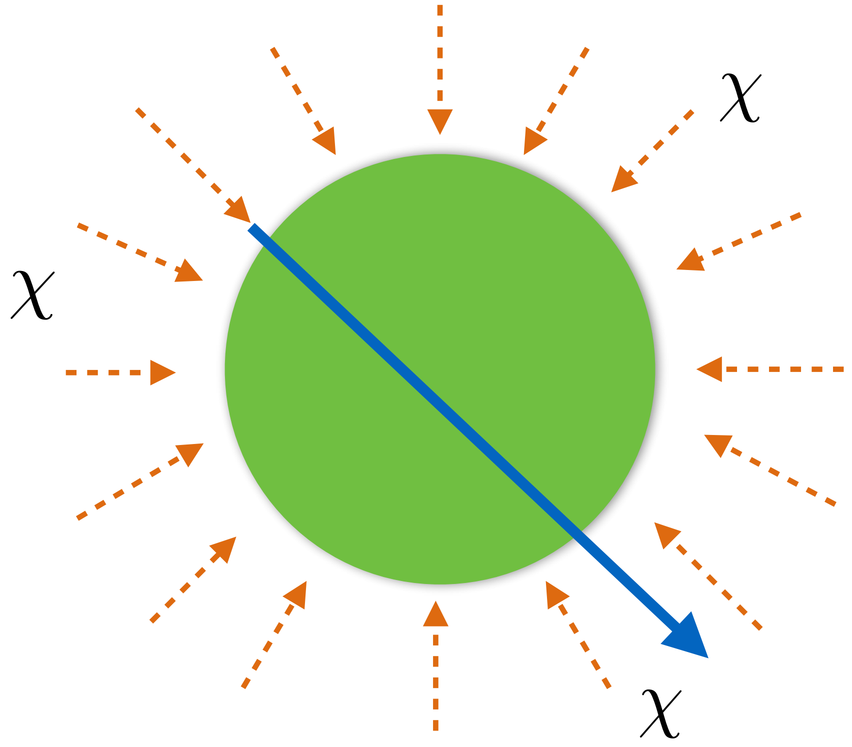
How hot can dark matter keep my soup?

$T = 0.003$ Kelvin

Need a better detector

better

$$\sigma_{\rm threshold} = \frac{m_{\rm molec}}{\rho_{\rm soup} R_{\rm soup}} = \boxed{10^{-29}\ {\rm cm}^2}$$

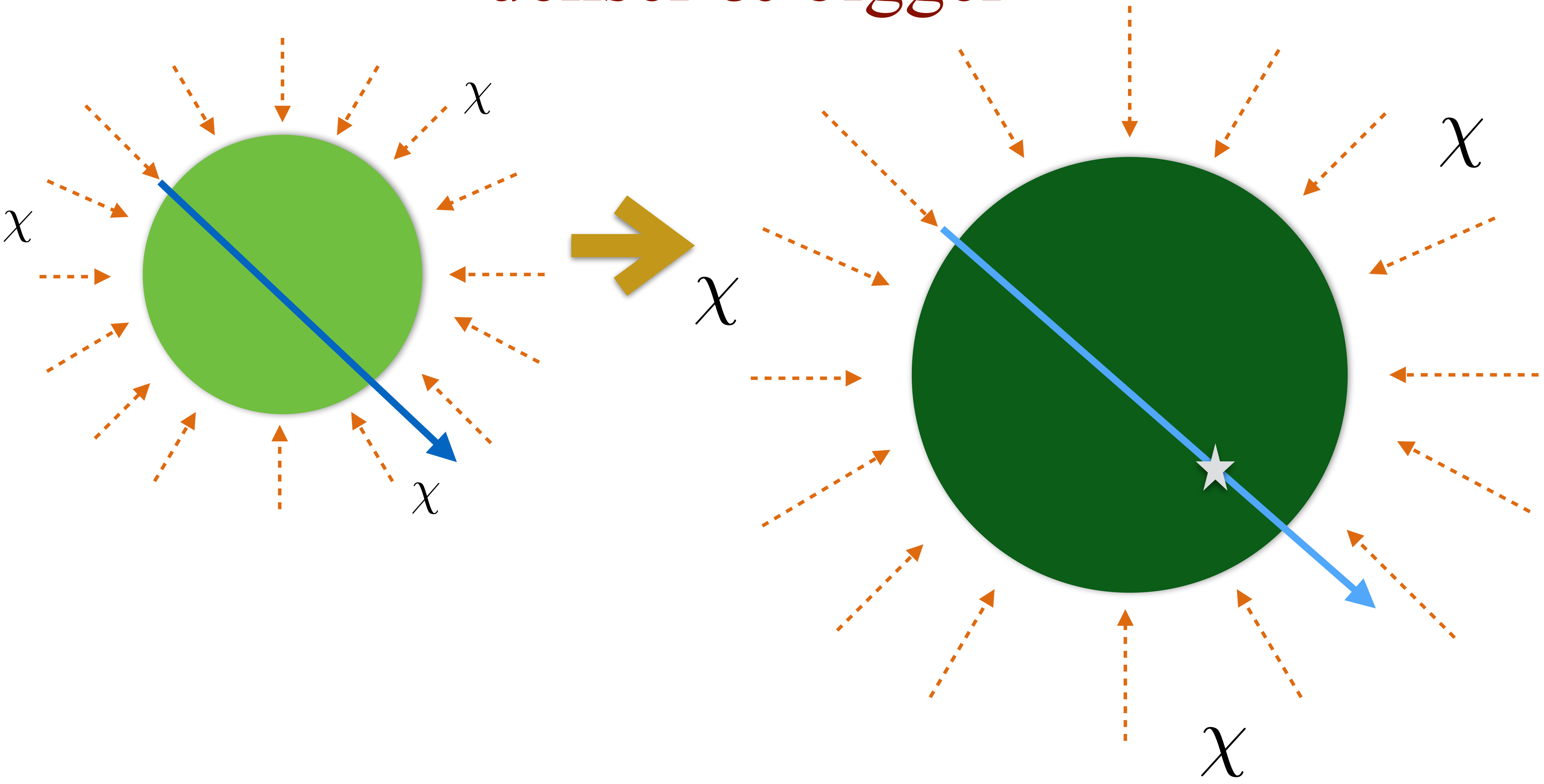


better

$$\sigma_{\text{threshold}} = \frac{m_{\text{molec}}}{\rho_{\text{soup}} R_{\text{soup}}}$$

$\boxed{<<10^{-29}\ \text{cm}^2}$

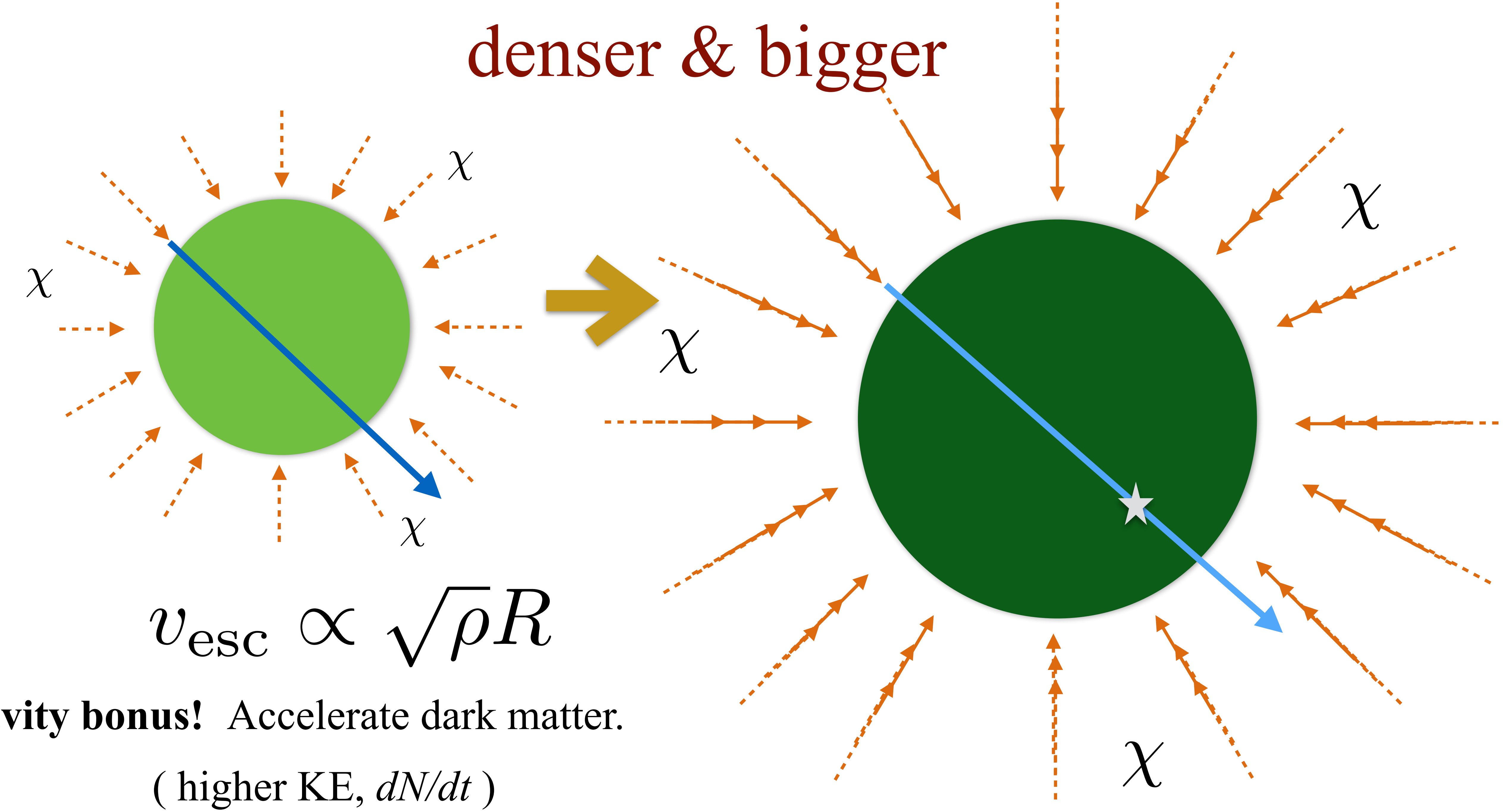
denser & bigger



better

$$\sigma_{\text{threshold}} = \frac{m_{\text{molec}}}{\rho_{\text{soup}} R_{\text{soup}}} \quad \boxed{<<10^{-29}\ \text{cm}^2}$$

denser & bigger



$$v_{\text{esc}} \propto \sqrt{\rho} R$$

Gravity bonus! Accelerate dark matter.

(higher KE, *dN/dt*)

better

left to itself	$T = 278$ Kelvin
dark fire-heated	$T = 0.003$ Kelvin

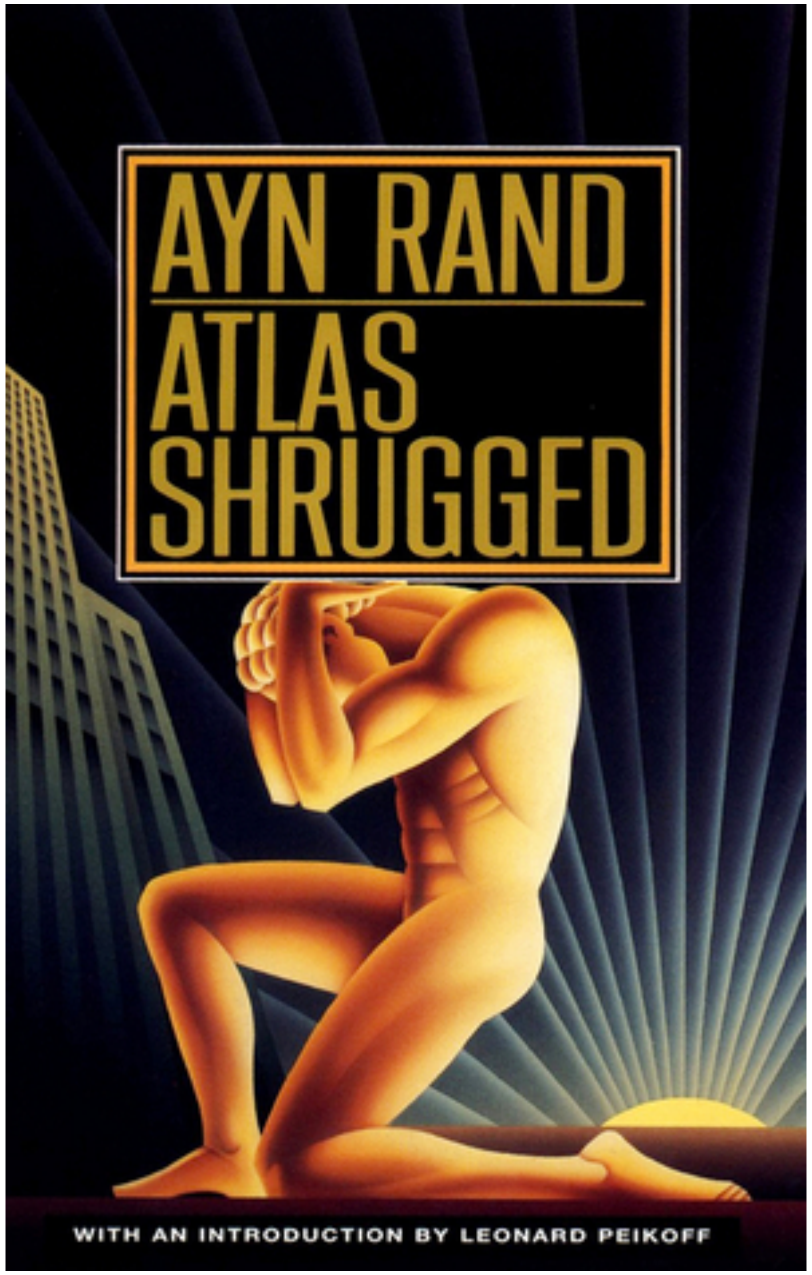
better

left to itself	$T = 278$ Kelvin	✗
dark fire-heated	$T = 0.003$ Kelvin	

be colder in natural state ✓

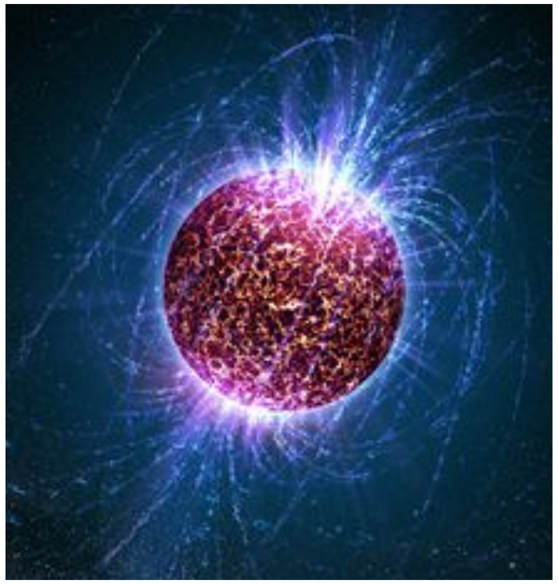
What’s dense, big, and cold?

What’s dense, big, and cold?

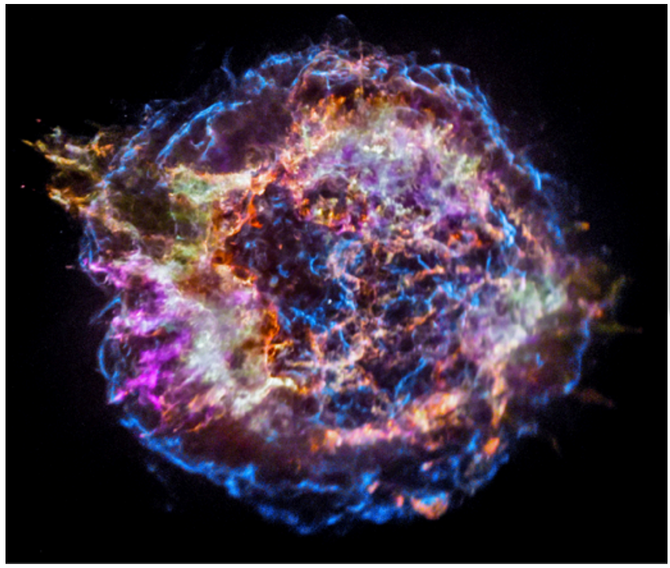


What’s dense, big, and cold?

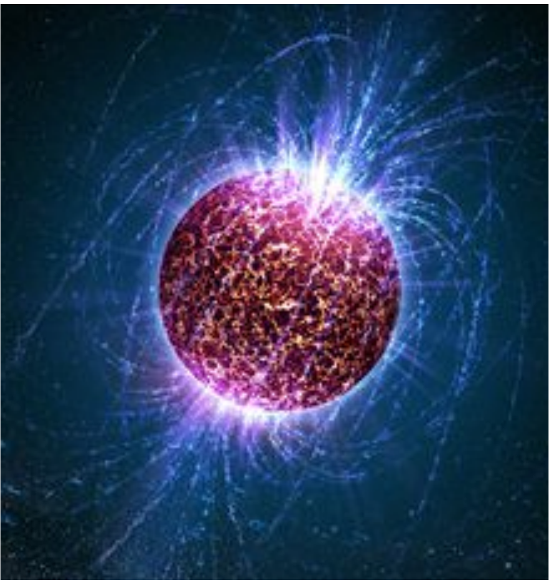
a neutron star!



1934



core-collapse
supernova

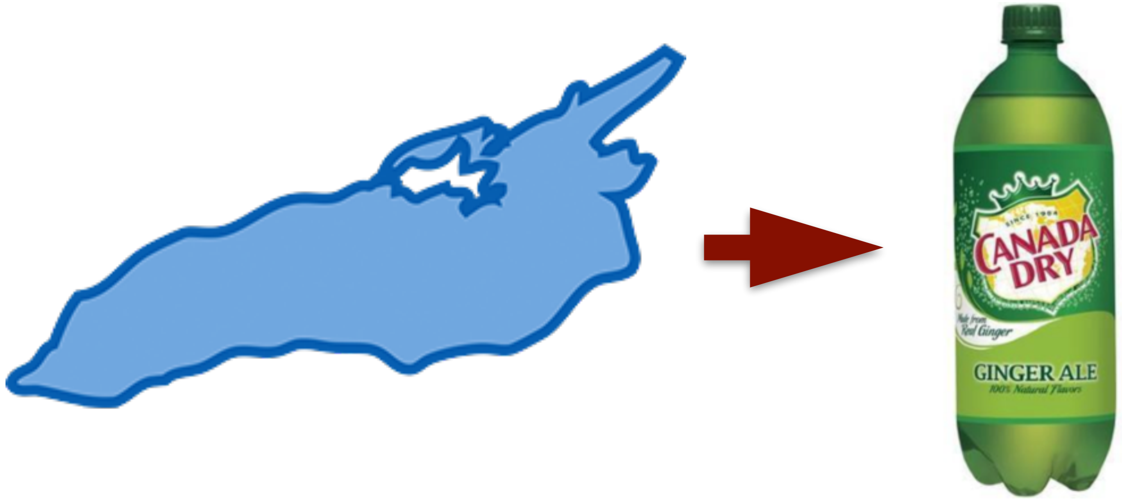


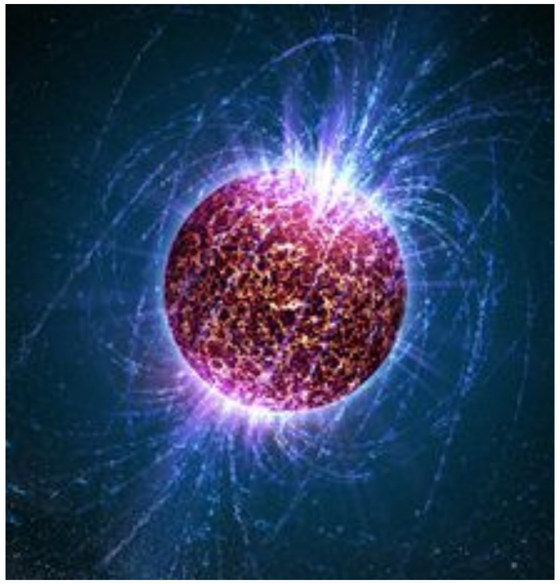
neutron star

Zwicky
again

dense

density: 7×10^{14} g/cm^3

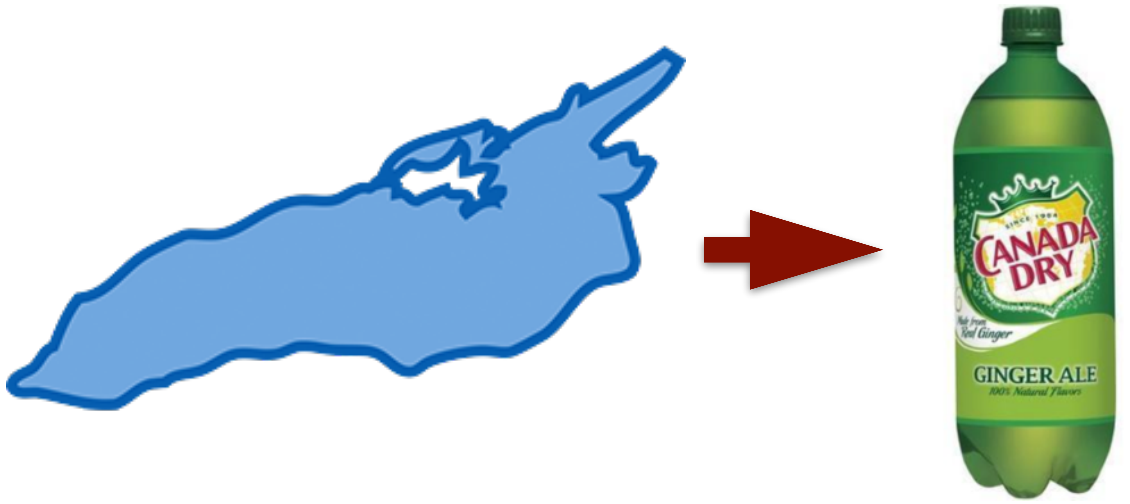


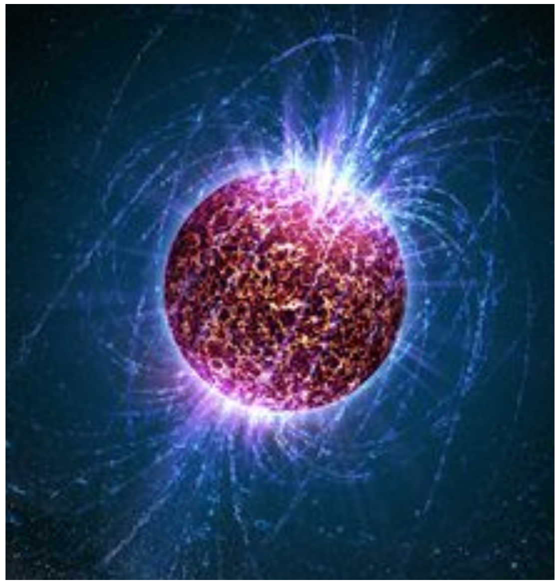


dense, big

density: 7 x 10^{14} g/cm^3

radius: 10 km

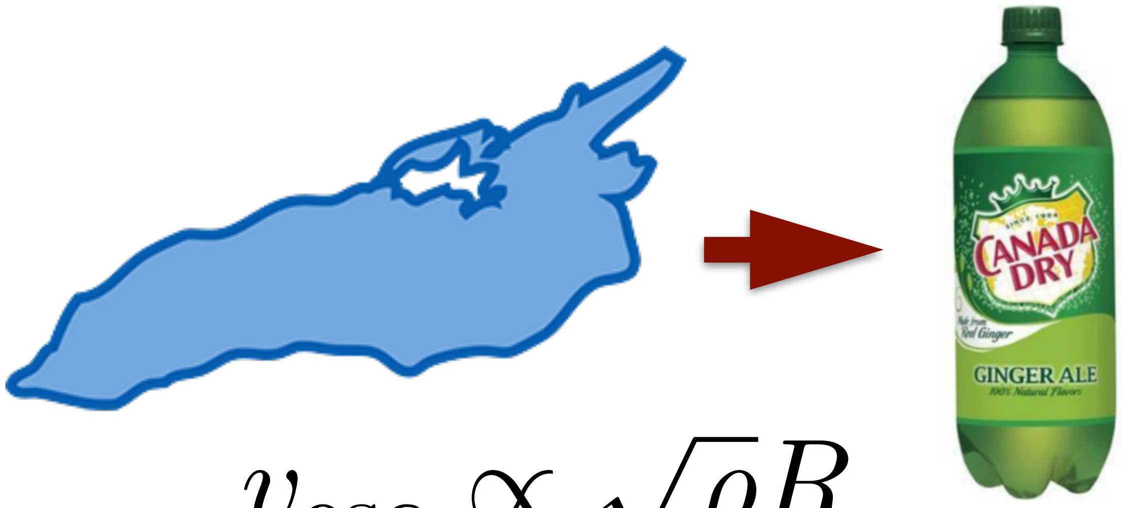


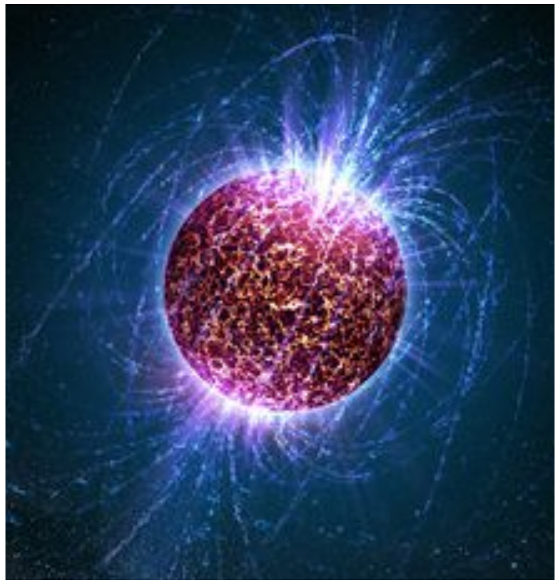


dense, big

density: 7 x 10^{14} g/cm^3

radius: 10 km



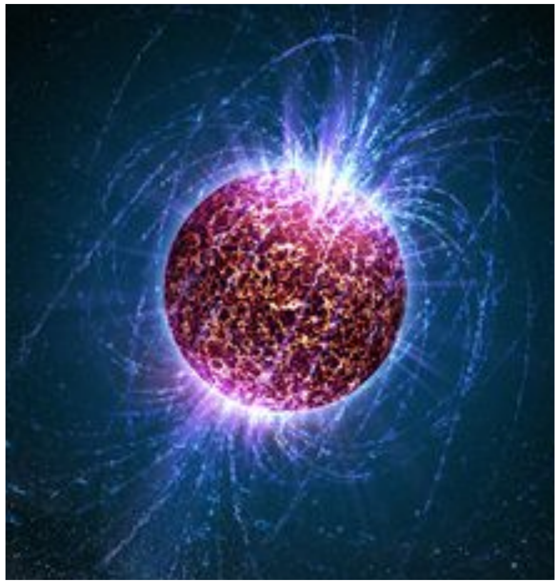


$$v_{\rm esc} \propto \sqrt{\rho}R$$
$$\simeq 0.7\,c$$

dense, big, and cold

density: 7 x 10^{14} g/cm^3

radius: 10 km



$$v_{\text{esc}} \propto \sqrt{\rho} R$$
$$\simeq 0.7\, c$$

$$\frac{20 \text{ x } 10^6 \text{ years}}{1 \text{ x } 10^9 \text{ years}}$$

$$\frac{T_{\text{effective}} \lesssim 1000 \text{ K}}{T_{\text{effective}} \sim 100 \text{ K}}$$

(compare with snowball)

Page, Lattimer, Prakash, Steiner (2004)
Yakovlev, Pethick (2004)

Detector properties

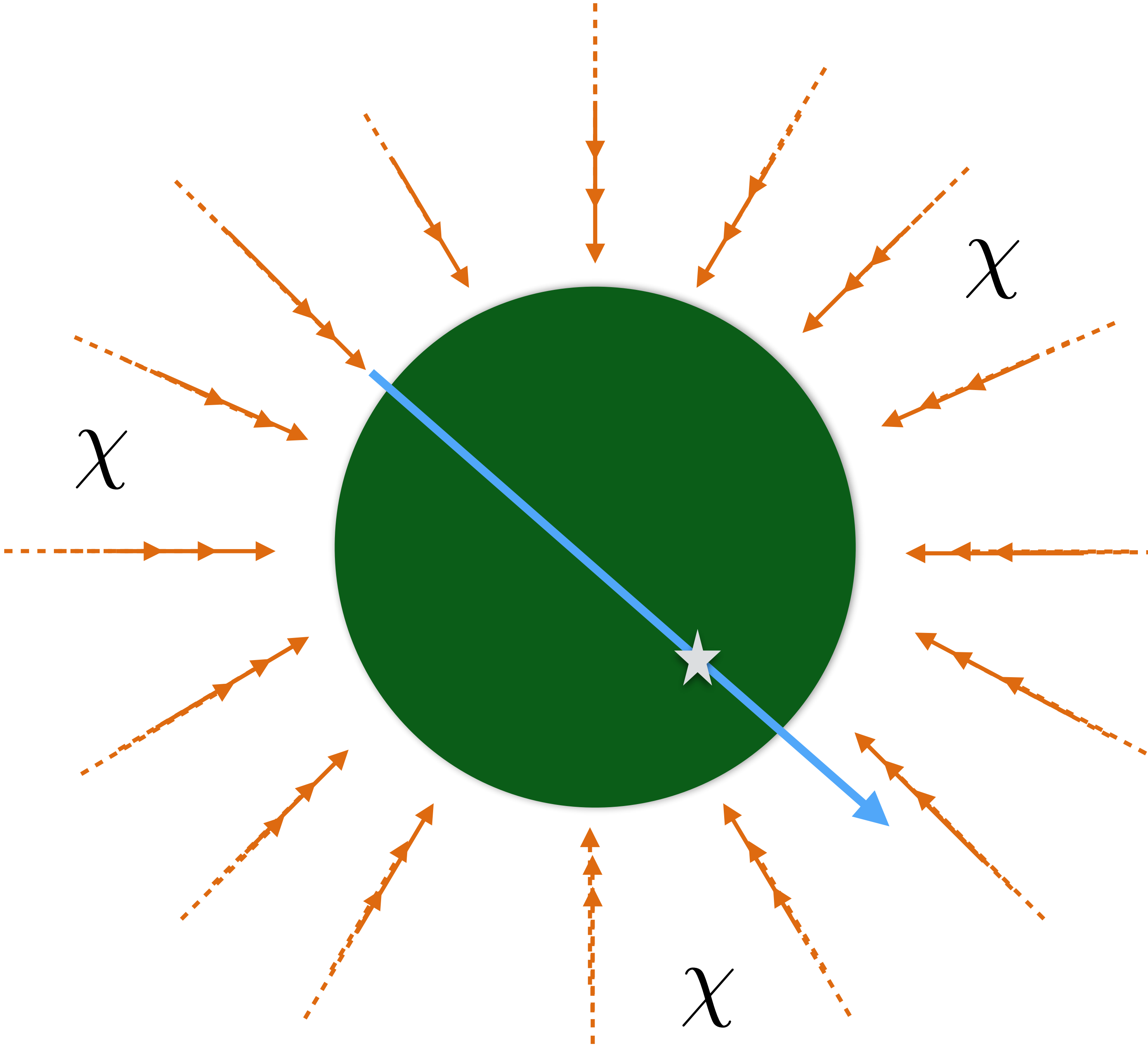
density: 7 x 10^{14} g/cm^3

radius: 10 km

$T_{\text{effective}} \sim 100$ K

Dark fires in the sky

$$\boxed{\mathrm{KE}_{\mathrm{DM}} \times \frac{dN_{\mathrm{DM}}}{dt}}$$



density: 7 x 10^{14} g/cm^3

radius: 10 km

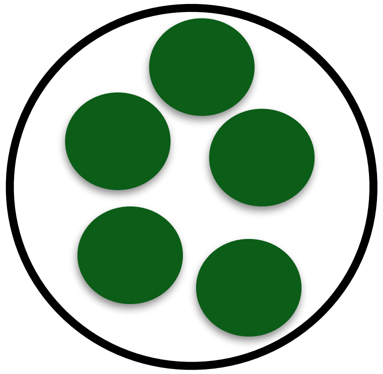
$T_{\mathrm{effective}} \sim 100$ K

M Baryakhtar, J Bramante, S Li, T Linden, **N. Raj**
Phys.Rev.Lett. *119, 131801 (2017)*
N. Raj, *P Tanedo, H-B Yu*
Phys.Rev.D. *97, 043006 (2017)*

How hot can dark matter keep my neutron star?

Zwicky misses the party

FROM LOCAL MEASUREMENTS



300 km/s

$0.3\ \mathrm{GeV/cm^3}$

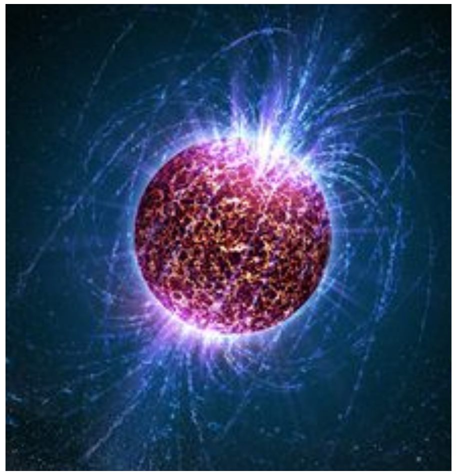
unknown to Zwicky

dark matter

1933

$$\boxed{\mathrm{KE_{DM}} \times \frac{dN_{\mathrm{DM}}}{dt}}$$

neutron star



1934

density: $7 \times 10^{14}\ \mathrm{g/cm^3}$

radius: 10 km

$T_{\mathrm{effective}} \sim 100$ K

estimable by Zwicky

How hot can dark matter keep my neutron star?

Dark fire: neutron star temperature

$$\text{heating rate} = \text{KE}_{\text{DM}} \times \frac{dN_{\text{DM}}}{dt}$$

$$\sigma_{\text{threshold}} = \frac{m_{\text{neutron}}}{\rho_{\text{NS}} R_{\text{NS}}}$$

$$\boxed{10^{-45}\ \text{cm}^2} << 10^{-29}\ \text{cm}^2$$

M Baryakhtar, J Bramante, S Li, T Linden, **N. Raj**
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Dark fire: neutron star temperature

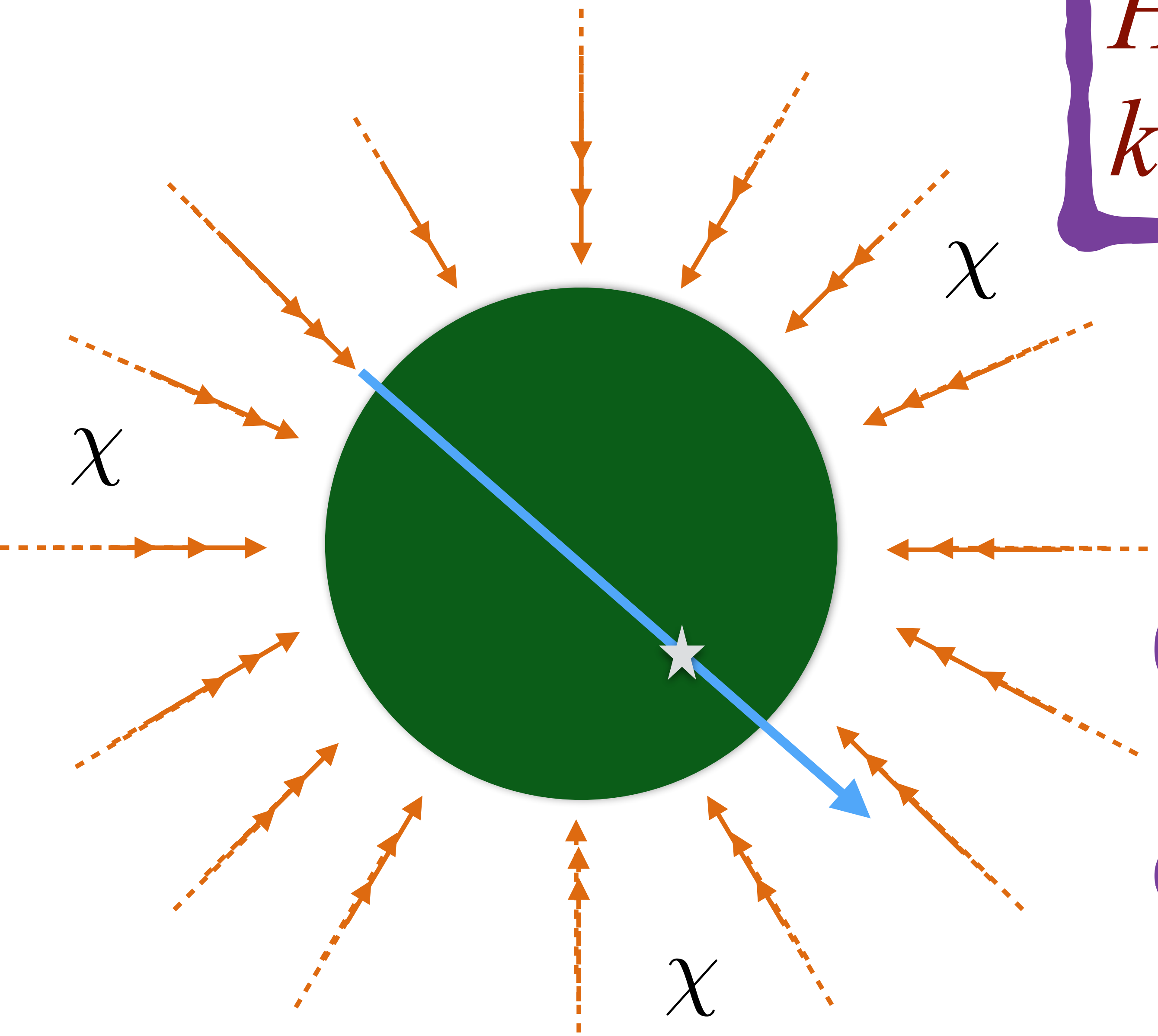
$$\text{heating rate} = \mathrm{KE}_{\mathrm{DM}} \times \frac{dN_{\mathrm{DM}}}{dt}$$

mass drops out!

$$0.35\ m_{\mathrm{DM}}$$

$$\frac{0.3\ \mathrm{GeV/cm^3}}{m_{\mathrm{DM}}}\pi R_{\mathrm{NS}}^2 \left(\frac{v_{\mathrm{esc}}}{300\ \mathrm{km/s}}\right)^2 (300\ \mathrm{km/s})$$

Dark fire: neutron star temperature



How hot can dark matter keep my neutron star?

$$T = 1750\ \text{Kelvin}$$

(infrared)

$$T = 1750 \text{ Kelvin}$$

snowball star

lava star

Coldest neutron star temperatures we were <u>able to measure</u>:

$$T = 10^5 \text{ Kelvin}$$

How to find dark fire-heated, lava-cold neutron stars?

Observation prospects

Radio telescopes
(design: pulsar discovery)

CHIME

FAST

100 old, cold neutron stars
in the local 50 pc.

O. Blaes, P. Madau (1993)

Observation prospects

Radio telescopes
(design: pulsar discovery)

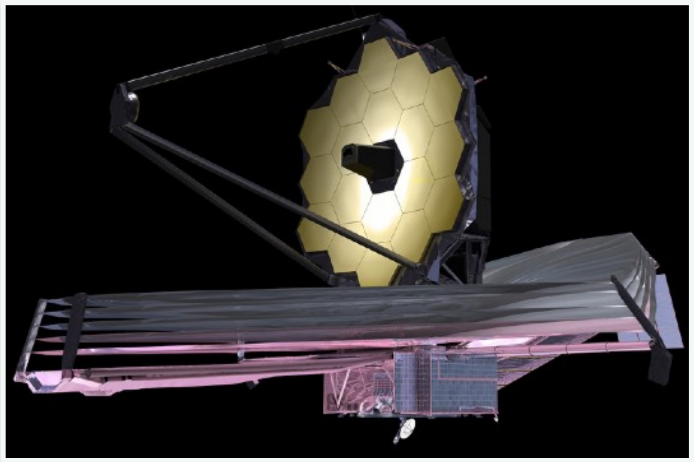
Infrared telescopes
(design: exoplanet atmosphere study)

CHIME

FAST

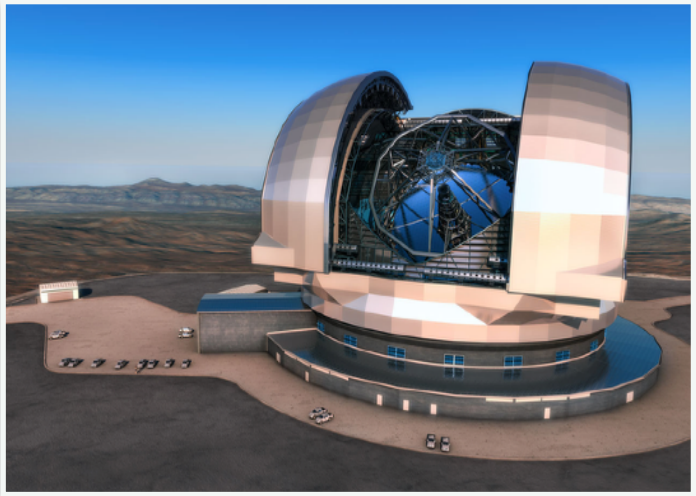
100 old, cold neutron stars
in the local 50 pc.

O. Blaes, P. Madau (1993)



James Webb

2021



European Extremely Large

2025

Thirty Meter

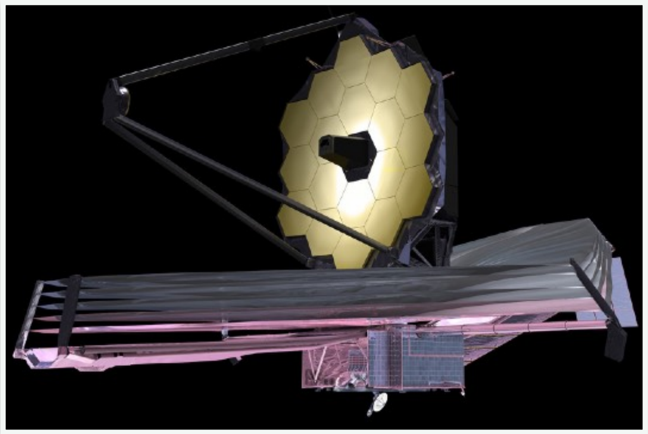
2027

Observation times

kinetic heating

(1.5 solar mass, 10 km star)

1750 K



James Webb

$$\boxed{10^5 \text{ sec} \left(\frac{d}{10\text{pc}}\right)^4}$$

Thirty Meter

$$\boxed{7 \times 10^4 \text{ sec} \left(\frac{d}{10\text{pc}}\right)^4}$$

for 2σ sensitivity

$$L \propto \boxed{(\gamma - 1) m_{\mathrm{DM}}}$$

kinetic heating

Minimum signature

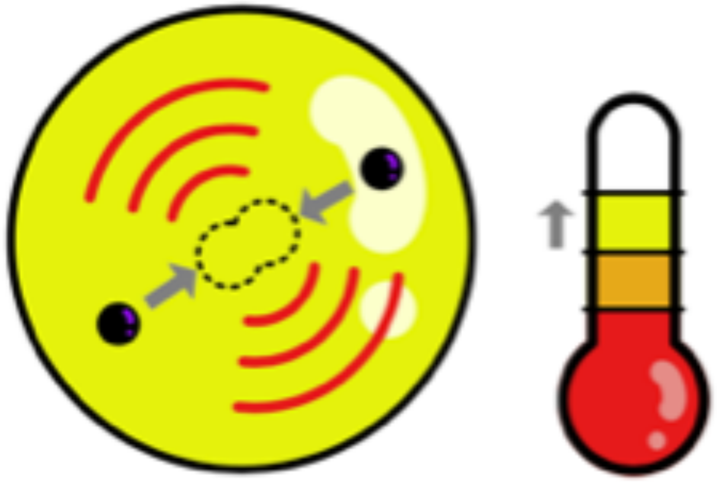
$$L \propto \boxed{\boxed{(\gamma - 1) m_{\rm DM}} + m_{\rm DM}}$$

kinetic heating

+ annihilation

Minimum signature

Possible bonus

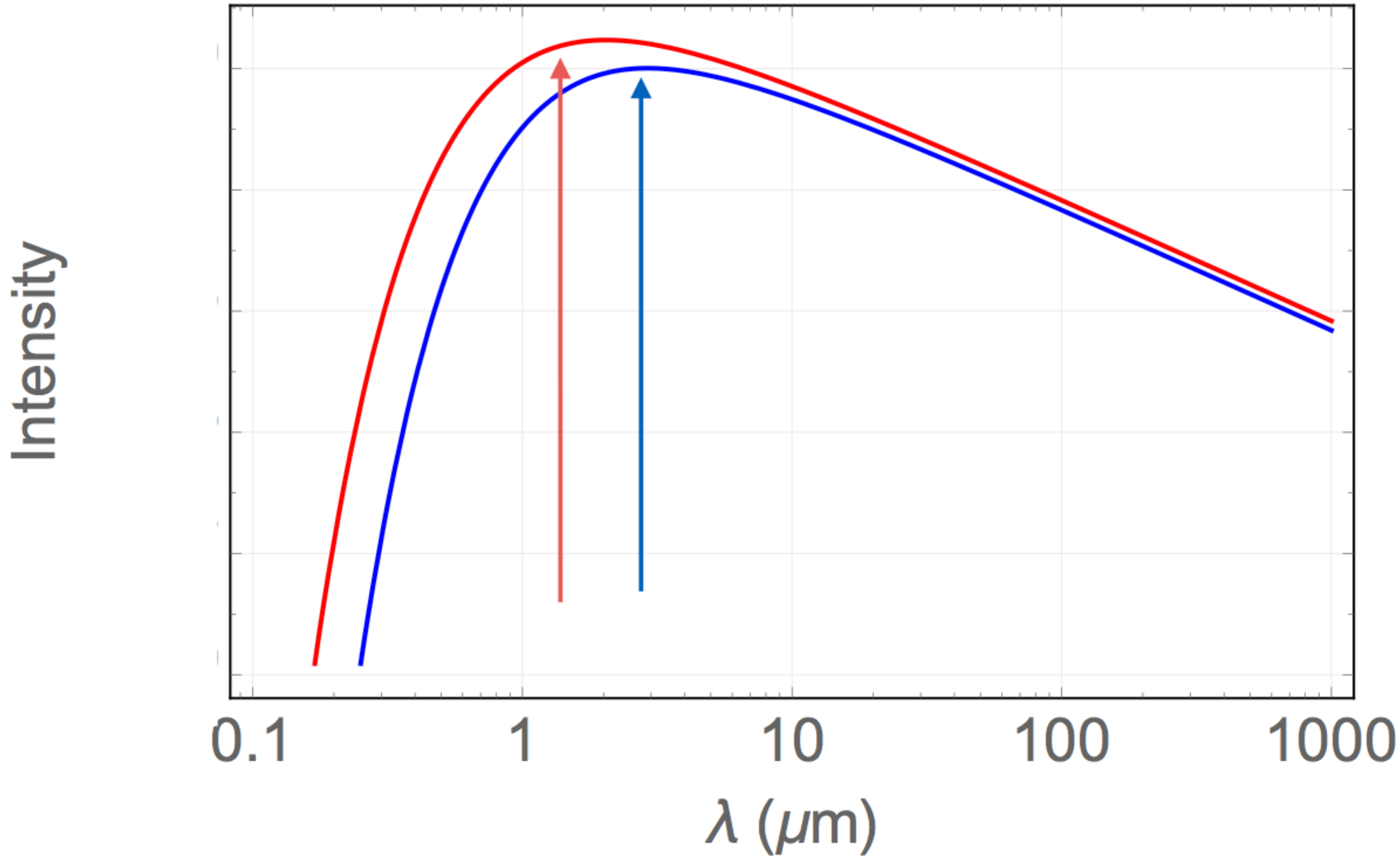


Brightness diagnosis

$$L \propto \boxed{\boxed{(\gamma - 1) m_{\rm DM}} + m_{\rm DM}}$$

kinetic heating

+ annihilation



Affects choice of filter, observation time.

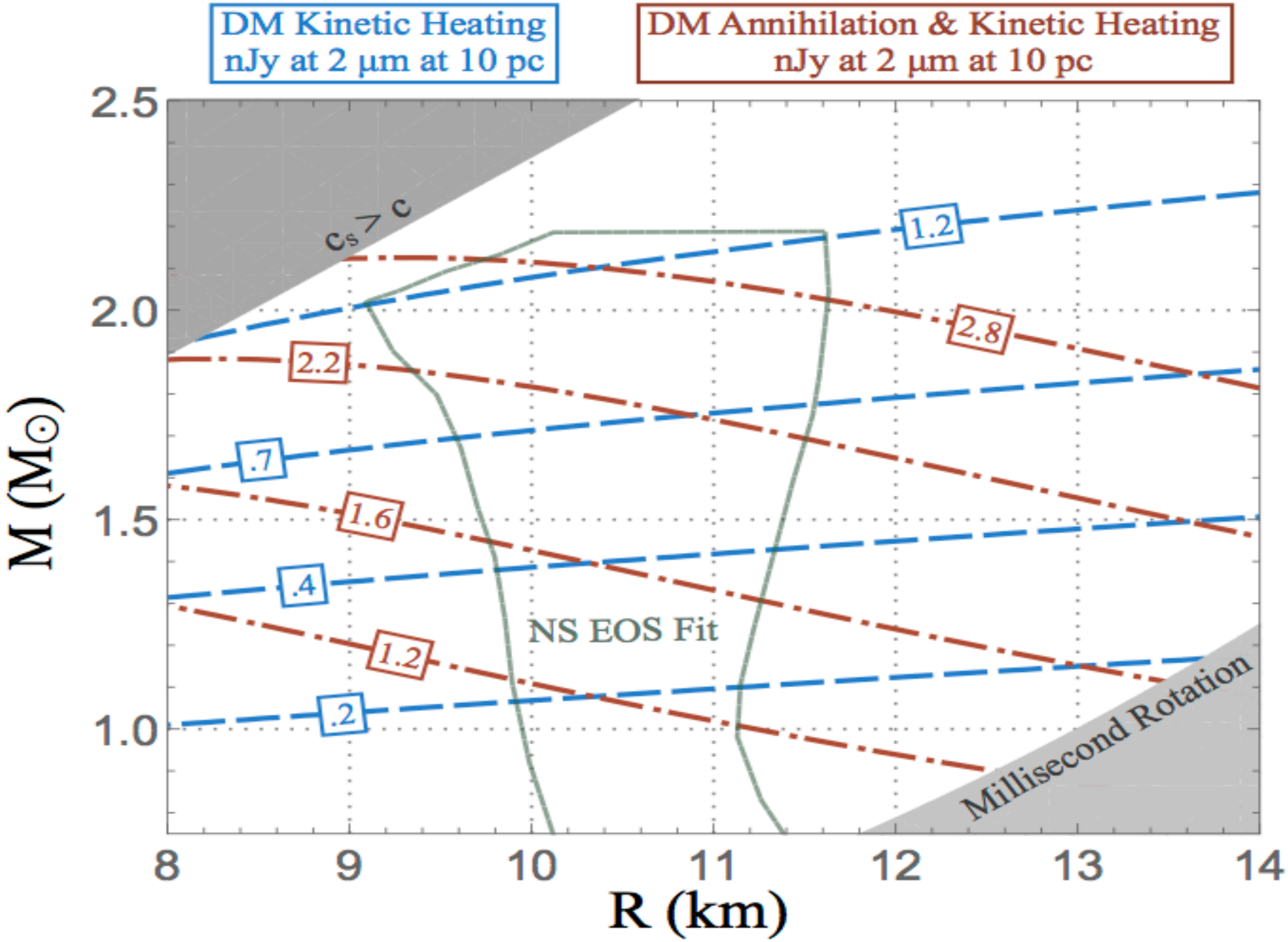
Brightness diagnosis

$$L \propto \boxed{\boxed{(\gamma - 1)m_{\rm DM}} + m_{\rm DM}}$$

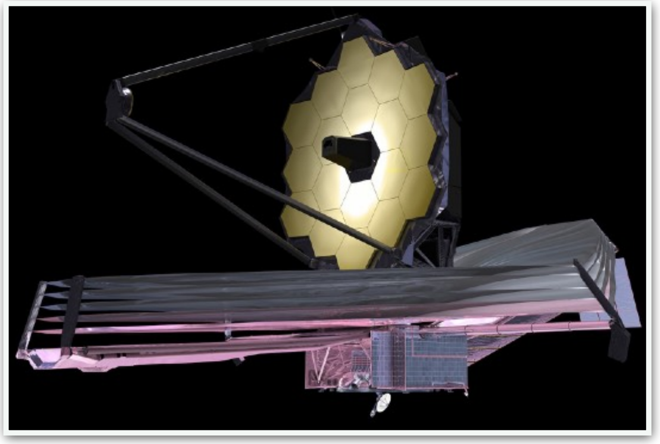
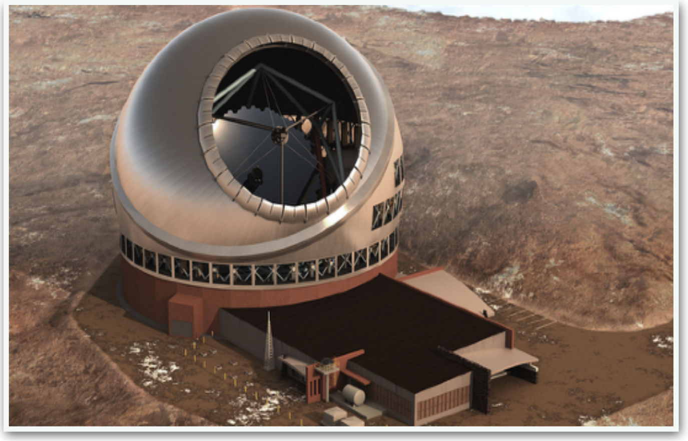
$$\left(\gamma = \frac{1}{\sqrt{1 - 2GM/R}}\right)$$

kinetic heating

\+ annihilation



Observation times

	kinetic heating	+ annihilation
(1.5 solar mass, 10 km star)	1750 K	2480 K
 James Webb	$\boxed{10^5 \text{ sec} \left(\frac{d}{10\text{pc}}\right)^4}$	$\boxed{9 \times 10^3 \text{ sec} \left(\frac{d}{10\text{pc}}\right)^4}$
 Thirty Meter	$\boxed{7 \times 10^4 \text{ sec} \left(\frac{d}{10\text{pc}}\right)^4}$	$\boxed{2 \times 10^3 \text{ sec} \left(\frac{d}{10\text{pc}}\right)^4}$

for 2σ sensitivity

Annihilation saves observation time (= $$)
by a factor of >10!

Difficulties

- Old, nearby neutron stars must turn up
- Internal backgrounds not fully understood
- Age determination not fully reliable

}

Claim of dark matter discovery
will be premature/exaggerated!
(As opposed to exclusion = clean.)

So what do you buy?

Rewards?

- Old, nearby neutron stars must turn up
- Internal backgrounds not fully understood
- Age determination not fully reliable

Claim of dark matter discovery will be premature/exaggerated! (As opposed to exclusion = clean.)

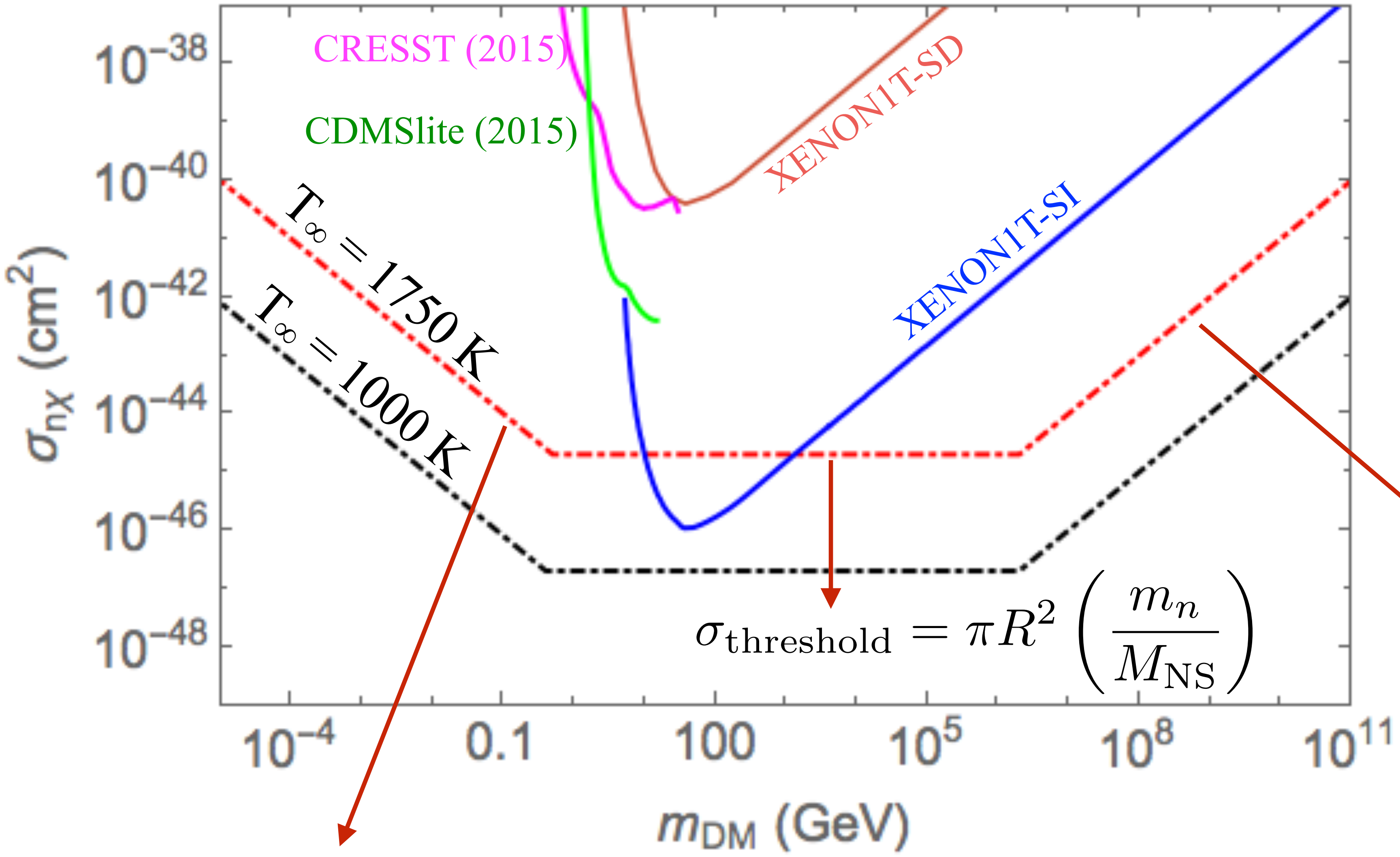
So what do you buy?

Does dark kinetic heating expand our direct detection frontiers?

(1) Low mass ☒
(2) High mass ☒
(3) Strongly interacting ☒
(4) Spin-dependent ☒
(5) v-suppressed ☒
(6) Inelastic ☒
(7) Neutrino floors ☒

Complementing terrestrial searches

M Baryakhtar, J Bramante, S Li, T Linden, **N R**;1704.01577



$$\sigma_{\rm threshold} = \pi R^2 \left(\frac{m_n}{M_{\rm NS}}\right)$$

(1) Low mass ✅

(2) High mass ✅

$$\sigma_{\rm threshold} \propto \text{ number of scatters}$$
$$= E_{\rm DM}/E_{\rm recoil}$$

$$E_{\rm recoil} \sim 2 m_n v_{\rm esc}^2$$

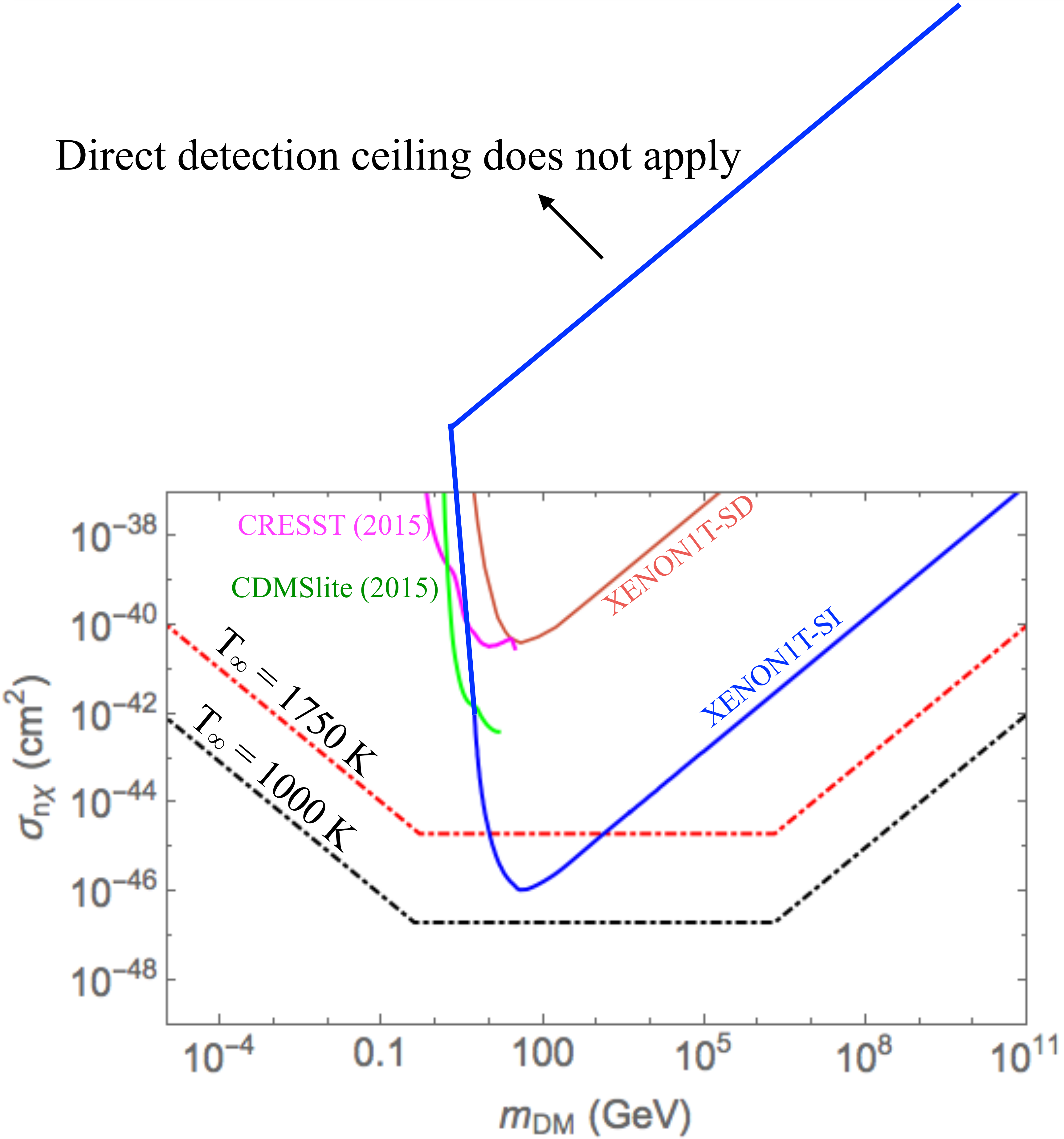
versus

$$E_{\rm DM} \sim \frac{1}{2} m_{DM} v_{\rm DM}^2$$

"Pauli blocking"

$\sigma^{-1}_{\rm threshold} \propto$ fraction of nucleons excitable to > Fermi momentum $\dfrac{\gamma m_{\rm DM} v_{\rm esc}}{p_{\rm Fermi}}$

Complementing terrestrial searches

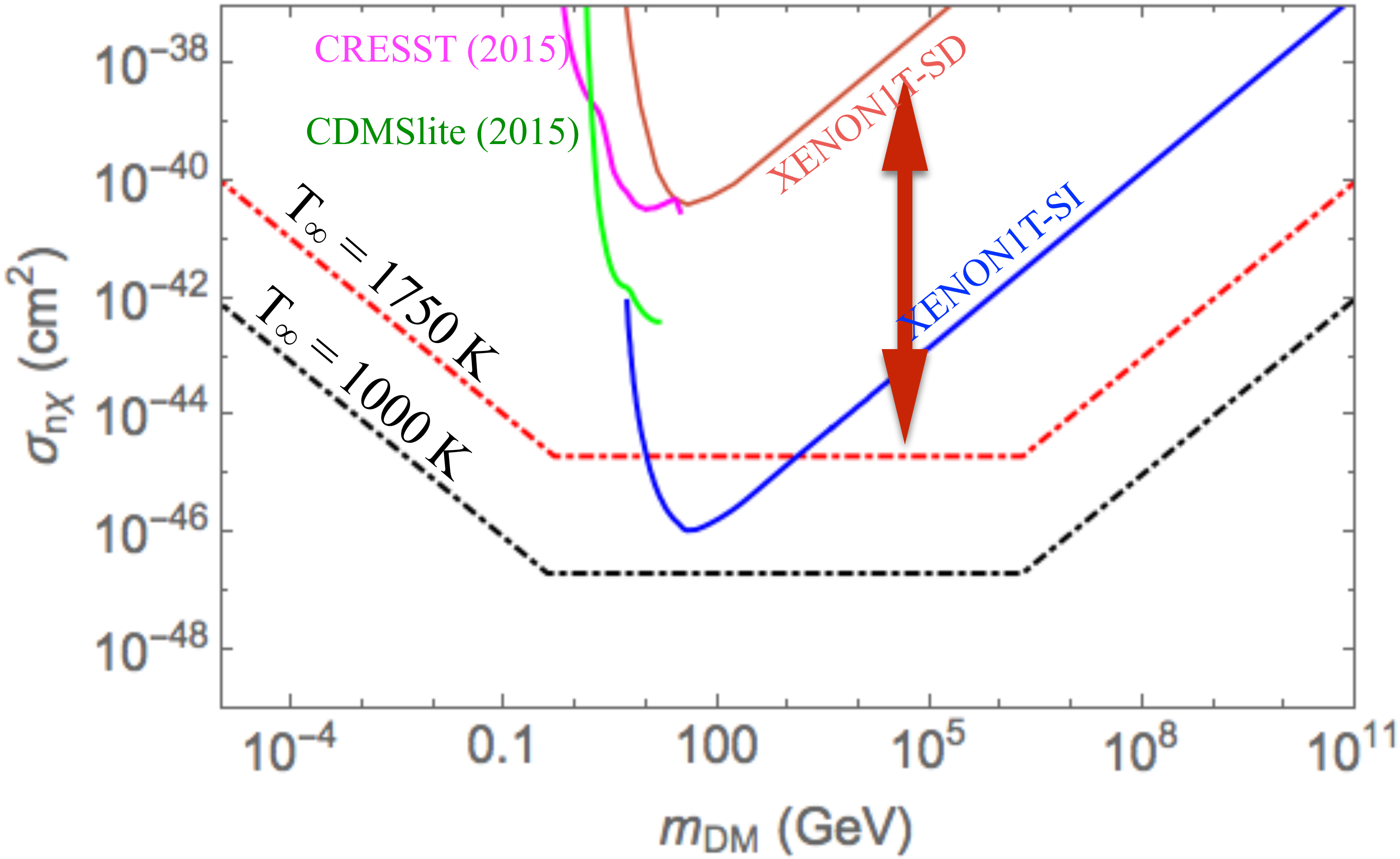


Direct detection ceiling does not apply

(1) Low mass ✓
(2) High mass ✓
(3) Strong ✓

Complementing terrestrial searches

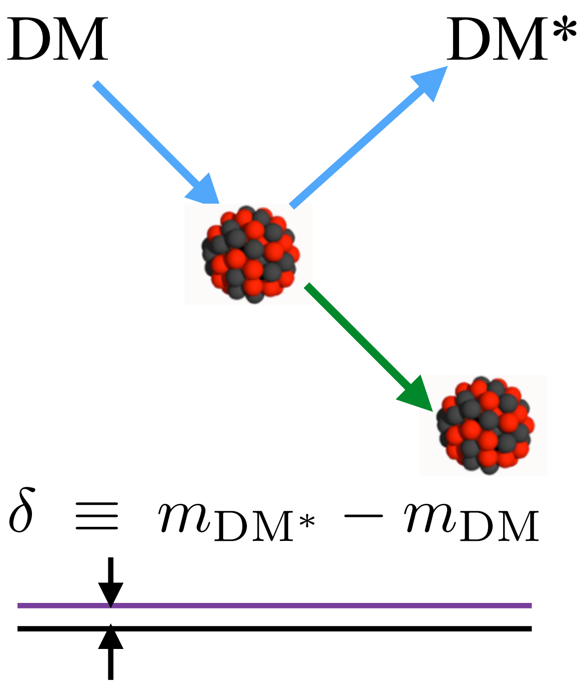
M Baryakhtar, J Bramante, S Li, T Linden, **N R**;1704.01577



(1) Low mass ✅
(2) High mass ✅
(3) Strong ✅
(4) Spin-dependent ✅

Scattering with neutrons:
apathy to nuclear coherence

Complementing terrestrial searches



$$\delta \equiv m_{\mathrm{DM}^*} - m_{\mathrm{DM}}$$

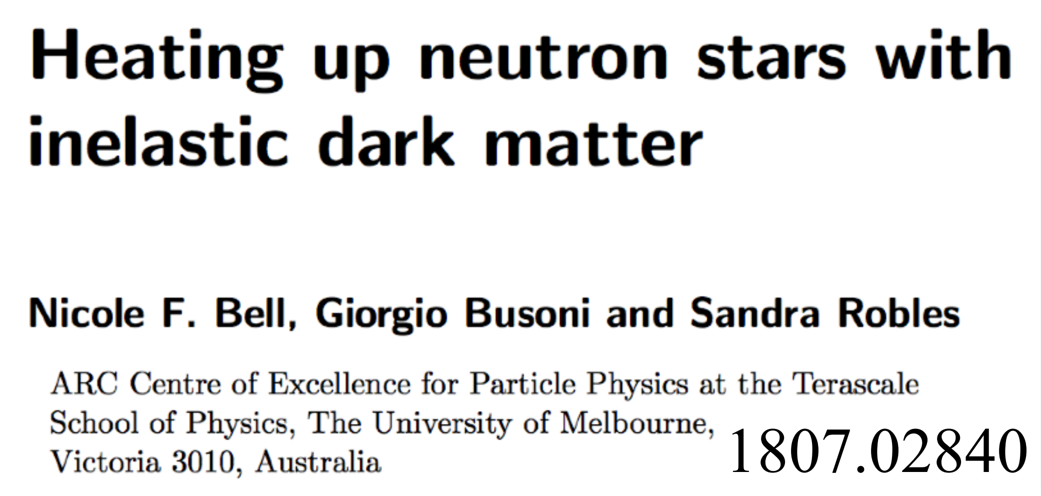
Scattering proceeds so long as mass splitting is below

$$\delta_{\max} = \frac{\mu_{n\chi} v^2}{2}$$

$$= 200 \text{ MeV}$$

(Direct detection $\delta_{\max}$: $O(100 \text{ keV})$)

(1) Low mass ☑
(2) High mass ☑
(3) Strong ☑
(4) Spin-dependent ☑
(5) Inelastic ☑

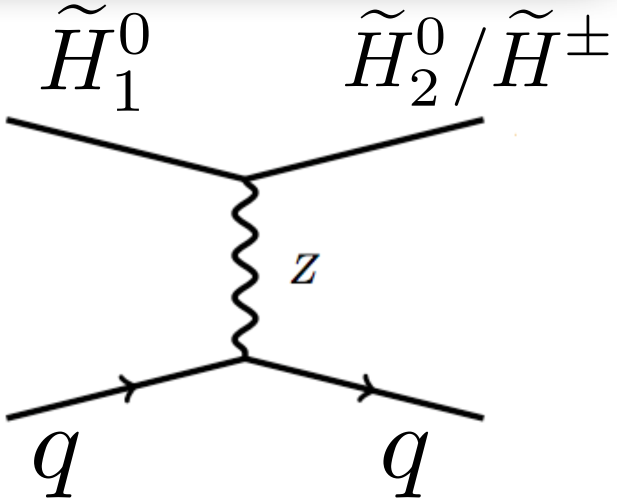


Heating up neutron stars with inelastic dark matter

Nicole F. Bell, Giorgio Busoni and Sandra Robles

ARC Centre of Excellence for Particle Physics at the Terascale
School of Physics, The University of Melbourne,
Victoria 3010, Australia

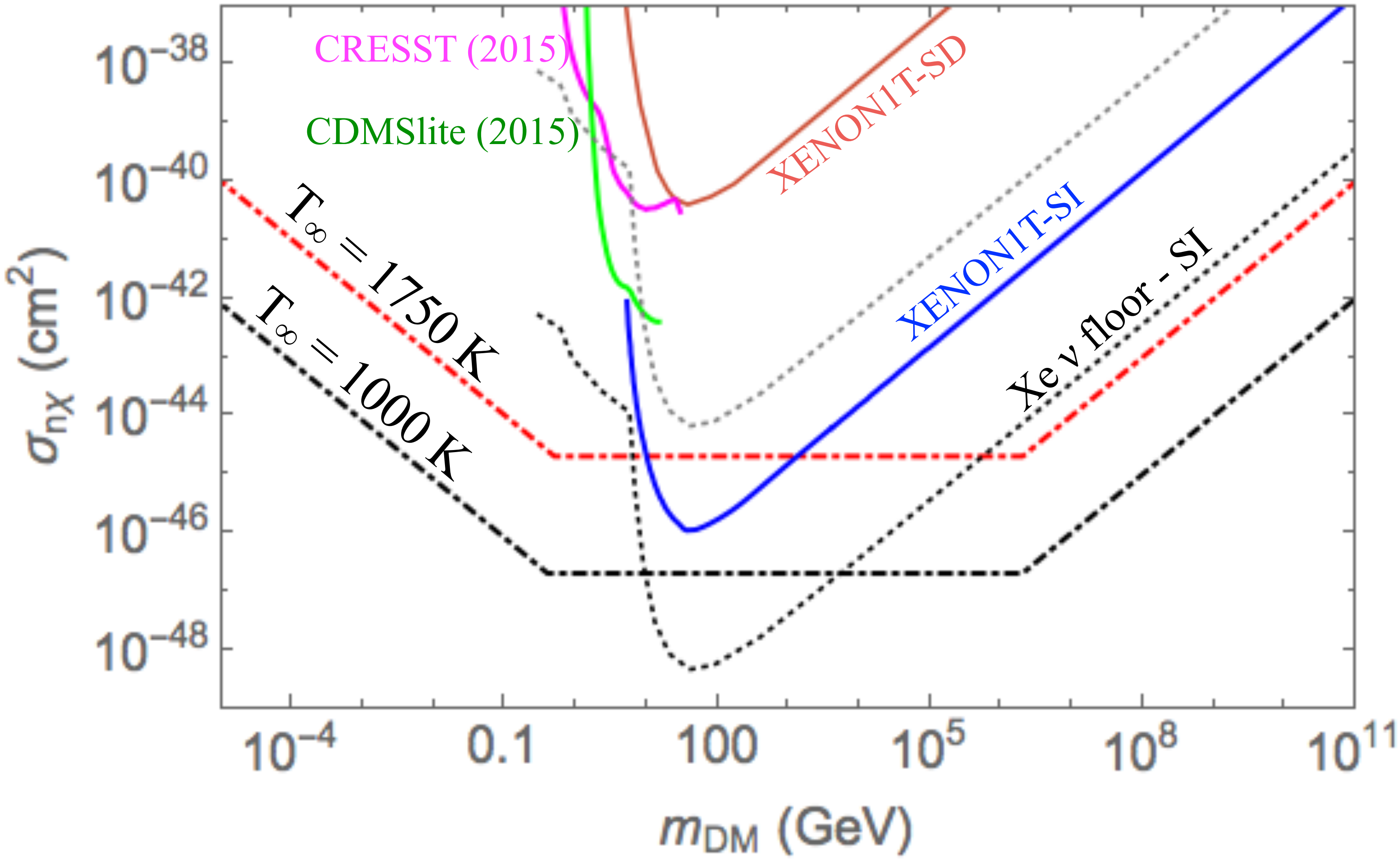
1807.02840



Great news for Higgsino lovers!

Complementing terrestrial searches

M Baryakhtar, J Bramante, S Li, T Linden, **N R**;1704.01577

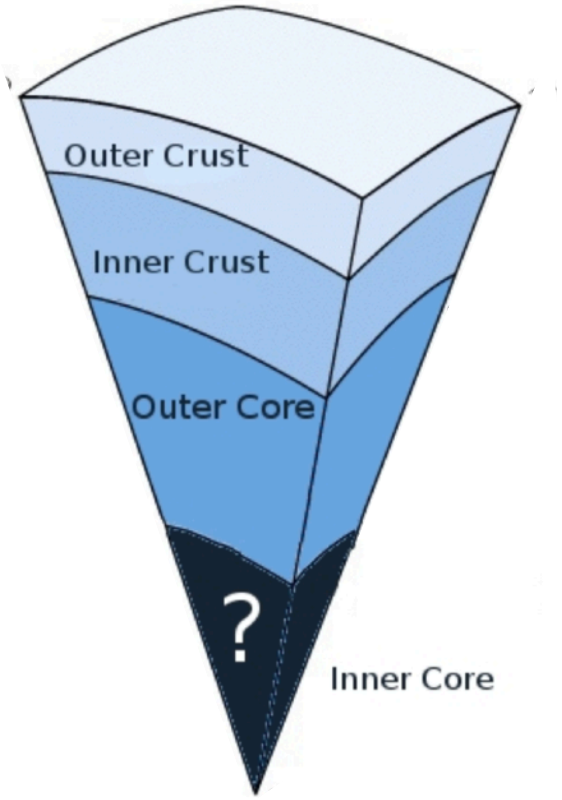


- (1) Low mass ✅
- (2) High mass ✅
- (3) Strong ✅
- (4) Spin-dependent ✅
- (5) Inelastic ✅
- (6) Neutrino floors ✅

Important variations on a theme

Are we barking down the wrong target particle?

A Joglekar, **N R**, P Tanedo, H-B. Yu; 1911.13293



Are we barking down the wrong stellar region?

J Acevedo, J Bramante, R Leane, **N R**; 1911.06334

Zippy electrons in the core

species	$\langle Y_\mathrm{T} \rangle$	mass (MeV)	$\langle p_\mathrm{F} \rangle$ (MeV)
e	0.06	0.51	146
μ	0.02	105.7	50
p	0.07	938.3	160
n	0.93	939.6	373

products of β equilibrium

speed ~c, random directions

frozen in the star

A Joglekar, **N R**, P Tanedo, H-B. Yu; 1911.13293

Trouble with zippy electrons

species	$\langle Y_{\rm T}\rangle$	mass (MeV)	$\langle p_{\rm F}\rangle$ (MeV)
e	0.06	0.51	146
μ	0.02	105.7	50
p	0.07	938.3	160
n	0.93	939.6	373

products of β equilibrium

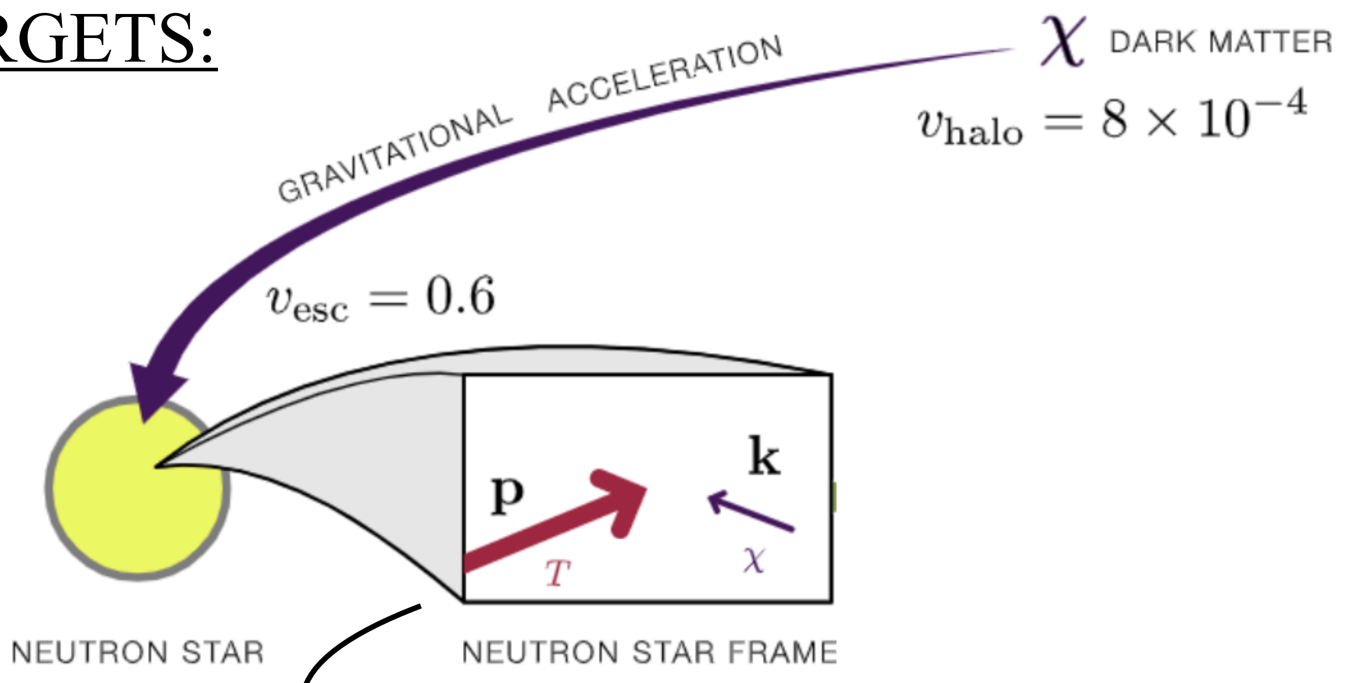
speed ~c, random directions

frozen in the star

NON-RELATIVISTIC TARGETS:

$$\text{capture probability } f = \frac{\text{scattering cross section}}{\text{geometric cross section}} = \sigma n_{\rm T} R_\star$$

ELECTRON TARGETS:



notion of "geometric cross section" breaks down!

A Joglekar, **N R**, P Tanedo, H-B. Yu; 1911.13293

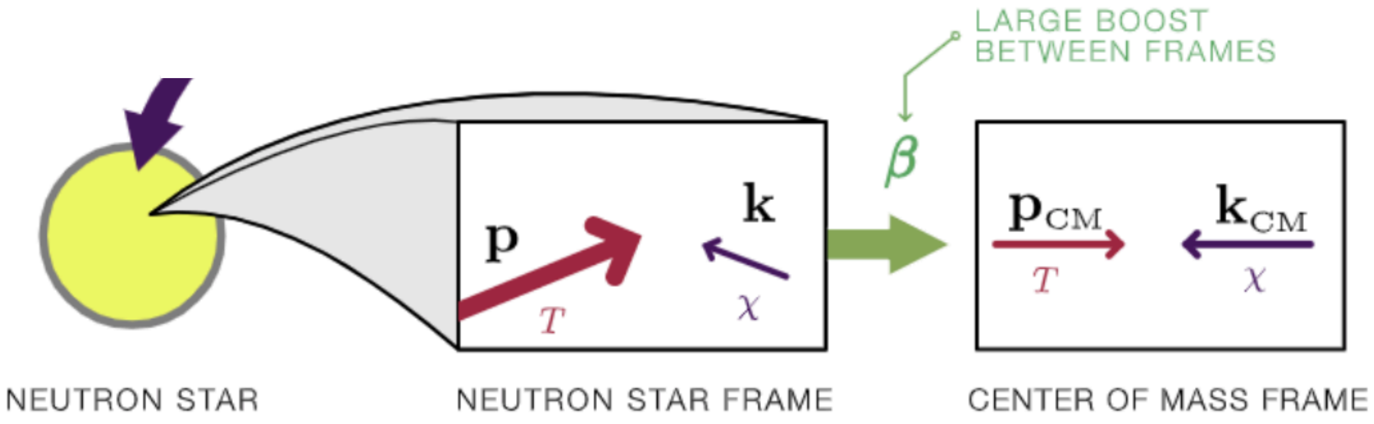
More trouble with zippy electrons

capture probability f must be Lorentz-invariant

But scattering ingredients aren’t

Fermi-Dirac distribution best known (to me) here

cross sections most conveniently expressed here



A Joglekar, **N R**, P Tanedo, H-B. Yu; 1911.13293

Putting frames together Lorentz-invariantly

DM's stellar transit time

Møller velocity

energy transfer to electron

electron Fermi momentum

$$\boxed{f = \sum_{N_{\text{hit}} \in \mathbb{Z}} \frac{\langle n_{\text{T}} \rangle \, \Delta t}{N_{\text{hit}}} \int d\Omega_{\text{NS}} \int_0^{p_{\text{F}}} d|\bar{p}| \frac{|\bar{p}|^2}{V_{\text{F}}} v_{\text{Møl}} \int d\Omega_{\text{CM}} \left(\frac{d\sigma}{d\Omega}\right)_{\text{CM}} \underbrace{\Theta\left(\Delta E + E_p - E_{\text{F}}\right)}_{\text{Pauli blocking}} \underbrace{\Theta\left(\frac{E_{\text{halo}}}{N_{\text{hit}} - 1} - \Delta E\right) \Theta\left(\Delta E - \frac{E_{\text{halo}}}{N_{\text{hit}}}\right)}_{\text{possible multiscatter capture}}}$$

electron orientation with respect to star

Fermi sphere volume

In limit of non-relativistic target (& $N_{\text{hit}} \to 1$):

$$f = \langle n_{\text{T}} \rangle \Delta t \ \text{x} \ 1 \ \text{x} \ v_{\text{DM}} \ \text{x} \ \sigma = \sigma n_{\text{T}} R_\star = \frac{\text{scattering cross section}}{\text{geometric cross section}}$$

A Joglekar, **N R**,
P Tanedo, H-B. Yu;
1911.13293

Putting frames together Lorentz-invariantly

DM's stellar transit time

Møller velocity

energy transfer to electron

electron Fermi momentum

$$f = \sum_{N_{\rm hit}\in\mathbb{Z}} \frac{\langle n_{\rm T}\rangle\,\Delta t}{N_{\rm hit}} \int d\Omega_{\rm NS} \int_0^{p_{\rm F}} d|\bar{p}|\,\frac{|\bar{p}|^2}{V_{\rm F}}\,v_{\rm Møl} \int d\Omega_{\rm CM} \left(\frac{d\sigma}{d\Omega}\right)_{\rm CM} \underbrace{\Theta\left(\Delta E + E_p - E_{\rm F}\right)}_{\text{Pauli blocking}} \underbrace{\Theta\left(\frac{E_{\rm halo}}{N_{\rm hit}-1} - \Delta E\right)\Theta\left(\Delta E - \frac{E_{\rm halo}}{N_{\rm hit}}\right)}_{\text{possible multiscatter capture}}$$

electron orientation with respect to star

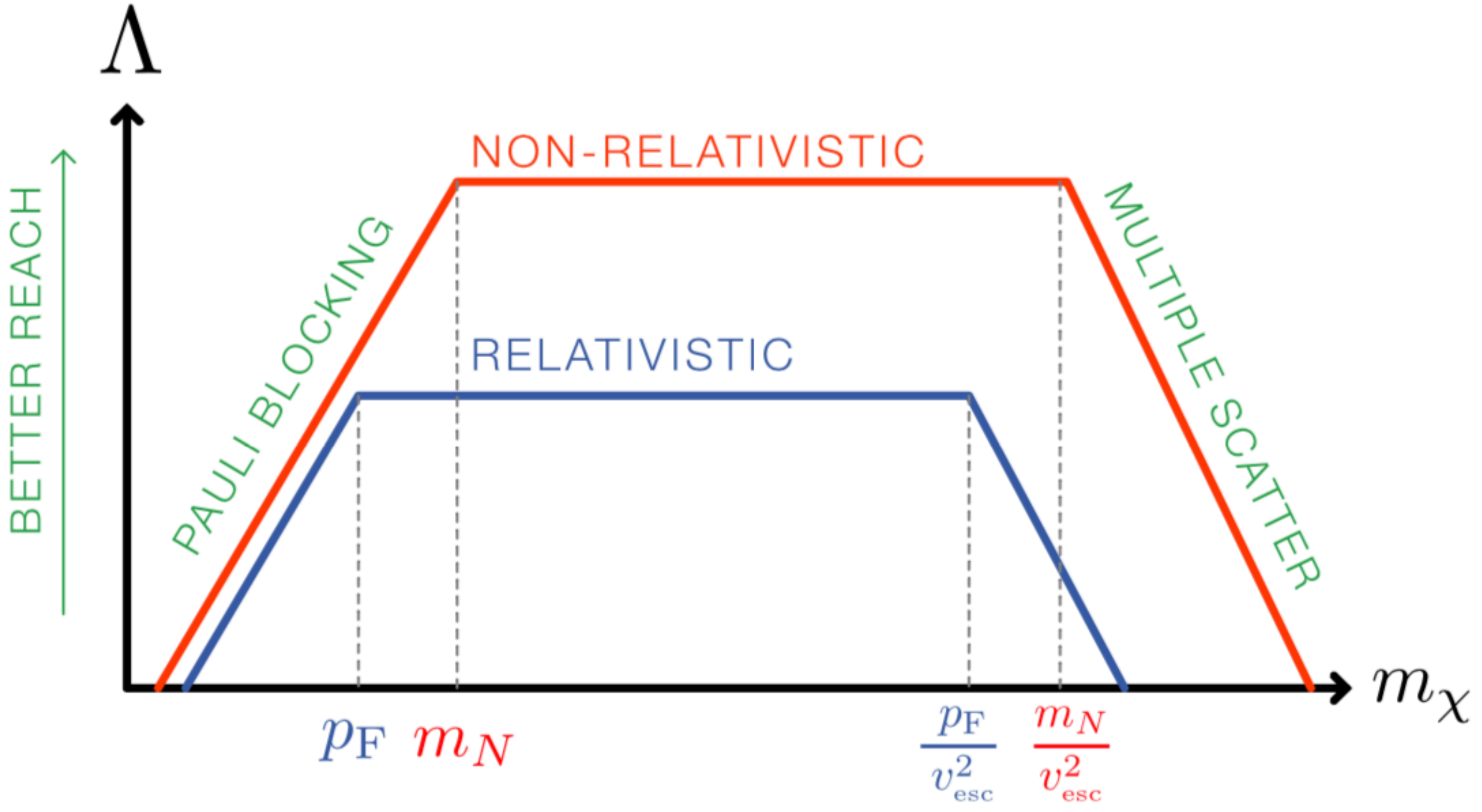
Fermi sphere volume

In limit of non-relativistic target (& $N_{\rm hit} \to 1$):

$$f = \langle n_{\rm T}\rangle\Delta t \;\text{x}\; 1 \;\text{x}\; v_{\rm DM} \;\text{x}\; \sigma \;=\; \sigma n_{\rm T} R_\star = \frac{\text{scattering cross section}}{\text{geometric cross section}}$$

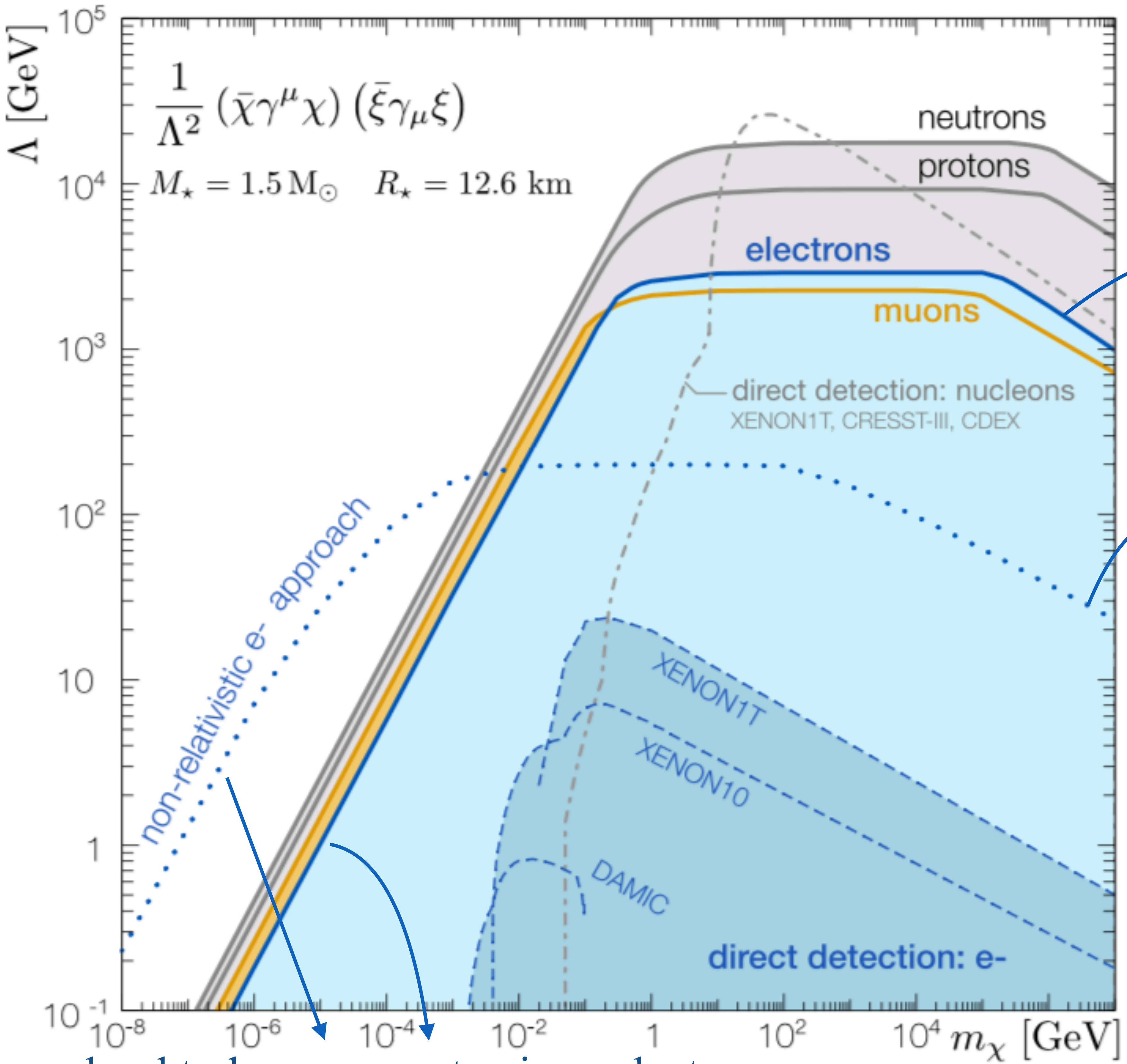
For short-distance interactions:

$$\frac{(\bar{\chi}\Gamma_\chi\chi)(\bar{f}\Gamma_f f)}{\Lambda^2}$$



A Joglekar, **N R**,
P Tanedo, H-B. Yu;
1911.13293

“Electron star” dark matter detection



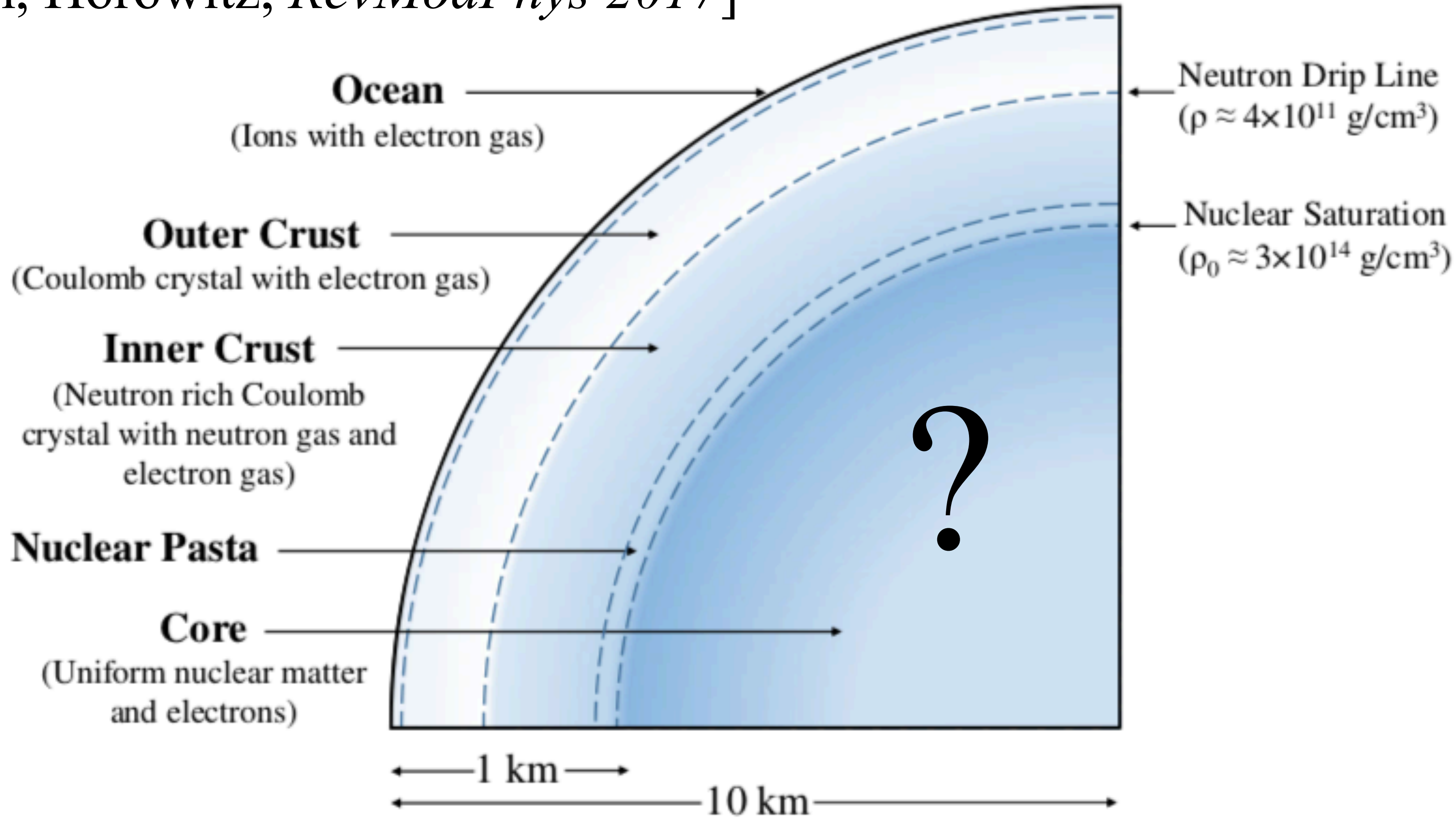
cross section ∝ Fermi energy^2
[(150 MeV)^2]

cross section ∝ electron mass^2
[(0.5 MeV)^2]

hard to lose energy to zippy electrons

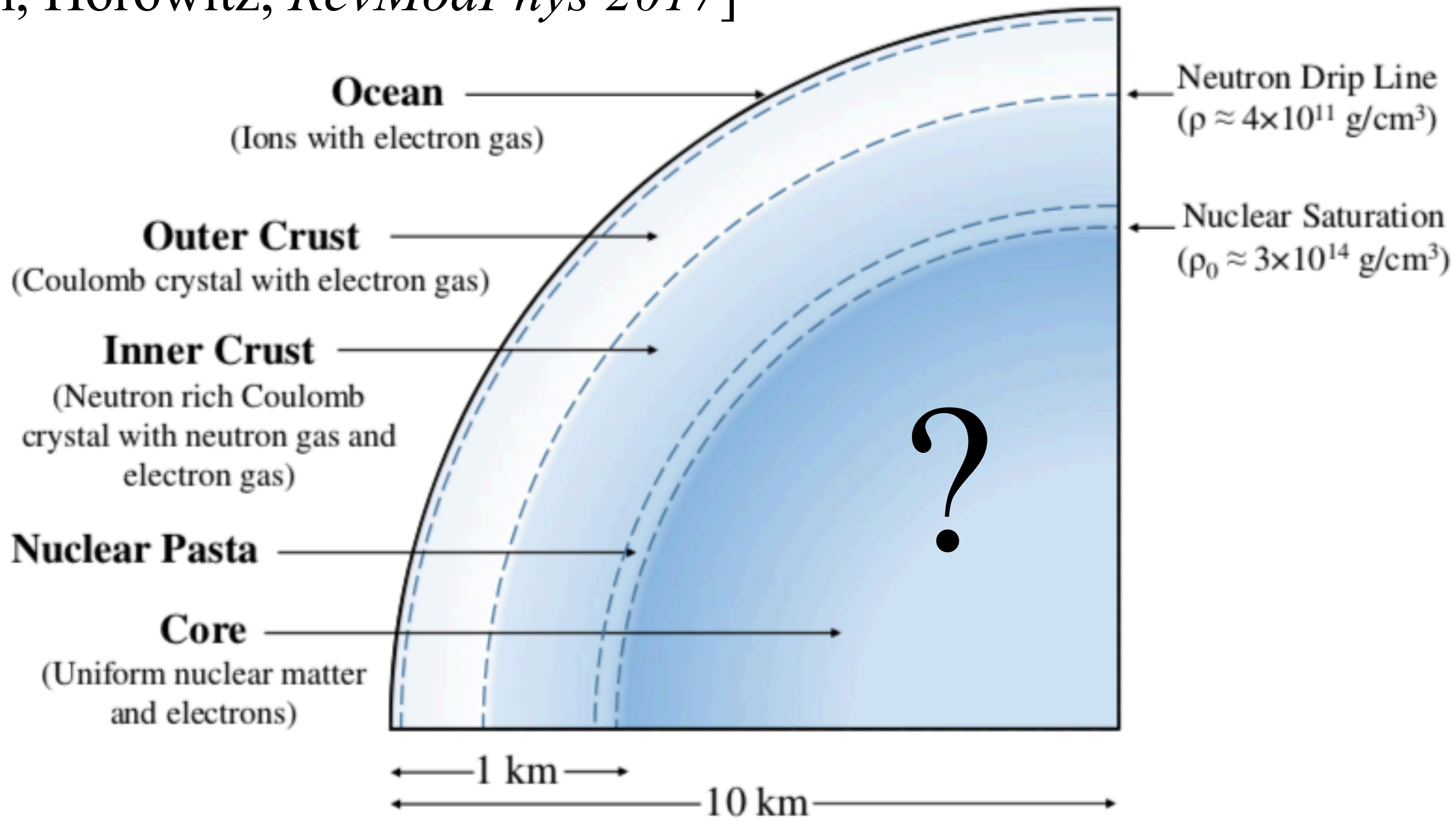
A Joglekar, **N R**, P Tanedo, H-B. Yu; 1911.13293

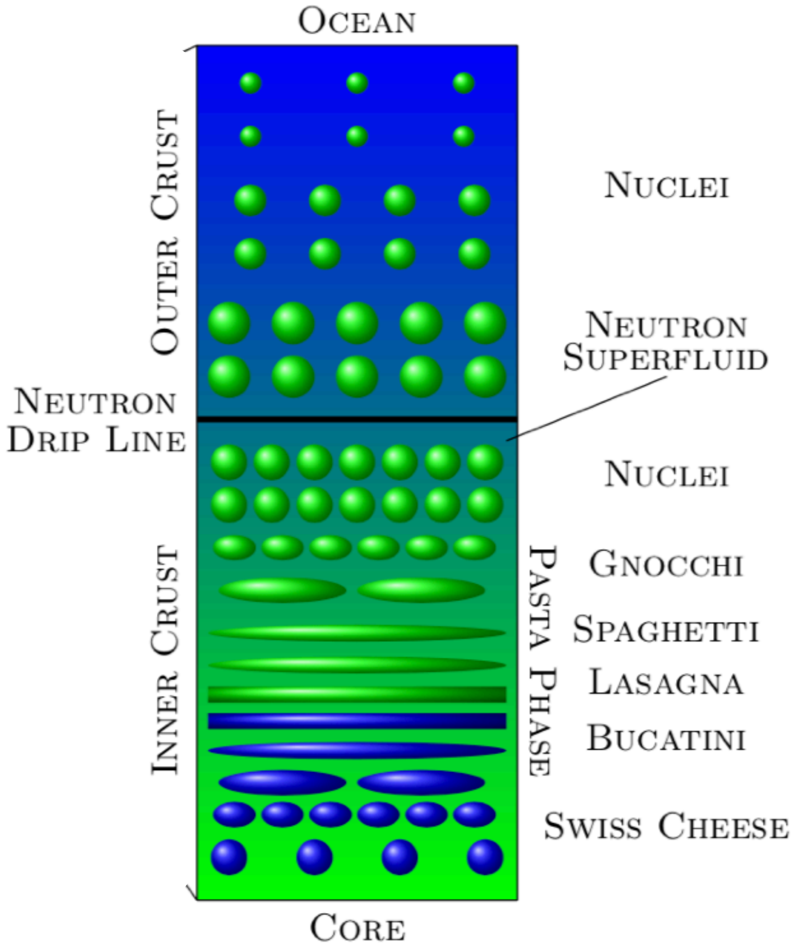
[Caplan, Horowitz, *RevModPhys 2017*]



Are we barking down the wrong region?

[Caplan, Horowitz, *RevModPhys 2017*]

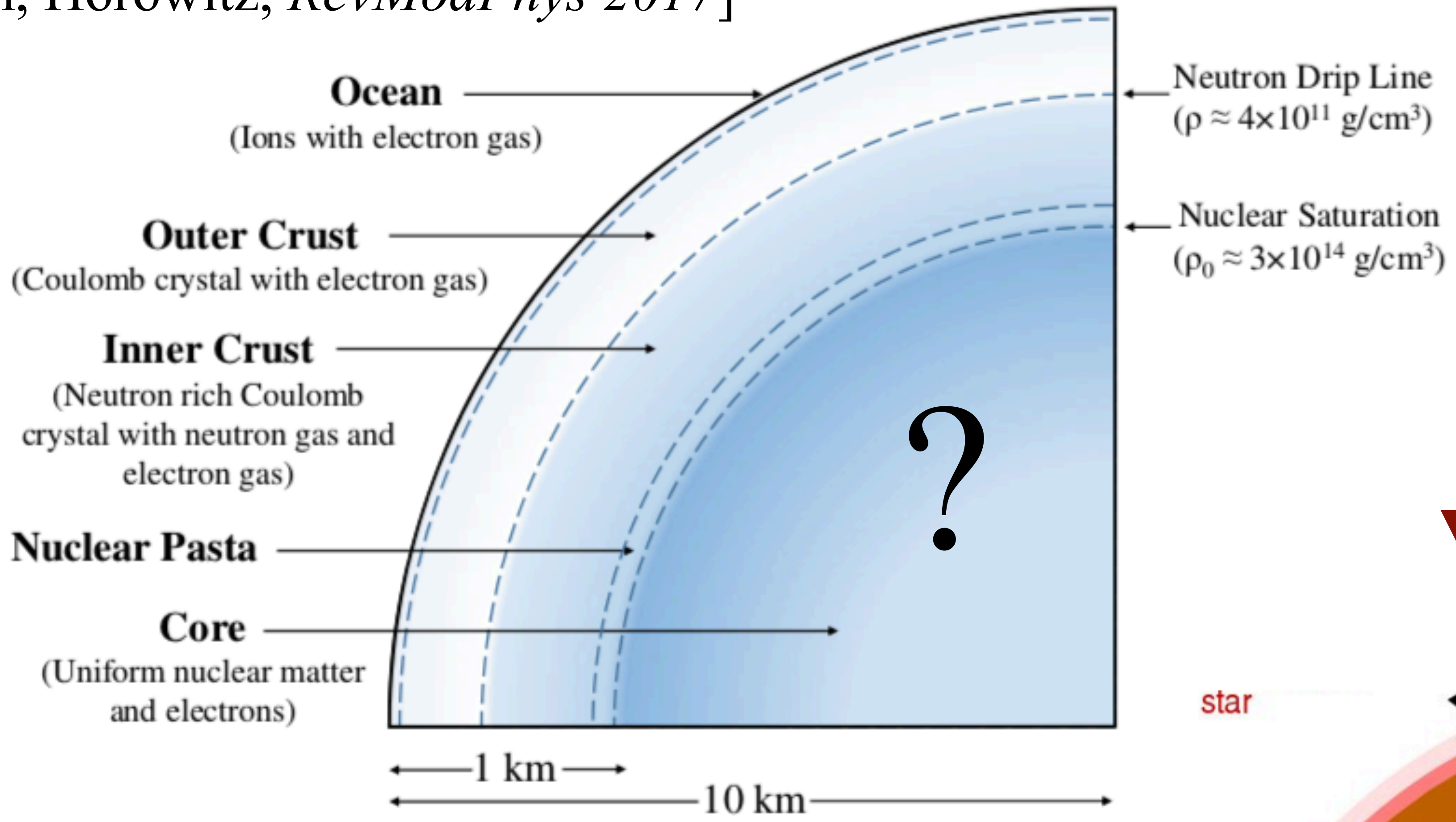




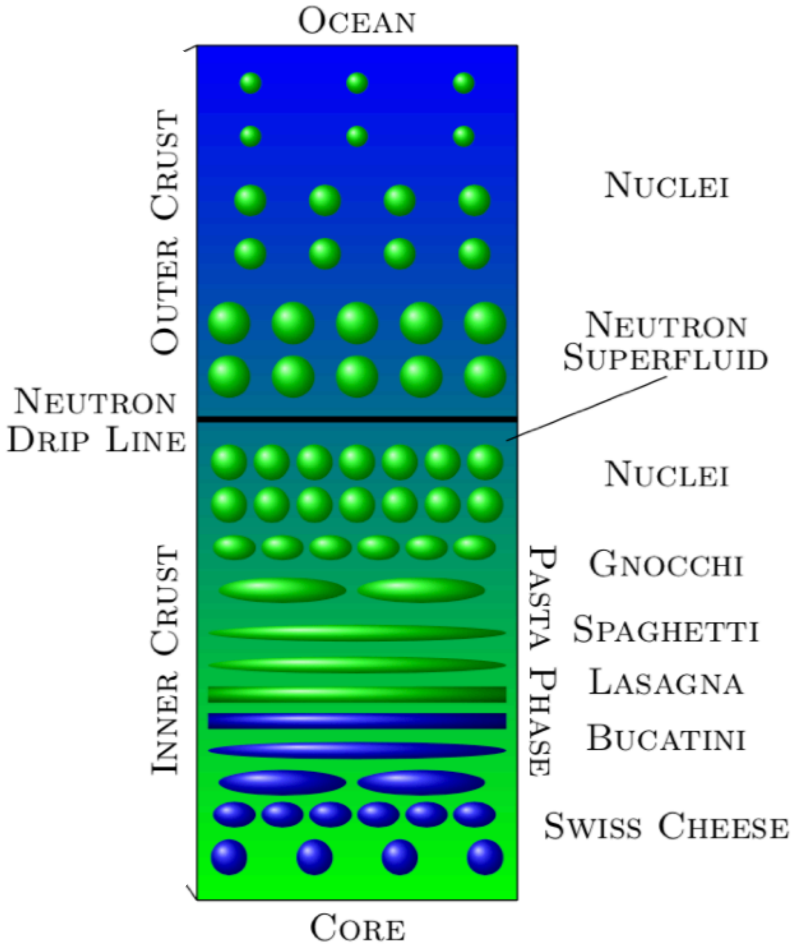
}

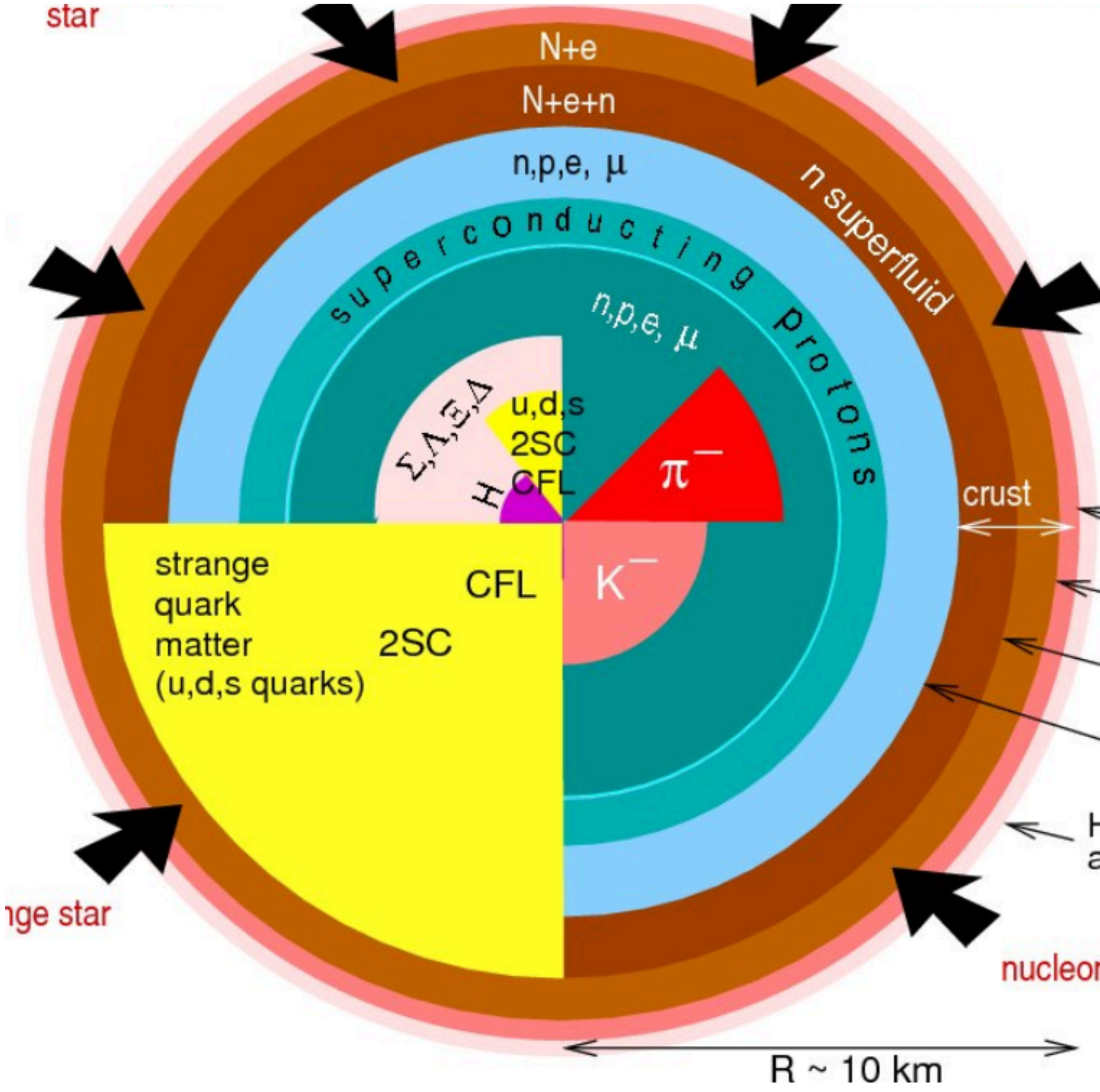
structure of the crust,
better understood than core

[Caplan, Horowitz, *RevModPhys 2017*]



deeper =>
knowledge of structure
more uncertain

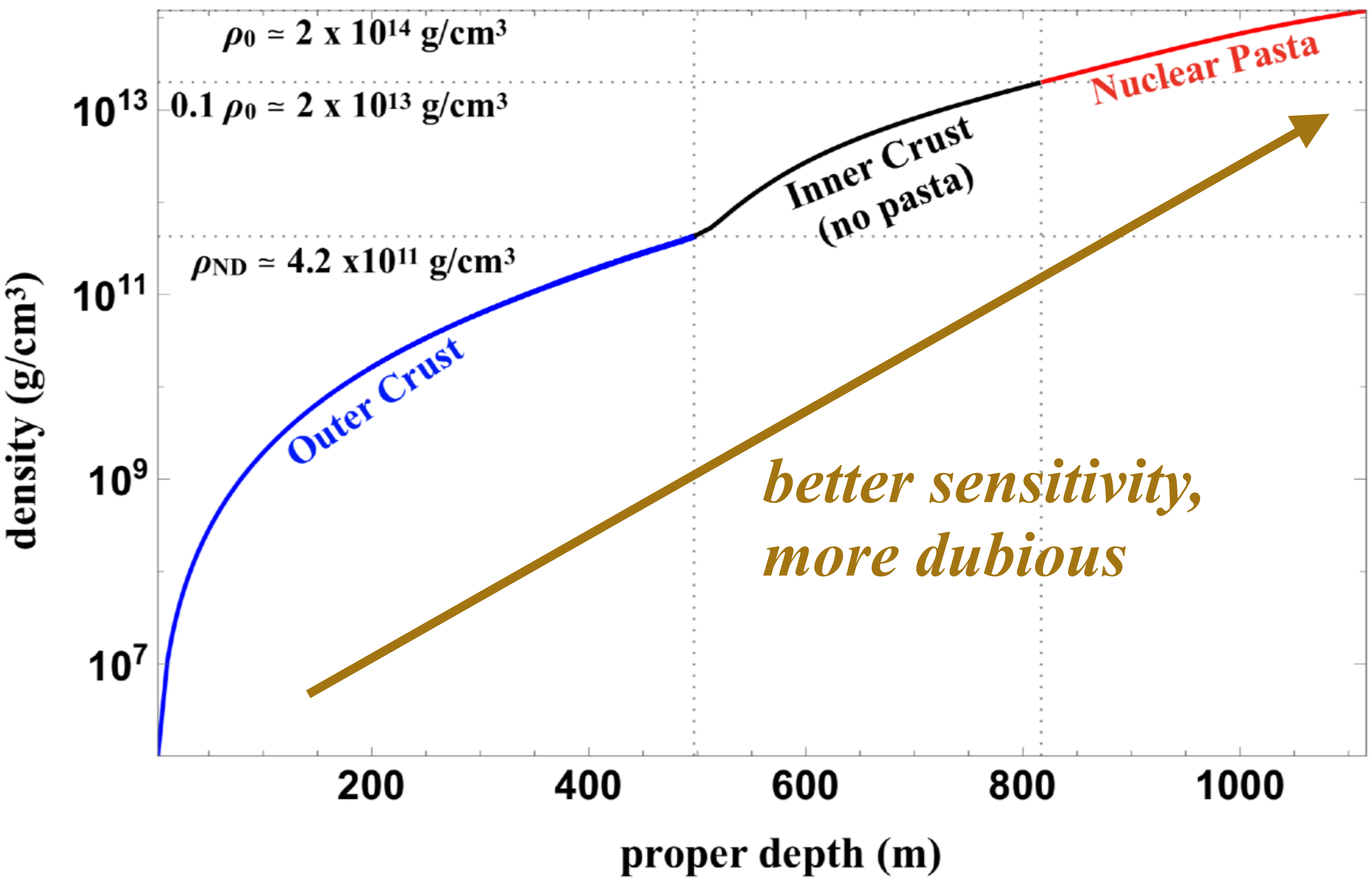




Climbing down the layers

deeper =>
knowledge of structure
more uncertain

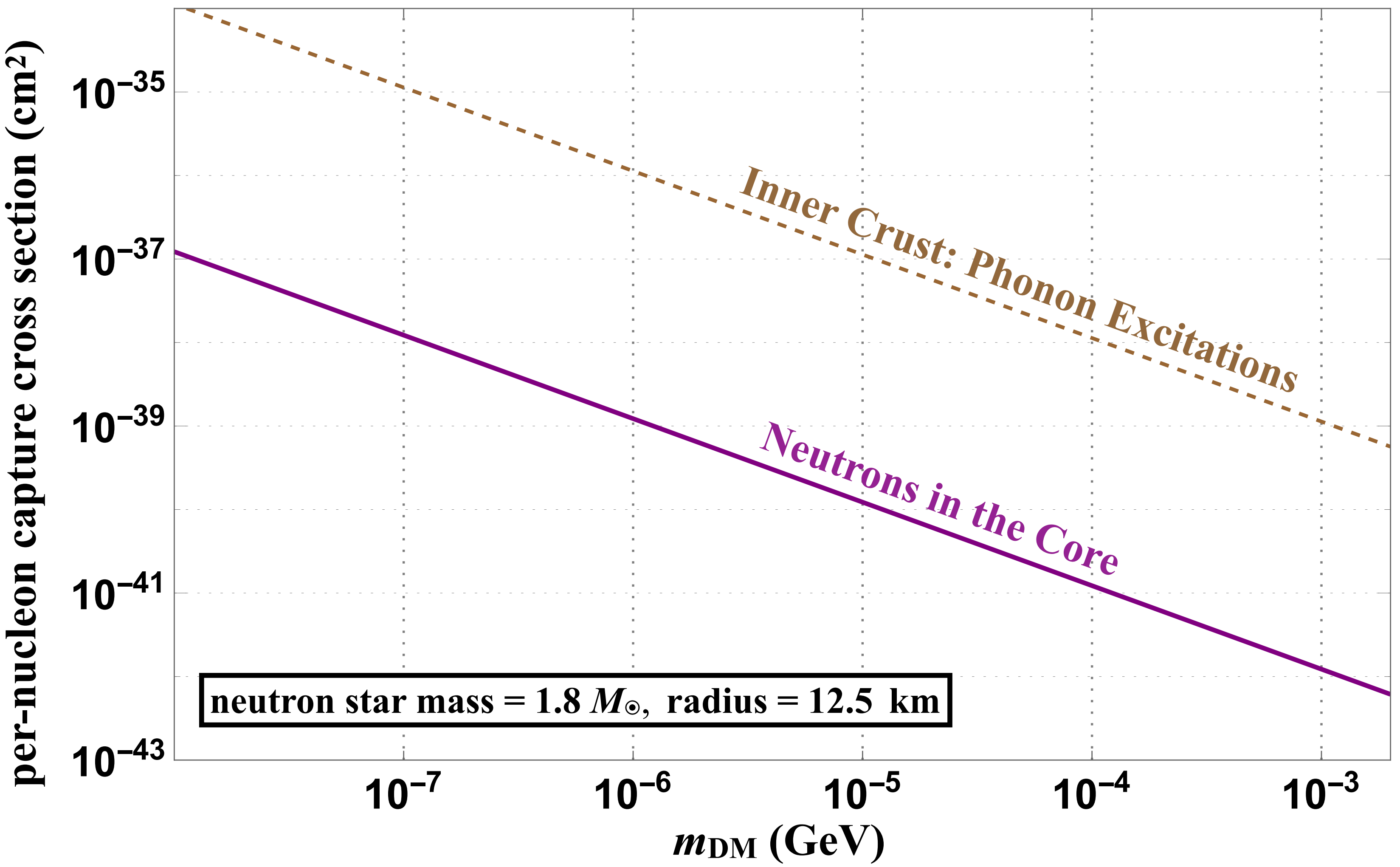
worthwhile to investigate capability of every layer to capture dark matter



better sensitivity,
more dubious

J Acevedo, J Bramante, R Leane, **N R**; 1911.06334

Crust vs low mass dark matter



capture by exciting single superfluid phonon:

energy deposited > halo KE

[q x phonon speed [m_{DM} (10^{-3} c)^2]

~ m_{DM} v_{esc} x 0.04 c]

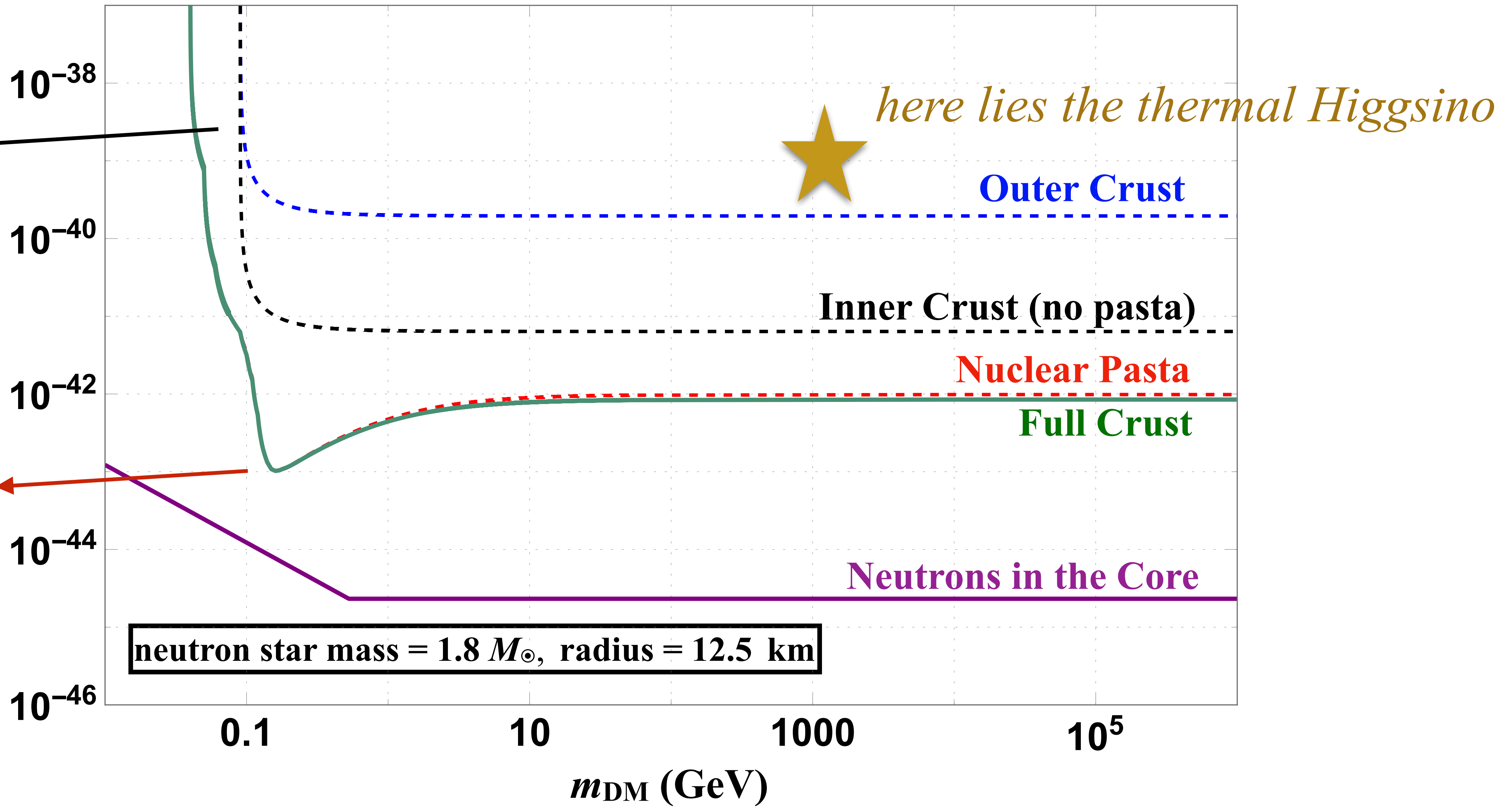
$$\sigma_{\mathrm{phonon}}(q) = \underset{\downarrow}{S_{\mathrm{phonon}}(q)}\sigma_{\mathrm{n}\chi}$$

q/ (2m_{n} x phonon speed)

J Acevedo, J Bramante, R Leane, **N R**; 1911.06334

Crust vs WIMPs & heavier dark matter

capture by (quasi-)elastic scattering on *nucleons*



energy transfer < nucleon binding energy ~ 10 MeV

response peak

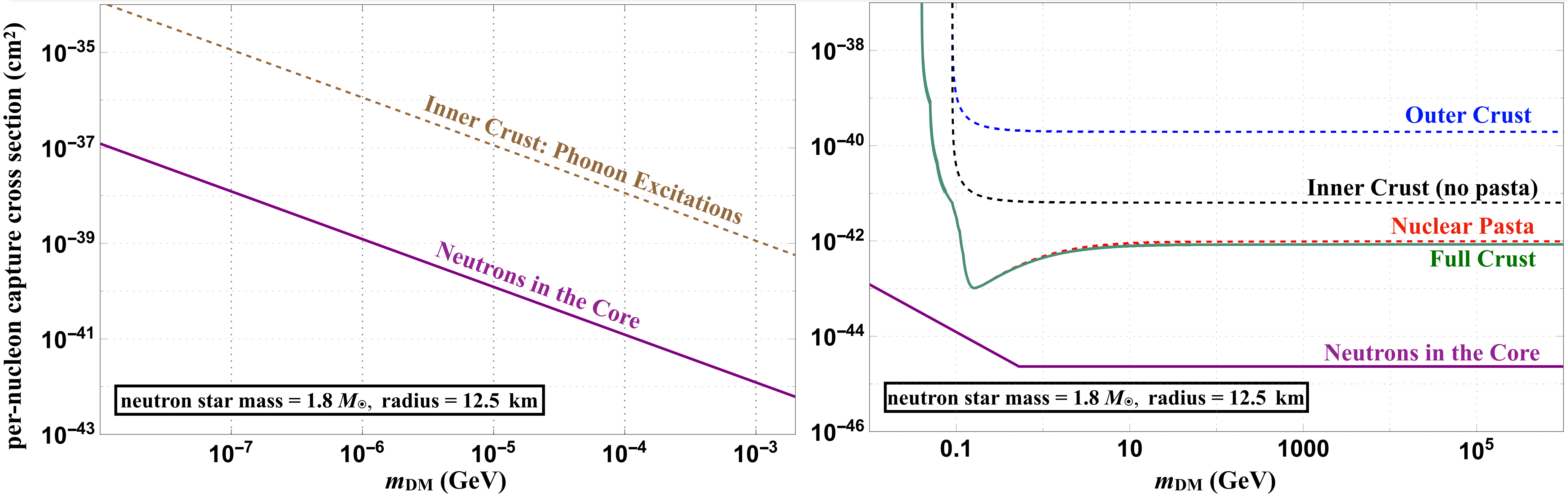
capture by pasta:

$$\sigma_{\rm pasta}(q) = \underbrace{S_{\rm pasta}(q)}\ \sigma_{\rm n\chi}$$

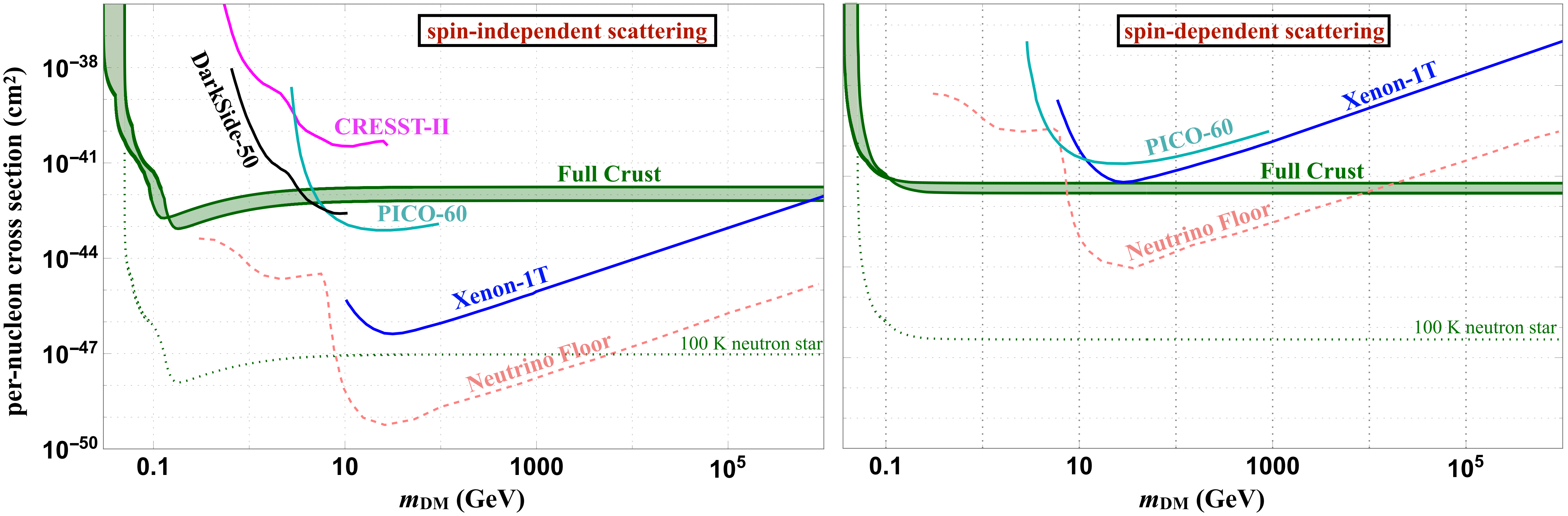
response function describing correlations among *nucleons* in pasta

J Acevedo, J Bramante, R Leane, **N R**; 1911.06334

Neutron star crust vs Earth crust



versus direct detection:



Takeaways

- Dark kinetic heating of neutron stars via scattering on *non-relativistic* **nucleonic** or *ultra-relativistic* **electronic** targets, in the *less-understood* **core** or *fail-safe* **crust**, seriously advances the direct detection frontiers of

  low mass (sub-GeV) ,
  high mass (> 100 GeV) ,
  spin-dependence ($\sigma_{SD} > 10^{-45}$ cm^2) ,
  velocity-dependence ,
  inelasticity ($<$ GeV splittings) , and
  sub-neutrino floors .

- Exoplanet observers like James Webb and Thirty Meter Telescope can unmask it with a day's worth of exposure.

Backup

The importance of being

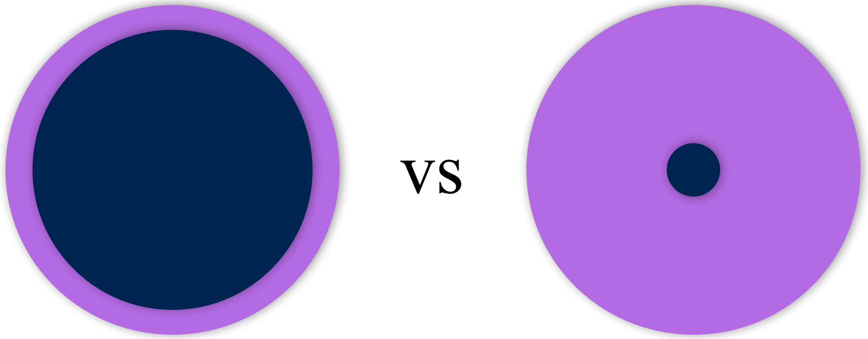
Annihilation saves observation time (= $$)
by a factor of >10!

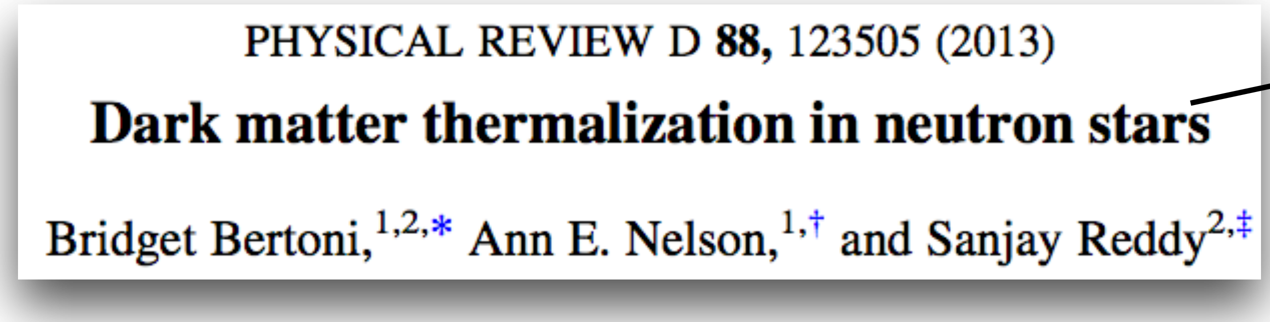
But how much annihilation is guaranteed?

Asymmetric (with Z_2-given stability) — none
p-wave — very suppressed

Does DM even thermalize with the star?

Affects DM spatial distribution,
hence annihilation rate:




PHYSICAL REVIEW D **88,** 123505 (2013)

Dark matter thermalization in neutron stars

Bridget Bertoni,[1,2,*] Ann E. Nelson,[1,†] and Sanjay Reddy[2,‡]

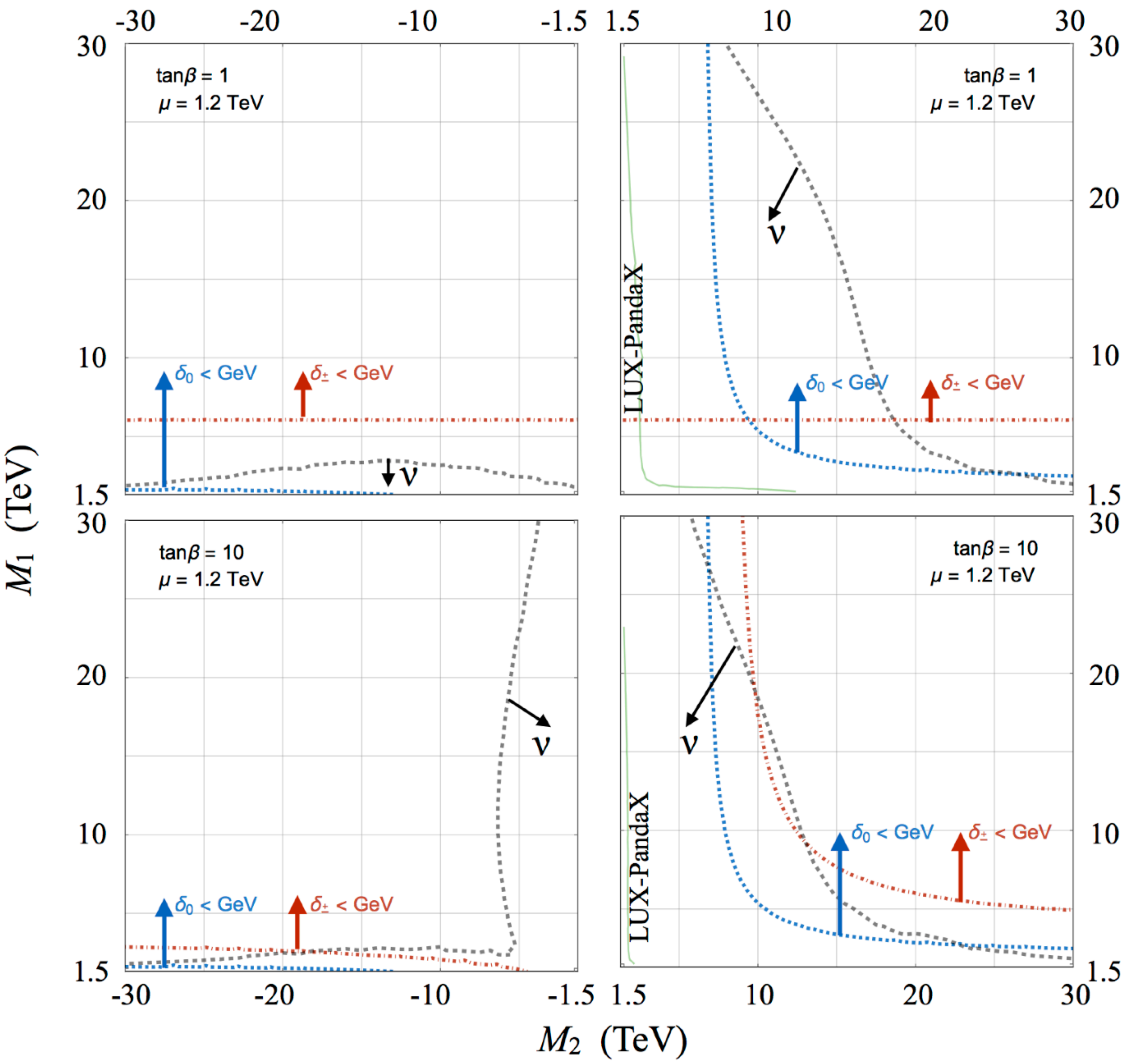
→ Spin-0 DM, vector interaction with quarks

What happens for other spins & interactions?
What if scattering is velocity-suppressed?
inelastic?

Investigation ongoing…

Complementing terrestrial searches

M Baryakhtar, J Bramante, S Li, T Linden, **N R**;1704.01577

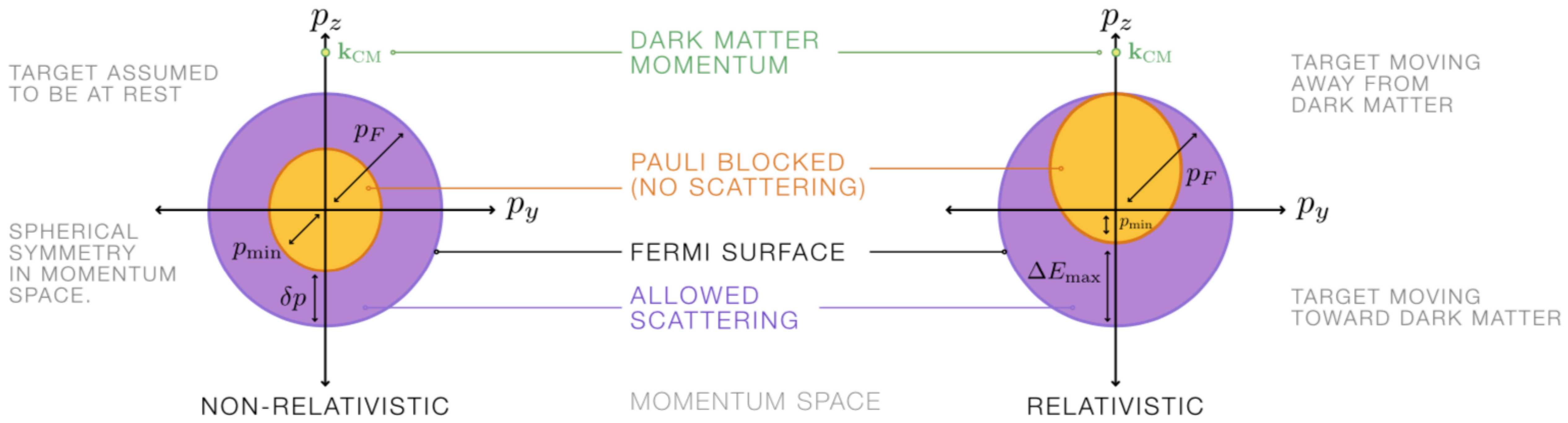


(1) Low mass ✅
(2) High mass ✅
(3) Strong ✅
(4) Spin-dependent ✅
(5) Inelastic ✅

$$\delta_0 \simeq \frac{v^2}{4}\left(\frac{g_1^2}{M_1}+\frac{g_2^2}{M_2}\right),$$

$$\delta_\pm^{\rm tree} \simeq \frac{v^2}{4}\left(\frac{g_1^2}{M_1}(1+\sin 2\beta)+\frac{g_2^2}{M_2}(1-\sin 2\beta)\right),$$

$$\delta_\pm^{\rm loop} \simeq \left(\frac{g_2}{4\pi}\right)^2 \mu \sin^2\theta_W\, f\left(\frac{m_Z}{\mu}\right),$$



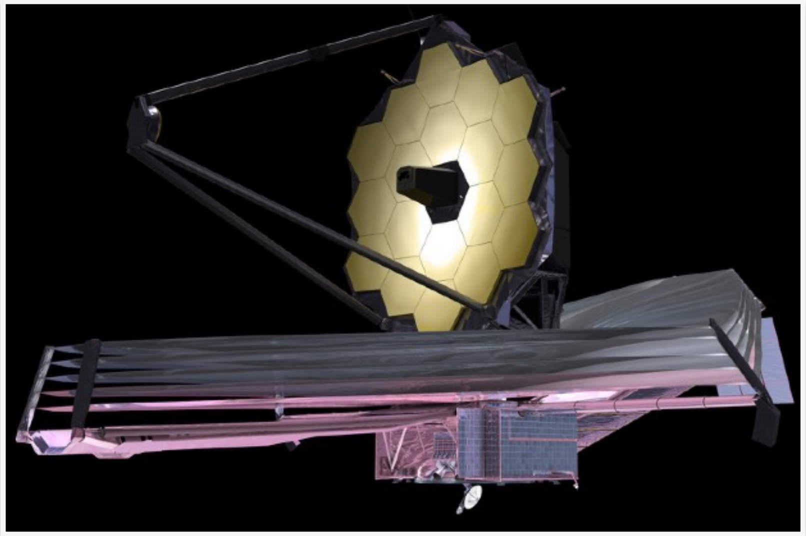
p_z
$\mathbf{k}_{\rm CM}$
DARK MATTER MOMENTUM
p_z
$\mathbf{k}_{\rm CM}$
TARGET ASSUMED TO BE AT REST
TARGET MOVING AWAY FROM DARK MATTER
p_F
PAULI BLOCKED (NO SCATTERING)
p_F
p_y
p_y
SPHERICAL SYMMETRY IN MOMENTUM SPACE.
$p_{\rm min}$
$p_{\rm min}$
FERMI SURFACE
$\Delta E_{\rm max}$
δp
ALLOWED SCATTERING
TARGET MOVING TOWARD DARK MATTER
NON-RELATIVISTIC
MOMENTUM SPACE
RELATIVISTIC

Detection: infrared telescopes

T = **1750 Kelvin** (infrared emission)

backup

Peak wavelength: $1.65\ \mu\mathrm{m}$

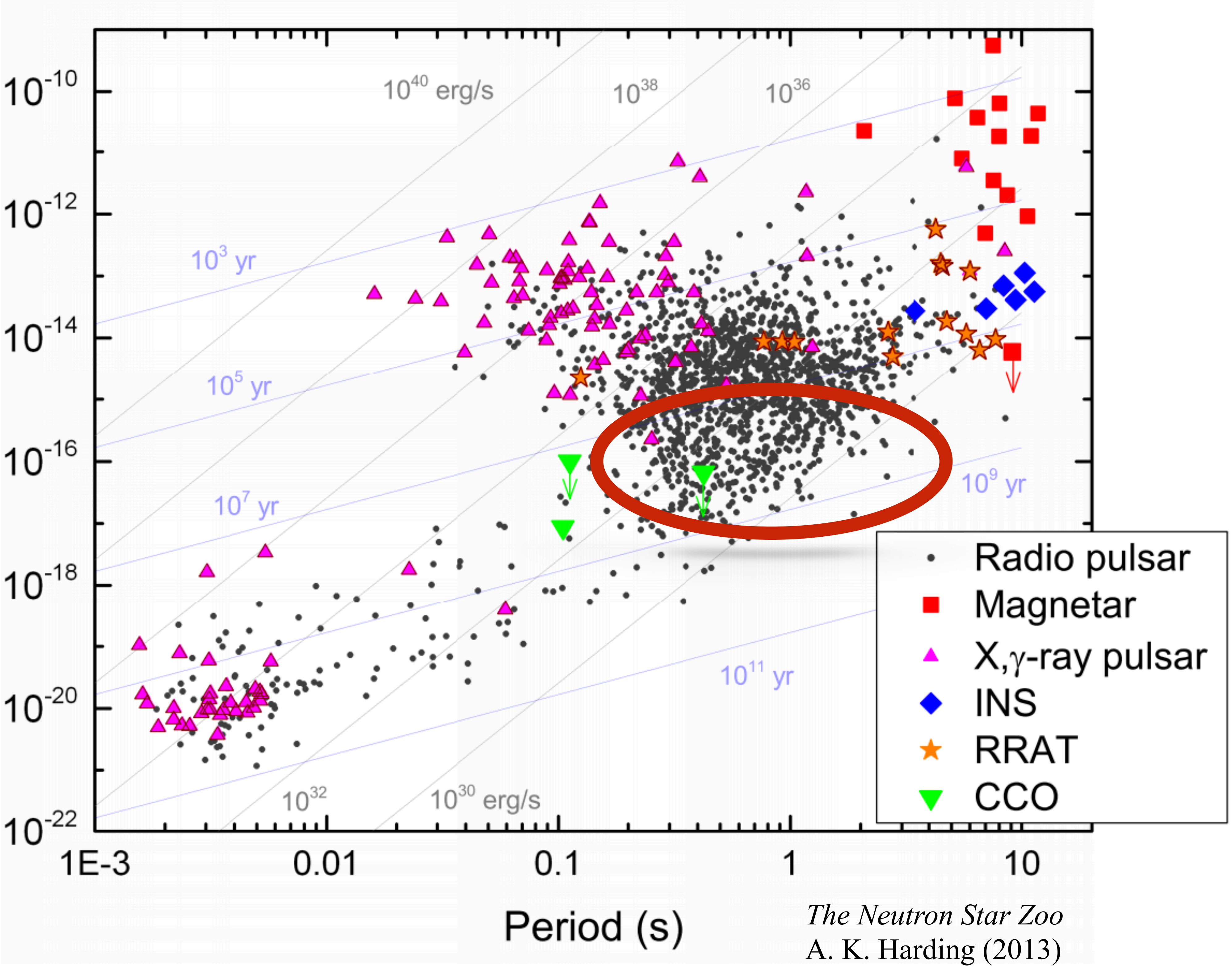


James Webb

Thirty Meter

Imager	NIRCam	IRIS
Filter	F200W	K-band
	$1.75 - 2.2\ \mu\mathrm{m}$	$2.0 - 2.4\ \mu\mathrm{m}$
Observ. time for 2σ sensitivity	$10^5\ \mathrm{sec} \left(\frac{d}{10\mathrm{pc}}\right)^4$	$7 \times 10^4\ \mathrm{sec} \left(\frac{d}{10\mathrm{pc}}\right)^4$

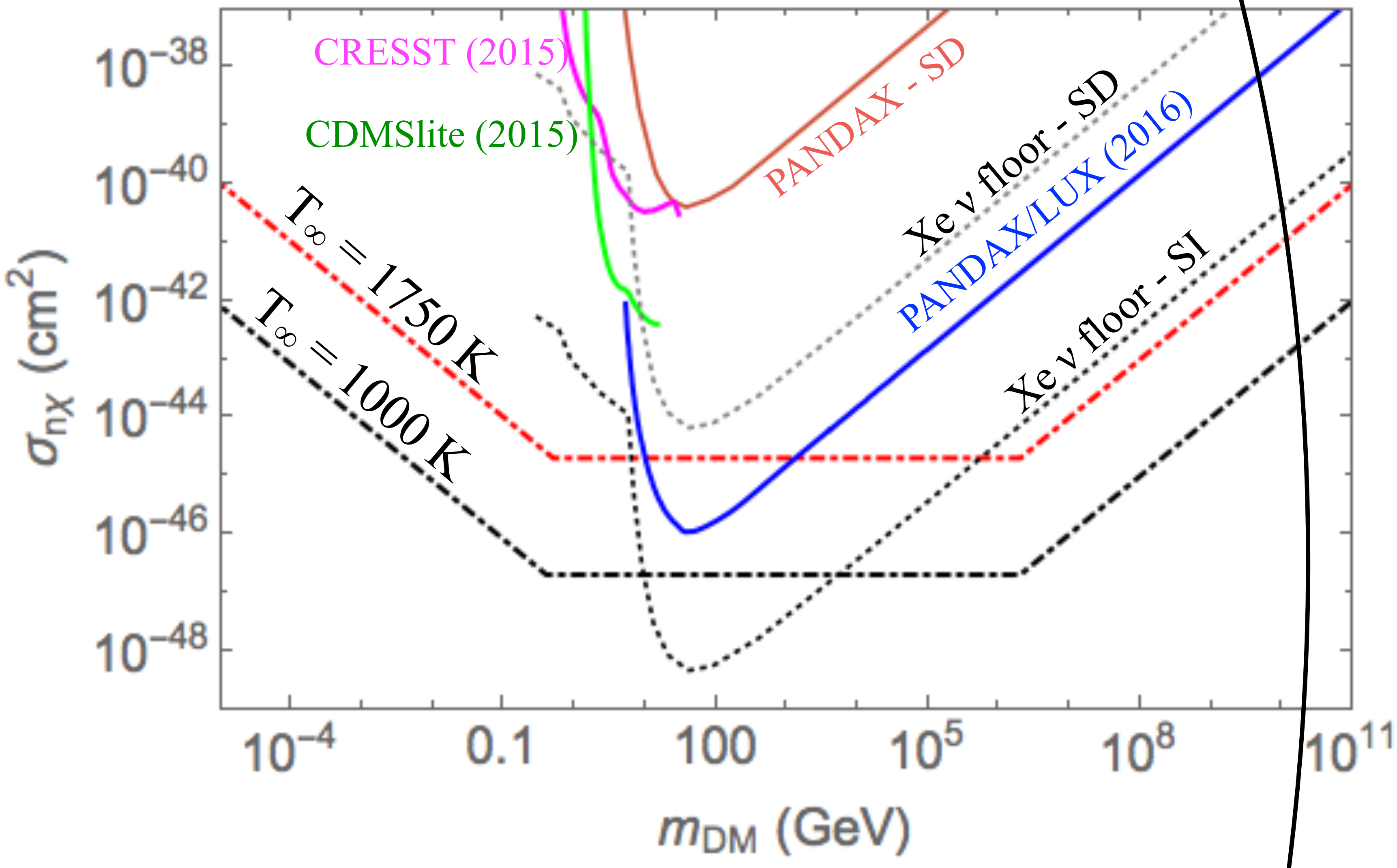
Detection: radio pulsing



The Neutron Star Zoo
A. K. Harding (2013)

Complementing

M Baryakhtar, J Bramante, S Li, T Linden, **N R**;1704.01577



(1) Low mass ✅
(2) High mass ✅
(3) Spin-dependent ✅
(4) v-suppressed ✅
(5) Inelastic ✅
(6) Neutrino floors ✅

Underground searches would terrifically complement us.