

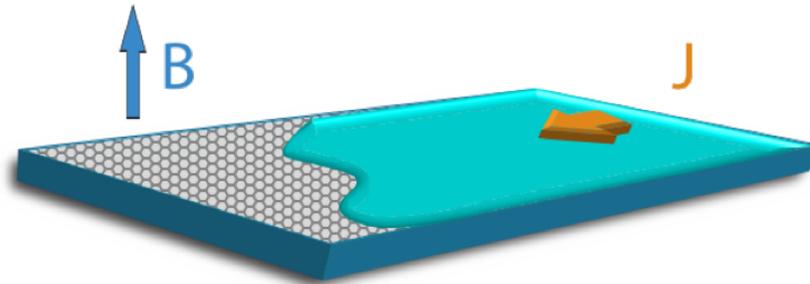
# Electron hydrodynamics

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University of Toronto

Department of Physics Colloquium

Sept 26th 2019



# Could it finally be true?

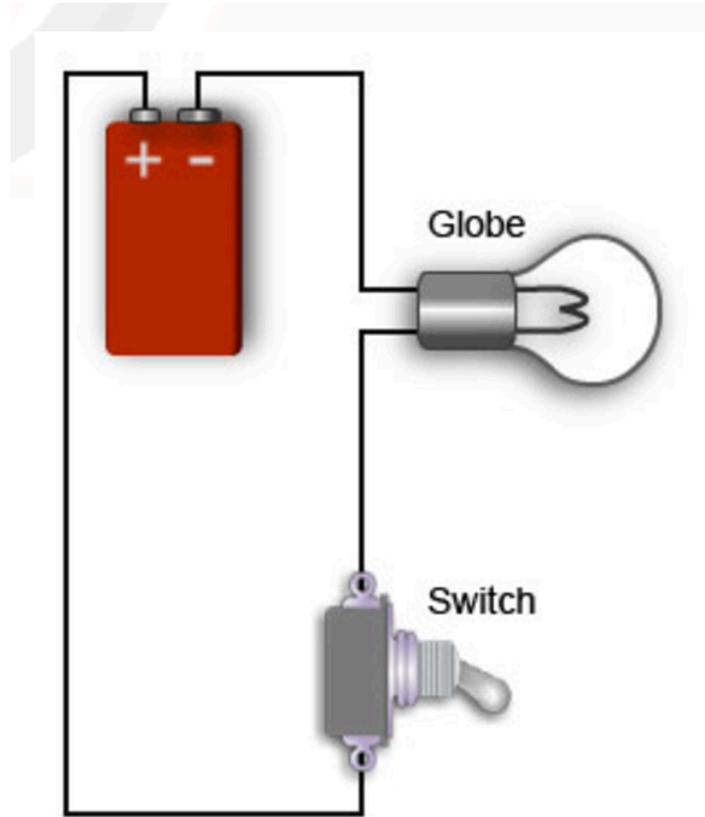


Figure 1: Electrical

?

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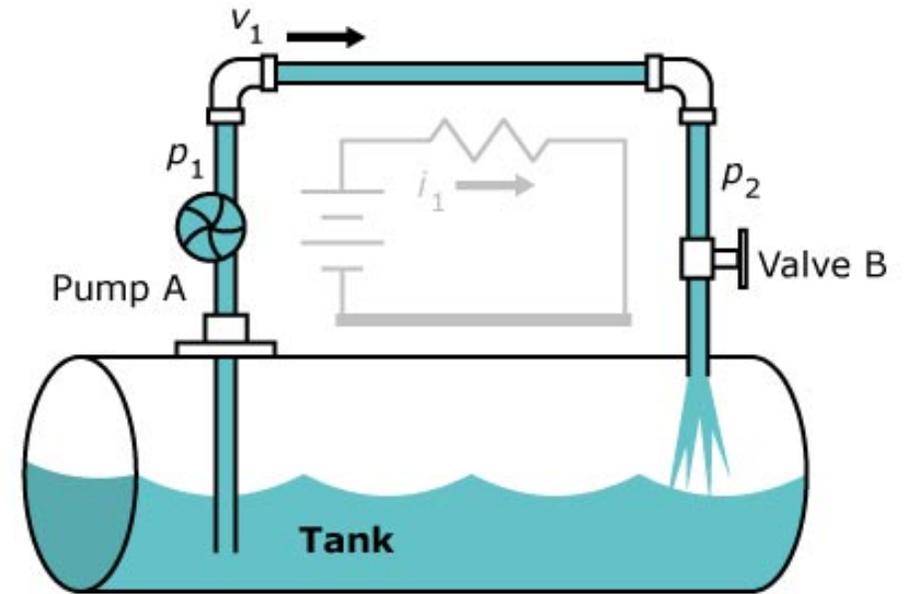
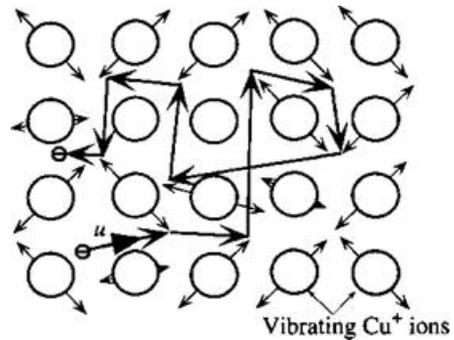


Figure 1—Water circulates endlessly around this closed loop much like electrical current.

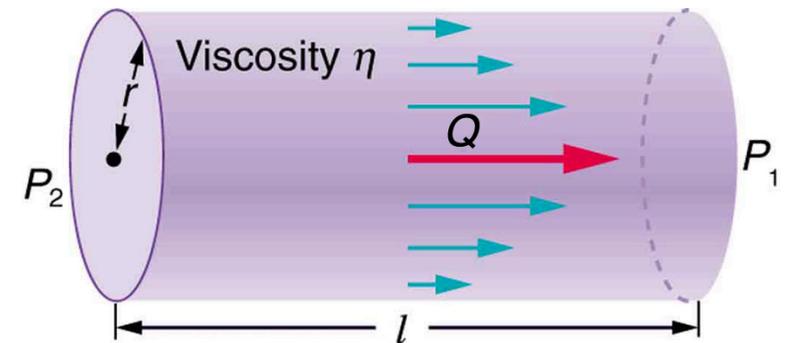
# Electric conduction versus water flow

## Metal



- Resistance arises through external scattering due to the lattice (impurities, phonons,...)

## Water

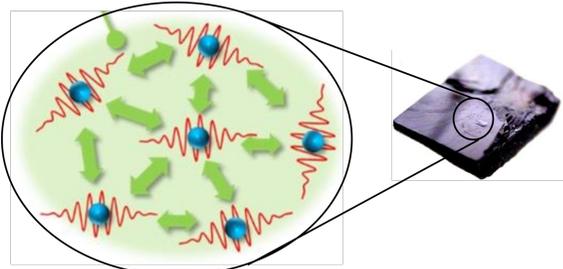


- Resistance arises through internal scattering (viscosity)

# Outline

- Electron hydrodynamics – What and why?
- Viscous Fermi liquids
- Using magnetic fields to detect viscous effects
- Conclusion and outlook

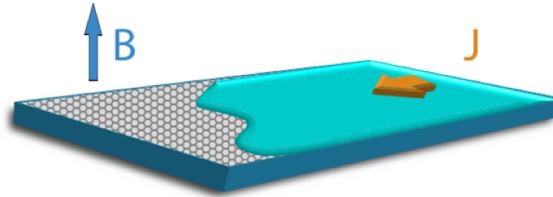
## Correlated electrons



## Hydrodynamics



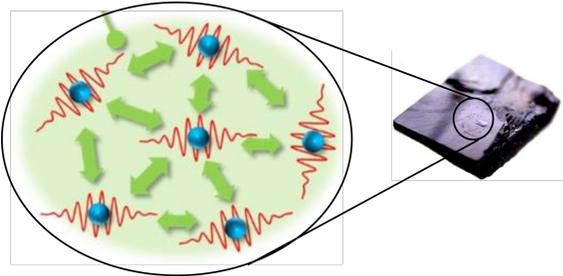
## Electron hydrodynamics



Can quantum mechanics  
*constrain* hydrodynamics?

Can quantum mechanics  
*enrich* hydrodynamics?

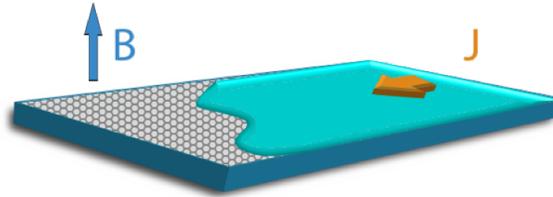
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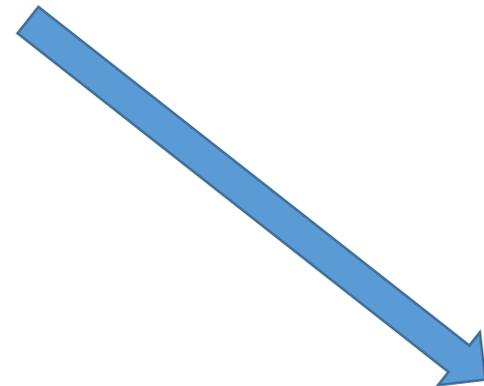
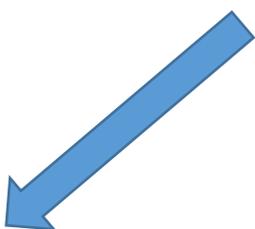
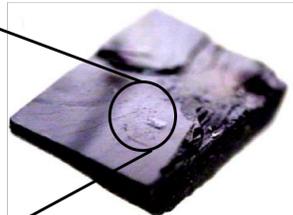
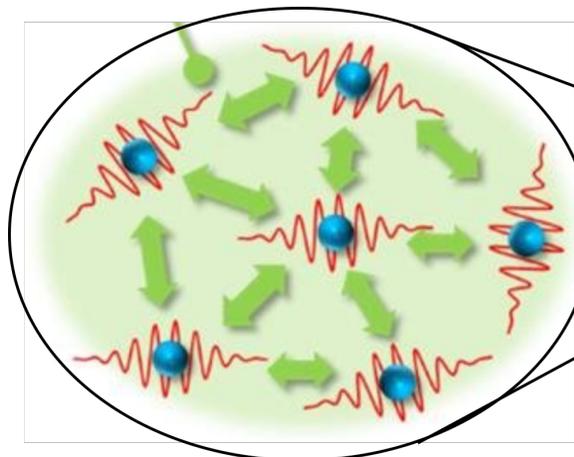
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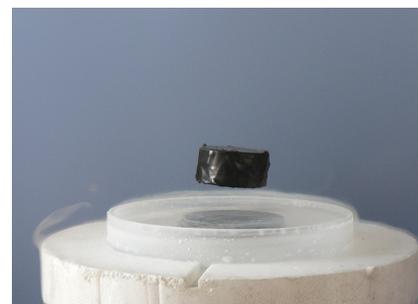
# “More is different” [Anderson '72]



Spin liquids



Topological phases



Superconductors

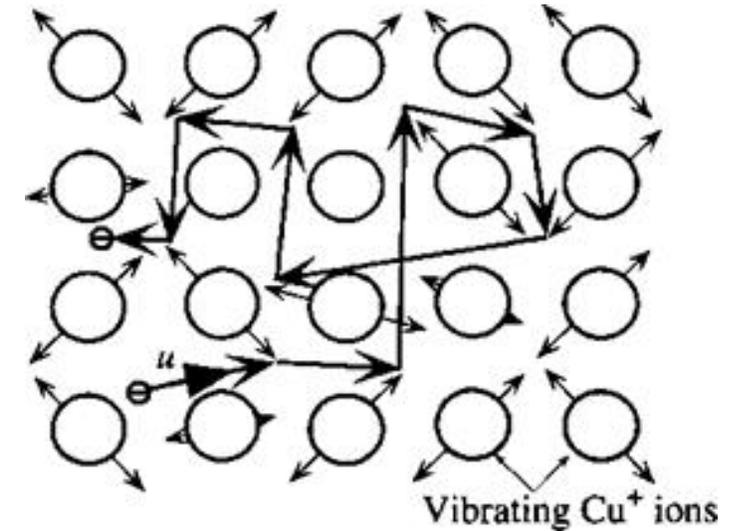
...

# Conventional metallic transport [Drude 1900]



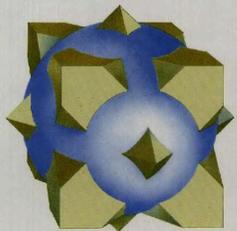
$$\rho = \frac{m}{ne^2} \left( \frac{1}{\tau} \right)$$

Mean free time of electrons, set by defects and vibrations of the lattice



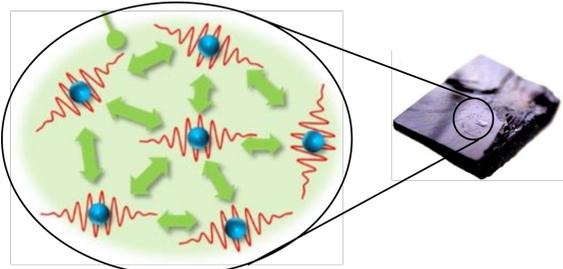
Electron-electron scattering<sup>5</sup> plays a relatively minor role in the theory of conduction in solids, for reasons to be described in Chapter 17.

ASHCROFT / MERMIN



SOLID STATE PHYSICS

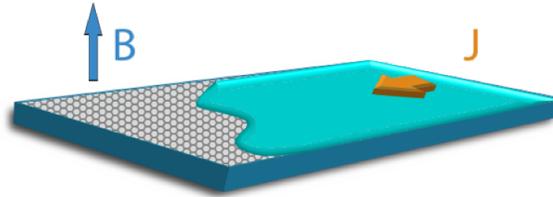
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## Hydrodynamics



## Electron hydrodynamics



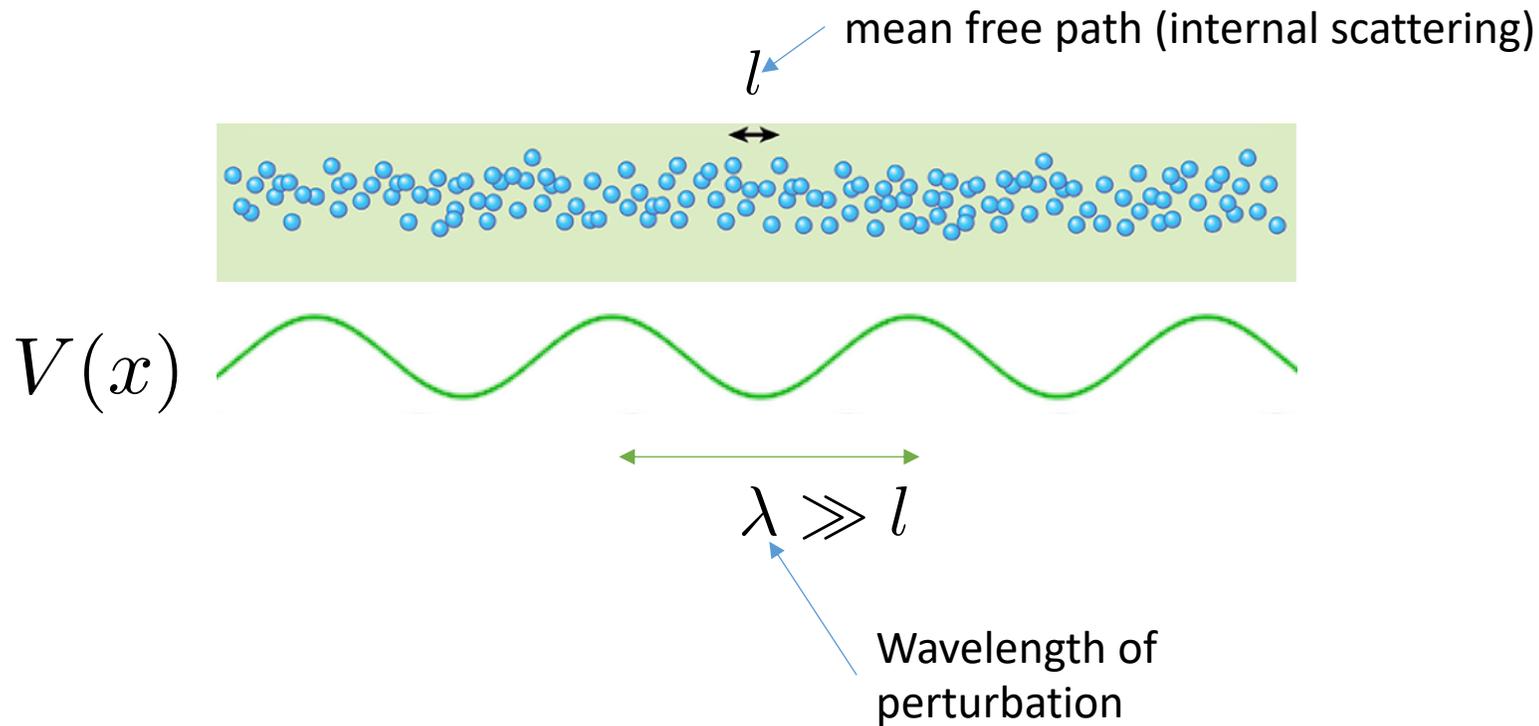
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*constrain* hydrodynamics?

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# Hydrodynamics

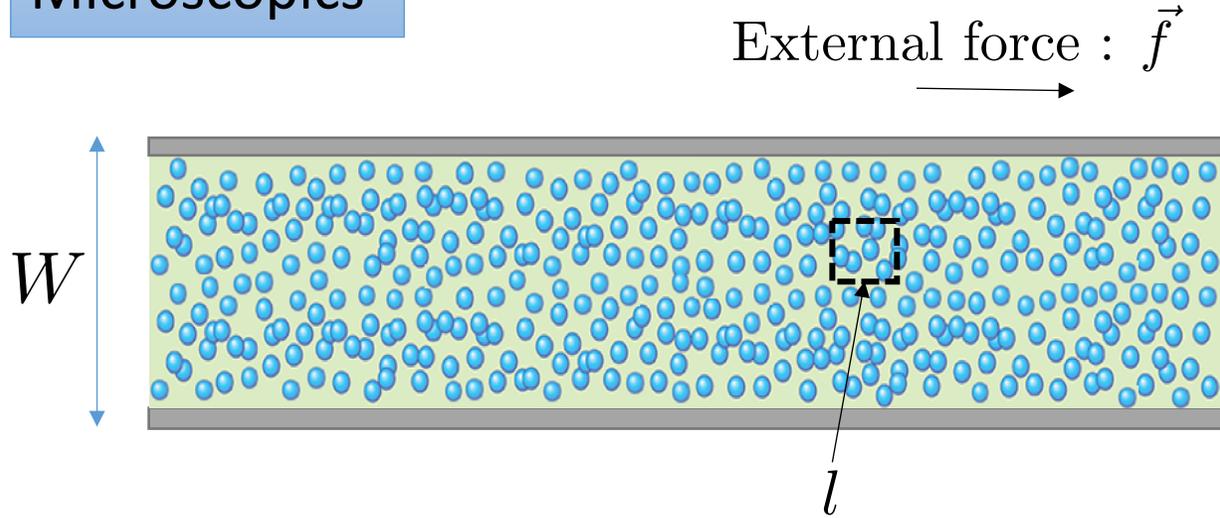


- Universal description of fluids based on conserved quantities: momentum, energy, charge,...
- Works at length/time scales much larger than the microscopic ones

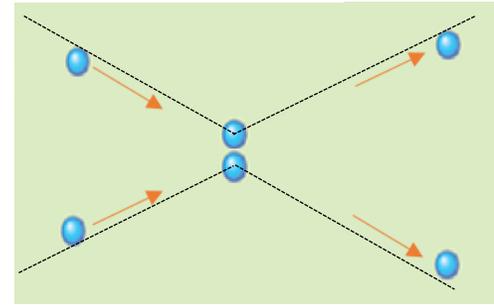


# Viscous fluid description based on momentum conservation

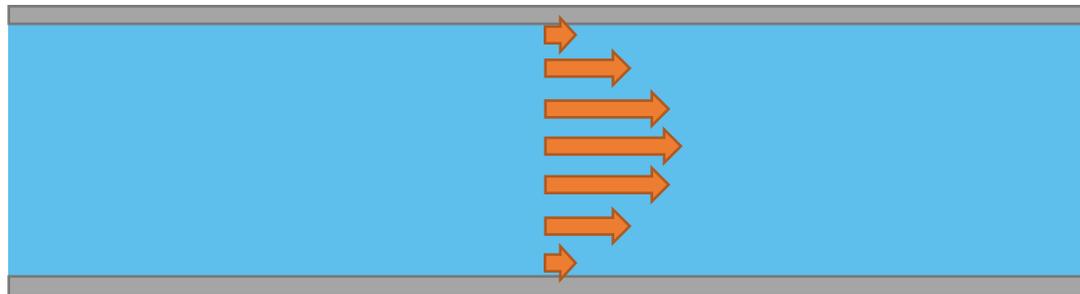
## Microscopics



$$\vec{v}_1 + \vec{v}_2 = \vec{v}'_1 + \vec{v}'_2$$



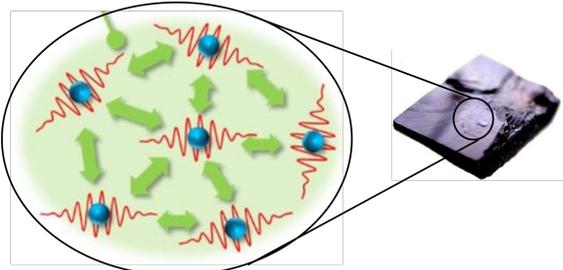
## Hydrodynamics



Navier-Stokes

$$\partial_t \vec{v} + \vec{v} \cdot \nabla \vec{v} = \eta \nabla^2 \vec{v} + \frac{\vec{f}}{m}$$

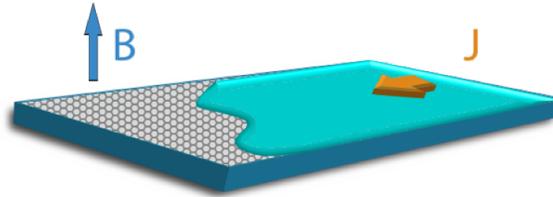
## Correlated electrons



## Hydrodynamics



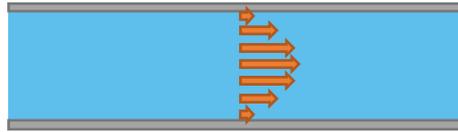
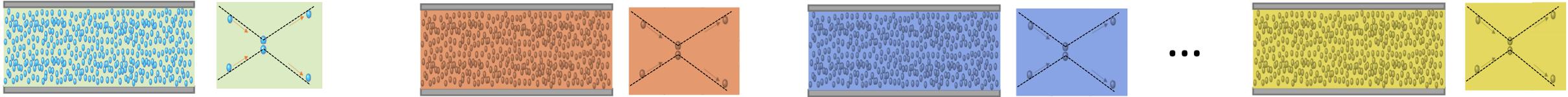
## Electron hydrodynamics



Can quantum mechanics  
*constrain* hydrodynamics?

Can quantum mechanics  
*enrich* hydrodynamics?

# Can quantum mechanics *constrain* hydrodynamics?



$$\partial_t \vec{v} + \vec{v} \cdot \nabla \vec{v} = \underbrace{\eta}_{\text{viscosity}} \nabla^2 \vec{v} + \frac{\vec{f}}{m}$$

## Classically

Viscosity can take any value

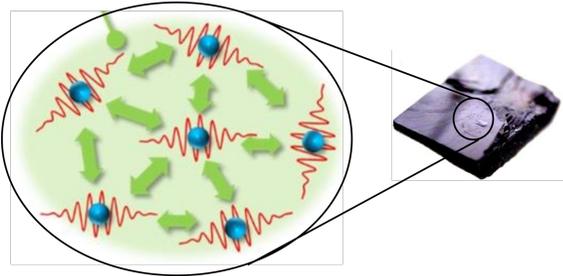
## Quantum mechanically

$$\frac{\eta}{s} \geq \frac{\hbar}{4\pi k_B}$$

[Kovtun, Son, Starinets, PRL 2005]

Conjectured fundamental bound based on holography.

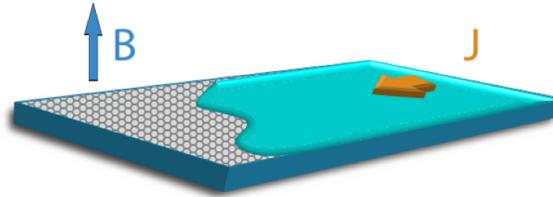
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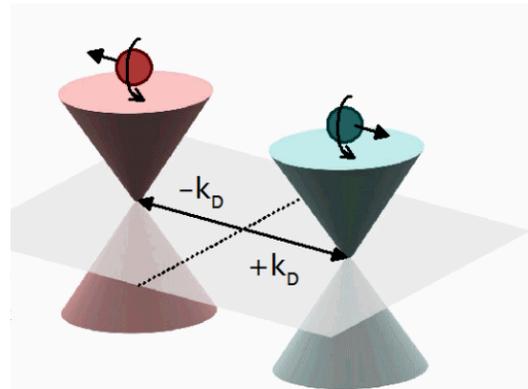
Can quantum mechanics  
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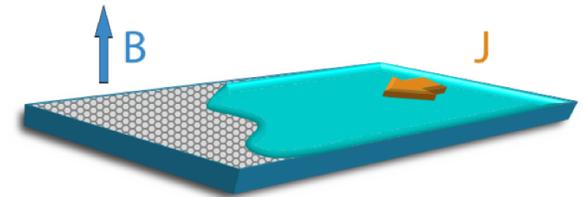
# Can quantum mechanics *enrich* hydrodynamics?

For electrons in a solid:  $H = \frac{p^2}{2m}$   $\longrightarrow$   $H(\mathbf{k})\psi(\mathbf{k}) = E(\mathbf{k})\psi(\mathbf{k})$

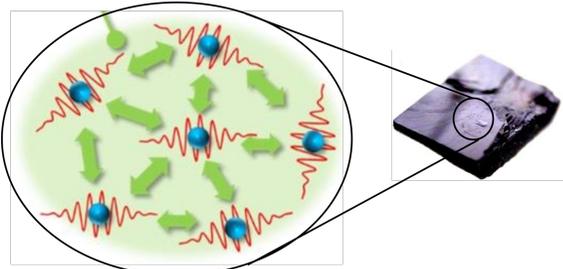
**Relativistic fluids:**  
graphene, Weyl  
semimetals,...



**Topological effects:**  
Berry phase, Hall  
viscosity,...



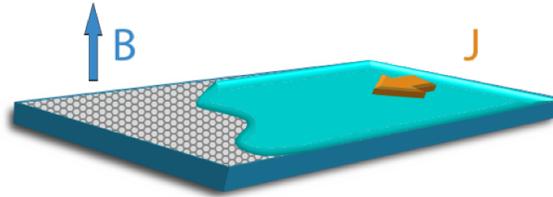
## Correlated electrons



## Hydrodynamics



## Electron hydrodynamics



Can quantum mechanics  
*constrain* hydrodynamics?

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# Outline

- Electron hydrodynamics – What and why?
- **Viscous Fermi liquids**
- Using magnetic fields to detect viscous effects
- Conclusion and outlook

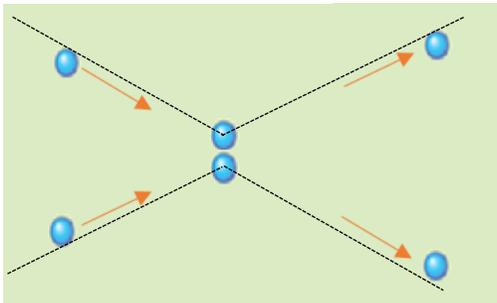
# Hydrodynamic flow of electrons possible if...

$$l_{MC} \ll W \ll l_{MR}$$

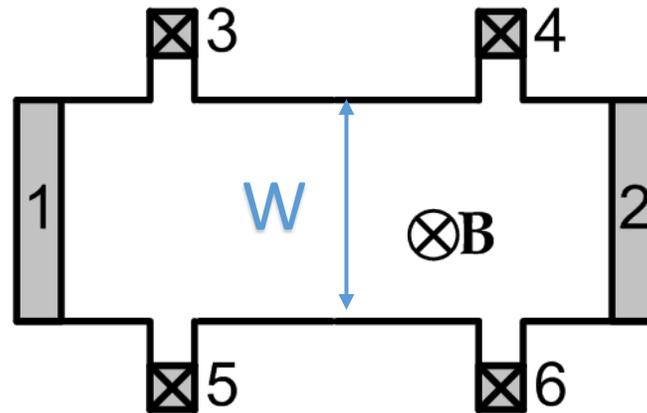
Momentum-conserving mean free path (“internal scattering”):

- e-e scattering

$$\vec{v}_1 + \vec{v}_2 = \vec{v}'_1 + \vec{v}'_2$$



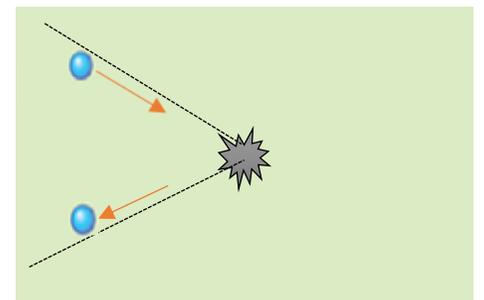
Sample size



Momentum-relaxing mean free path (“external scattering”):

- Impurities
- Phonons

$$\vec{v} \neq \vec{v}'$$



# Which materials?

$$l_{MC} \ll W \ll l_{MR}$$

Strong interactions,  
not too low T

~ 1  $\mu\text{m}$

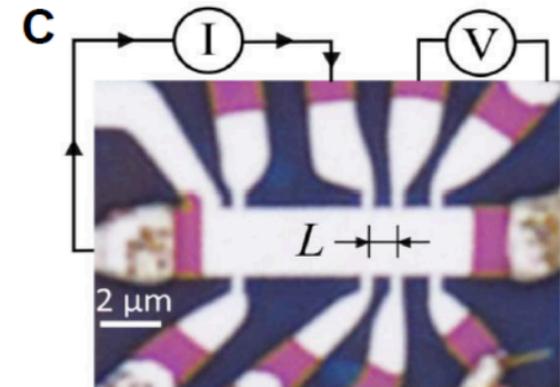
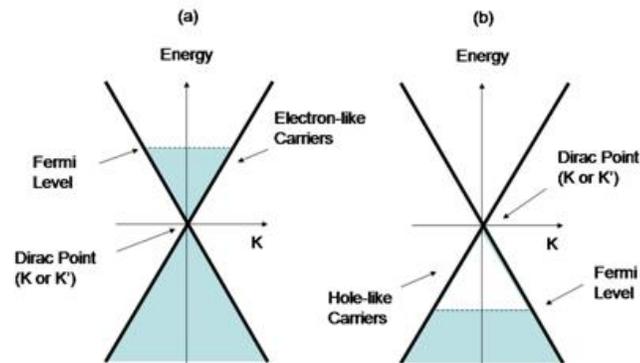
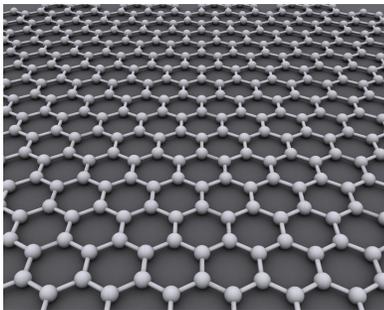
Mesoscopic samples

~ 10  $\mu\text{m}$

Clean samples,  
not too high T

~ 100  $\mu\text{m}$

- 2D electron gases [de Jong, Molenkamp] (1994)
- PdCoO<sub>2</sub> (high mobility layered metal) (2016)
- **Graphene** [Bandurin et al] [Crossno et al] (2016)

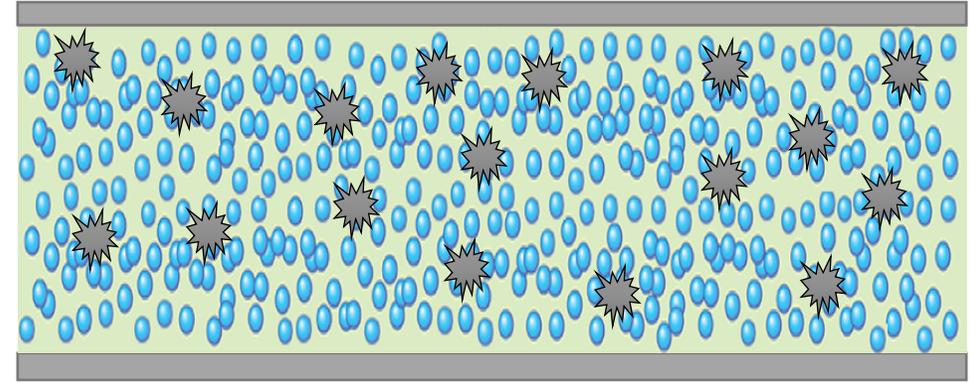


Bandurin et al, Science 351,  
1055-1058 (2016)

# Three regimes

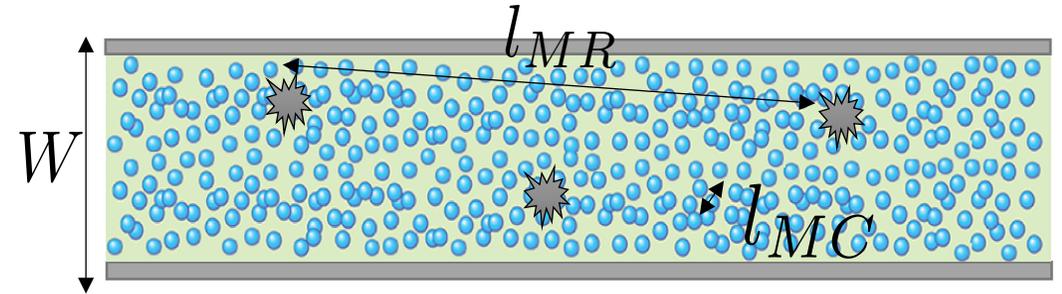
Ohmic

$$l_{MR}, l_{MC} \ll W$$



Viscous

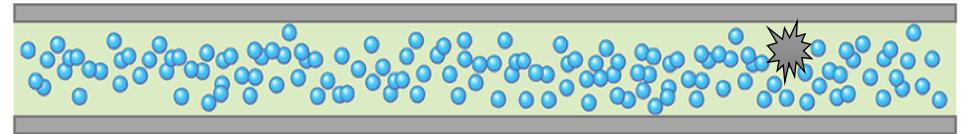
$$l_{MC} \ll W \ll l_{MR}$$



Diffuse-Ballistic

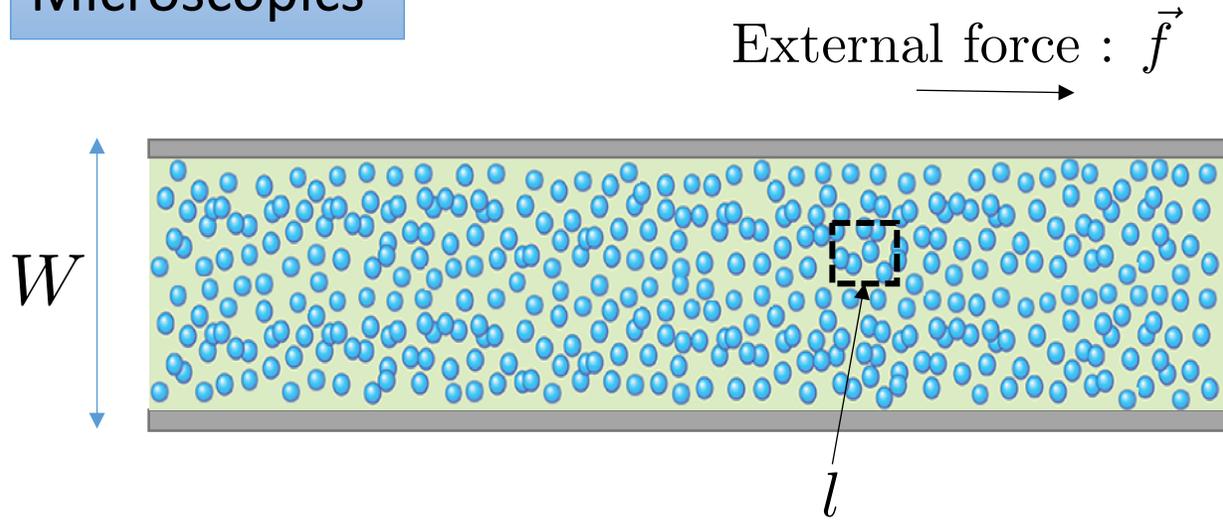
$$W \ll l_{MR}, l_{MC}$$

“Knudsen flow”

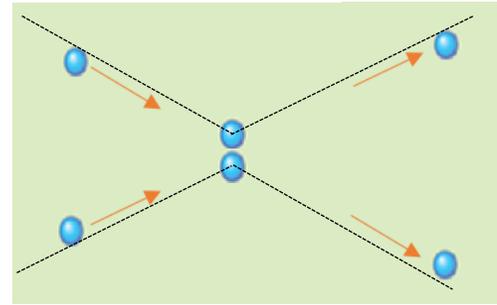


# Viscous fluid

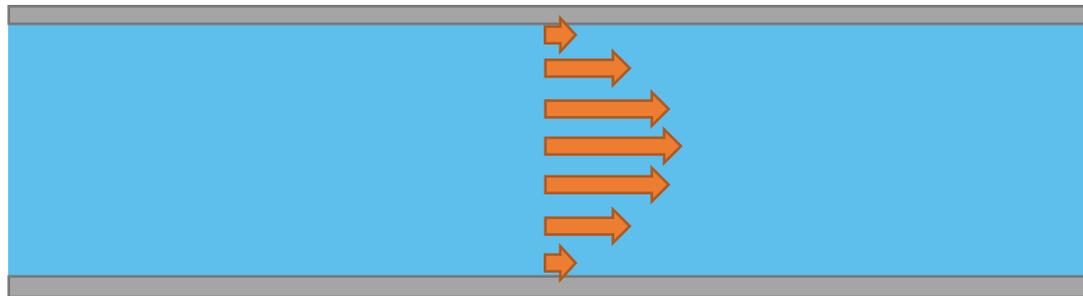
## Microscopics



$$\vec{v}_1 + \vec{v}_2 = \vec{v}'_1 + \vec{v}'_2$$



## Hydrodynamics



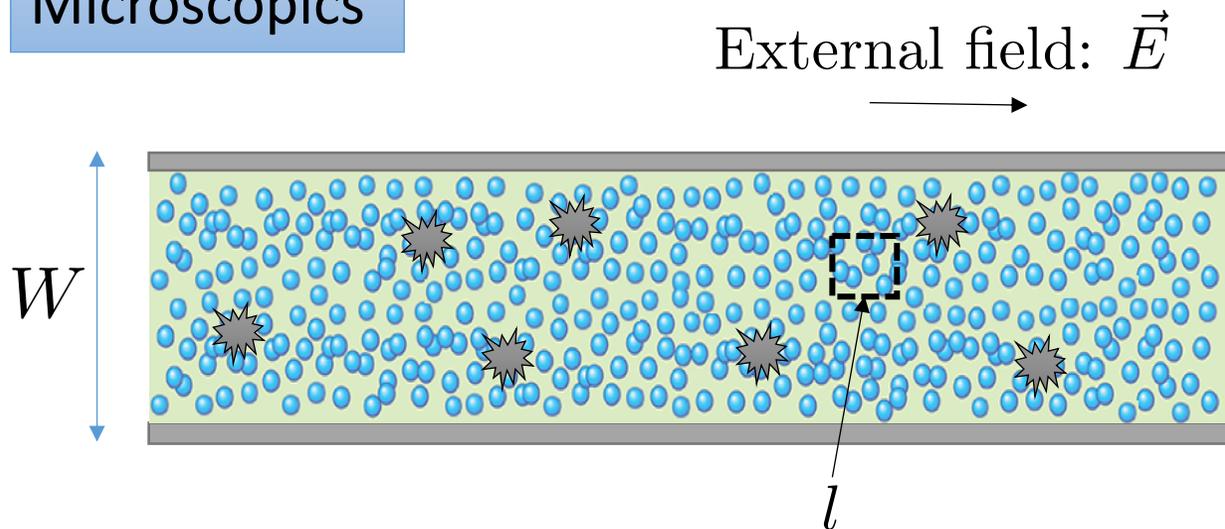
Navier-Stokes

$$\partial_t \vec{v} + \vec{v} \cdot \nabla \vec{v} = \eta \nabla^2 \vec{v} + \frac{\vec{f}}{m}$$

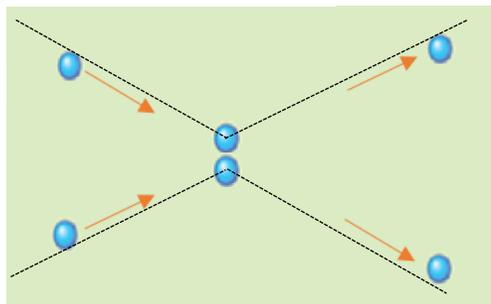
# Viscous electronic fluid

$$\vec{j} = ne\vec{v}$$

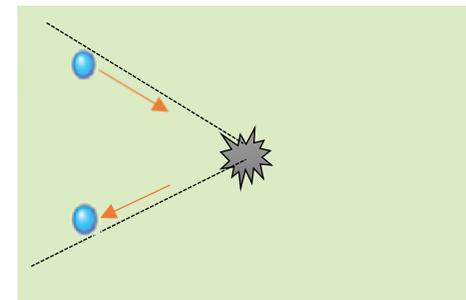
Microscopics



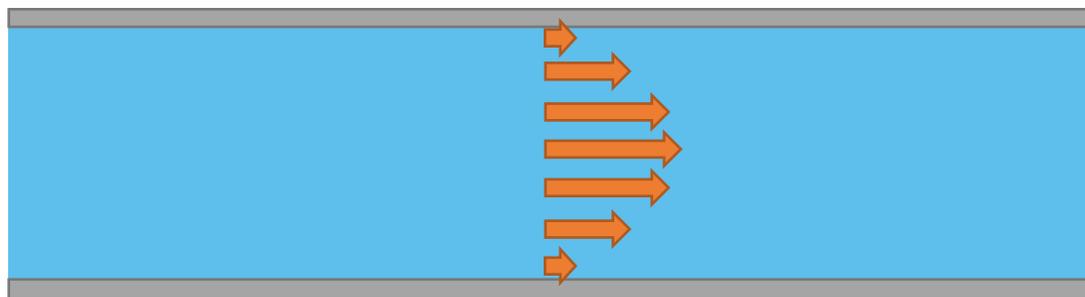
$$\vec{v}_1 + \vec{v}_2 = \vec{v}'_1 + \vec{v}'_2$$



$$\vec{v} \neq \vec{v}'$$



Hydrodynamics



$$\partial_t \vec{v} + \vec{v} \cdot \nabla \vec{v} = \eta \nabla^2 \vec{v} - \frac{1}{\tau_{MR}} \vec{v} + \frac{e\vec{E}}{m}$$

# Two regimes

$$\partial_t \vec{v} = \eta \nabla^2 \vec{v} - \frac{1}{\tau_{MR}} \vec{v} + \frac{e \vec{E}}{m}$$

$$\vec{j} = ne\vec{v}$$

If term dominates

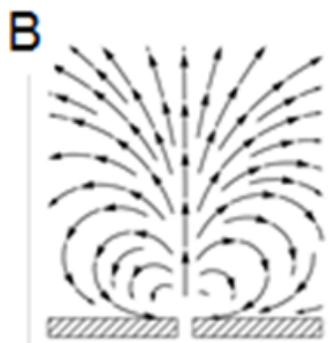
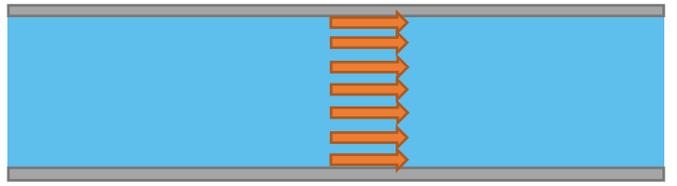
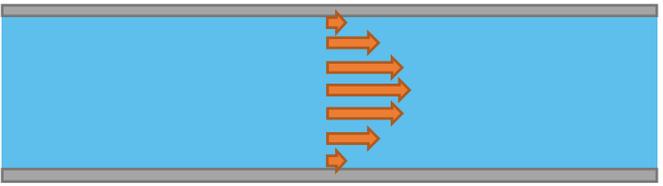
If term dominates

Viscous

Ohmic

$$\sigma = \frac{ne^2}{m} \tau_{MR}$$

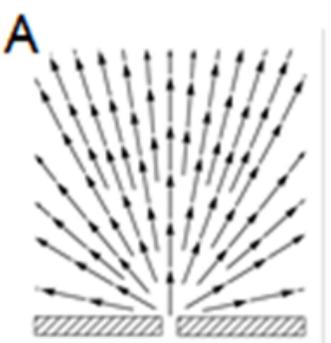
Non-local relation between current and electric field!



$$\eta \nabla^2 \vec{j} = -\frac{e^2 n}{m} \vec{E}$$

$$\nabla \times \vec{j} \neq 0$$

$$\bar{\rho} = \frac{m}{e^2 n} \eta \frac{12}{W^2}$$



$$\vec{j} = \sigma \vec{E}$$

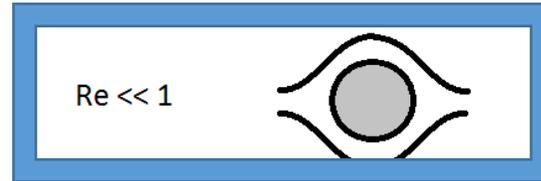
$$\nabla \times \vec{j} = 0$$

$$\bar{\rho} = \frac{m}{e^2 n} \frac{1}{\tau_{MR}}$$

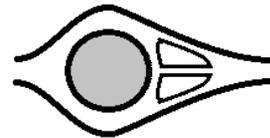
# Where did the convection term go?

$$\partial_t \vec{v} + \vec{v} \cdot \nabla \vec{v} = \eta \nabla^2 \vec{v} - \frac{1}{\tau_{MR}} \vec{v} + \frac{e\vec{E}}{m}$$

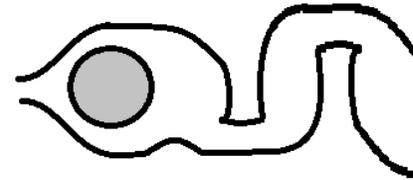
Today:



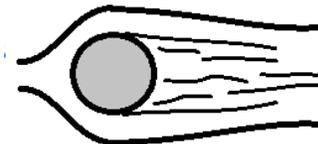
Re ~ 10



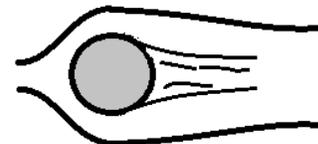
Re > ~90



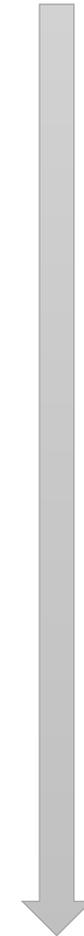
Re ~ 10<sup>4</sup> - ~ 10<sup>5</sup>



Re > ~ 10<sup>5</sup>



$$Re = \frac{\vec{v} \cdot \nabla \vec{v}}{\eta \nabla^2 \vec{v}}$$

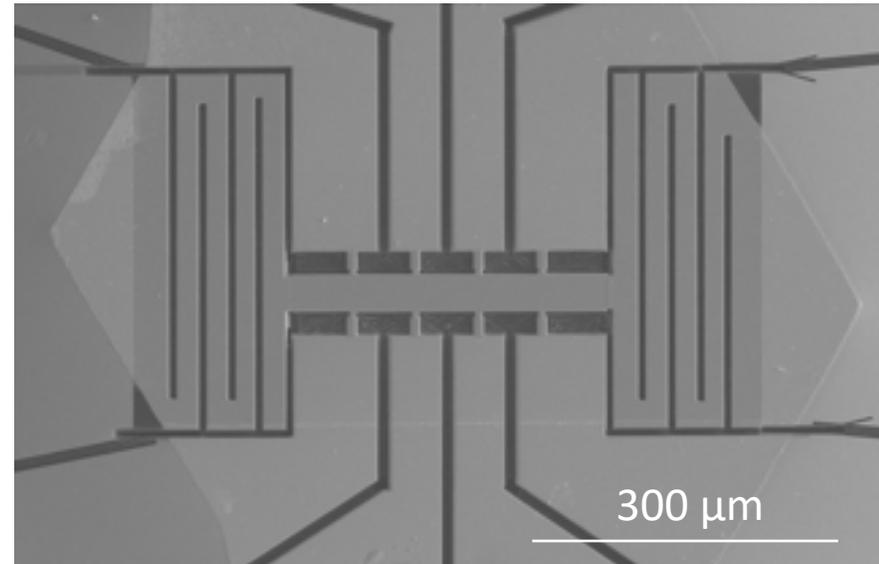
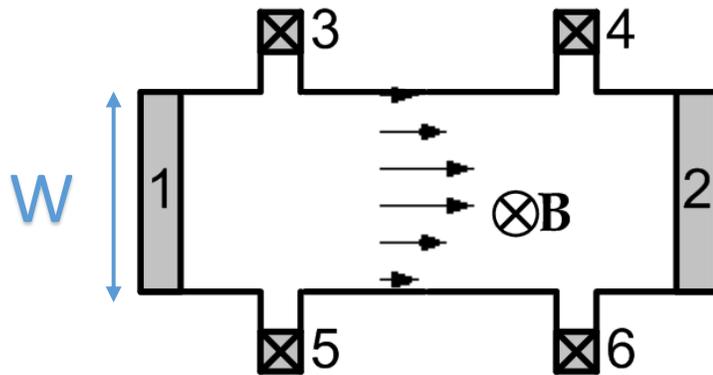


# Outline

- Electron hydrodynamics – What and why?
- Electron hydrodynamics in good metals
- Using magnetic fields to detect hydro effects
- Conclusion and outlook

## How to identify hydro effects?

- Idea: Look at finite size corrections to transport coefficients
- Problem: how to distinguish from ballistic effects?
- Solution: use magnetic field



# Navier-Stokes under magnetic field

$$\partial_t \vec{v} = \eta_{xx} \nabla^2 \vec{v} - \frac{1}{\tau_{MR}} \vec{v} + \frac{e}{m} \vec{E}$$



Add B field

$$\partial_t \vec{v} = \eta_{xx} \nabla^2 \vec{v} - \frac{1}{\tau_{MR}} \vec{v} + \frac{e}{m} (\vec{E} + \vec{v} \times \vec{B}) + \eta_{xy} \nabla^2 \vec{v} \times \hat{z}$$

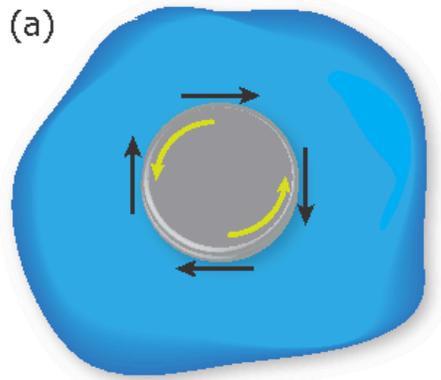
Lorentz force

Hall viscosity

# A tale of two viscosities

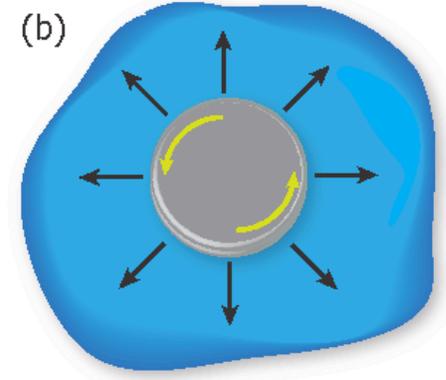
$$\partial_t \vec{v} = \eta_{xx} \nabla^2 \vec{v} - \frac{1}{\tau_{MR}} \vec{v} + \frac{e}{m} (\vec{E} + \vec{v} \times \vec{B}) + \eta_{xy} \nabla^2 \vec{v} \times \hat{z}$$

Shear viscosity



Probed by  $\rho_{xx}$

Hall viscosity

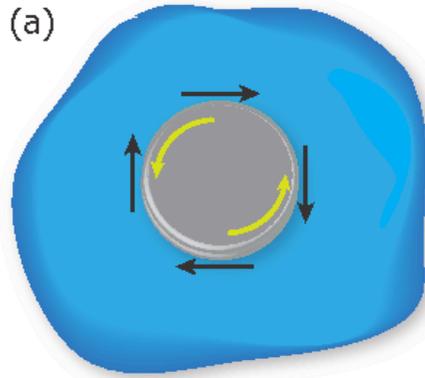


Probed by  $\rho_{xy}$

# A tale of two viscosities

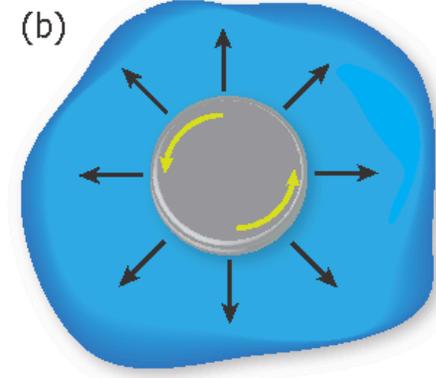
$$\partial_t \vec{v} = \eta_{xx} \nabla^2 \vec{v} - \frac{1}{\tau_{MR}} \vec{v} + \frac{e}{m} (\vec{E} + \vec{v} \times \vec{B}) + \eta_{xy} \nabla^2 \vec{v} \times \hat{z}$$

## Shear viscosity



Probed by  $\rho_{xx}$

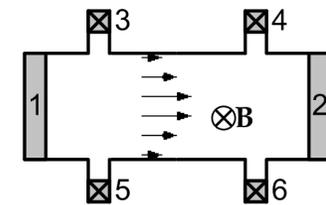
## Hall viscosity



Probed by  $\rho_{xy}$

# Hydrodynamic solution

[TS et al, PRL '17]



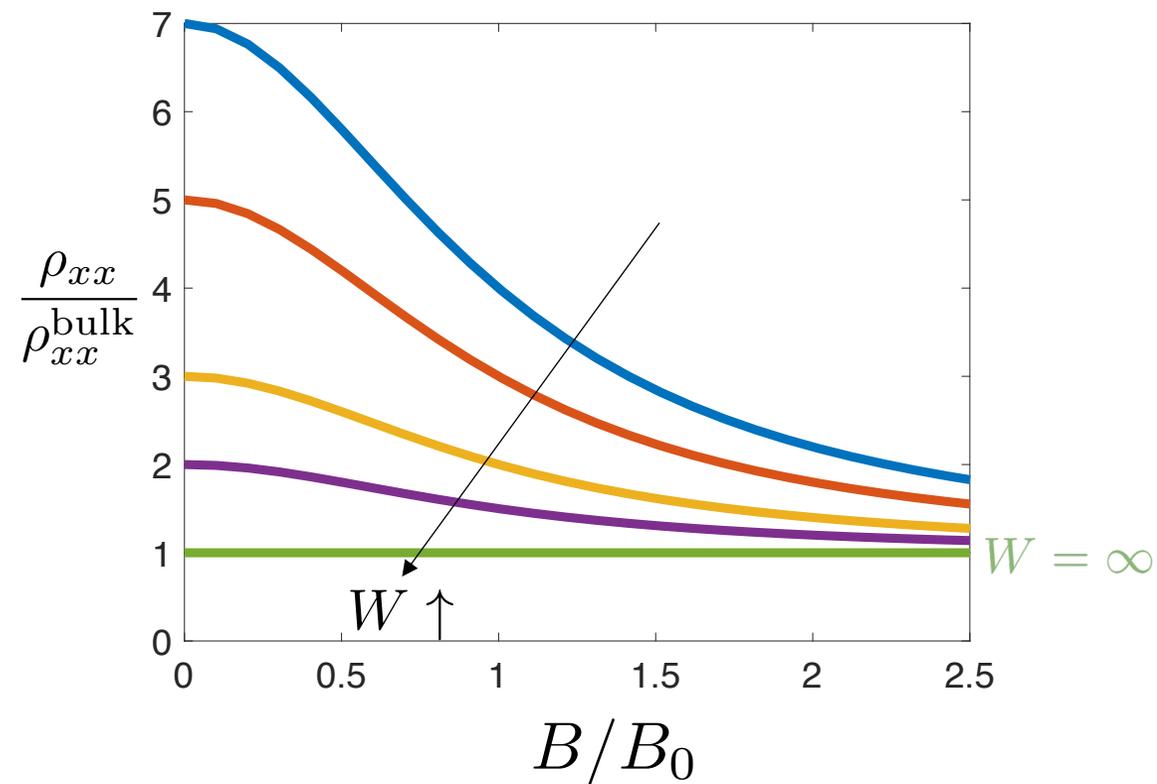
$$\partial_t \vec{v} = \eta_{xx} \nabla^2 \vec{v} + \eta_{xy} \nabla^2 \vec{v} \times \vec{z} + \frac{e}{m} (\vec{E} + \vec{v} \times \vec{B}) - \frac{1}{\tau_{MR}} \vec{v}$$

$$\rho_{xx} \simeq \rho_{xx}^{bulk} + \frac{m}{e^2 n} \eta_{xx} \frac{12}{W^2}$$

$$\rho_{xx}^{bulk} = \frac{m}{e^2 n} \frac{1}{\tau_{MR}}$$

Result from Boltzmann theory for a charged Fermi liquid:

$$\eta_{xx}(B) \sim \eta_{xx}(B=0) \frac{1}{1 + (B/B_0)^2}$$



# Navier-Stokes is not enough

In theory:

$$l_{MC} \ll W \ll l_{MR} \quad \longrightarrow \quad \text{Navier-Stokes}$$

In practice:

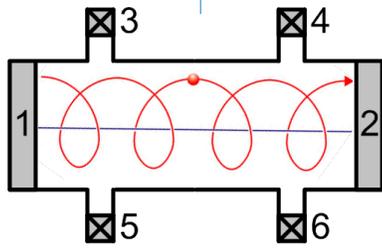
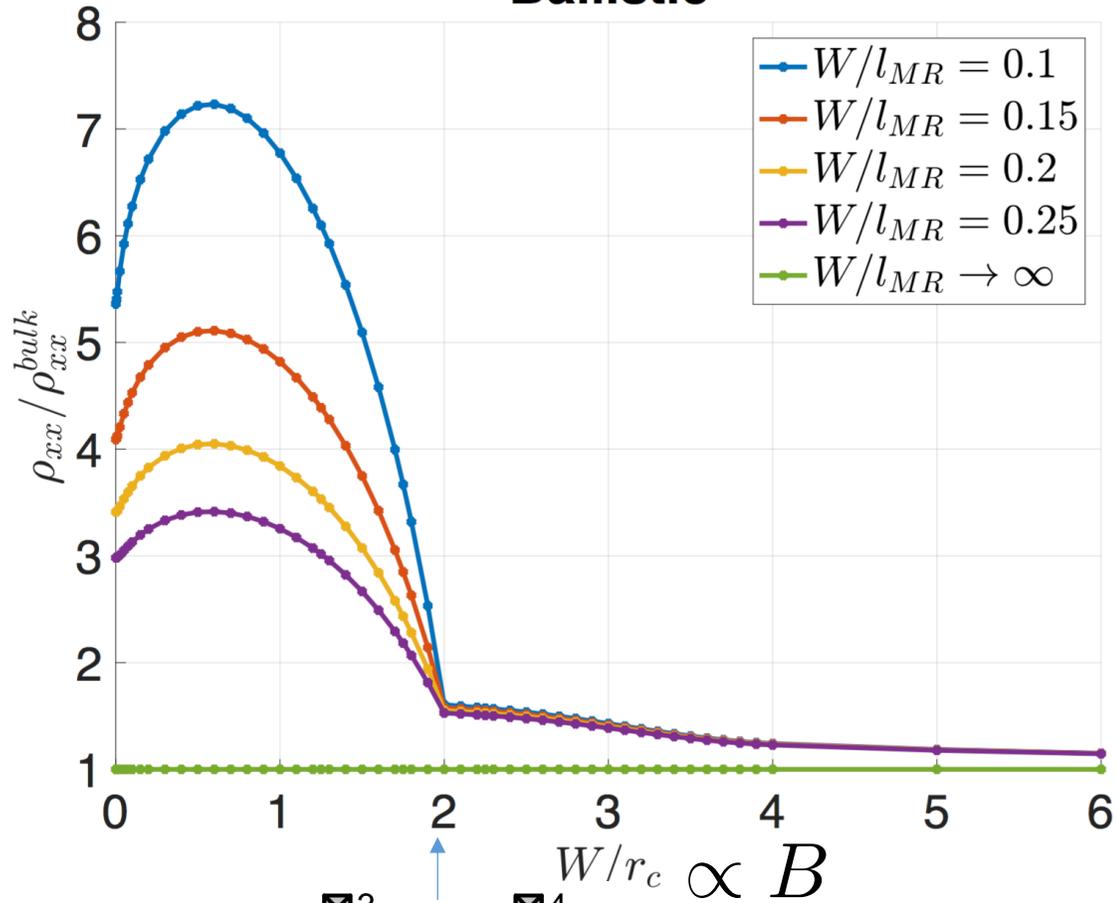
$$l_{MC} \lesssim W \lesssim l_{MR} \quad \longrightarrow \quad \text{Kinetic theory}$$

$$\partial_t f + \vec{v} \cdot \nabla_{\vec{r}} f + \frac{e}{m} (\vec{E} + \vec{v} \times \vec{B}) \cdot \nabla_{\vec{v}} f = -\frac{1}{\tau_{MR}} f - \frac{1}{\tau_{MC}} I[f]$$

# Results of Boltzmann-hydro: magnetoresistance

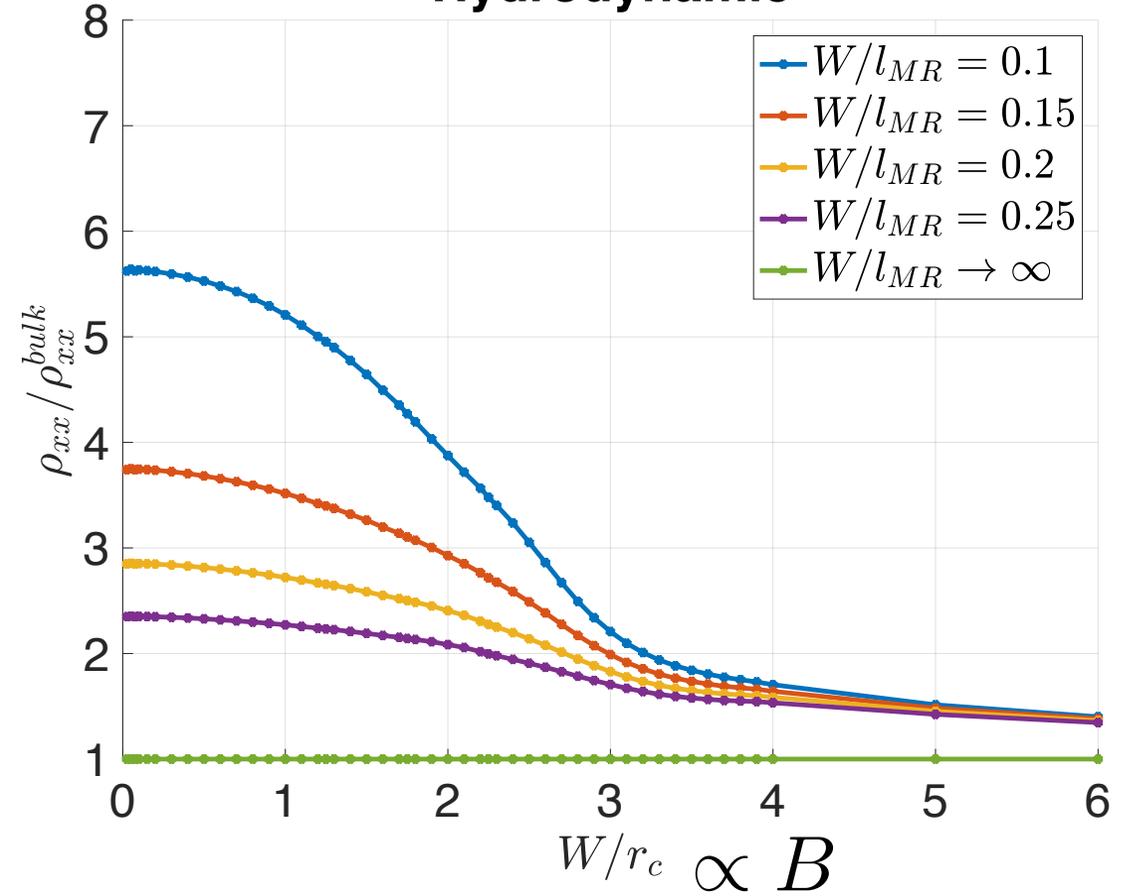
$$W \ll l_{MR}, l_{MC}$$

**Ballistic**



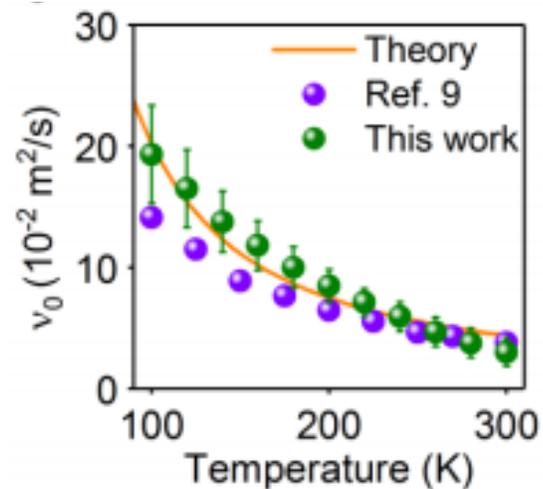
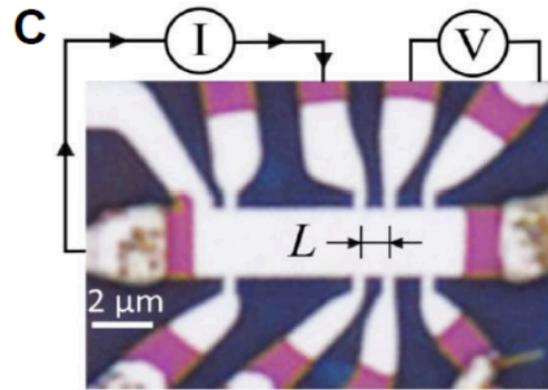
$$l_{MC} \ll W \ll l_{MR}$$

**Hydrodynamic**



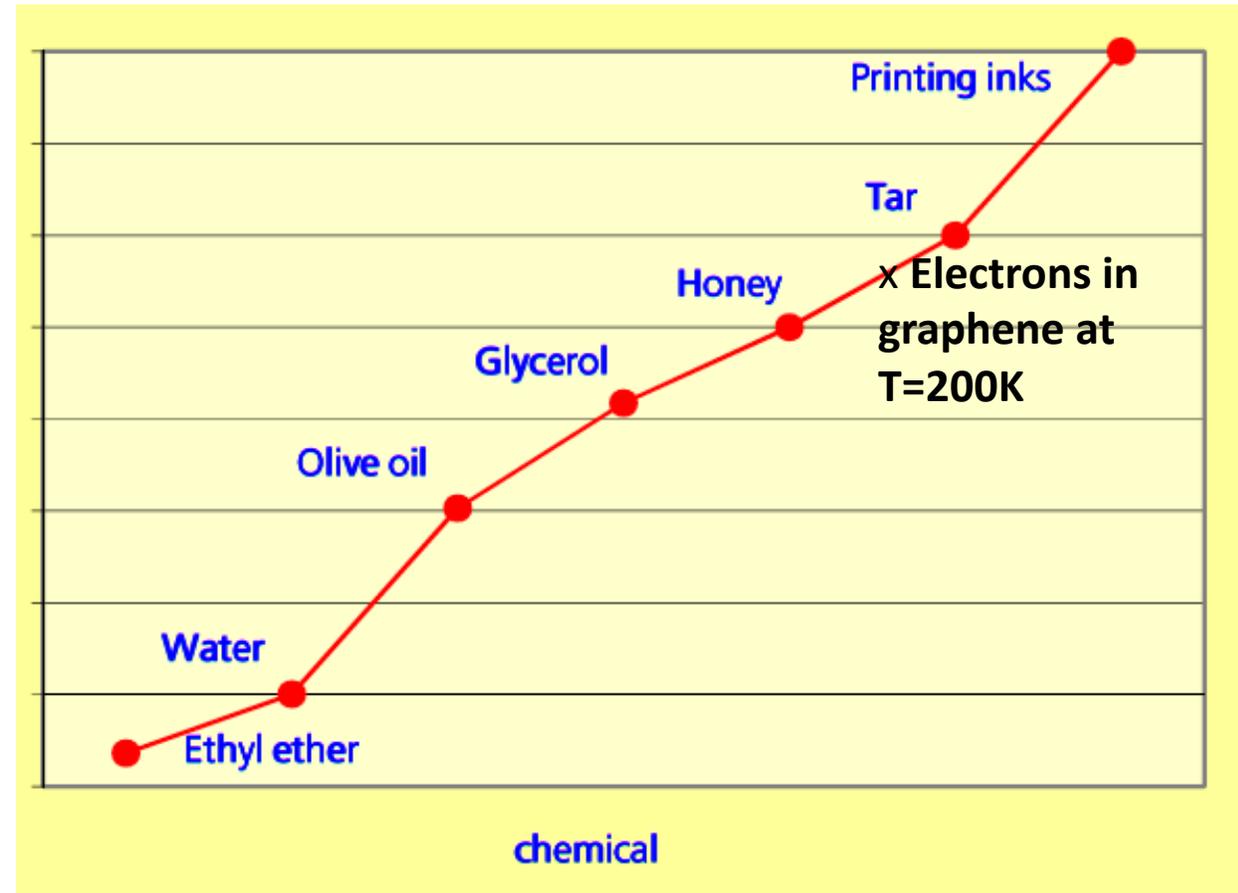
# Measurement of shear viscosity in graphene

[Bandurin et al Science (2016)] [Berdyugin et al, Science 2018]



Kinematic viscosity ( $\text{m}^2 \cdot \text{s}^{-1}$ )

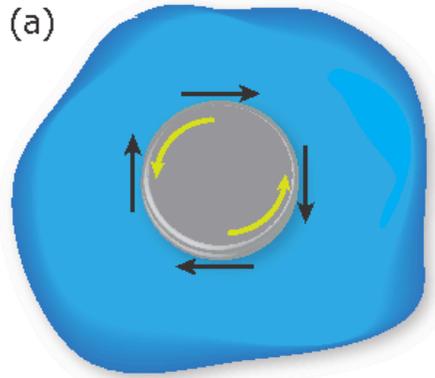
1  
 $10^{-1}$   
 $10^{-2}$   
 $10^{-3}$   
 $10^{-4}$   
 $10^{-5}$   
 $10^{-6}$



# A tale of two viscosities

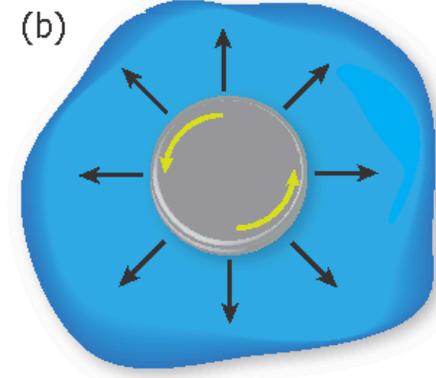
$$\partial_t \vec{v} = \eta_{xx} \nabla^2 \vec{v} - \frac{1}{\tau_{MR}} \vec{v} + \frac{e}{m} (\vec{E} + \vec{v} \times \vec{B}) + \eta_{xy} \nabla^2 \vec{v} \times \hat{z}$$

Shear viscosity



Probed by  $\rho_{xx}$

Hall viscosity



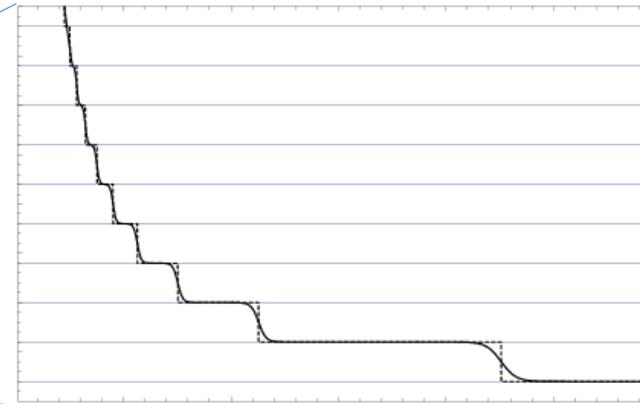
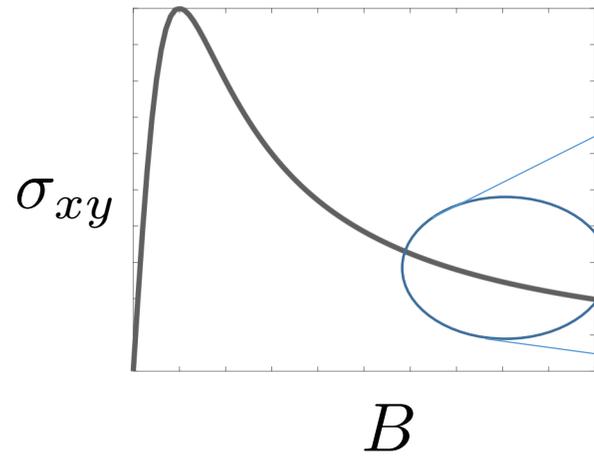
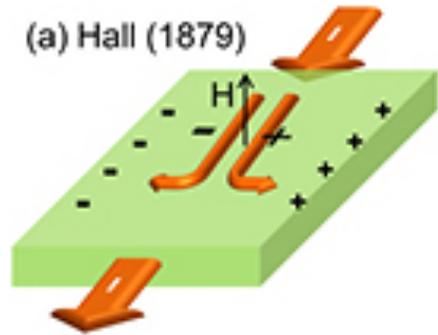
Probed by  $\rho_{xy}$

# Historical perspective on Hall effect



## Classical Hall Effect [Hall 1879]

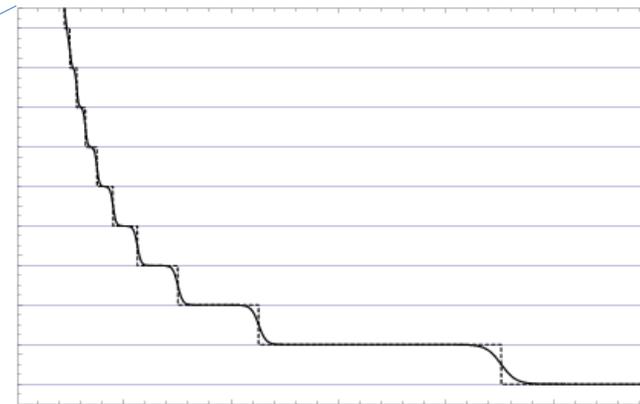
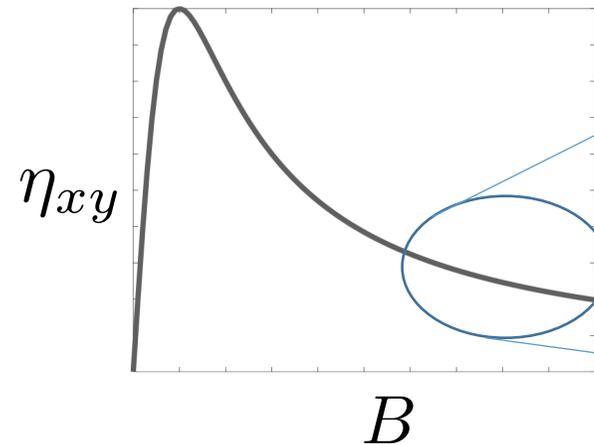
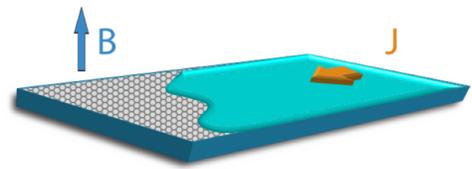
## Quantum Hall Effect [von Klitzing 1980]



$$\sigma_{xy} = \nu \frac{e^2}{h}$$

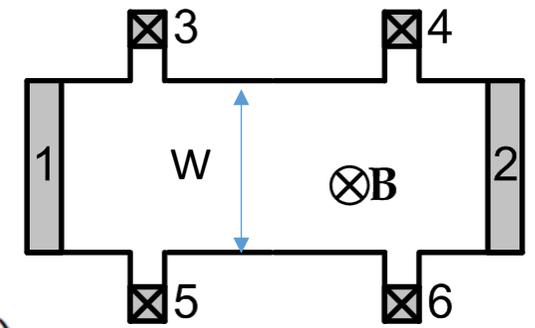
## Classical Viscous Hall Effect

## Quantum Viscous Hall Effect



$$\eta_{xy} = \nu n \frac{\hbar}{4}$$

## Hall resistivity =&gt; Hall viscosity



$$\partial_t \vec{v} = \eta_{xx} \nabla^2 \vec{v} + \eta_{xy} \nabla^2 \vec{v} \times \vec{z} + \frac{e}{m} (\vec{E} + \vec{v} \times \vec{B})$$



$$\rho_{xy} = \rho_{xy}^{bulk} \left( 1 - \eta_{xy} \frac{12}{W^2} \frac{1}{\omega_c} \right)$$

$$\rho_{xy}^{bulk} = \frac{B}{ne}$$

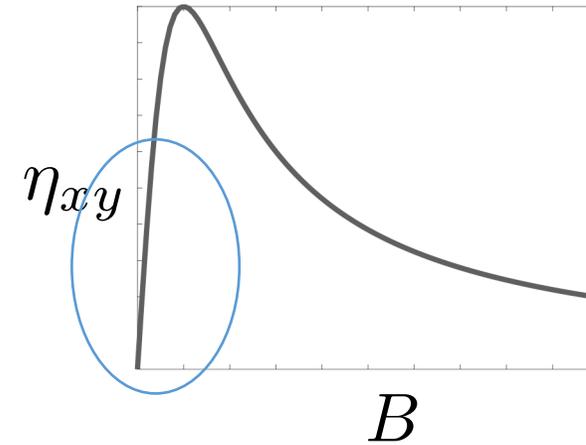
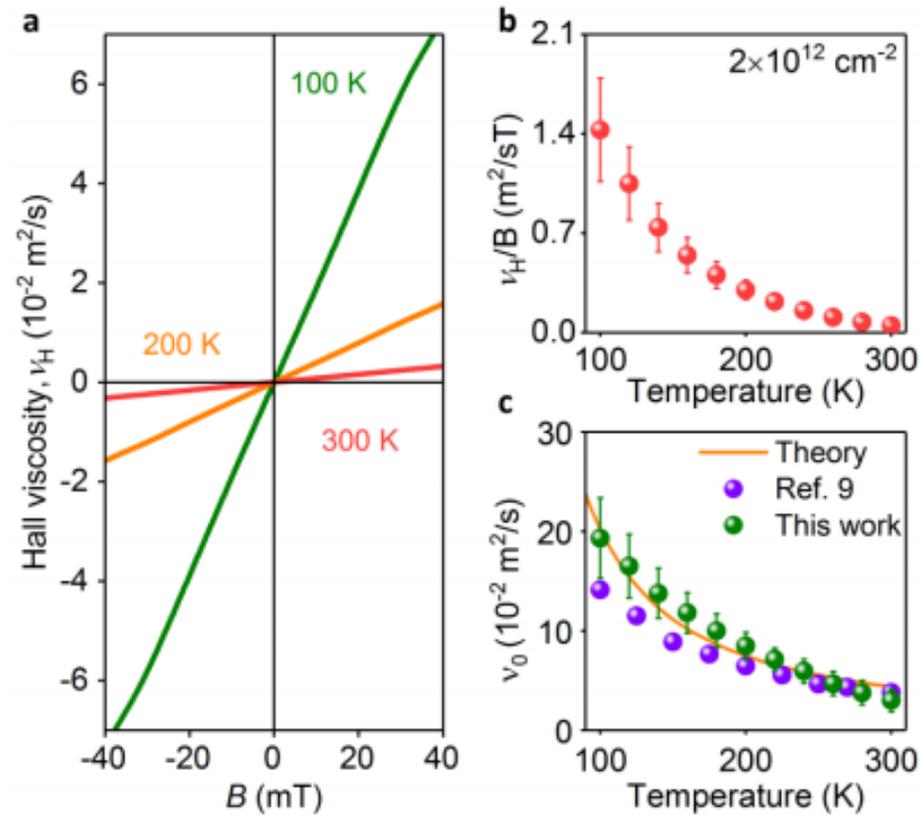


Hall viscosity can be measured by looking at finite-size effects in Hall resistivity

# Measuring Hall Viscosity of Graphene's Electron Fluid Science 2018

A. I. Berdyugin, S. G. Xu, F. M. D. Pellegrino, R. Krishna Kumar, A. Principi, I. Torre, M. Ben Shalom, T. Taniguchi, K. Watanabe, I. V. Grigorieva, M. Polini, A. K. Geim, D. A. Bandurin

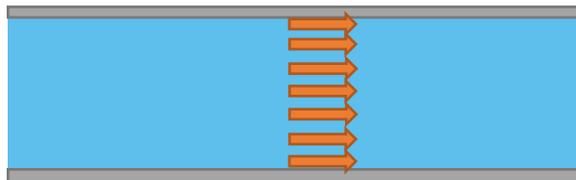
(Submitted on 5 Jun 2018)



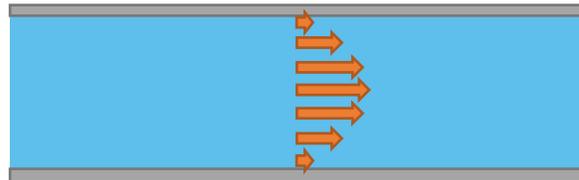
# Last part of the talk: local properties

- So far we've only discussed quantities that are averaged over the cross section of the channel
- What about their spatial dependence? Any smoking gun features of hydro?
- Naïvely, yes:

Ohmic



Hydro

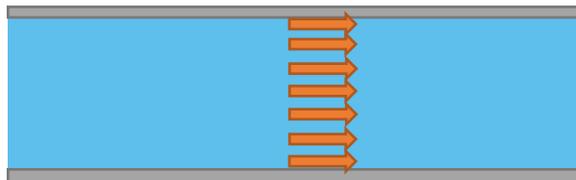


# Last part of the talk: local properties

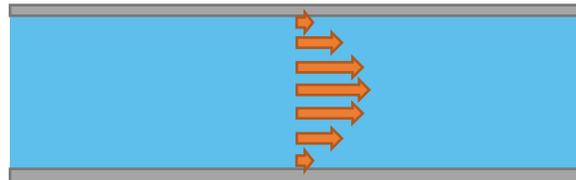
- So far we've only discussed quantities that are averaged over the cross section of the channel
- What about their spatial dependence? Any smoking gun features of hydro?
- Naïvely, yes:

But:

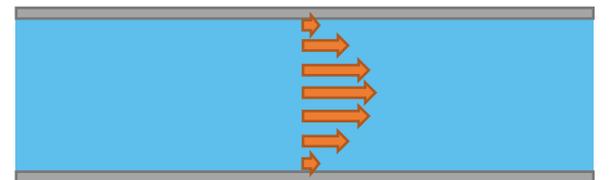
Ohmic



Hydro



Ballistic

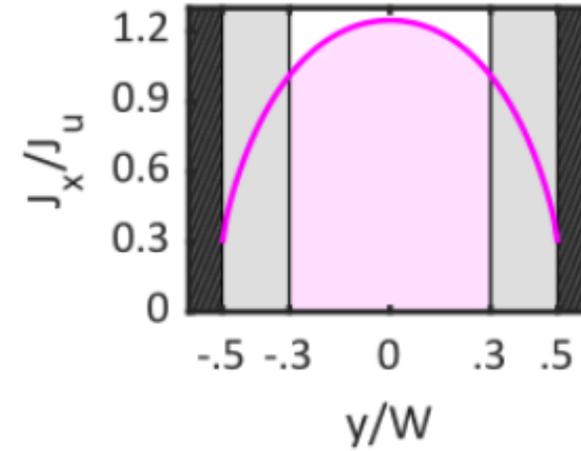
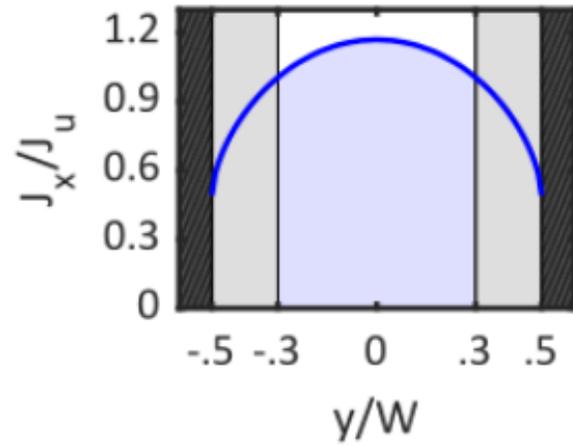


Magnetic fields are useful again:

Ballistic

Hydro

Current density

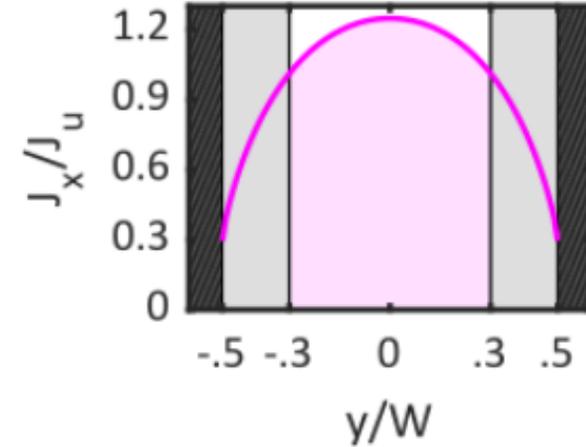
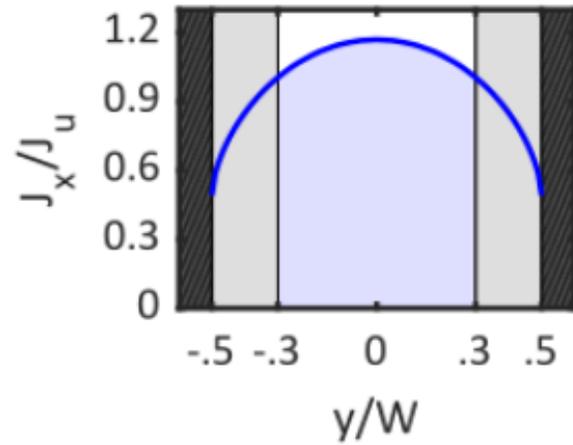


Magnetic fields are useful again:

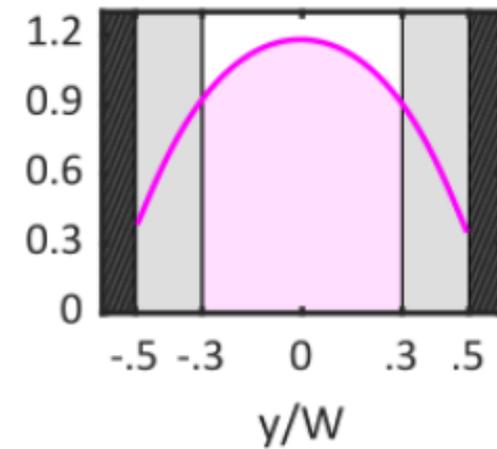
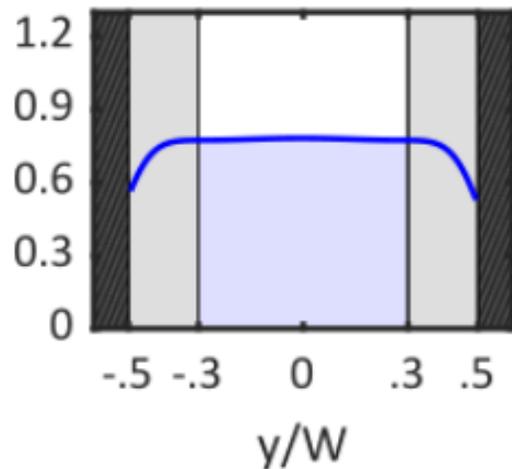
Ballistic

Hydro

Current density



Hall electric field

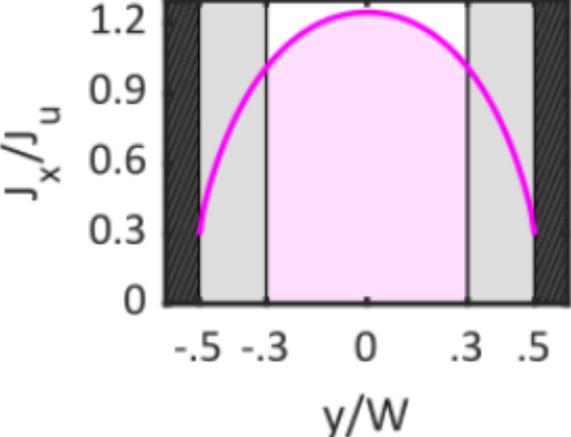
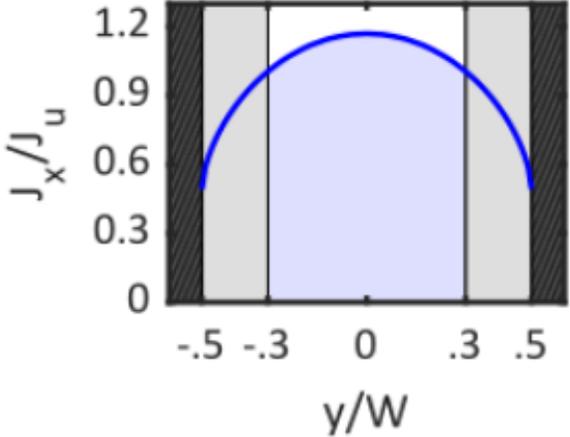


Magnetic fields are useful again:

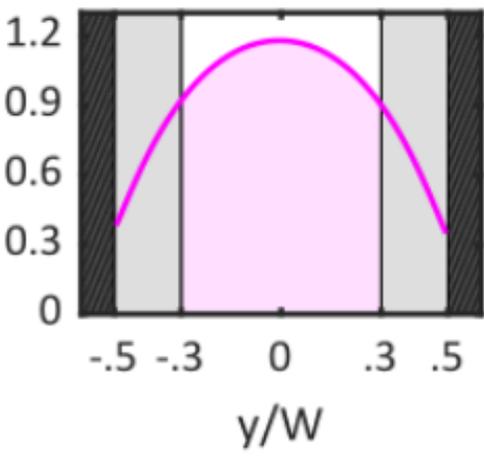
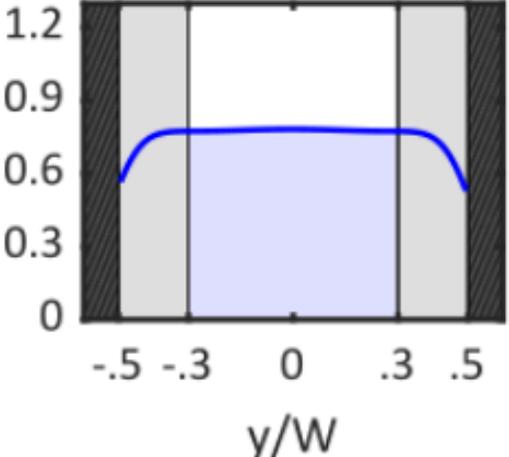
Ballistic

Hydro

Current density



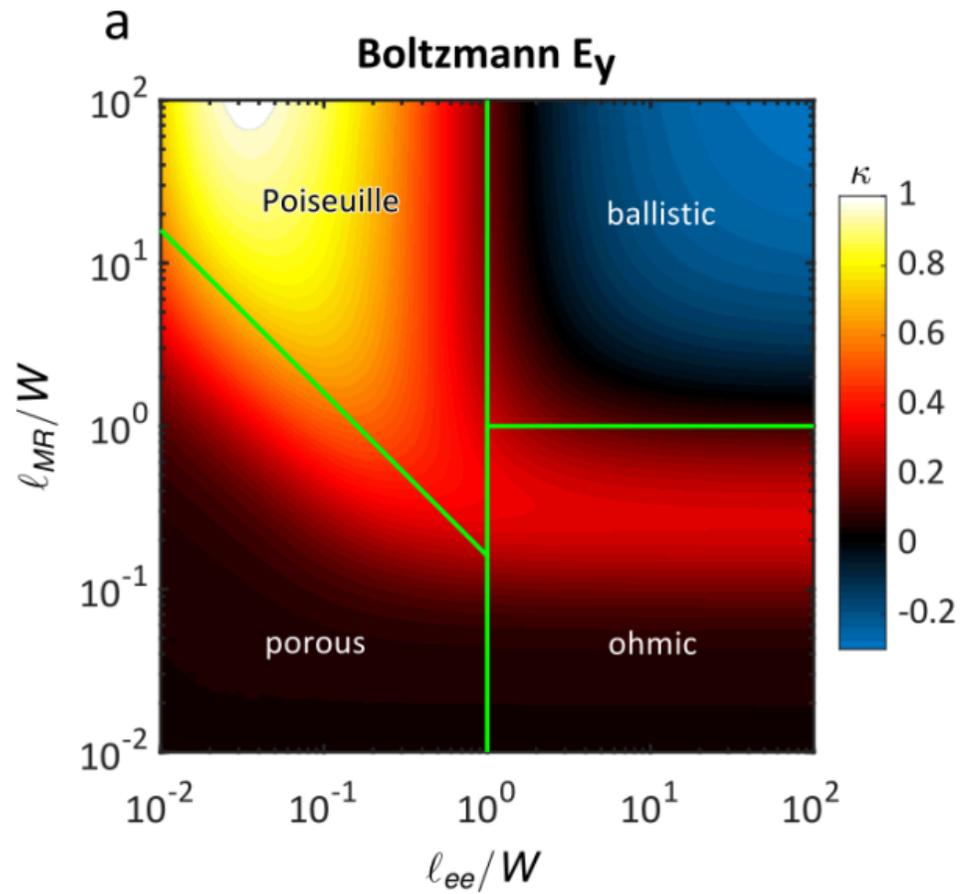
Hall electric field



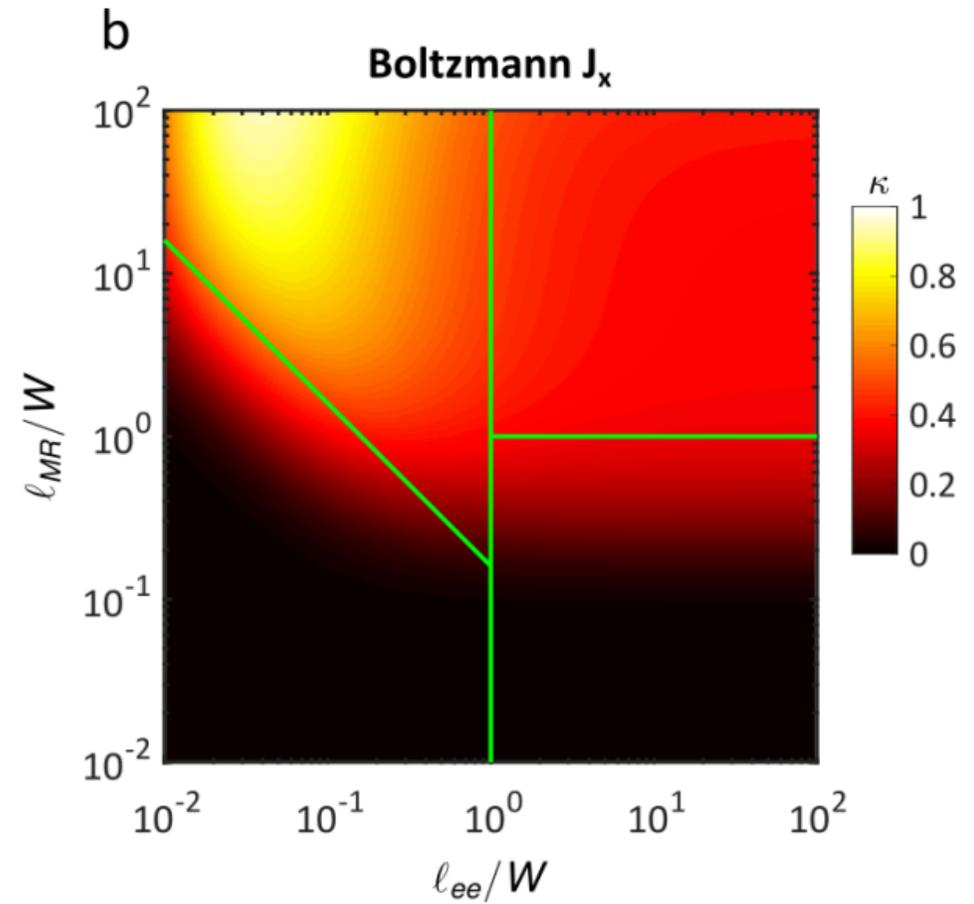
$$E_y = \frac{B}{ne} \left( j_x + \frac{1}{2} l_{ee}^2 \partial_y^2 j_x \right)$$

# “Phase diagram” based on profile curvature $\kappa$

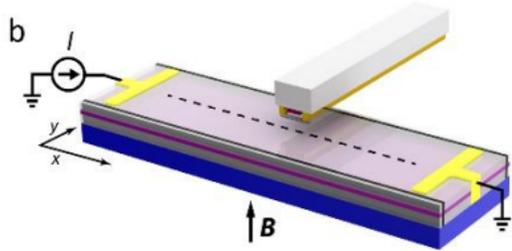
## Hall electric field



## Current

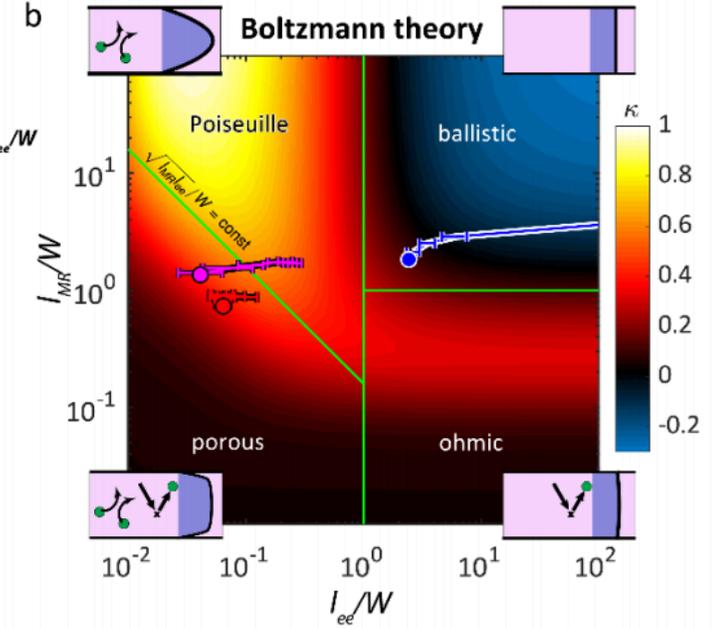
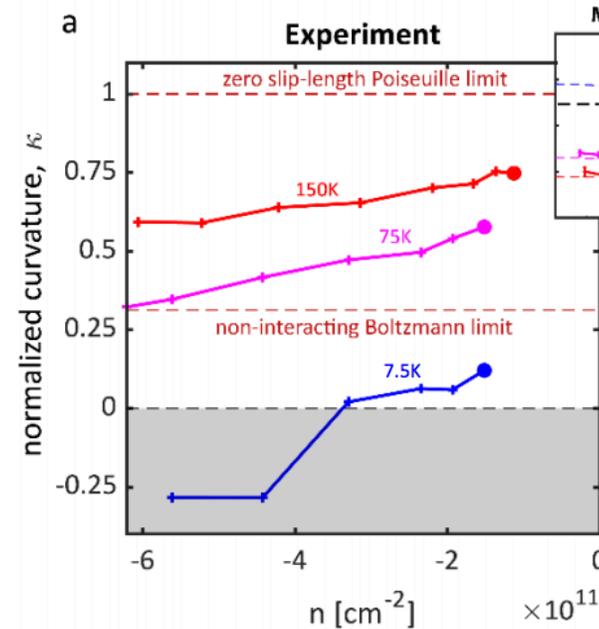
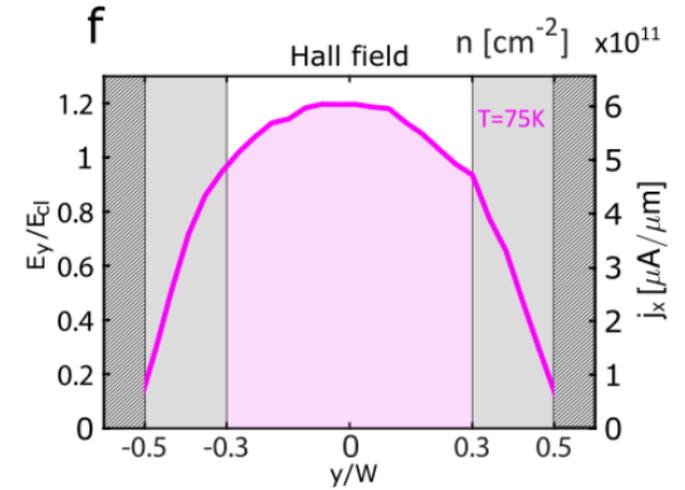
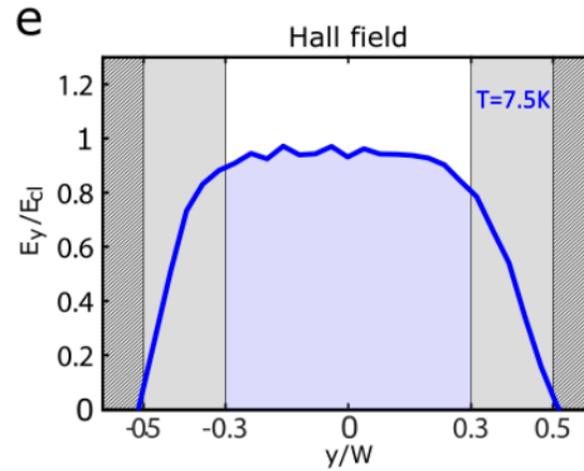


# First visualization of hydrodynamic flow of electrons



Single-electron transistor

[J.A. Sulpizio<sup>1†</sup>, L. Ella<sup>1†</sup>, A. Rozen<sup>1†</sup>, J. Birkbeck<sup>2,3</sup>, D.J. Perello<sup>2,3</sup>, D. Dutta<sup>1</sup>, M. Ben-Shalom<sup>2,3,4</sup>, T. Taniguchi<sup>5</sup>, K. Watanabe<sup>5</sup>, T. Holder<sup>1</sup>, R. Queiroz<sup>1</sup>, A. Stern<sup>1</sup>, TS, A.K. Geim<sup>2,3</sup>, and S. Ilani, [arXiv:1905.11662](https://arxiv.org/abs/1905.11662), to appear in Nature (2019) ]



# References

[TS, Nandi, Schmidt, Mackenzie, and Moore, **Phys. Rev. Lett.** **118**, 226601]

[Holder, Queiroz, TS, Silberstein, Rozen, Sulpizio, Ella, Ilani, Stern, **arXiv:1901.08546**]

[J.A. Sulpizio<sup>1†</sup>, L. Ella<sup>1†</sup>, A. Rozen<sup>1†</sup>, J. Birkbeck<sup>2,3</sup>, D.J. Perello<sup>2,3</sup>, D. Dutta<sup>1</sup>, M. Ben-Shalom<sup>2,3,4</sup>, T. Taniguchi<sup>5</sup>, K. Watanabe<sup>5</sup>, T. Holder<sup>1</sup>, R. Queiroz<sup>1</sup>, A. Stern<sup>1</sup>, TS, A.K. Geim<sup>2,3</sup>, and S. Ilani, **arXiv:1905.11662**, to appear in Nature (2019) ]

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- UC Berkeley: Joel Moore
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- Weizmann institute: Tobias Holder, Raquel Queiroz, Navot Silberstein, Asaf Rozen, Joseph A. Sulpizio, Lior Ella, Shahal Ilani, Ady Stern

Thanks for your attention!



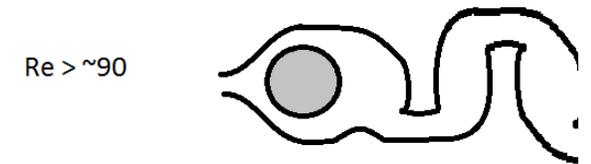
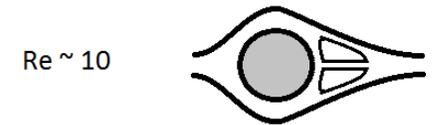
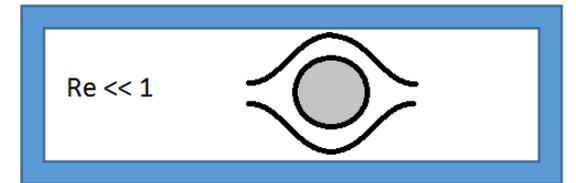
UNIVERSITY OF  
TORONTO

GORDON AND BETTY  
MOORE  
FOUNDATION

# Reynolds number?

$$\partial_t \vec{v} + \vec{v} \cdot \nabla \vec{v} = \eta \nabla^2 \vec{v} - \frac{1}{\tau_{MR}} \vec{v} + \frac{e\vec{E}}{m}$$

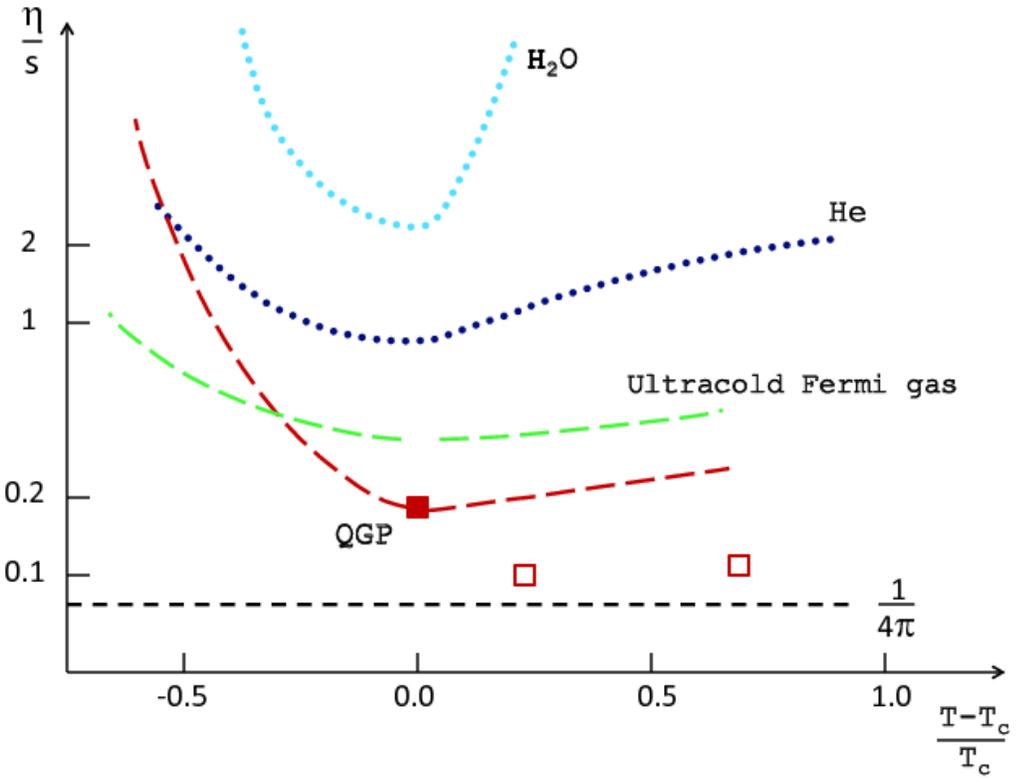
$$Re = \frac{v_{\text{drift}} W}{\eta} \leq \frac{v_{\text{drift}}}{v_F} \times \frac{l_{e-imp}}{l_{ee}} \ll 1$$



# Bound on viscosity

$$\frac{\eta}{s} \geq \frac{1}{4\pi} \frac{\hbar}{k_B}$$

- Three arguments
  - Class of strongly interacting QFTs with gravity dual saturate this bound
  - Empirical evidence:
    - Extrapolation of kinetic theory

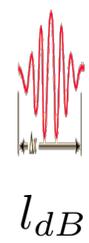
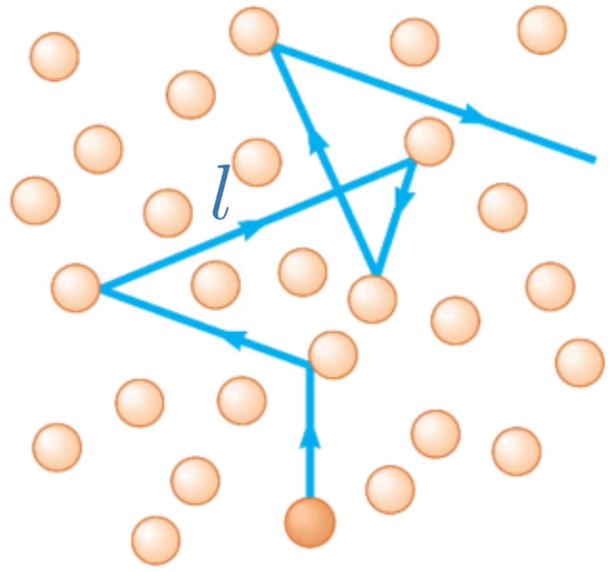


# Bound on viscosity

$$\frac{\eta}{s} \sim \frac{mv l}{k_B} \sim \frac{\hbar}{k_B} \frac{l}{l_{dB}}$$

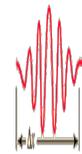
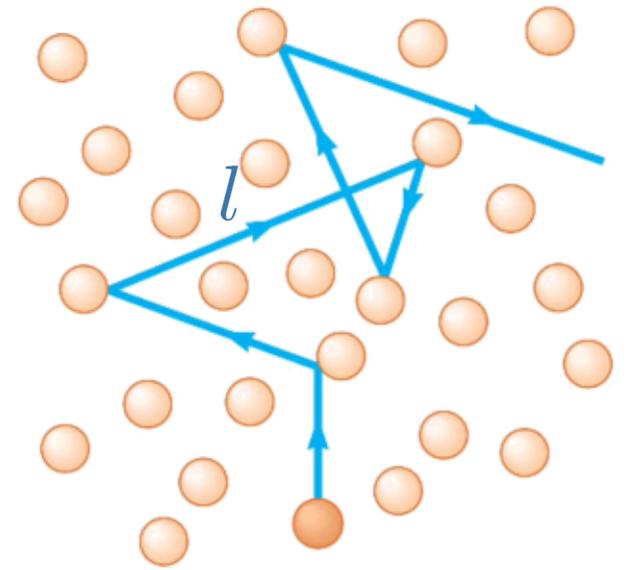
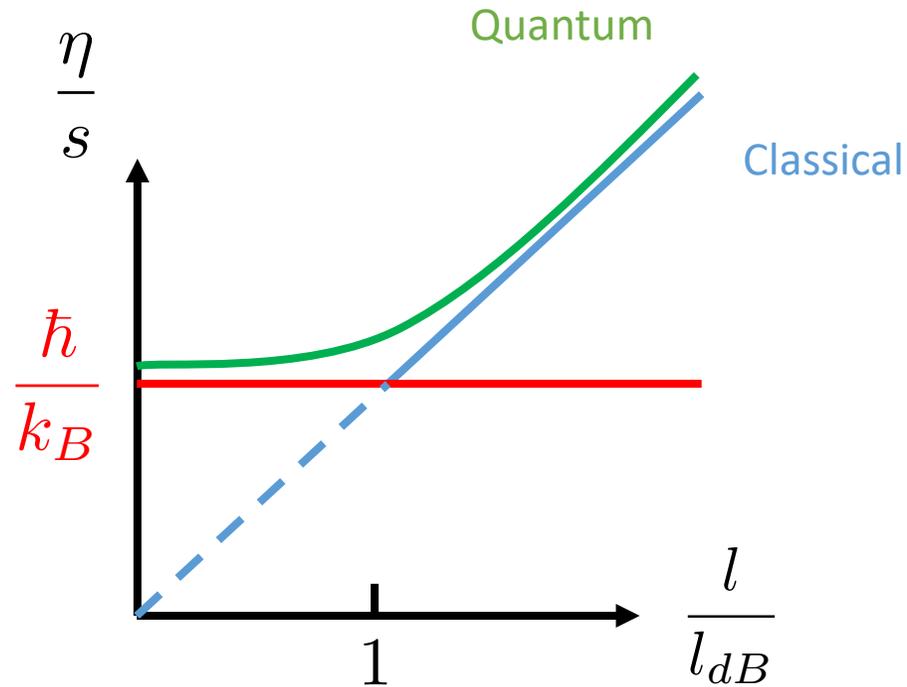
$$\eta \sim nmvl$$
$$s \sim k_B n$$

$$p = mv \sim \frac{\hbar}{l_{dB}}$$



# Bound on viscosity

$$\frac{\eta}{s} \sim \frac{\hbar}{k_B} \frac{l}{l_{dB}}$$



$l_{dB}$