

# Quantum critical dynamics: CFT, Monte Carlo & holography

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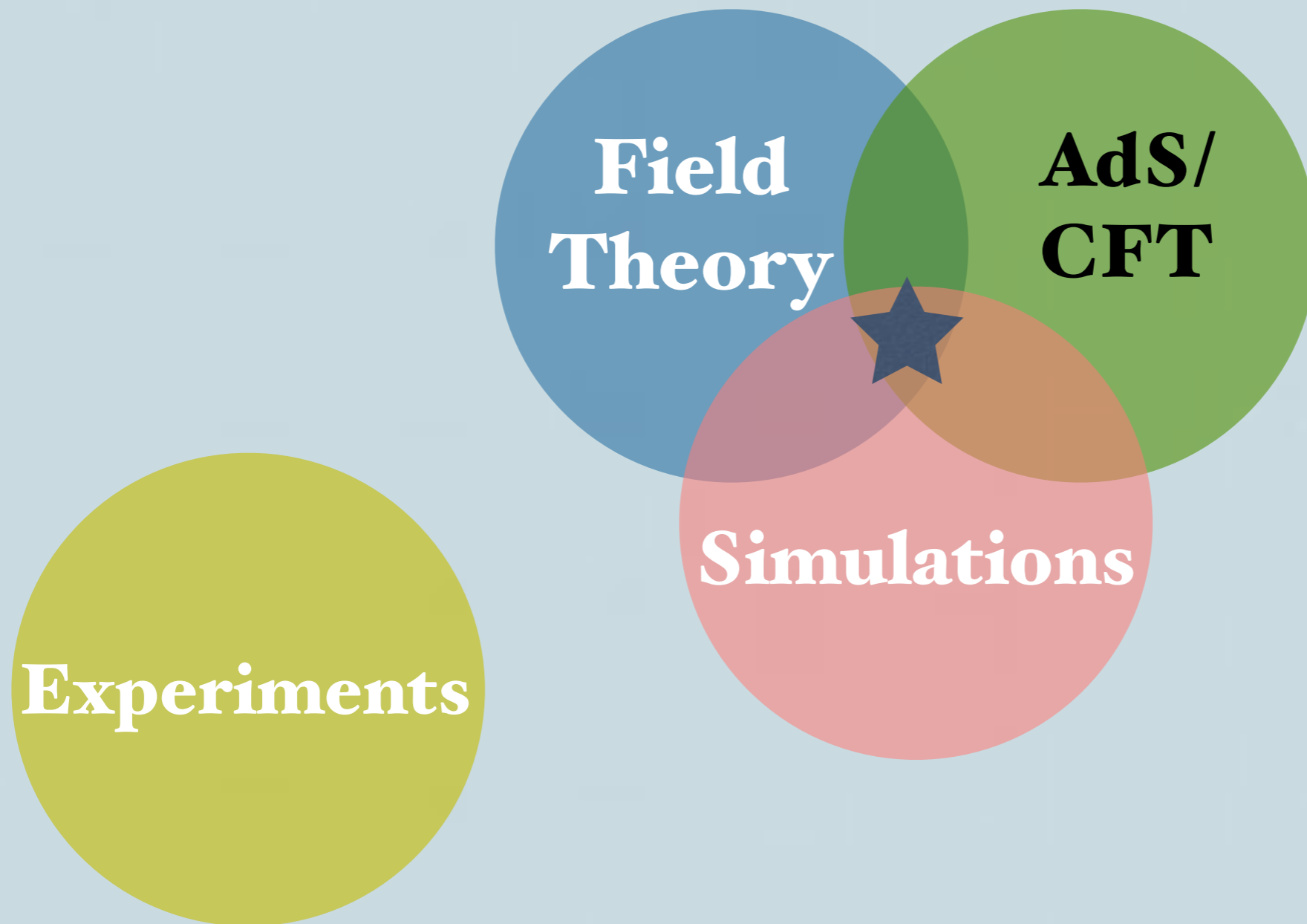
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@Boston U.

- ❖ E. Katz, S. Sachdev, E. Sørensen, WWK, PRB **90**, 2014 [Ed. Sugg.]
- ❖ WWK, E. Sørensen, S. Sachdev, Nat. Phys. **10**, 2014
- ❖ WWK, S. Sachdev, PRB **87**, 2013 [Ed. Sugg.]
- ❖ WWK, S. Sachdev, PRB **86**, 2012

# ★ Real time **finite $T$** response of strongly correlated quantum fluids

- (quasi)particles:  $\infty$  lifetime (poles)
- If have sharp (quasi)particles, get dynamics from Boltzmann Eq.
- Systems *without* quasiparticles:
  - ▶ Quantum critical phase transition
  - ▶ Strongly coupled gauge theories, etc.

# Quantum critical dynamics

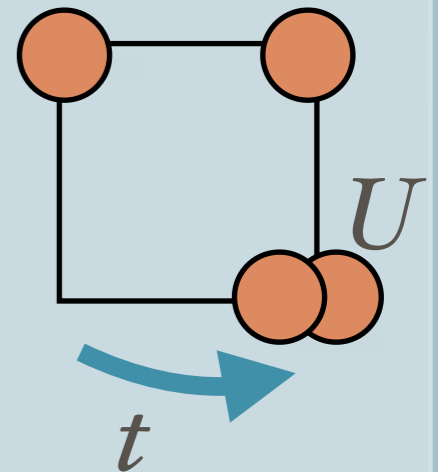
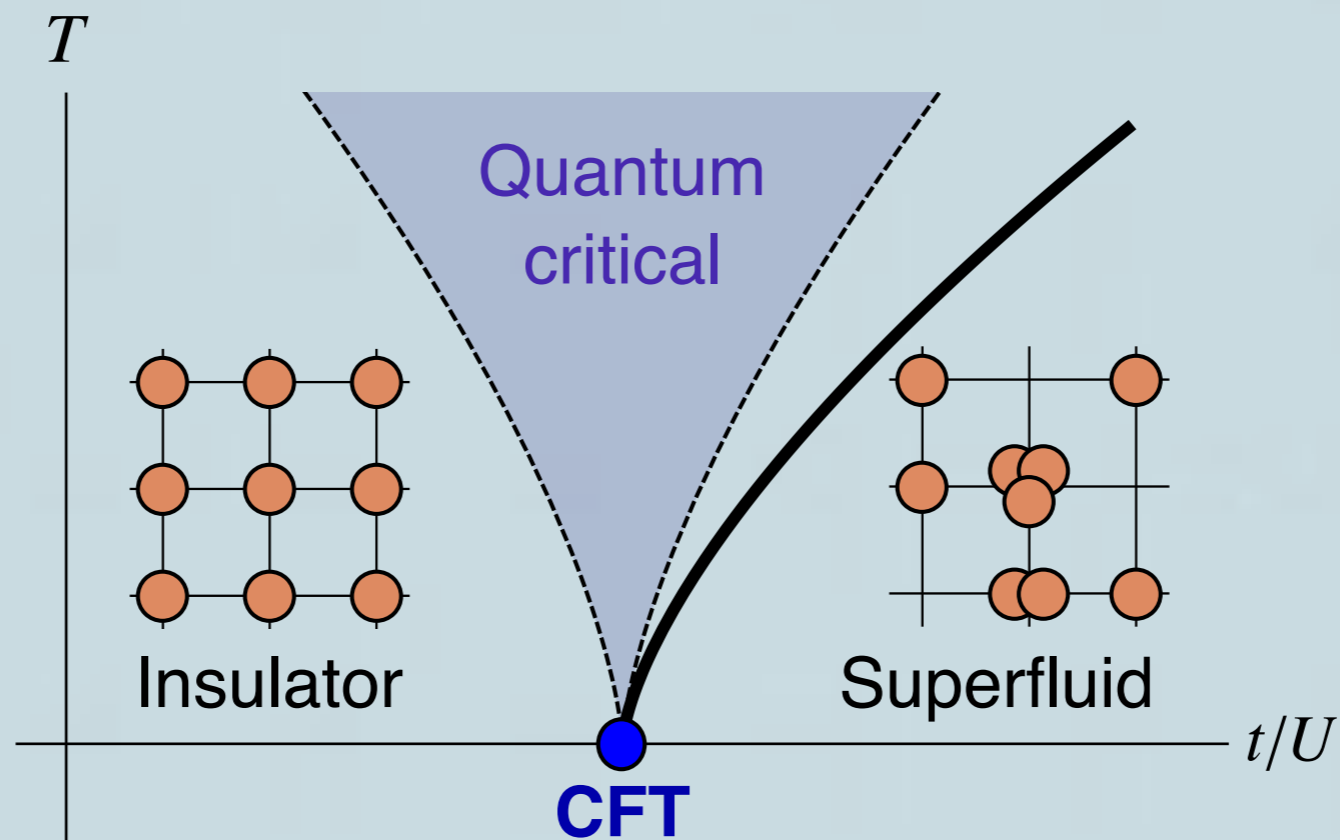


# Conformal Field Theory

- ❖ Scale + Lorentz invariant QFT
- ❖ Best characterized quantum fluid w/out (quasi)particles
  - Many open questions: dynamics, etc
- ❖ Why care?
  - ❖ Experiments!
  - ❖ Realistic models: numerics
  - ❖ Gravity = CFT? AdS/CFT [Maldacena, etc]

# Bosons in 2+1D

$$H = -t \sum_{\langle i,j \rangle} b_i^\dagger b_j + U \sum_i n_i(n_i - 1) - \mu \sum_i n_i$$



- ❖  $O(2)$  universality class: Bose-Hubbard, quantum rotors, XY spins, etc

❖ SSB of O(2) order parameter

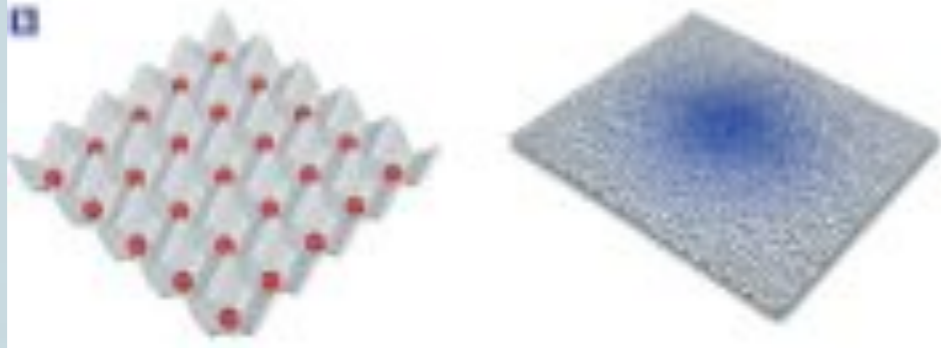
$$\mathcal{L} = (\partial_t \vec{\phi})^2 + (\nabla \vec{\phi})^2 + m^2 \phi^2 + u \phi^4$$

Strongly coupled fixed point  
in 2d (Wilson-Fisher)

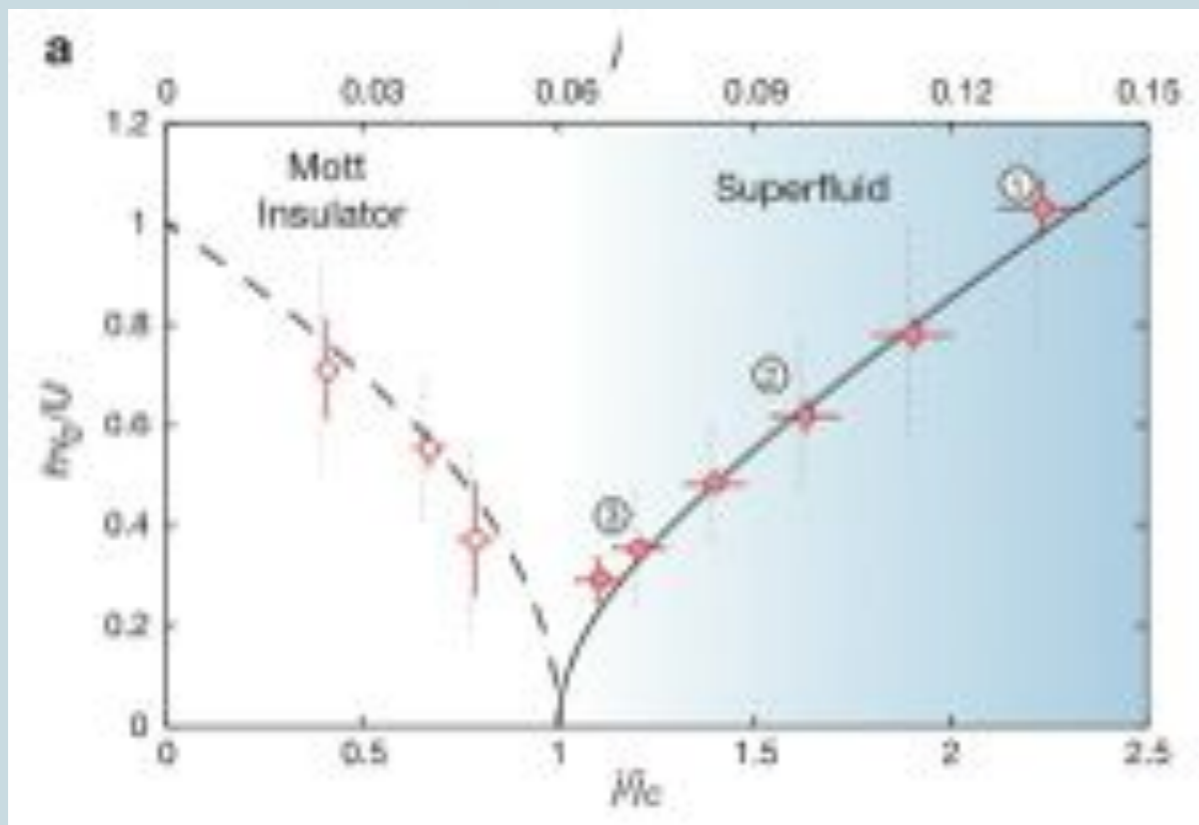


Superfluid

[Bloch]



Insulator



$^{87}\text{Rb}$

[Endres *et al*]



# *Break* conformal symmetry

- ❖ Characteristic time scale:  $1 / T$
- ❖ **Short** times  $\omega \gg T$ : probing near vacuum
- ❖ **Long** times  $\omega \ll T$ : excitations interact strongly with thermal background

# Operator Product Expansion

- ❖ A scalar primary op  $O(x)$ , w/ scaling dim  $\Delta$ :

$$\langle O(x) O(0) \rangle = 1 / x^{2\Delta}$$

- ❖ OPE:

$$\mathcal{O}(x) \mathcal{O}(0) \rightarrow \sum_n \frac{\mathcal{O}_n(0) + \text{desc.}}{x^{2\Delta - \Delta_n}}$$

[Wilson; Polyakov; Ferrara *et al*; etc]

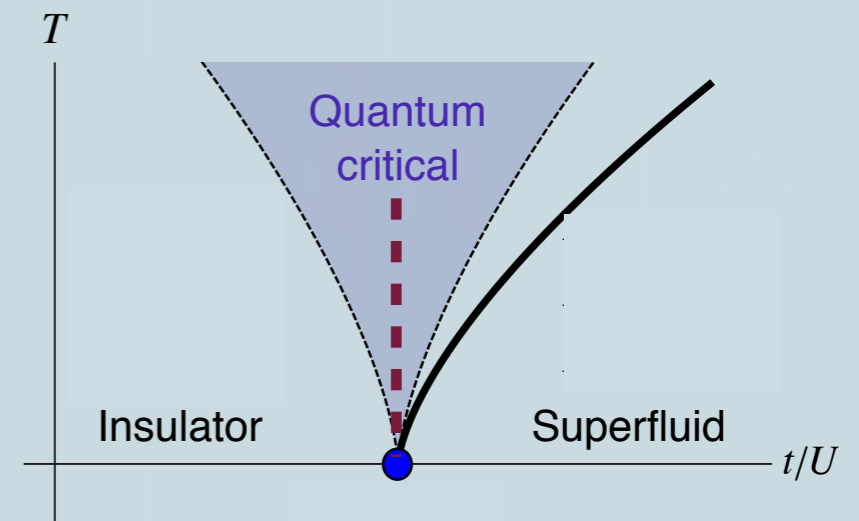
# Current correlators

$$\partial_\mu J^\mu = 0$$

$$\sigma(\omega) = \frac{1}{i\omega} \langle J_x(\omega) J_x(-\omega) \rangle$$


**Universal** scaling function:


$$\sigma(\omega) = \Phi\left(\frac{\omega}{T}\right)$$



# Current OPE

$$J_\mu(x) J_\nu(0) = \frac{I_{\mu\nu} \mathbf{1}}{x^{2\cdot 2}} + \mathcal{C}_{JJ\mathcal{O}} \frac{x_\mu x_\nu \mathcal{O}(0)}{x^{6-\Delta}} + \mathcal{C}_{JJT} \frac{T_{\mu\nu}(0)}{|x|} + \dots$$

  
O(N) CFT:  
**Relevant**  
scalar  $\mathcal{O} \sim \phi^2$

  
Stress tensor

- ❖ Get OPE coefficients from  $\langle JJ\mathcal{O} \rangle$

# Conductivity via OPE

$$\lim_{|q| \gg |p|} J_x(q) J_x(-q + p) = -|q| \sigma_\infty \delta(p) - \mathcal{C}_{JJO} \frac{\mathcal{O}_g(p)}{|q|^{\Delta-1}} \\ + \frac{\mathcal{C}_{JJT}}{|q|^2} [T_{xx} - T_{yy} - 12\gamma(T_{xx} + T_{yy})] \Big|_p + \dots$$

Thermal average



$$\langle \mathcal{O} \rangle_T = BT^\Delta$$

$$\sigma(i\omega_n) \stackrel{\omega_n \gg T}{=} \sigma_\infty + b_1 \left( \frac{T}{\omega_n} \right)^\Delta + b_2 \left( \frac{T}{\omega_n} \right)^3 + \dots$$

# Dominant op in OPE?

- ❖ For  $O(N)$  Wilson-Fisher, it's **SCALAR**:

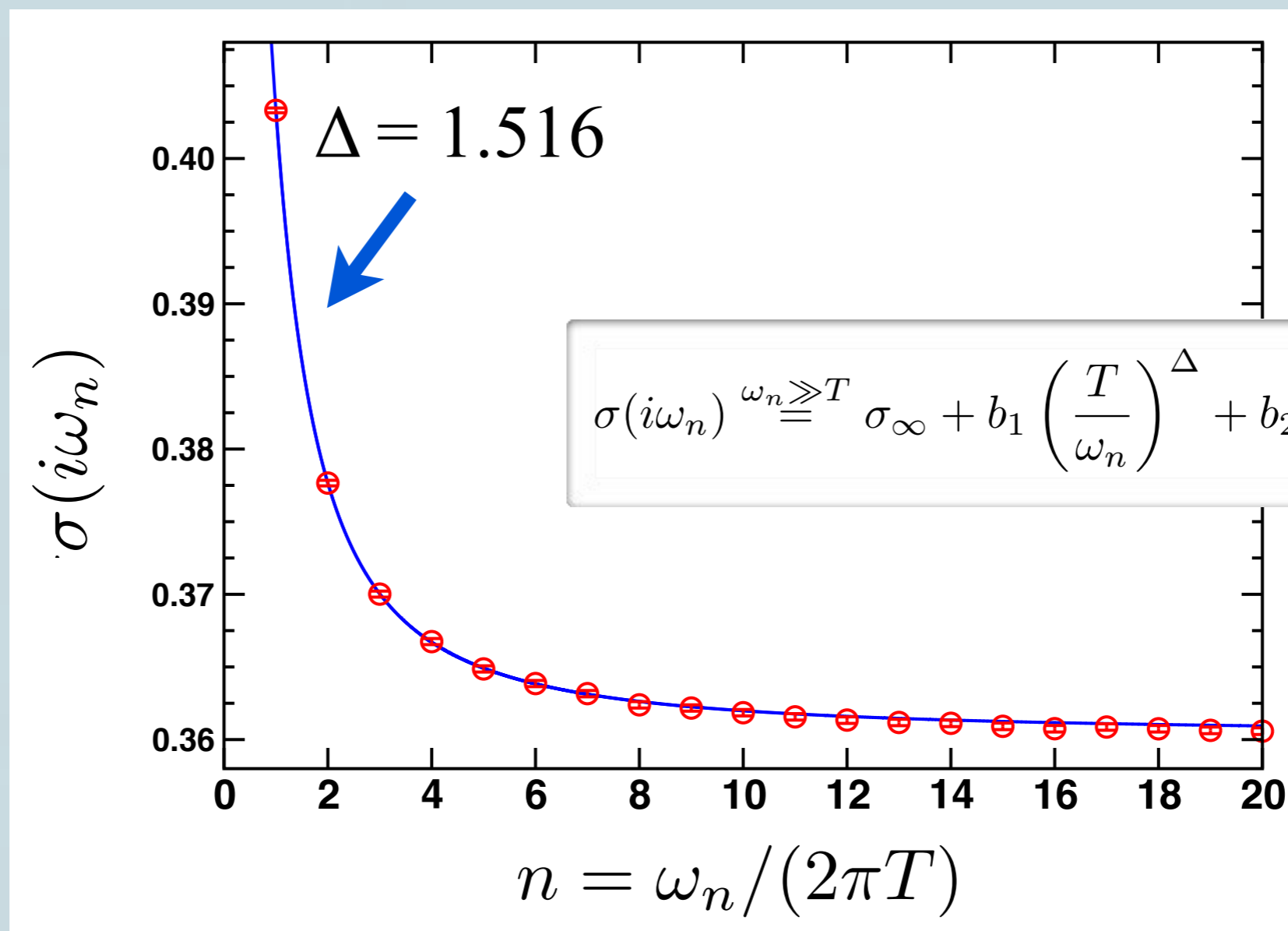
$$O(x) = \varphi^2(x) \quad \text{since} \quad \Delta = 3 - 1/\nu$$

- ❖  $N = \infty$ :  $\Delta = 2$   
 $N = 2$ :  $\Delta = 1.5106$
- ❖ **All** CFTs have  $1/\omega^3$  term because *there's always a stress tensor*
  - ❖ Sometimes it dominates (Dirac CFT, QED3,  $\mathcal{N}=4$  super-Y-M, etc)

# Quantum Monte Carlo

[WWK, Sorensen, Sachdev, Katz]

- ❖ Simulate lattice model for O(2) QCP → **imaginary** frequencies





*Long times &  
analytic continuation*

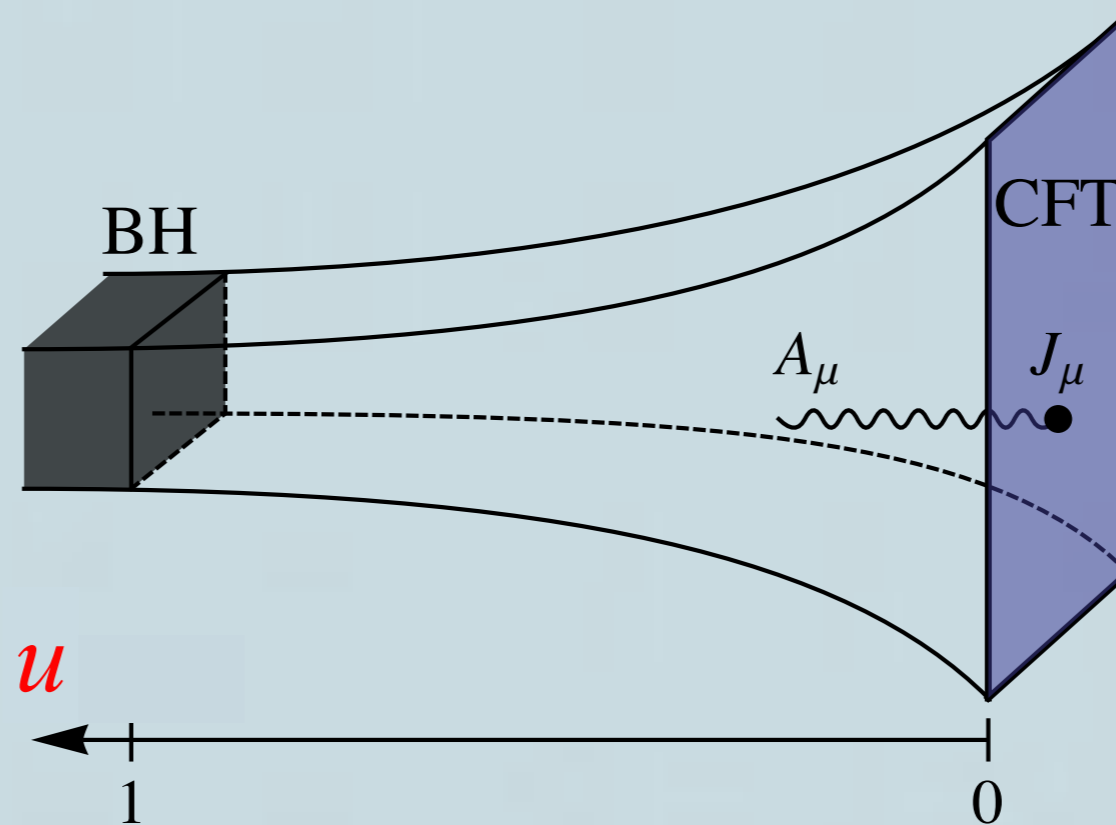


★ Use **AdS/CFT** to generate a family of physically motivated (constrained) variational functions  $\sigma(\omega/T)$

# $\sigma$ via AdS/CFT

[Maldacena *etc*]

spacetime:  
**B-Hole in  
AdS<sub>4</sub>**



$$A_\mu(t, x, y; u) \leftrightarrow J_\mu^{\text{CFT}}(t, x, y)$$

❖ Solve **classical** EoM for  $A$   $\rightarrow$  get  $J$ -correlator in **CFT**

$$\sigma(\omega/T) = \frac{T}{\omega} \left. \frac{\partial_u A_y(\omega, \vec{0}; u)}{A_y(\omega, \vec{0}; u)} \right|_{u=0}$$

# 1st try

$$S_{\text{bulk}}[A_\mu] = \int d^4x \sqrt{-g} \frac{1}{g_4^2} [F^2 + \gamma C^{abcd} F_{ab} F_{cd}]$$

$C$  = traceless part of Riemann

❖ Derivative expansion

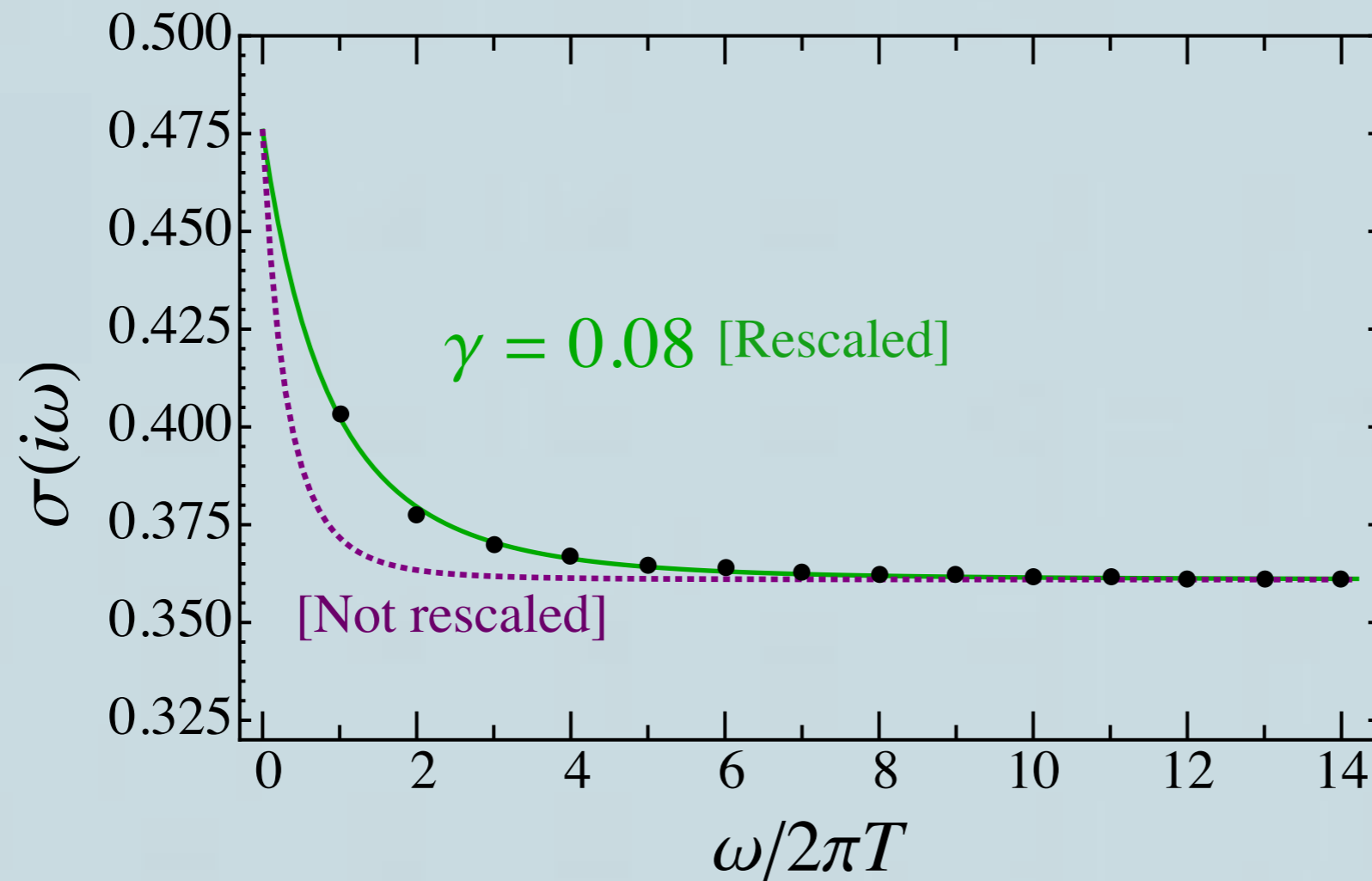
❖ Asymptotics [WWK]:  $\sigma \sim \sigma_\infty + (T/\omega_n)^3 + \dots$

<b>CFT</b>	<b>AdS</b>
$T_{\mu\nu}$	$g_{\mu\nu}$

[Herzog, Kovtun, Sachdev, Son; Ritz, Ward; Myers, Sachdev, Singh]

# Holographic fit: take 1

[WWK, Sorensen, Sachdev]



$$\sigma\left(\kappa \frac{\omega}{T}; \gamma\right)$$

# 2nd try

$$S_{\text{bulk}}[A_\mu] = \int d^4x \frac{1}{g_4^2} [1 + \alpha \varphi(x)] F_{ab} F^{ab}$$

- ❖ Add **scalar field** (dilaton)

<b>CFT</b>	<b>AdS</b>
$\mathcal{O}$	$\varphi$

- ❖ Simplest **Ansatz**:

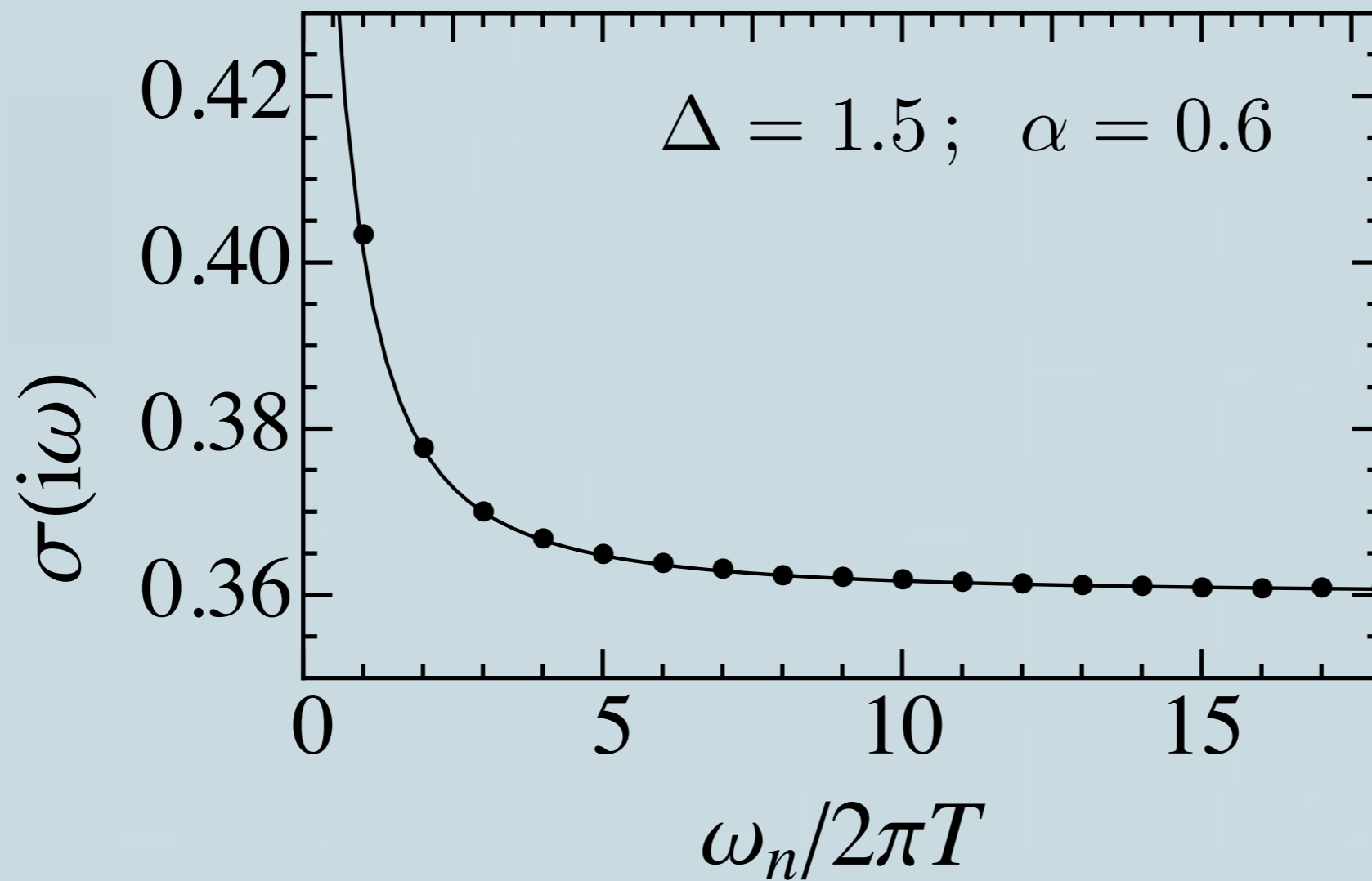
fix profile using OPE of  $O(2)$  Wilson-Fisher CFT

$$\varphi(u) = u^\Delta$$

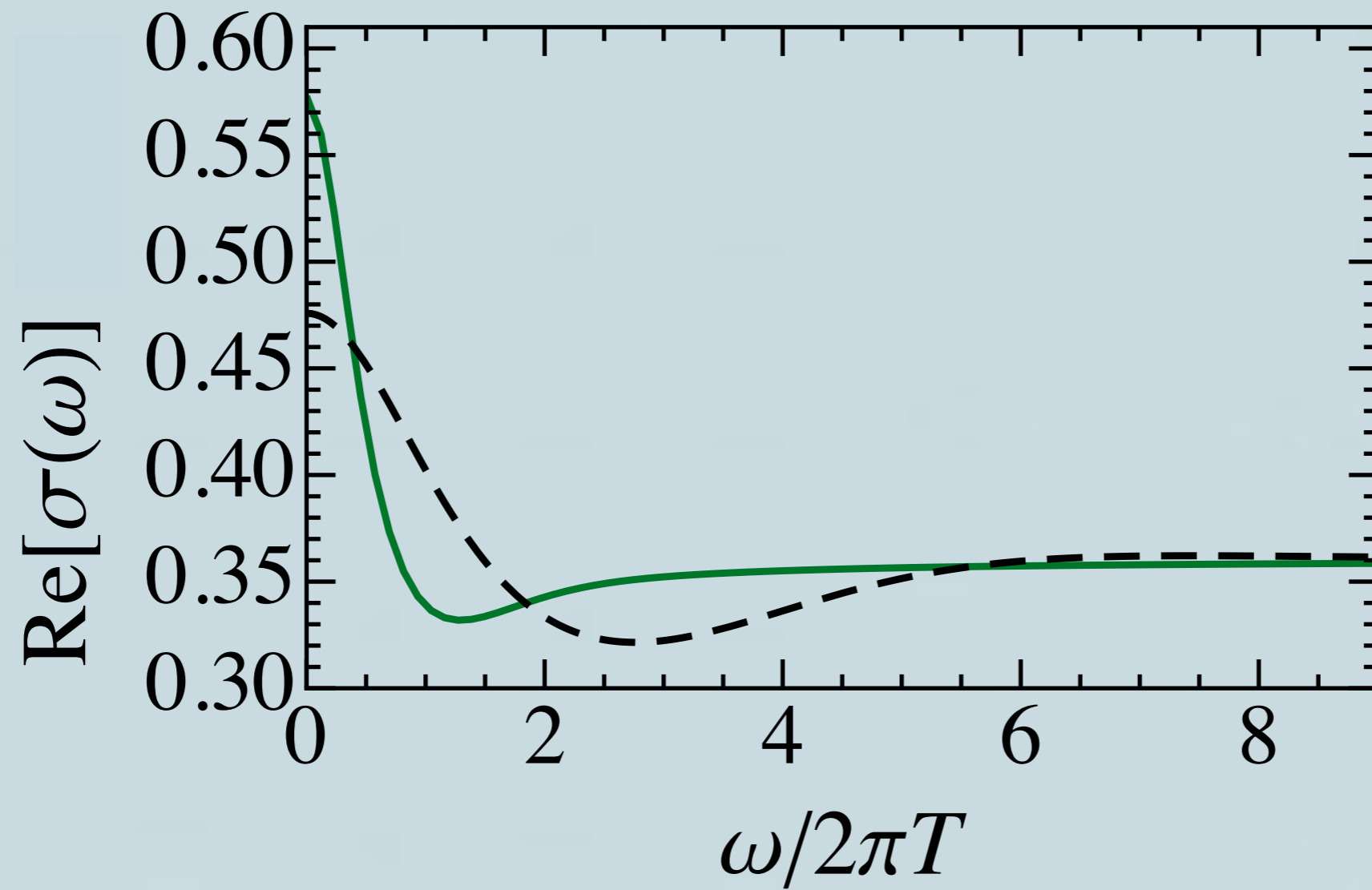
[Katz, Sachdev, Sorensen, WWK]

# Holographic fit: take 2

[Katz, Sachdev, Sorensen, WWK]



# Real frequencies



# Why dabble with black holes?

## ❖ **Physical** properties:

◆ ***Tailored*** holographic description to match asymptotics

◆ **Sum rules** 
$$\int_0^\infty d\omega \operatorname{Re}[\sigma(\omega/T) - \sigma(\infty)] = 0$$

◆ **Analytic properties** in common with field theory results

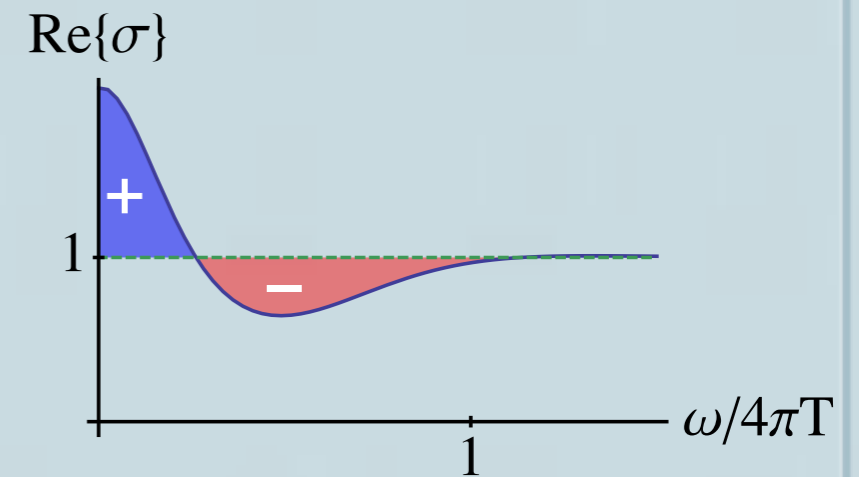
◆ **Practical:** real-time, small number of params, fast numerics



# Sum rules

[WWK, Sachdev; Gulotta, Herzog, Kaminski]

$$\int_0^\infty d\omega [\operatorname{Re} \sigma(\omega/T) - \sigma(\infty)] = 0$$



❖ S-dual version:

$$\int_0^\infty d\omega \left[ \operatorname{Re} \left\{ \frac{1}{\sigma(\omega/T)} \right\} - \frac{1}{\sigma(\infty)} \right] = 0$$

- ✓  $\mathcal{N}=8$  supersymmetric ABJM
- ✓  $O(N)$  CFT @  $N=\infty$
- ✓ Dirac CFT

# Sum rule proof via OPE

$$\sigma(i\omega_n) \stackrel{\omega_n \gg T}{\approx} \sigma_\infty + b_1 \left(\frac{T}{\omega_n}\right)^\Delta + b_2 \left(\frac{T}{\omega_n}\right)^3 + \dots$$

- ❖  $\sigma$  analytic in upper half-plane
- ❖  $\Delta > 1 \Rightarrow$  Integrable  $\Rightarrow$  Close contour in UHP ✓

# Conclusions

- ❖ Quantum critical dynamics (CFTs) in 2+1D
- ❖ OPE to constrain *short time* physics

- ❖ Large  $\omega$  conductivity

$$\sigma(i\omega_n) \stackrel{\omega_n \gg T}{\approx} \sigma_\infty + b_1 \left(\frac{T}{\omega_n}\right)^\Delta + b_2 \left(\frac{T}{\omega_n}\right)^3 + \dots$$

- ❖ Input OPE data of CFT into simple holographic ansatz
- ❖ Can match Monte Carlo data of O(2) CFT w/out unphysical tweaks

# Outlook

- ❖ Apply this program (OPE, sum rules) to other correlators & CFTs [WWK, arXiv:1501.xxxxx]
- ❖ Go beyond simplest holographic Ansatz [in progress with T. Sierens & R. Myers]
- ❖ Unitarity in proof of sum rules?

# Thanks!

- ❖ E. Katz, S. Sachdev, E. Sørensen, W.Witczak-Krempa, PRB **90**, 2014 [Ed. Sugg.]
- ❖ W.Witczak-Krempa, E. Sørensen, S. Sachdev, Nat. Phys. **10**, 2014
- ❖ W.Witczak-Krempa, S. Sachdev, PRB **87**, 2013 [Ed. Sugg.]
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