

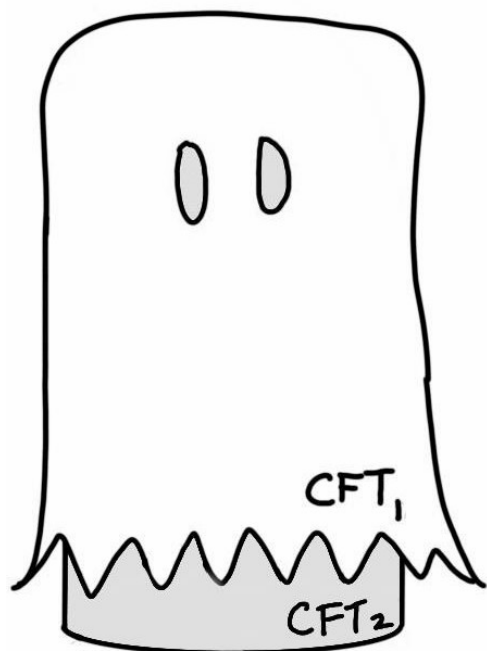
# HOLO-WEEN

Mark Van Raamsdonk, U.B.C.

Toronto HEP Theory Seminar

based on work w. Petar Simidzija

2006.13943



## MOTIVATION:

Deep connection between entanglement structure of holographic states & dual spacetime geometries.

How much info about dual spacetime is captured by "universal" properties of state?

matter fields?

internal space geometry?

## MOTIVATION:

Deep connection between entanglement structure of holographic states & dual spacetime geometries.

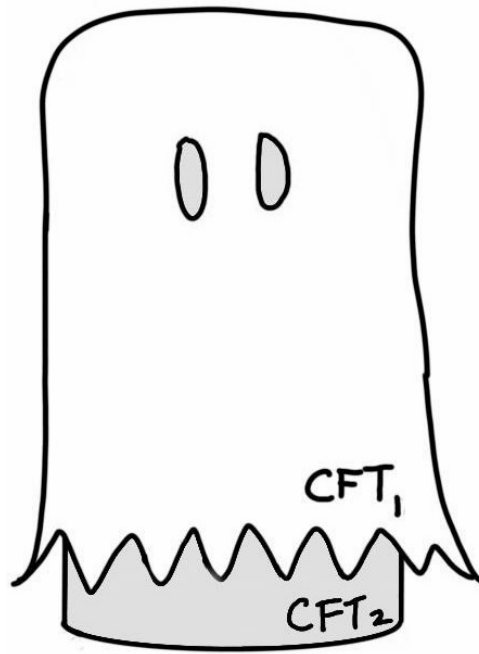
How much info about dual spacetime is captured by "universal" properties of state?

matter fields?

internal space geometry?

extreme possibility: precise microscopic d.o.f. & Hamiltonian unimportant

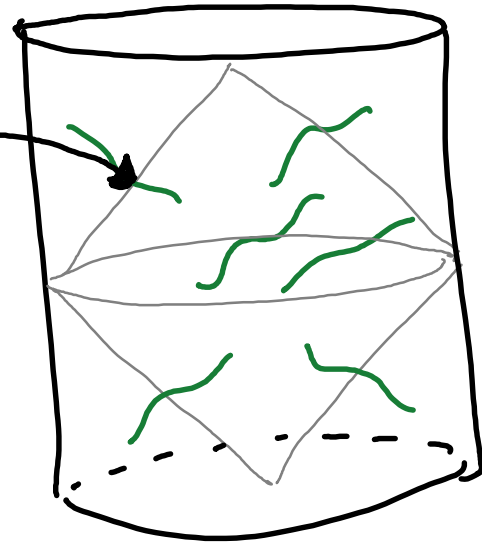
- Can a suitably chosen state of one holographic CFT faithfully encode spacetime dual to state of a completely different holographic CFT?



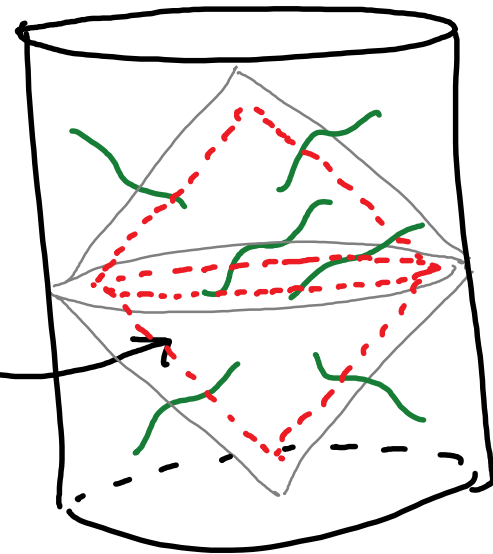
Can we "dress up"  $CFT_2$  to look like a state of  $CFT_1$ ?

$$|\Psi\rangle_{\text{CFT}_1}^{t=0}$$

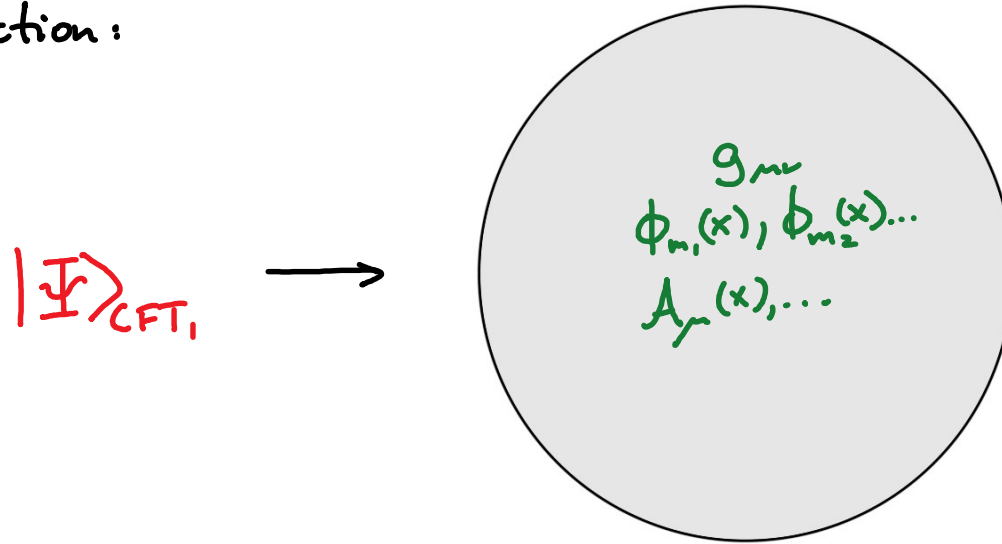
dual to



Can we find  $|\Psi\rangle_{\text{CFT}_2}$  whose dual spacetime includes



obvious objection:

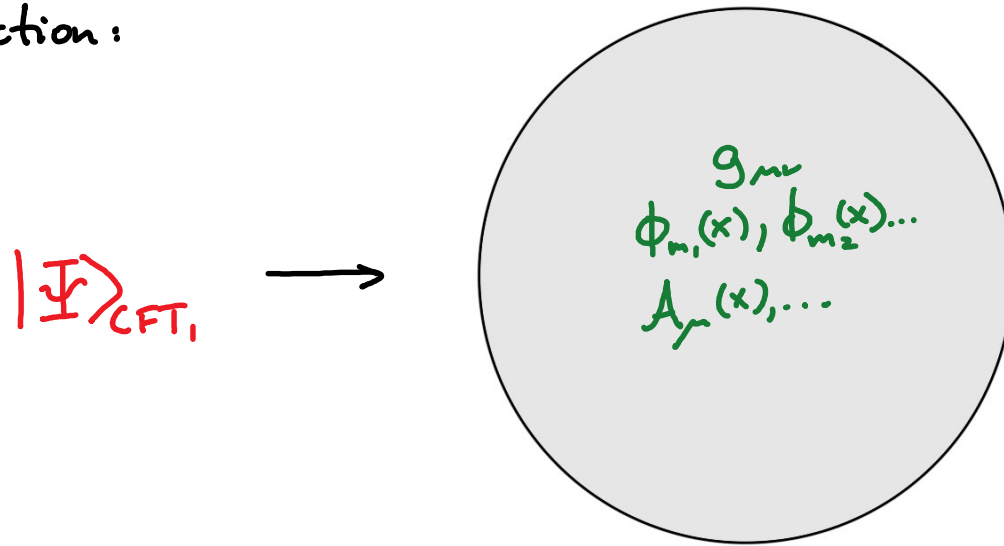


specific fields (incl. KK modes) associated w. specific CFT ops.

$$\begin{array}{ccc} \phi_{m_1} & \longleftrightarrow & \mathcal{O}_{\Delta_1} \\ \phi_{m_2} & \longleftrightarrow & \mathcal{O}_{\Delta_2} \\ & & \vdots \end{array}$$

asymptotic behavior of spacetime related to short-distance correlators of these operators.

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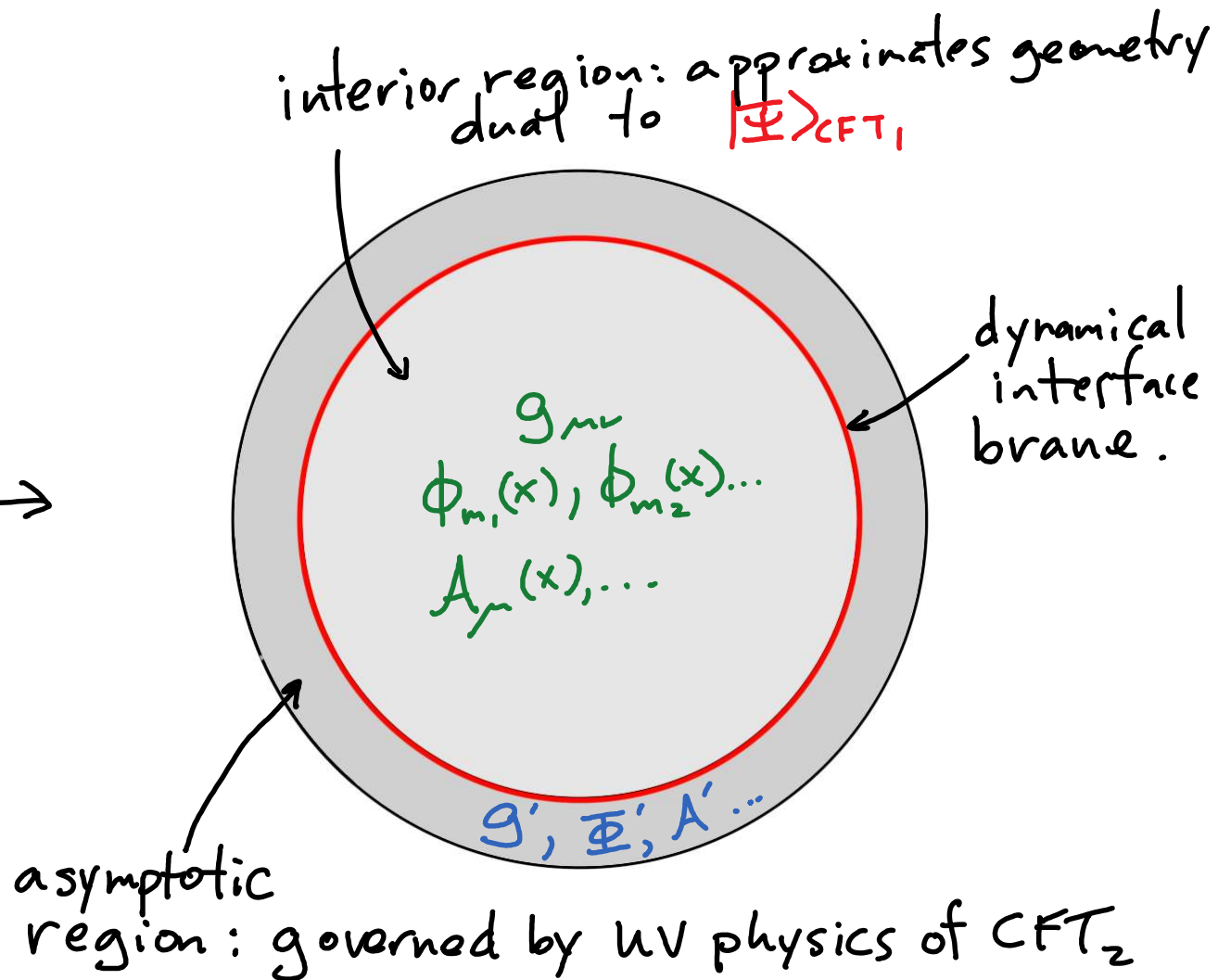
asymptotic behavior of spacetime related to short-distance correlators of these operators.

How can we hope to describe same spacetime using  $\text{CFT}_2$  w. completely different operator spectrum?

Key point: different semiclassical theories of gravity can be part of same non-perturbative theory.

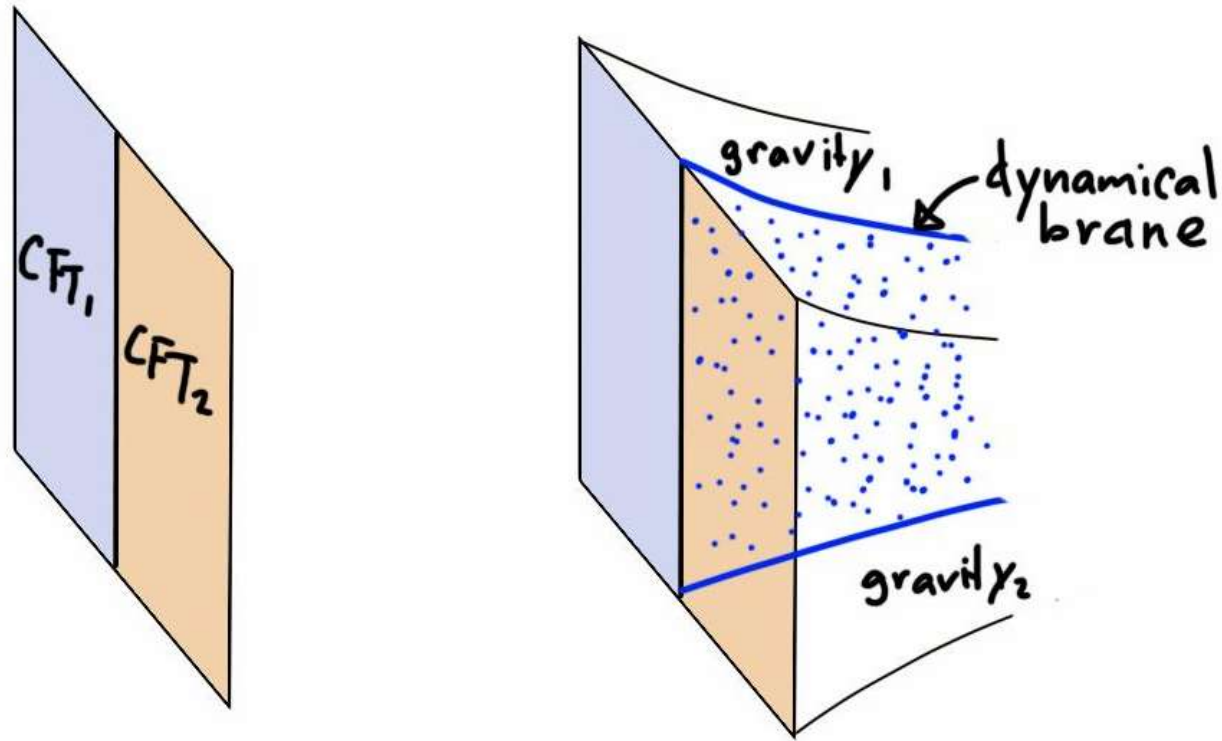
Could have:

$|\Psi\rangle_{\text{CFT}_2}$



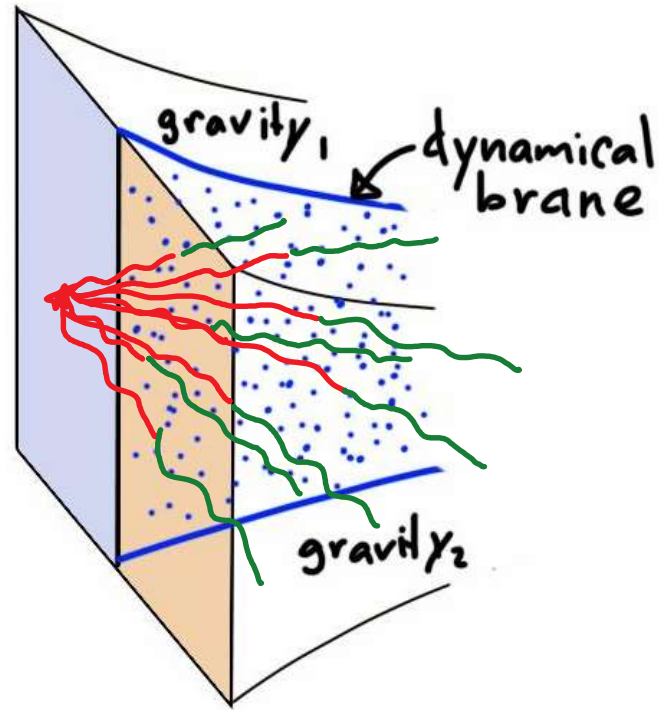
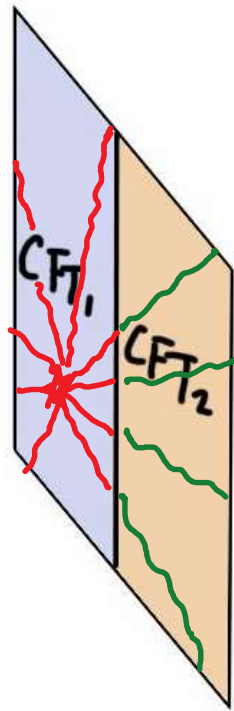


Which holographic CFTs / semiclassical gravity theories can be connected in this way?



Suggestion:

Theories of gravity dual to different CFTs must be part of the same theory (with non-perturbative interface brane) provided the CFTs can be coupled non-trivially at an interface.



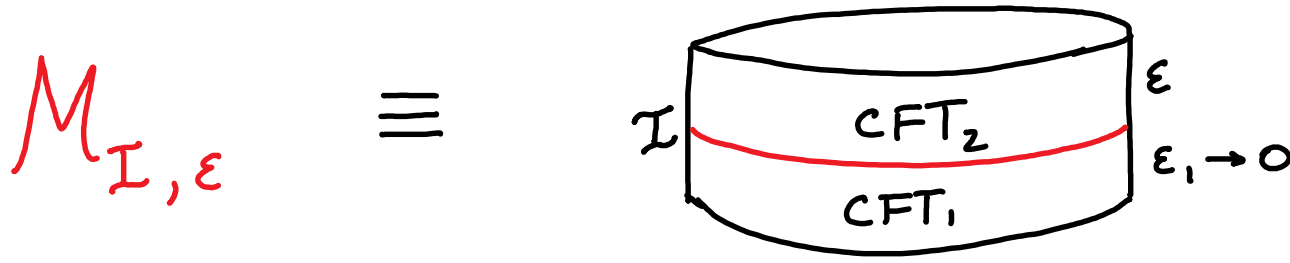
Suggestion:

Theories of gravity dual to different CFTs must be part of the same theory (with non-perturbative interface brane) provided the CFTs can be coupled non-trivially at an interface.

THIS TALK: assume  $CFT_1$  and  $CFT_2$  can be coupled in this way

- Construct specific states of  $CFT_2$  that approximate  $CFT_1$  states
- Explore gravity dual. Do we realize the bubble picture?

A map from  $\mathcal{H}_{\text{CFT}_1}$  to  $\mathcal{H}_{\text{CFT}_2}$



$$= \lim_{\epsilon_1 \rightarrow 0} e^{-\epsilon H_2} \hat{Q}_{\mathcal{I}} e^{-\epsilon_1 H_1}$$

Euclidean evolution w.r.t.  $H_{\text{CFT}_2}$ : results in finite energy state

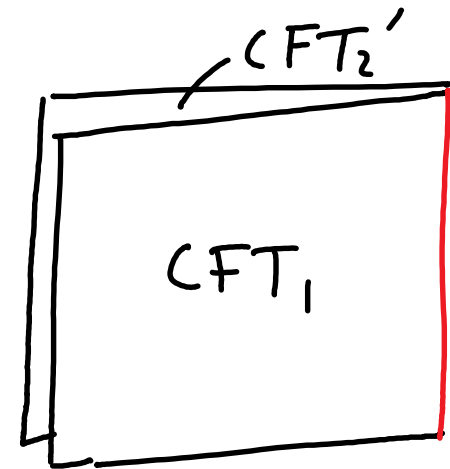
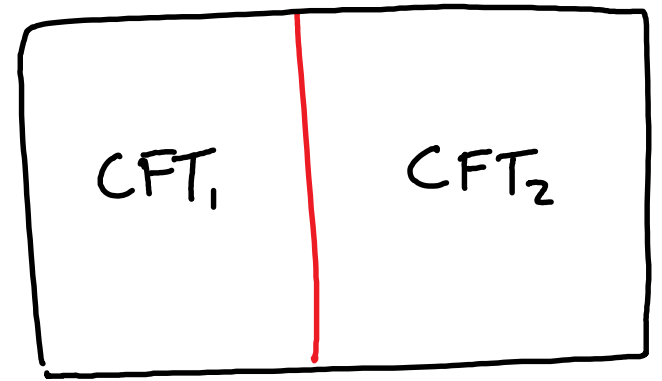
Quench operator (switches from  $H_{\text{CFT}_1}$  to  $H_{\text{CFT}_2}$ )

↓  
defined by choice of interface

Interface between  $CFT_1$  and  $CFT_2$

"

Boundary for  $CFT_1 \otimes CFT_2'$

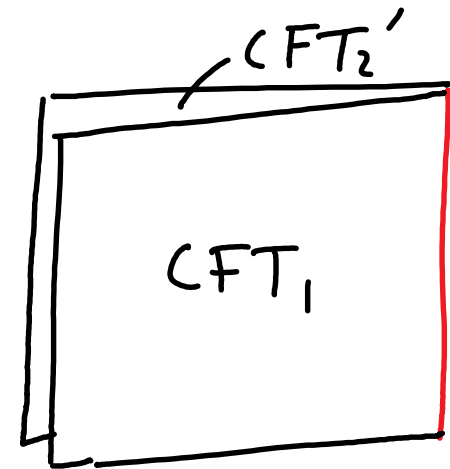
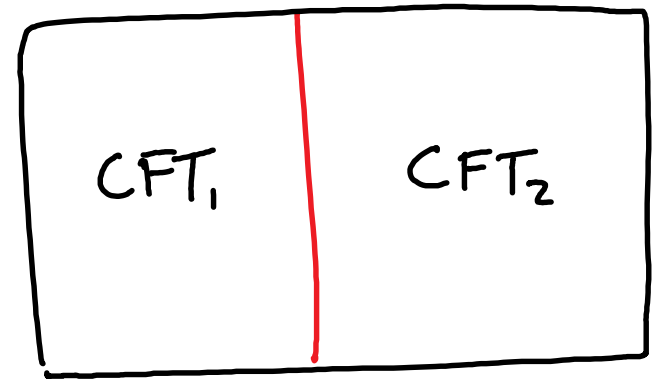


Usually, many possible choices (determined in principle by boundary bootstrap)

Interface between  $CFT_1$  and  $CFT_2$

"

Boundary for  $CFT_1 \otimes CFT_2'$



Usually, many possible choices (determined in principle by boundary bootstrap)

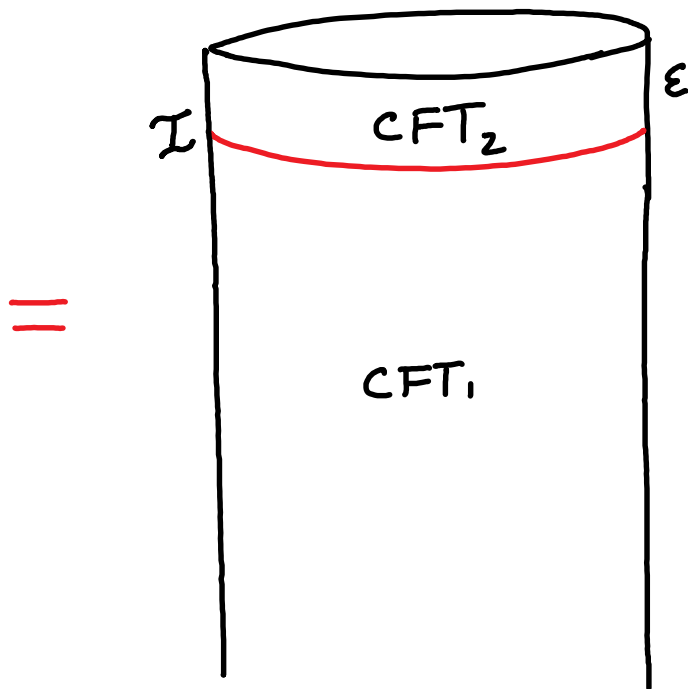
Useful parameter: interface entropy



$$S = \frac{c_1 + c_2}{6} \log\left(\frac{2L}{\epsilon}\right) + \log(g)$$

Focus on CFT<sub>1</sub> vacuum

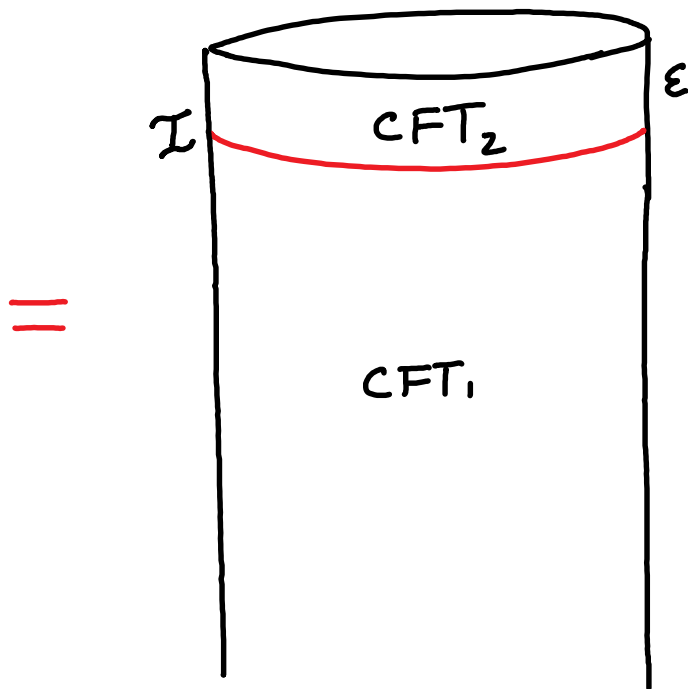
Take:  $|\Psi_2\rangle = M_{I,\epsilon} |vac\rangle_{CFT_1}$





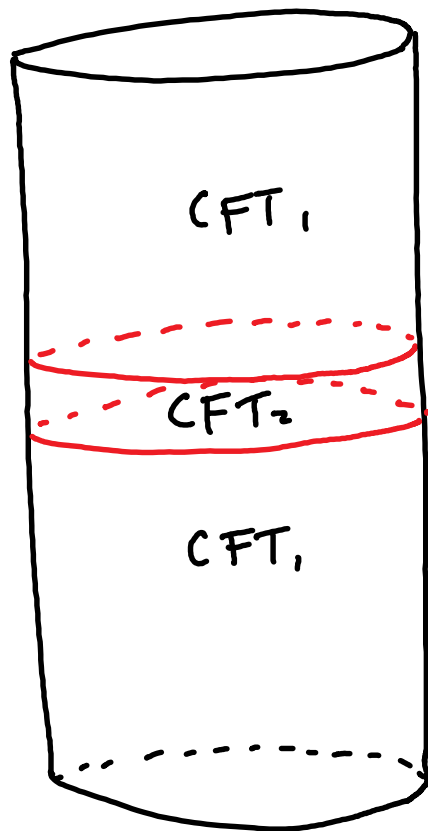
Focus on CFT<sub>1</sub> vacuum

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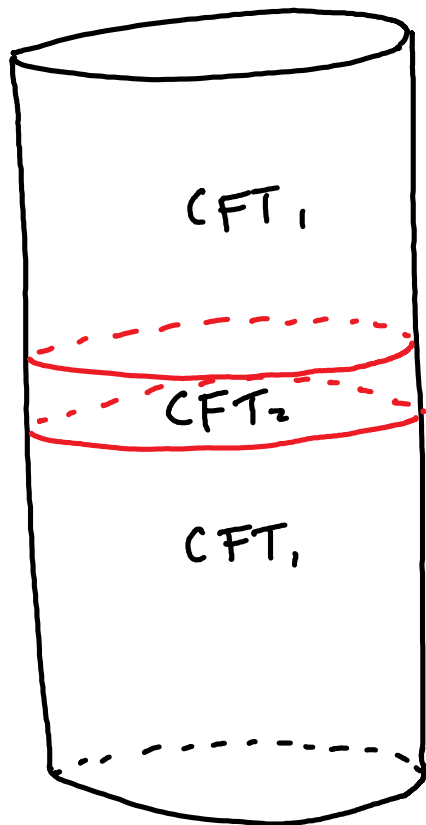


What is the dual geometry?

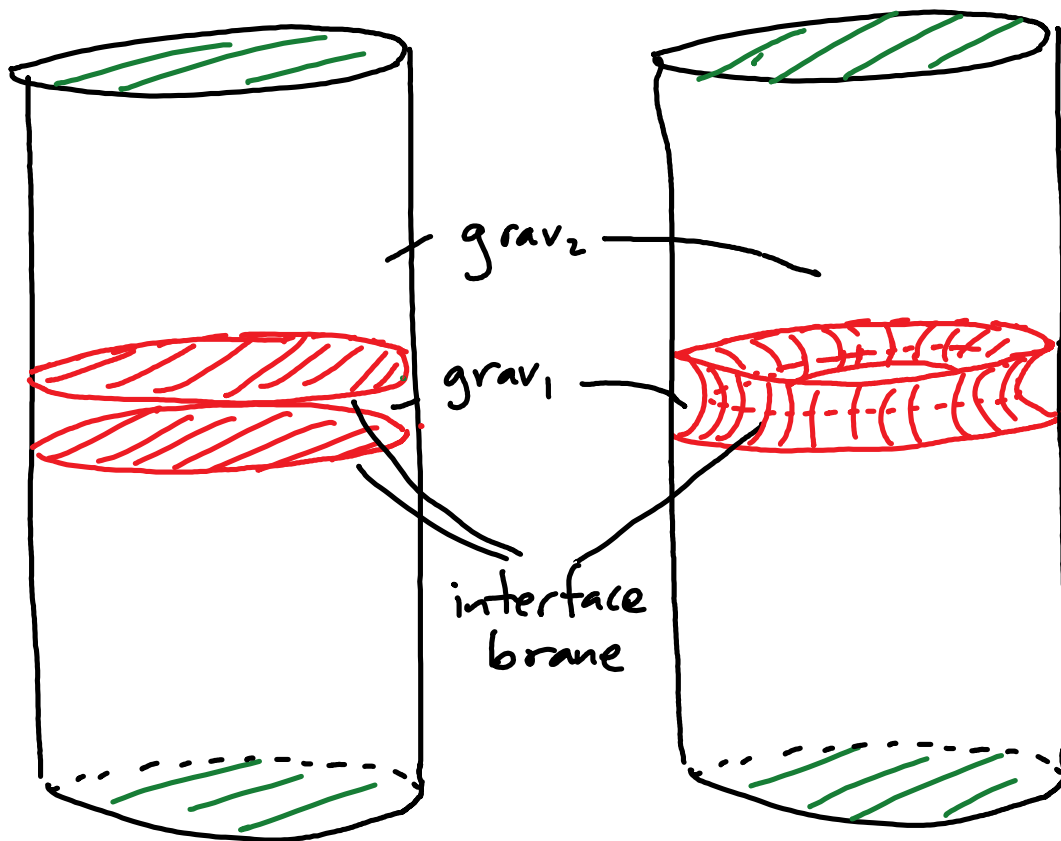
Path integral  
for  $\langle \Psi_2 | \dots | \Psi_1 \rangle$



Path integral  
for  $\langle \Psi_2 | \dots | \Psi_1 \rangle$

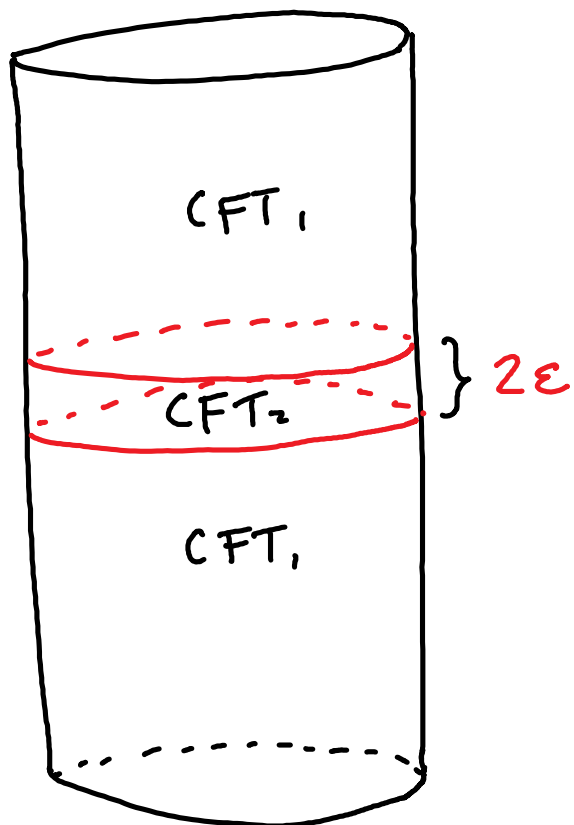


Possible dual Euclidean geometries:

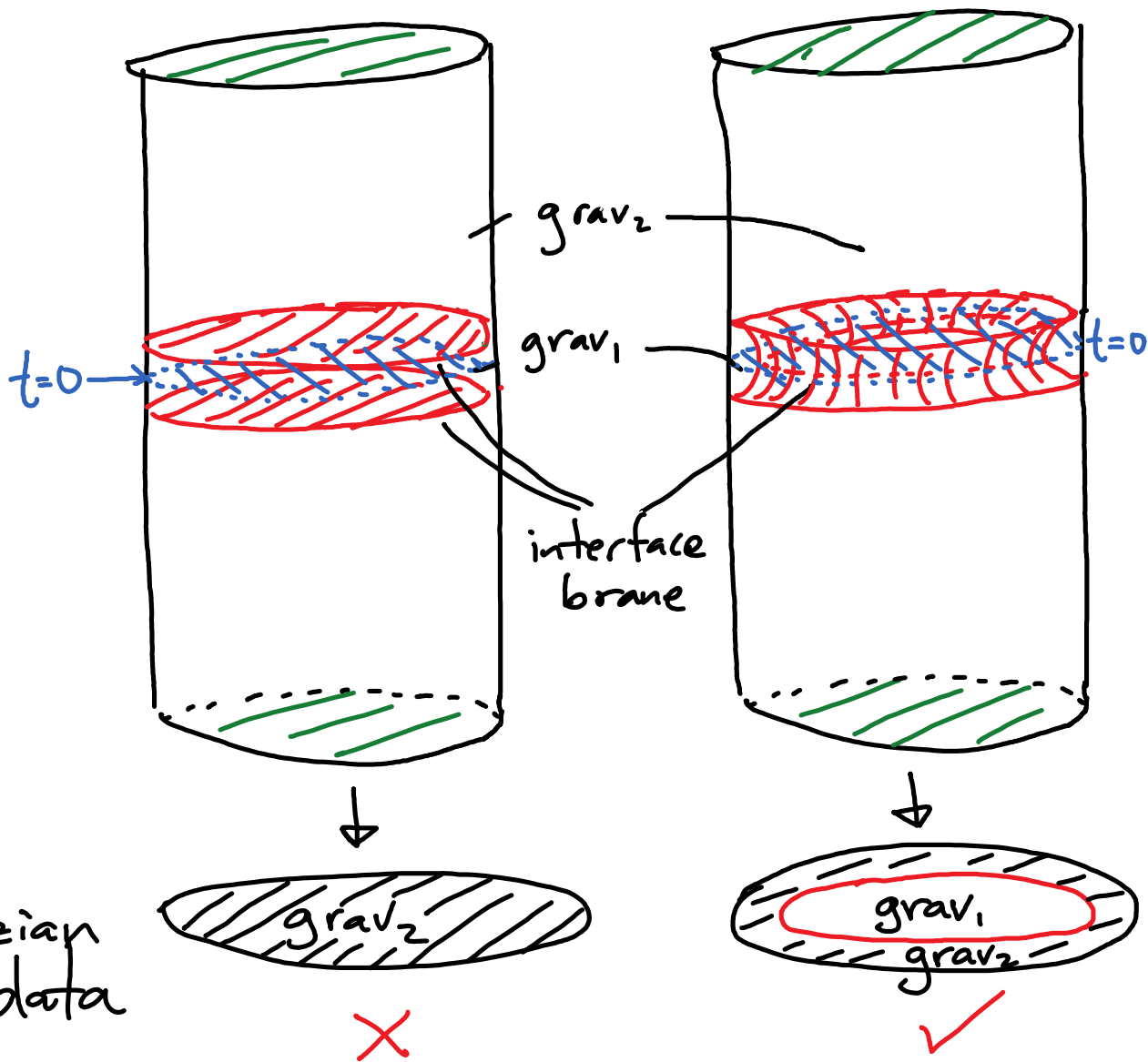


Path integral  
for  $\langle \Psi_2 | \dots | \Psi_1 \rangle$

Possible dual Euclidean geometries:

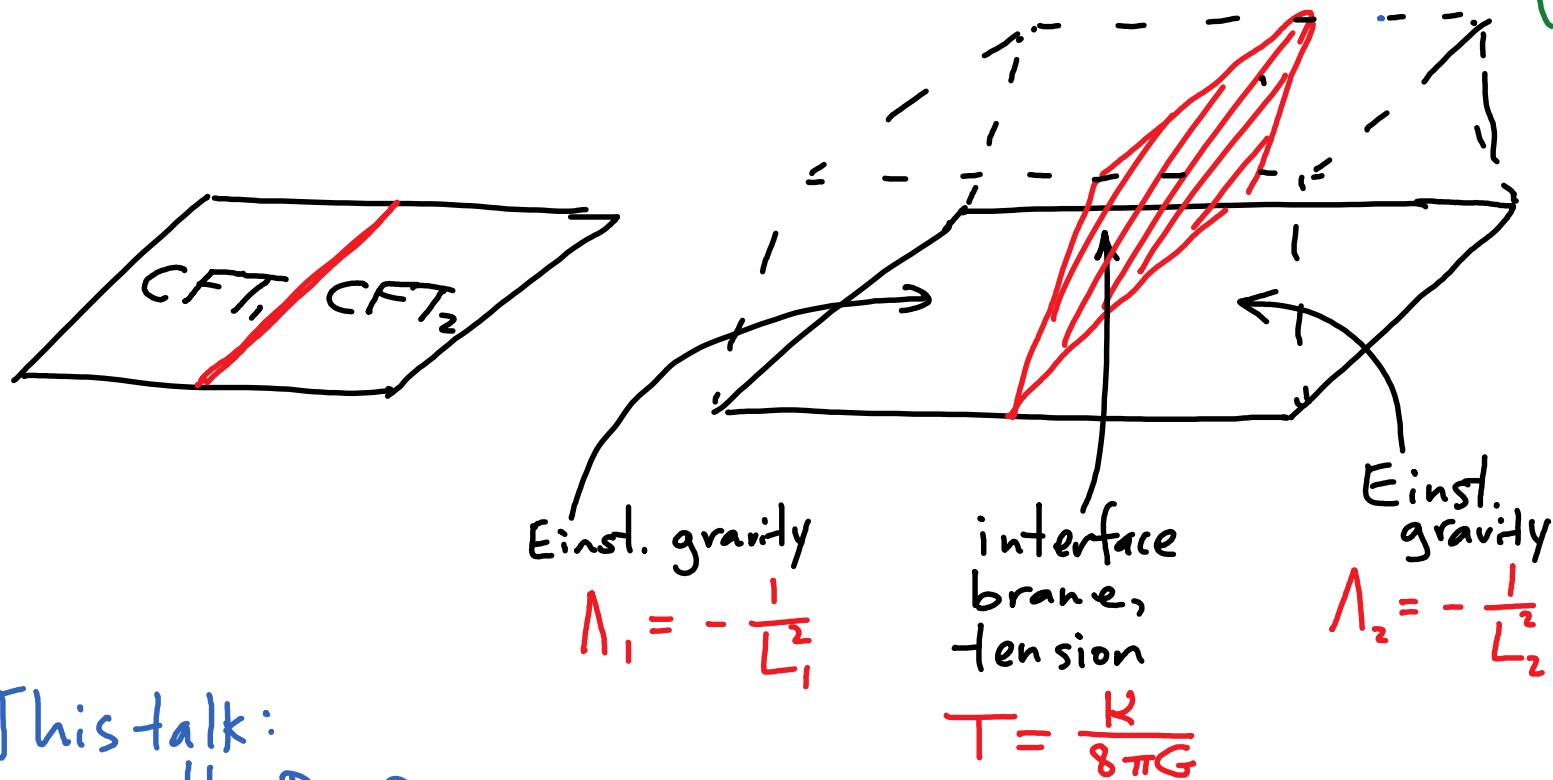


Lorentzian  
initial data



Does the "bubble" geometry have lower action when  $\varepsilon \rightarrow 0$ ?

Investigate in context of a simple model.

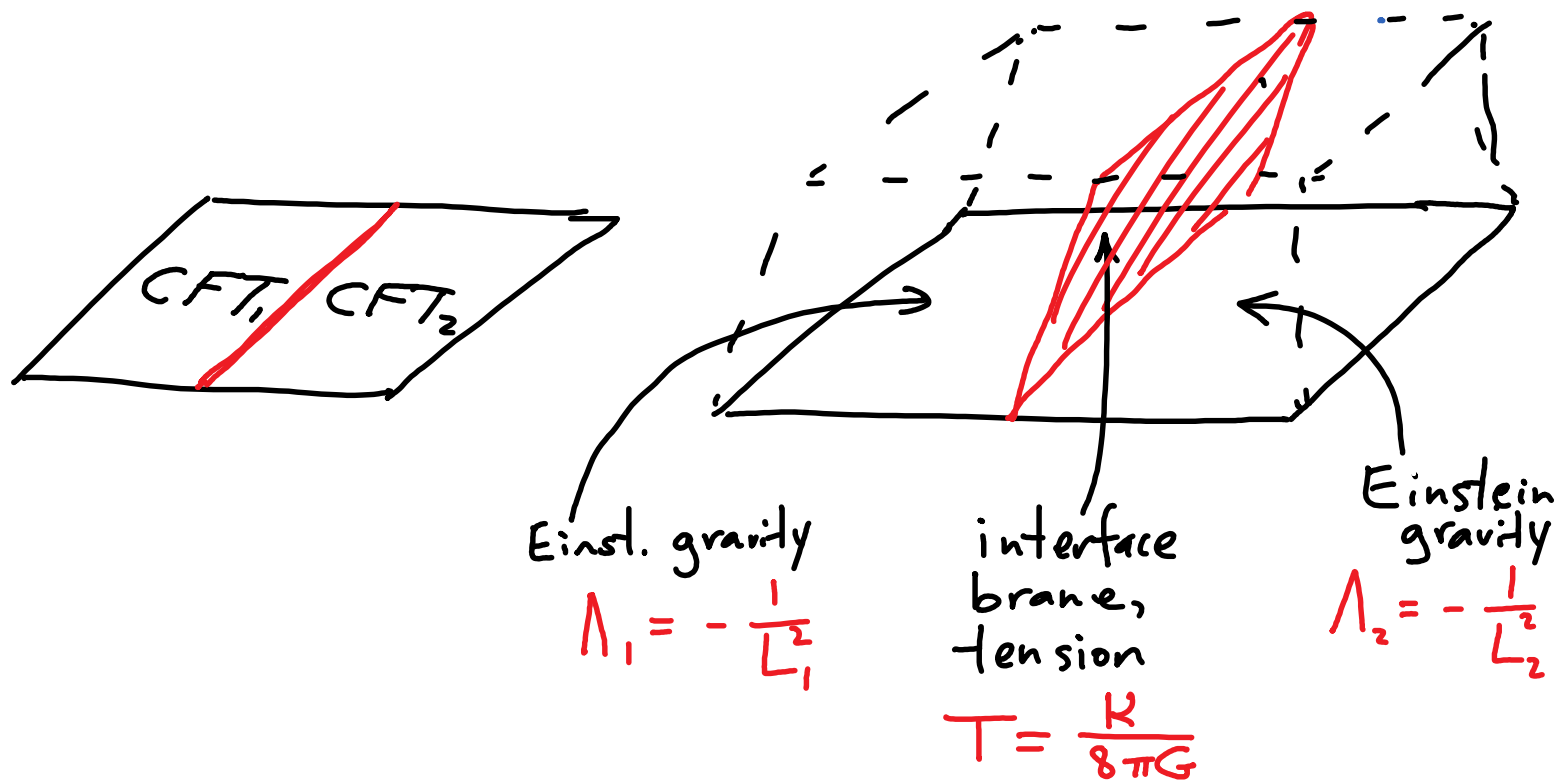


(similar to BCFT model of Karch/Randall, Takayanagi)

This talk:  
mostly  $D=3$

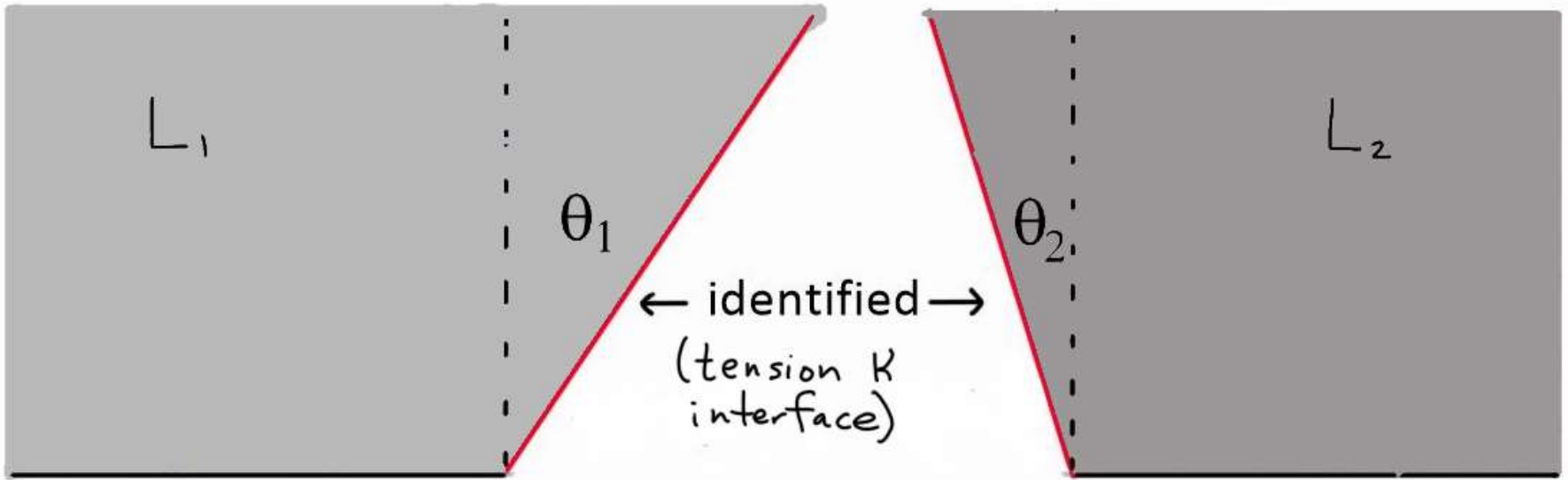
Does the "bubble" geometry have lower action when  $\varepsilon \rightarrow 0$ ?

Investigate in context of a simple model.



Microscopically: interface brane could represent region of 10D geometry where internal space has transition

# Single Interface



interface brane trajectory determined via  
junction conditions  $\rightarrow$  fixed angle in Poincaré coords.

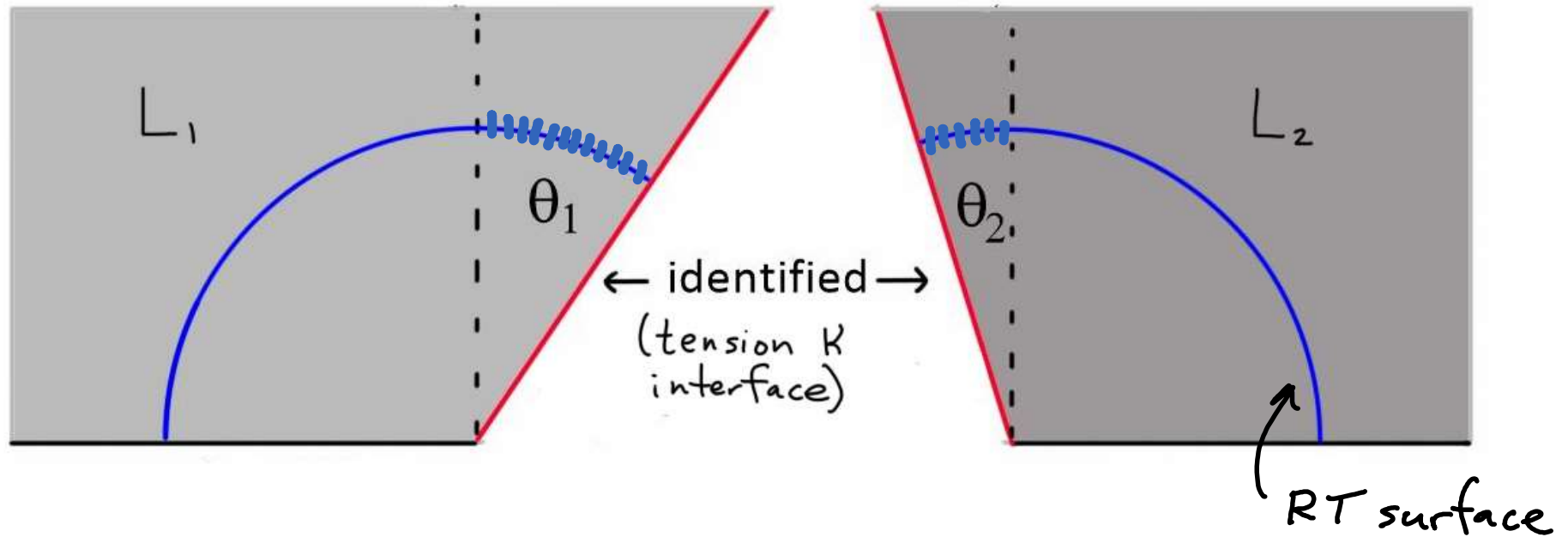




# Single Interface

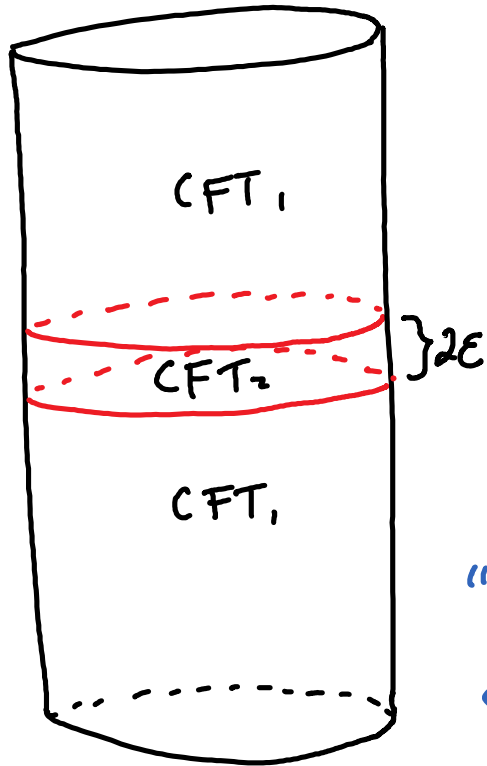
$$\frac{\log g}{c_1 \quad c_2}$$

$$\log g = \frac{1}{4G} (\text{area of } \text{---})$$



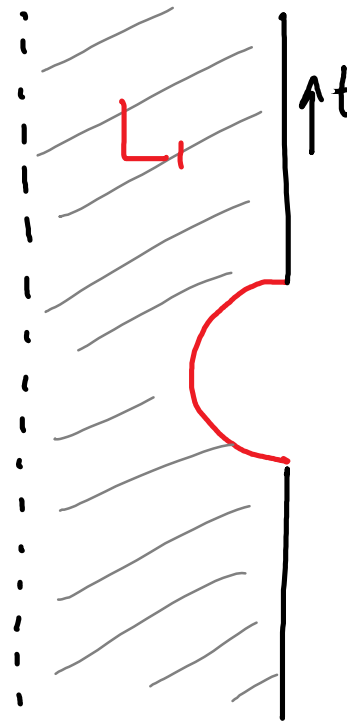
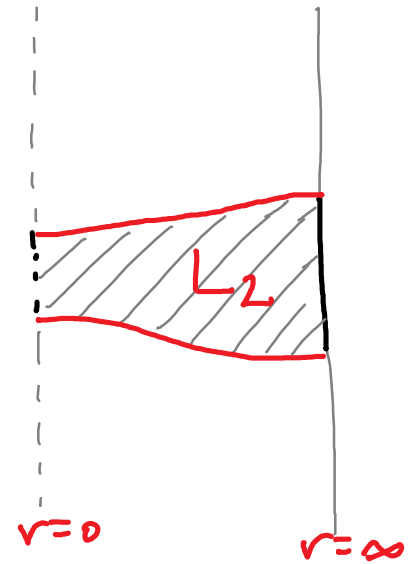
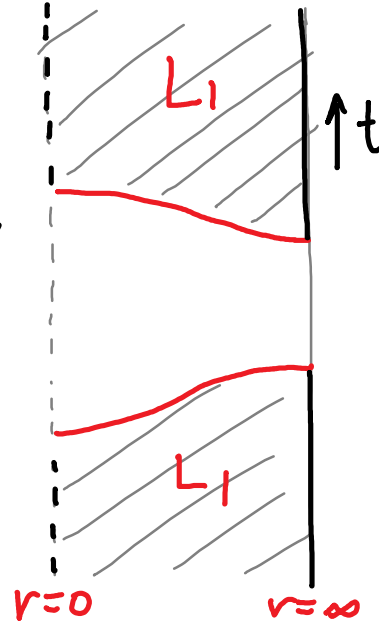
Brane tension related to interface entropy:

$$K \in \left( \left| \frac{1}{L_1} - \frac{1}{L_2} \right|, \frac{1}{L_1} + \frac{1}{L_2} \right) \longleftrightarrow \log(g) \in (-\infty, \infty)$$

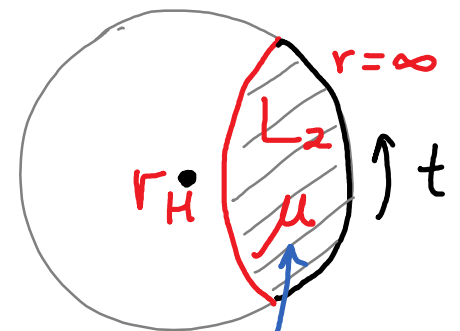


"pure AdS solutions"

"black hole solutions"



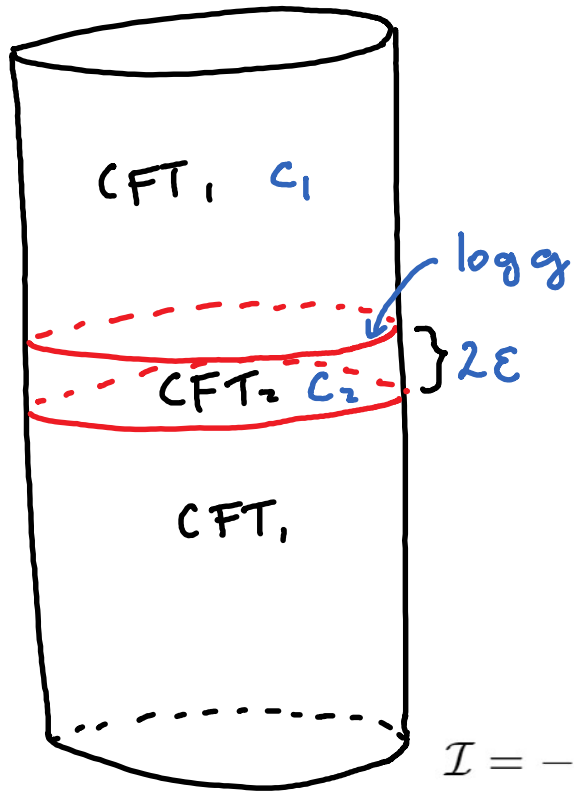
AdS-Schwarzschild



black hole mass parameter (fixed by ε)

see also: Fu, Marolf

For CFT parameters  $c_1, c_2, \log g, \varepsilon$ :



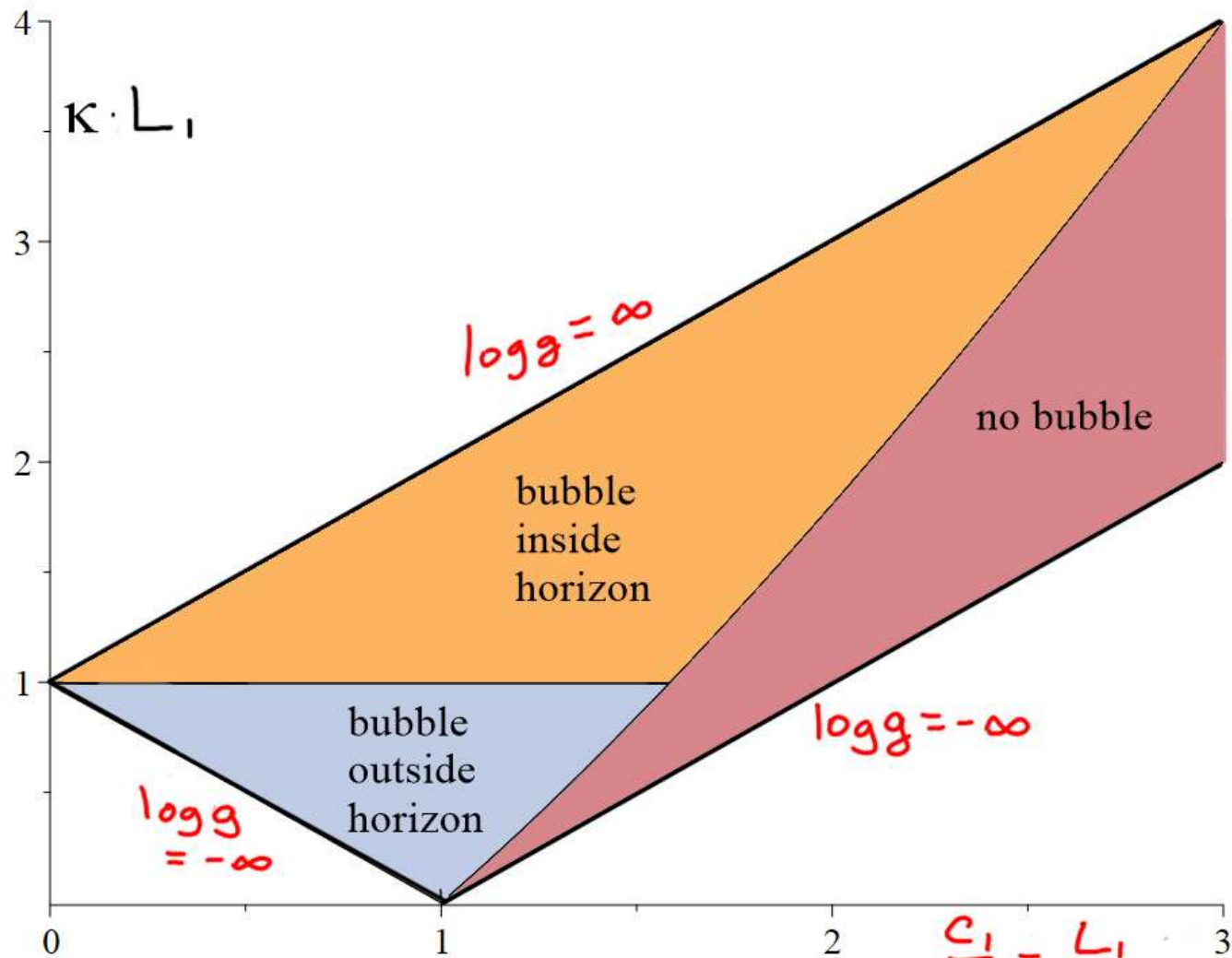
- find allowed solutions with gravity parameters  $\frac{L_i}{G} = \frac{2}{3} c_i, K(\log g)$

$$\mathcal{E}_{\text{grav}}(\mu) = \mathcal{E}_{\text{CFT}}$$

- compare actions to find least action solution

$$\mathcal{I} = -\frac{1}{16\pi G_D} \left[ \int_{\mathcal{M}_1} d^D x \sqrt{g_1} (R_1 - 2\Lambda_1) + \int_{\mathcal{M}_2} d^D x \sqrt{g_2} (R_2 - 2\Lambda_2) + 2 \int_S d^{D-1} y \sqrt{h} (K_1 - K_2) - 2(D-2) \int_S d^{D-1} y \sqrt{h} \kappa \right],$$

# Results for small $\epsilon$ :

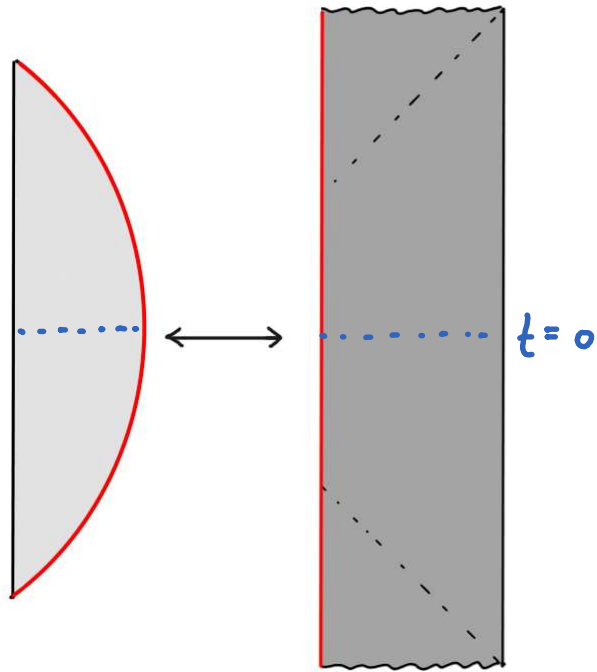


←  
CFT<sub>2</sub> has  
more d.o.f.

→  
CFT<sub>2</sub> has  
less d.o.f.

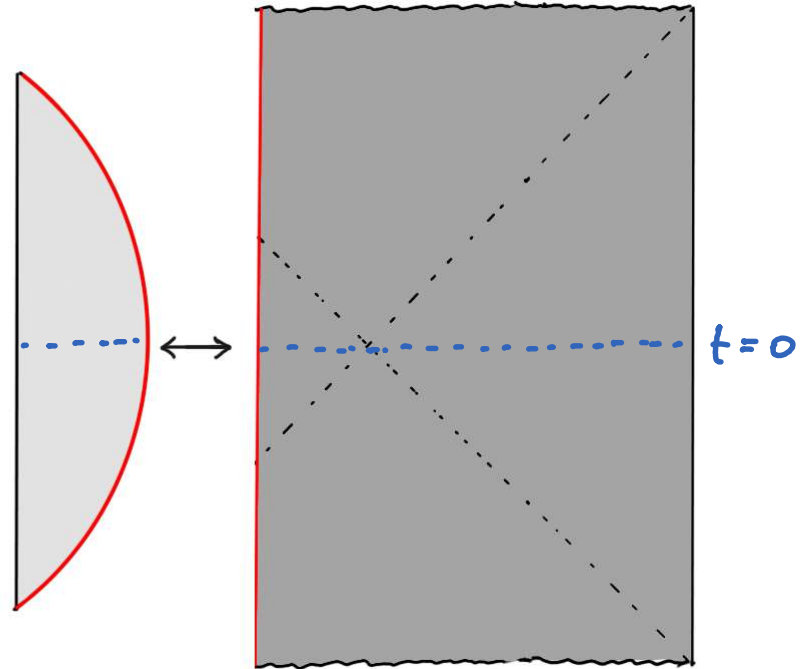
$$\frac{c_1}{c_2} = \frac{L_1}{L_2}$$

# Lorentzian Solutions:



bubble outside  
horizon:

$$c_2 > \frac{1}{3} c_1$$

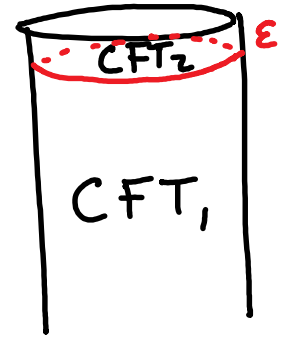


bubble behind  
horizon:

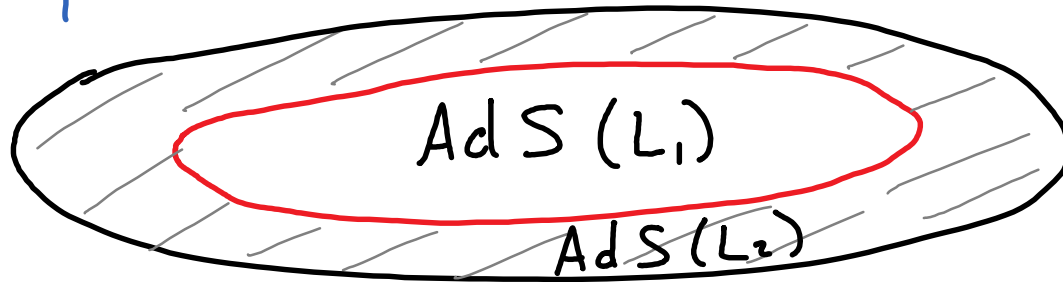
$$c_2 > \frac{1}{3} c_1, \text{ large enough } K / \log g$$

Summary so far:

$$\text{State } |\Phi\rangle_{\text{CFT}_2} = e^{-\epsilon H_2} Q_I |vac\rangle_{\text{CFT}_1}$$



faithfully encodes a bubble of geometry dual to CFT<sub>1</sub> vacuum:



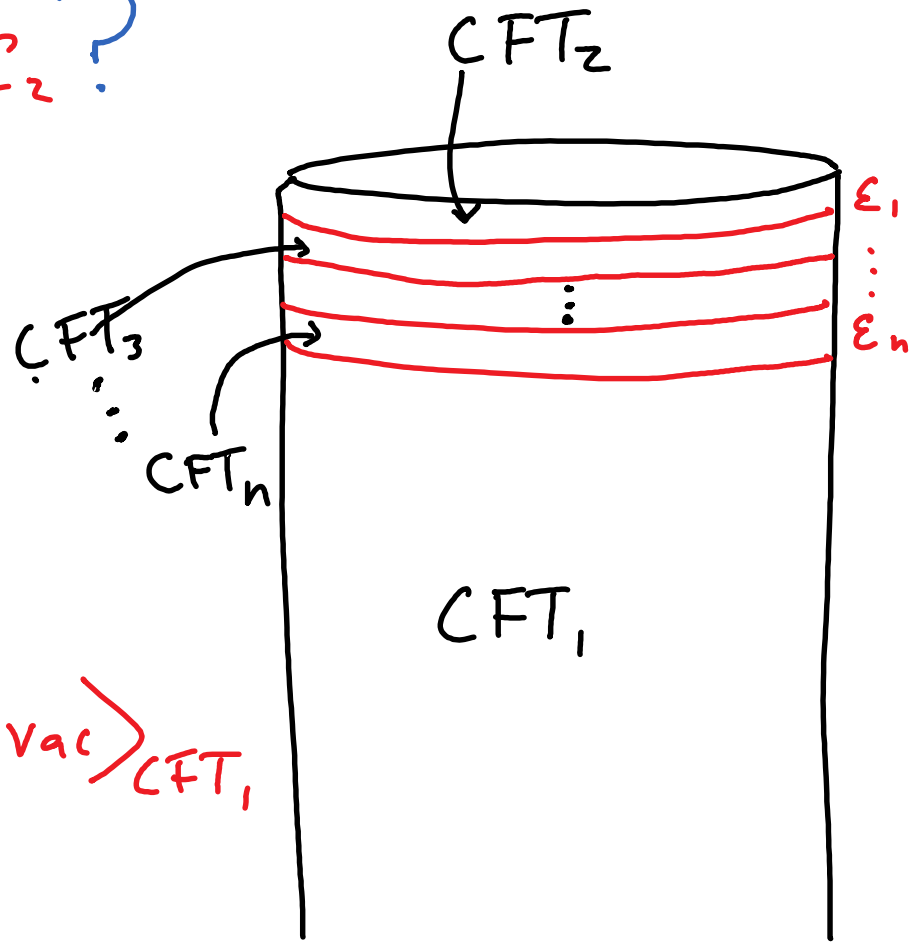
for small  $\epsilon$  if  $c_2 > c_1$

or  $\frac{1}{3}c_1 < c_2 < c_1$  and  $\log g$   
large enough.

What about smaller  $c_2$ ?

Try:

$$\begin{aligned} |\Psi_2\rangle_{\text{CFT}_2} &= e^{-\epsilon H_2} Q_{I_{23}} \dots e^{-\epsilon H_n} Q_{I_{n1}} |vac\rangle_{\text{CFT}_1} \end{aligned}$$



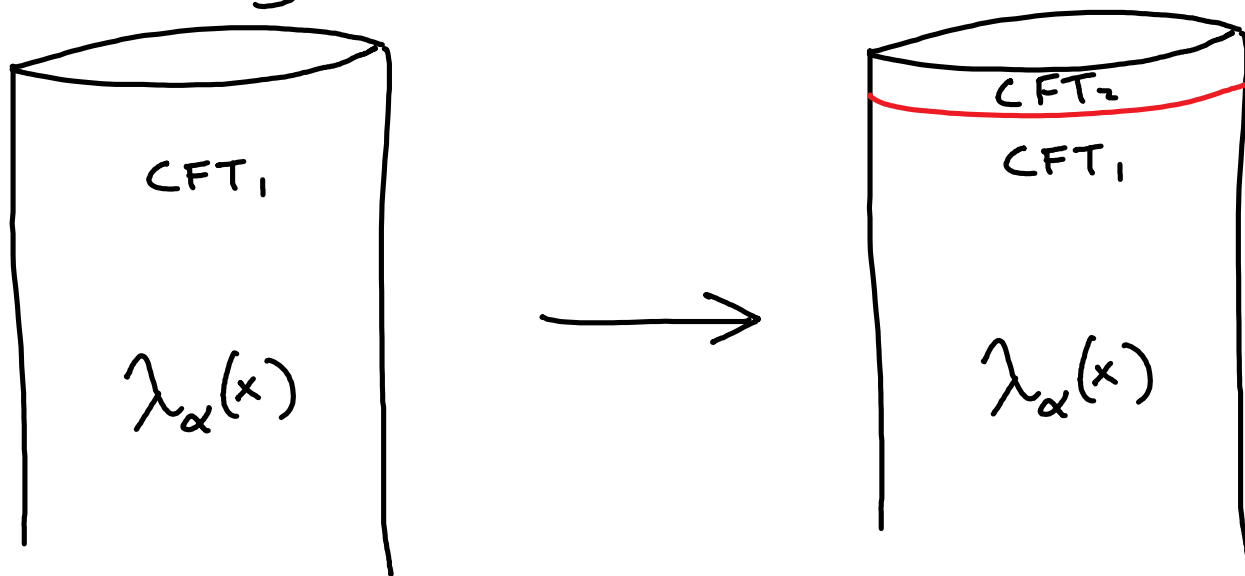
Preliminary results: for small  $(\frac{\Delta c}{c})_i$ , small  $K_i$ , nested bubble solution is preferred!

In the context of the model:

Any CFT can encode spacetime  
dual to vacuum of any other CFT



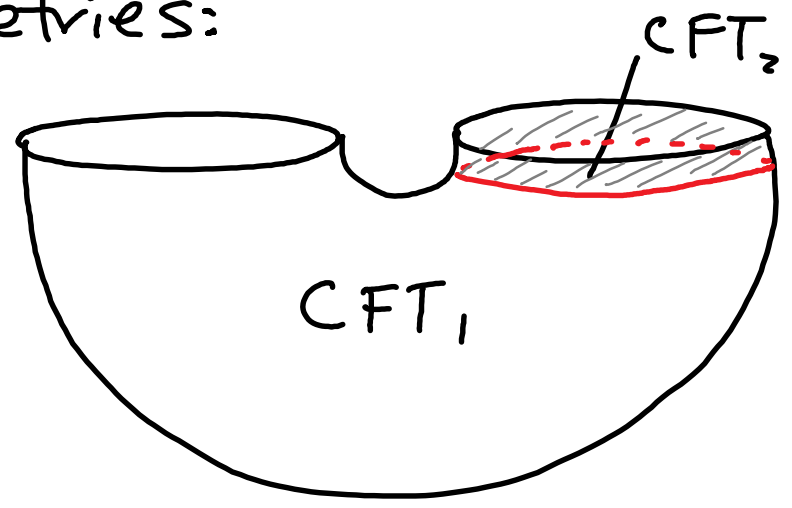
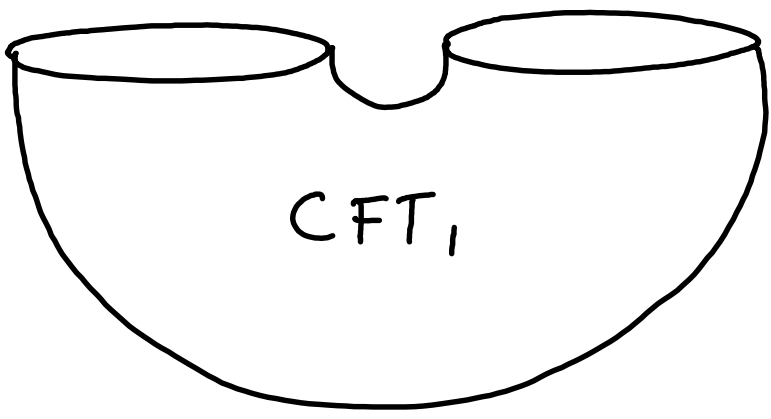
Approximating other states:



Add sources  $\int \lambda_\alpha(x) \mathcal{O}_\alpha(x)$  to Euclidean action in path integral to get CFT<sub>1</sub> states dual to very general perturbations of AdS.

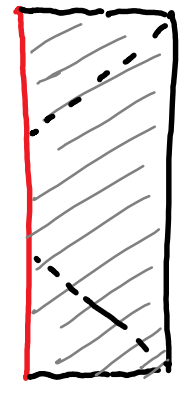
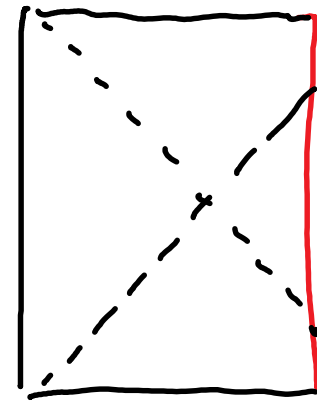
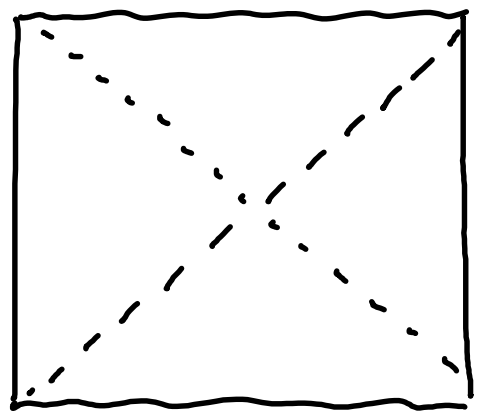
Same mapping should give CFT<sub>2</sub> state approximating this geometry (at least perturbatively)

Approximating black hole geometries:

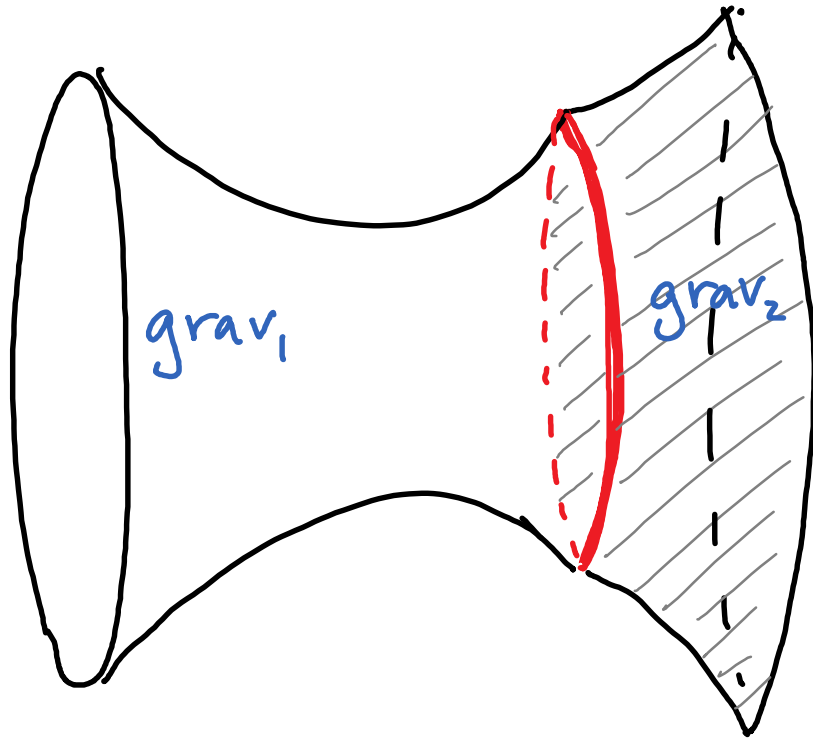
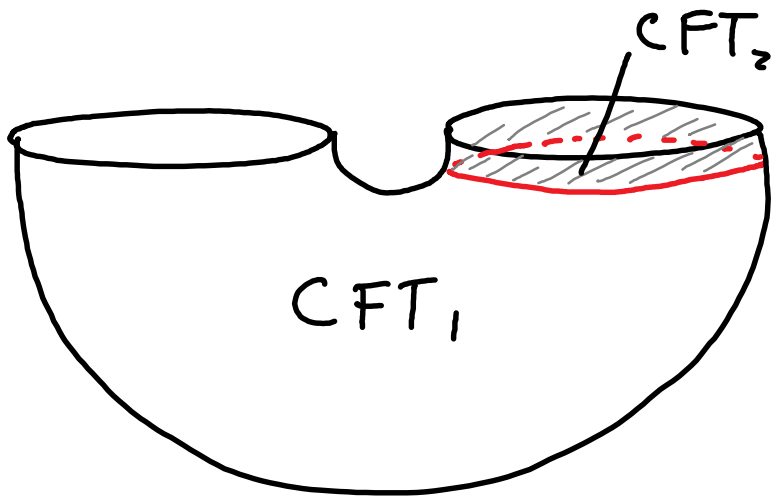


$$|\Psi_{\text{TFD}}\rangle_{\text{CFT}_1 \otimes \text{CFT}_1}$$

$$\mathbb{1}^L \otimes M_{I,\epsilon}^R |\Psi_{\text{TFD}}\rangle \in \mathcal{H}_{\text{CFT}_1} \otimes \mathcal{H}_{\text{CFT}_2}$$



CFT<sub>2</sub><sup>R</sup> encoded behind-the-horizon physics of B.H. dual to CFT<sub>1</sub><sup>L</sup>



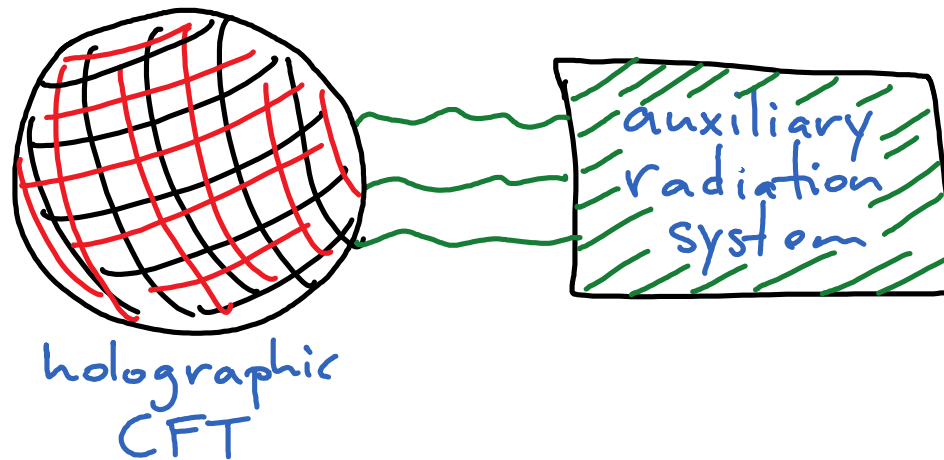
Entangled states of  
 $CFT_1 \rightarrow CFT_2$



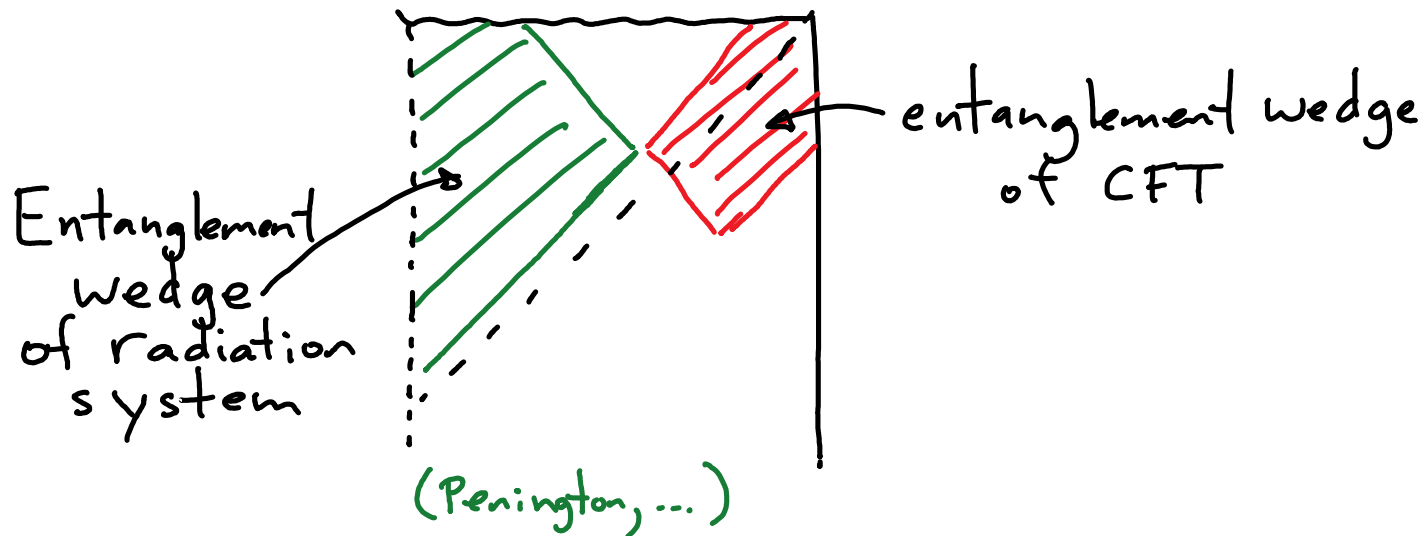
Wormhole connecting  
asymptotic regions  
described by different  
low-energy gravity  
theories

Auxiliary radiation systems can encode late-time black hole interiors.

- Maldacena-Susskind
  - MVR
  - Penington
- Almheiri, Englehardt, Marolf, Maxfield
- Almheiri, Mahajan, Maldacena, Zhao

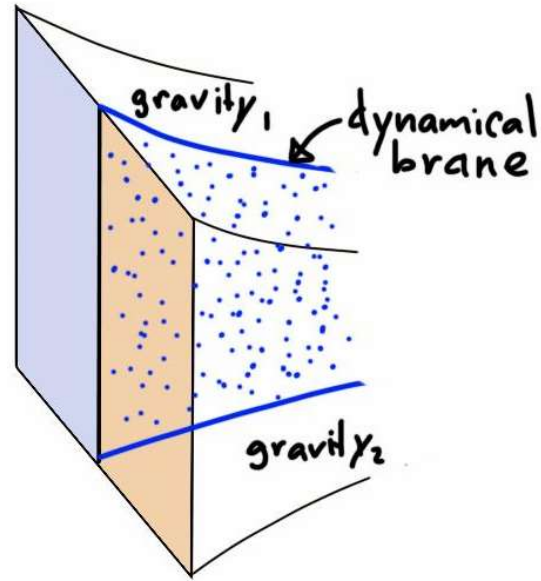
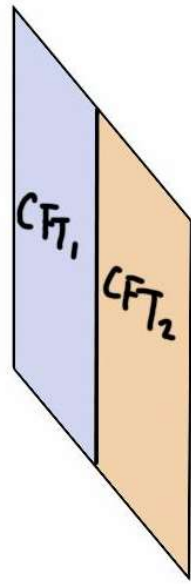


After Page time:



# SUMMARY

we  
have  
suggested:



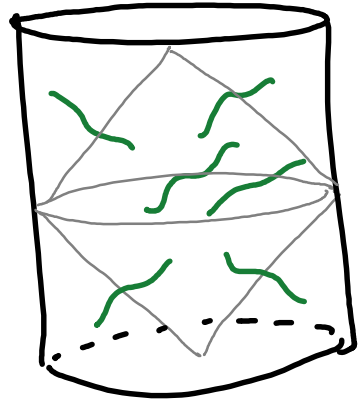
① Theories of gravity dual to CFTs that can be coupled non-trivially at an interface are part of same non-perturbative theory.

different theories  $\longleftrightarrow$  different equivalence classes of CFTs

Is there more than one class?

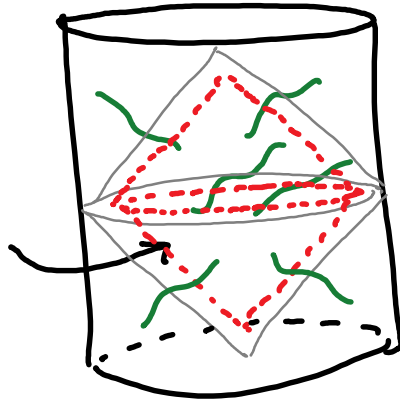
②

Given  $|\Psi\rangle_{\text{CFT}_1}$  dual to



For  $\text{CFT}_2$  in same class, can find  $|\Psi\rangle_{\text{CFT}_2}$

dual to spacetime containing



$$\text{e.g. } |\Psi\rangle_{\text{CFT}_2} = e^{-\epsilon H_2} Q_I |\Psi\rangle_{\text{CFT}_1}$$



Consistent with the idea that:

holographic CFTs



allowed asymptotically  
AdS behaviors for  
quantum gravity

but:

quantum gravity



universal physics of  
quantum systems  
e.g. entanglement  
structure & dynamics

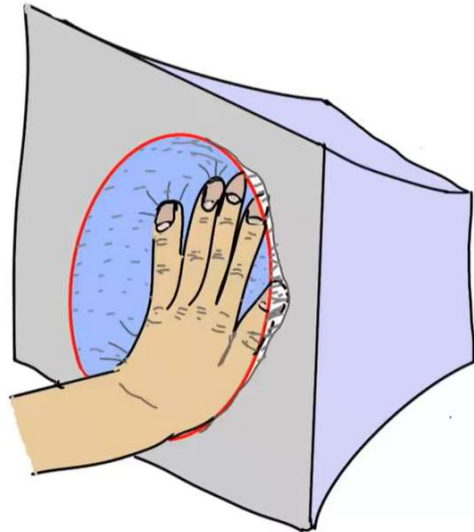
THANKS!





# PUSHING BOUNDARIES IN HOLOGRAPHY

Mark Van Raamsdonk, UBC



QGI seminar  
April 2020



QGI Virtual Seminar: Mark van Raamsdonk "Pushing boundaries in holography"

299 views • Apr 30, 2020

6 0 SHARE SAVE ...

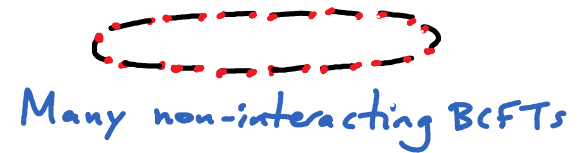
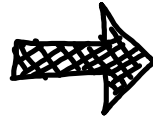


Gravity, Quantum Fields and Information  
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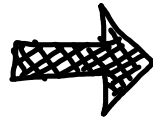
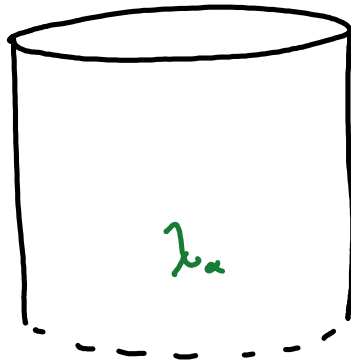
Discrete versions of holographic CFT states can faithfully encode causal patches of the original geometry

Theory:

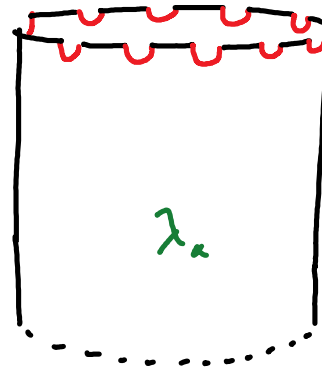


State:

$|\Psi\rangle =$



$|\Psi\rangle =$



Lorentzian  
Geometry:

