

Phase Diagram of QCD

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QCD – the theory of strong interactions

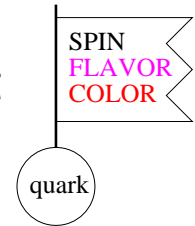
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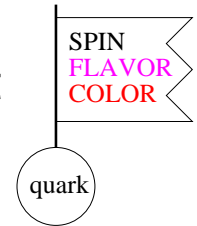
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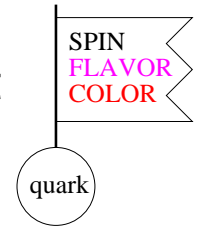
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- What force holds quarks together in hadrons? It weakens if the interaction is brief or if quarks are close to each other ($r \ll R_{\text{proton}} \sim 10^{-13}$ cm \equiv 1 fermi).



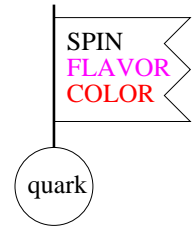
$$\alpha_{\text{strong}}(r) = \frac{Fr^2}{\hbar c} \xrightarrow{r \rightarrow 0} 0$$

Compare to E.M.: $\alpha = 1/137$ (and only *grows* as $r \rightarrow 0$).

In contrast $\alpha_{\text{strong}}(1\text{fermi}) \sim 1$ and $\rightarrow 0$ as $r \rightarrow 0$.

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ASYMPTOTIC FREEDOM

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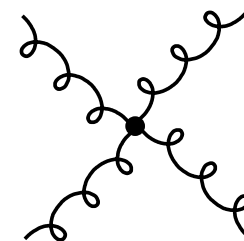
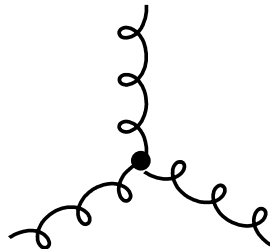
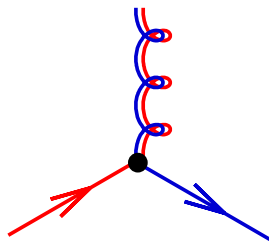
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QCD – asymptotically free field theory



- Only one such class of theories: Non-abelian gauge theories. (QED is abelian gauge theory.)
- Bonus: it involves a hidden symmetry (in QED – U(1) – phase rotation). In QCD it is the **color** symmetry – SU(3).
- QCD is a quantum field theory:

$$S = \int d^4x \left[\sum_{f=1}^{N_f} \bar{q}_f (i\not{\partial} + g\not{A} - m_f) q_f - \frac{1}{2} \text{Tr} F_{\mu\nu} F^{\mu\nu} \right];$$



- Predictive power

QCD thermodynamics

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 - Neutron stars (large density, low T)
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 - Neutron stars (large density, low T)
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- QCD allows first principle calculations
- Questions: phases, phase diagram, as function of T, μ_B, \dots

● Early expectations \Rightarrow

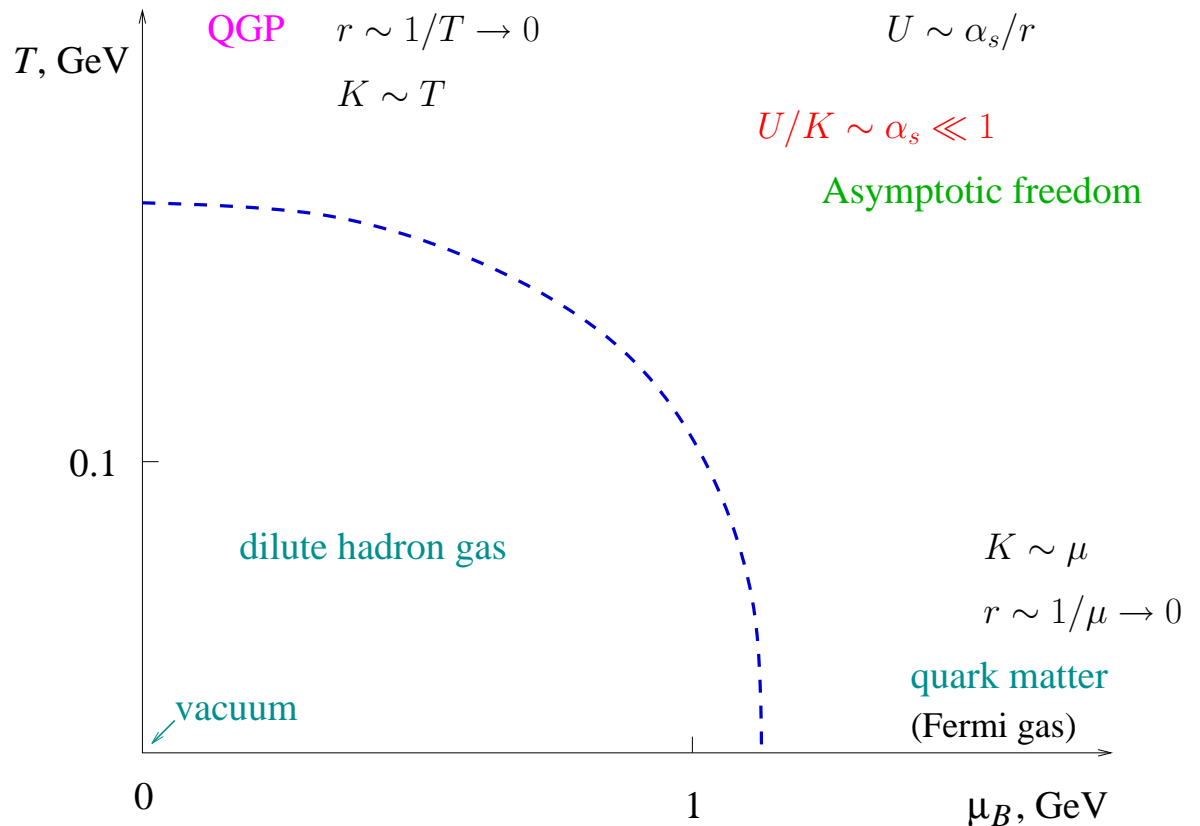
Natural scale:

$$kT \sim \frac{\hbar c}{1\text{fm}^2} = 0.2 \text{ GeV}.$$

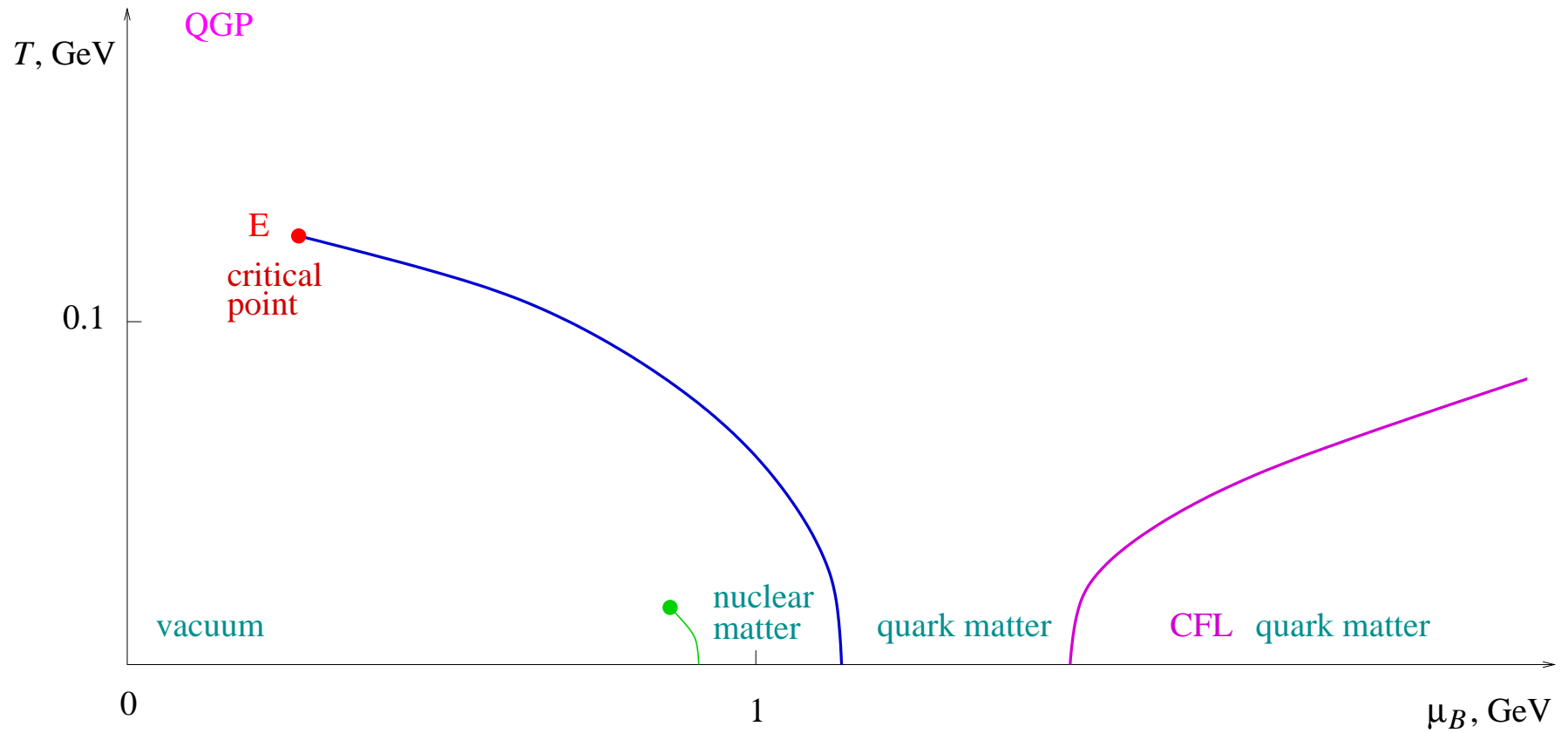
$$(T \sim 10^{12} \text{ K})$$

or

$$\rho_B \sim 1\text{fm}^{-3}.$$

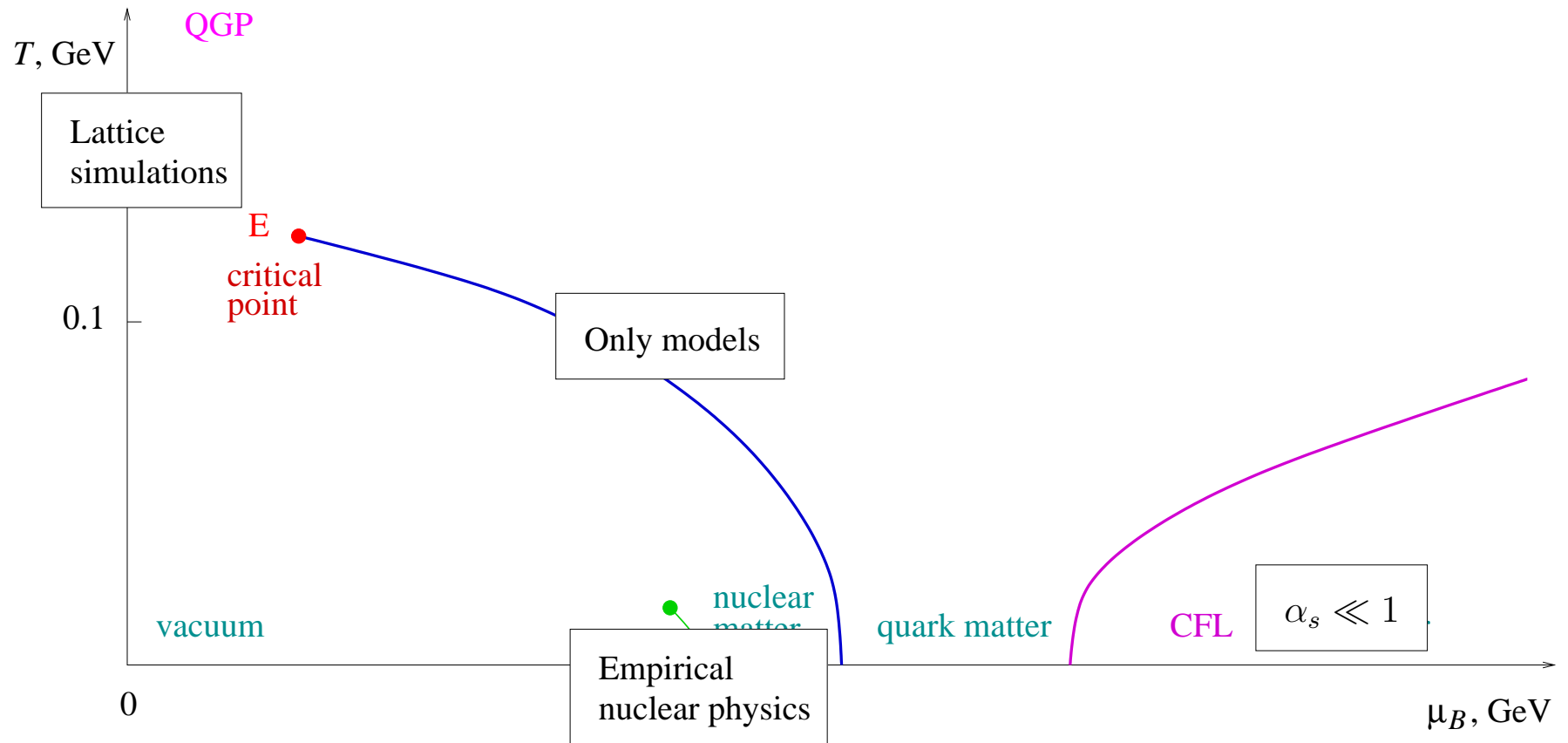


Contemporary view



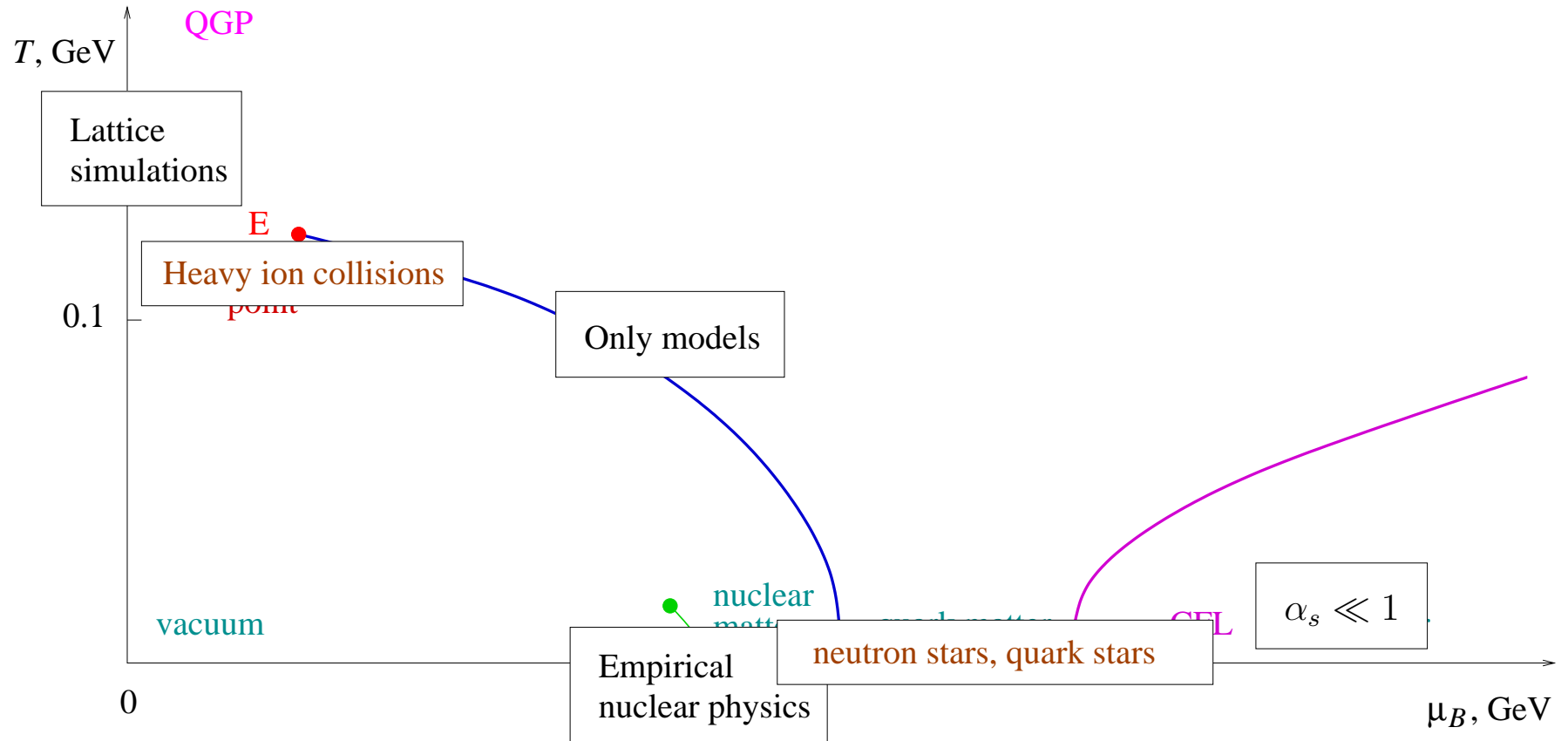
“Minimal” phase diagram

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Role of chiral symmetry ($m_q = 0$)

$$q \equiv \begin{pmatrix} u \\ d \end{pmatrix} \quad \text{Chiral:} \quad \begin{pmatrix} u_L \\ d_L \end{pmatrix} \rightarrow e^{+i\alpha \cdot \tau} \begin{pmatrix} u_L \\ d_L \end{pmatrix}, \quad \begin{pmatrix} u_R \\ d_R \end{pmatrix} \rightarrow e^{-i\alpha \cdot \tau} \begin{pmatrix} u_R \\ d_R \end{pmatrix}$$

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Order parameter:

$$\langle \bar{q}q \rangle = \langle q_L^\dagger q_R + \text{h.c.} \rangle$$

- non-zero in one phase (in vacuum),
- *exactly* zero in another phase — protected by symmetry!

Compare ferromagnet : for $T < T_c$ $\langle M \rangle \neq 0$, while for $T > T_c$ $\langle M \rangle = 0$.

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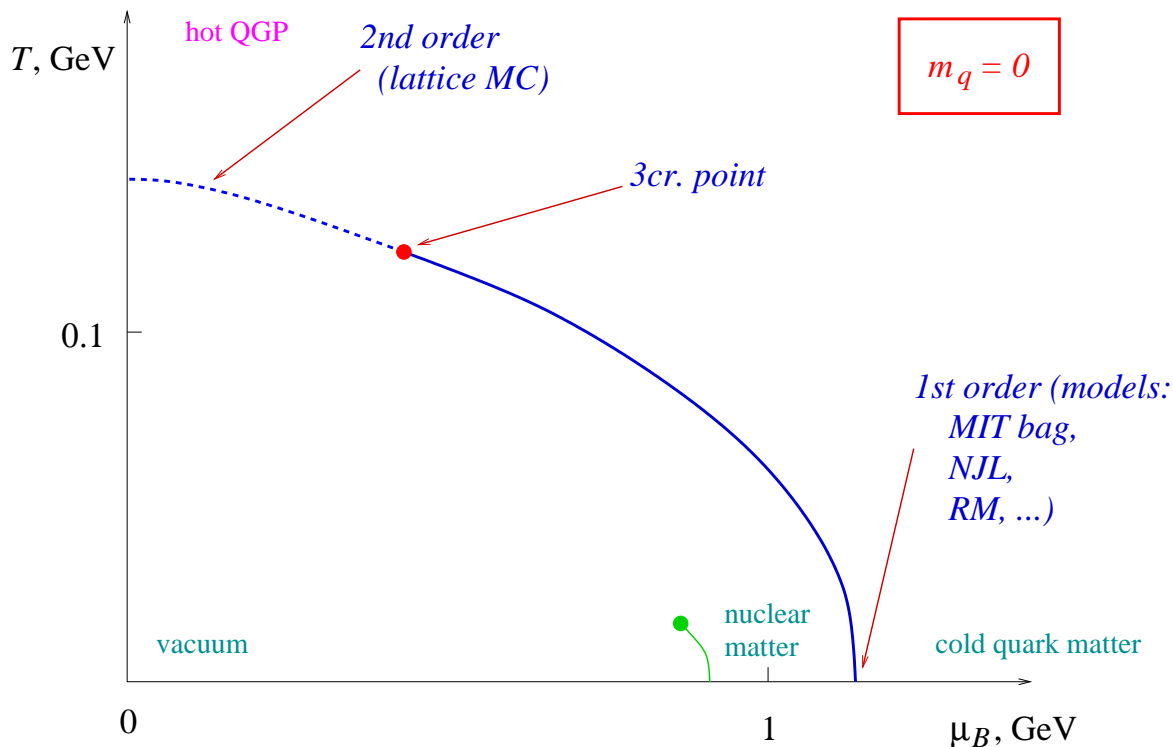
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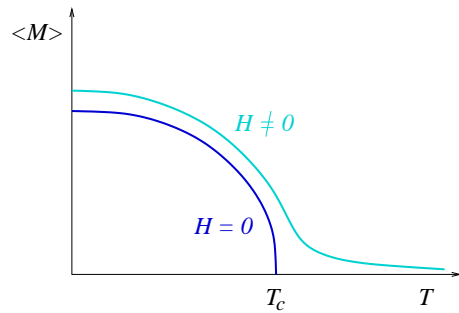
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This requires singularity \equiv phase transition

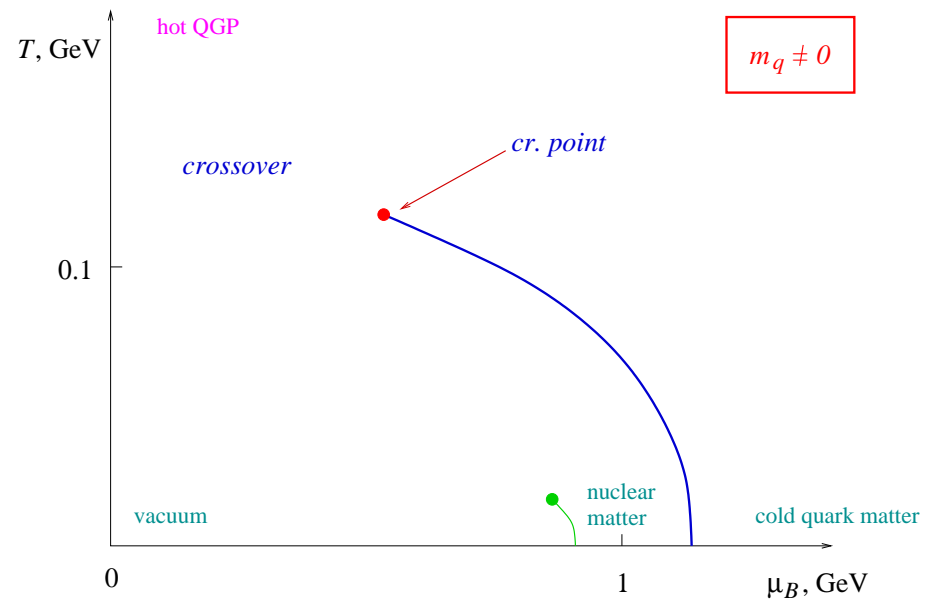


$$m_q \neq 0$$

Compare ferromagnet
at $H \neq 0$:

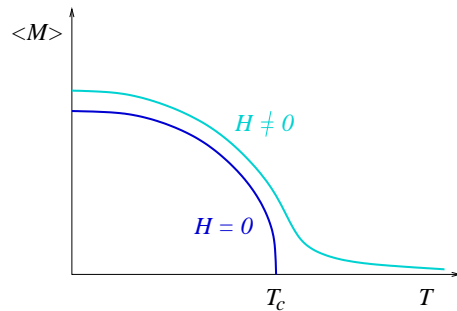


no phase transition, but
crossover

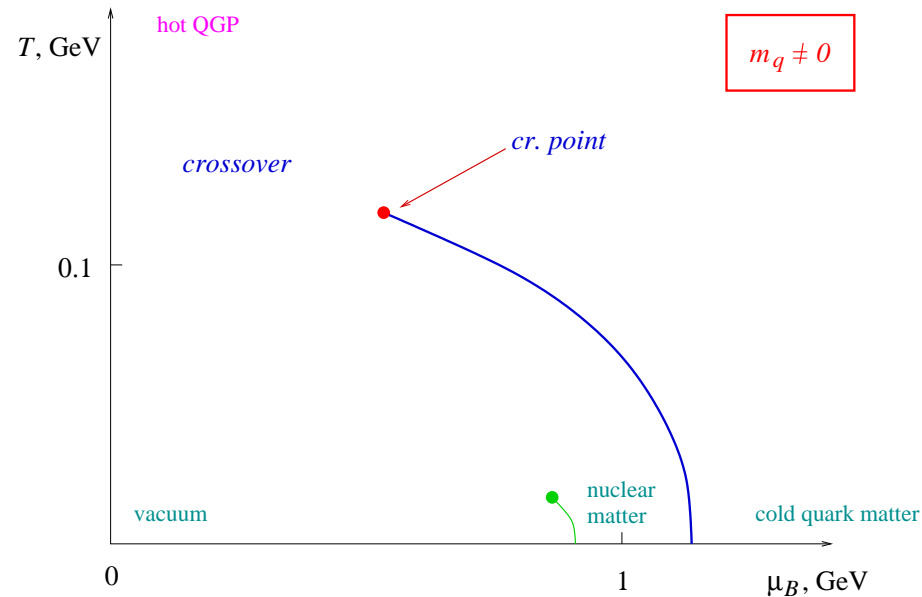


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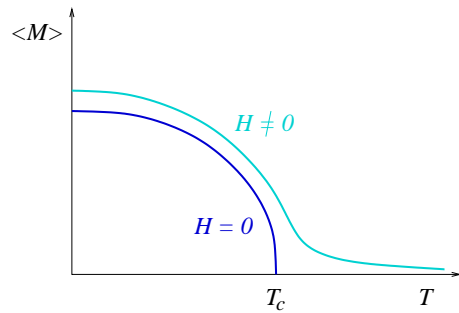
- Why all phases (can be) connected? Because \nexists order parameters: no chiral symmetry ($m_q \neq 0$);

...

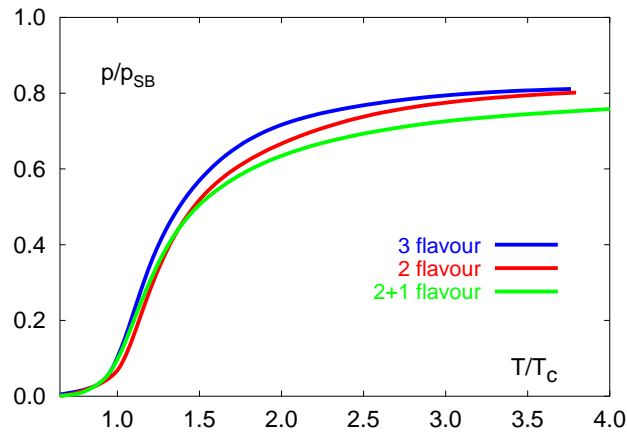
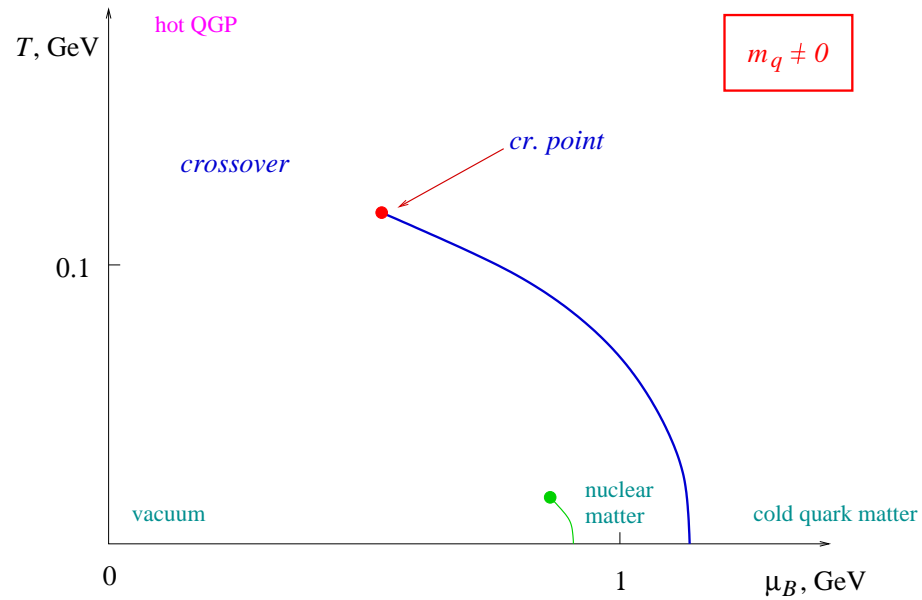
No phase boundary does not mean same physics, of course – (compare water-vapor, gas-plasma, ...) Critical point in many liquids – critical opalescence

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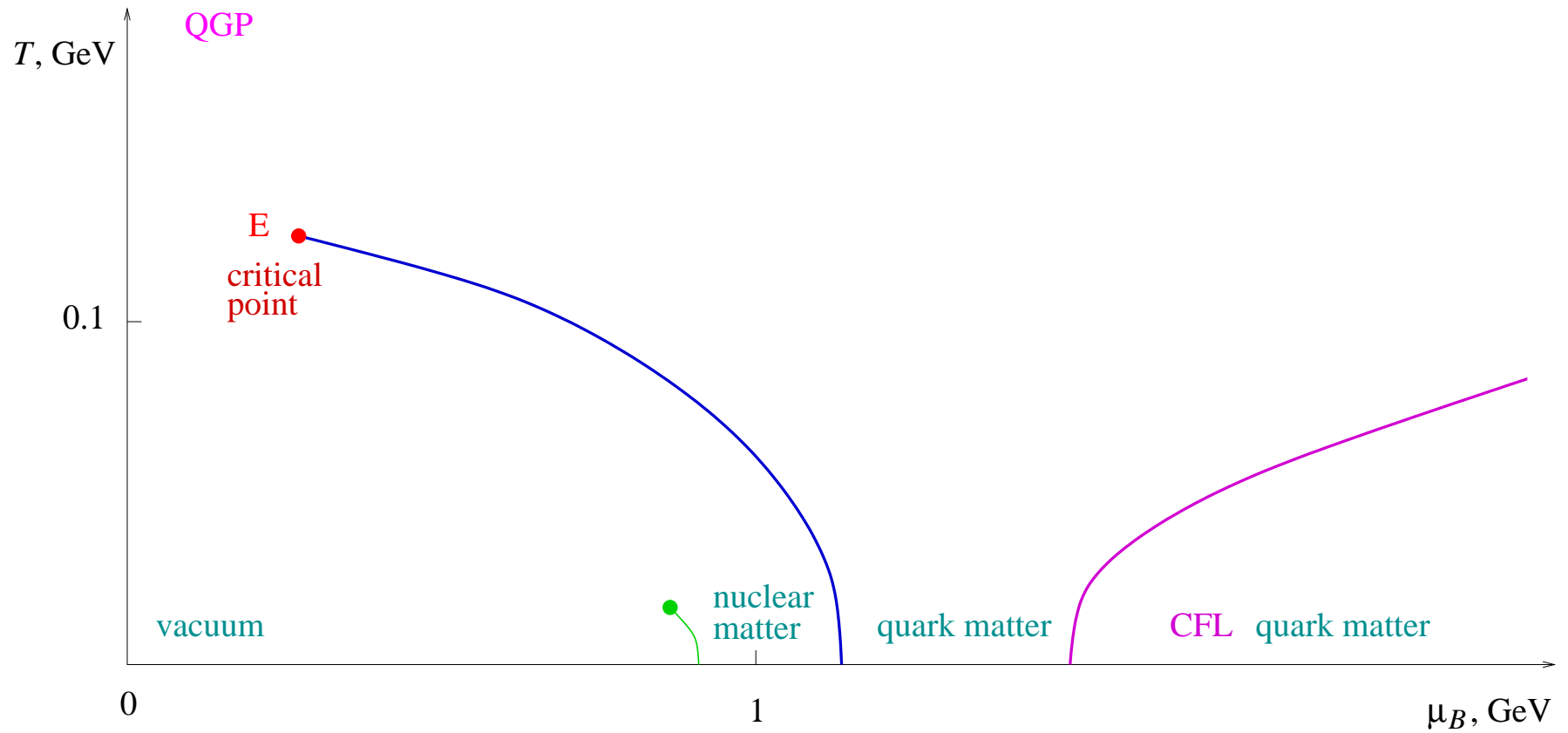


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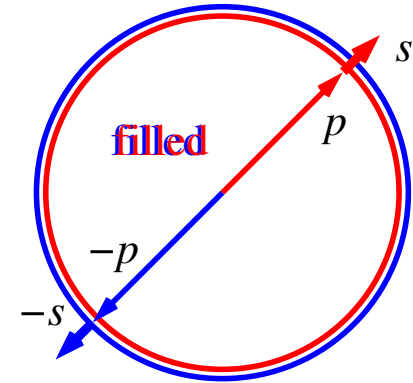
(Bielefeld group, lattice Monte Carlo)

Phase Diagram of QCD



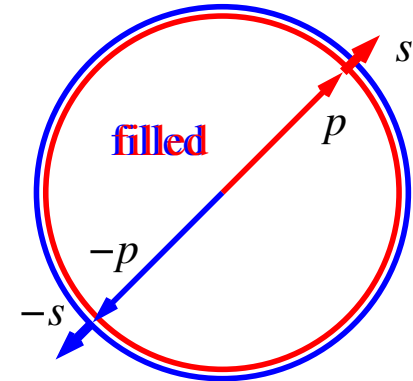
Large μ , color superconductivity and CFL

- Asymptotic freedom $\Rightarrow \alpha_s(\mu) \rightarrow 0$.
- Quarks of “different color” (color antisymmetric state) attract. Fermi sphere is unstable towards condensation of quark pairs (Cooper).
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- Does not break chiral symmetry (unlike $\bar{q}q = q_L^\dagger q_R + \text{h.c.}$).



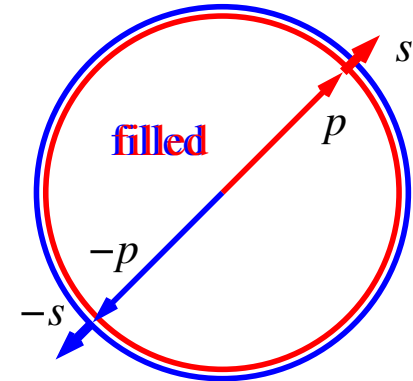
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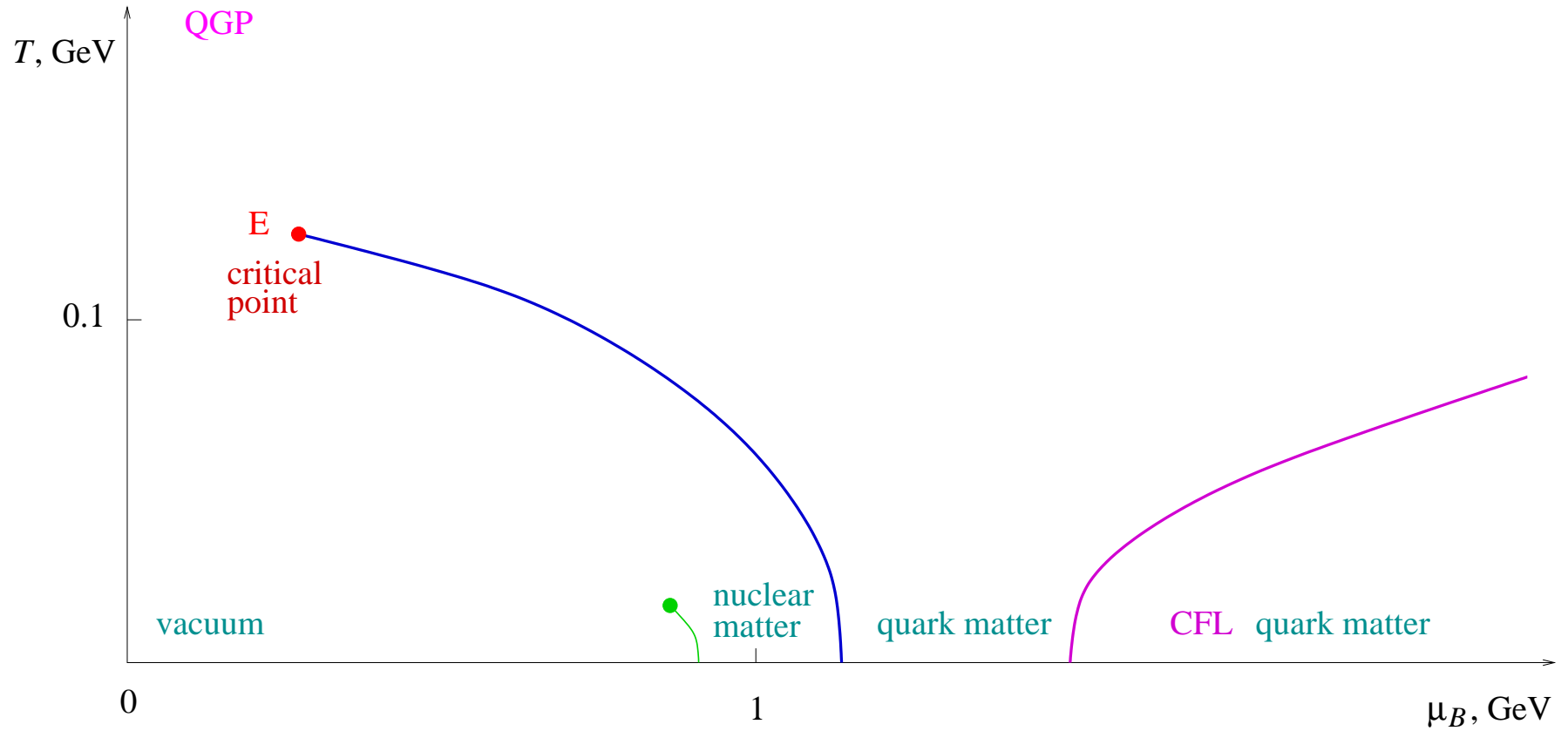
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- But flavor transformation on X can be undone by a color transformation, which is “invisible”. However, there is also $Y = u_L d_L + \dots$, and although it can also be undone by a color rotation...
- the simultaneous transformations of q_L and q_R do change the ground state if they are not equal. This is the broken chiral symmetry.

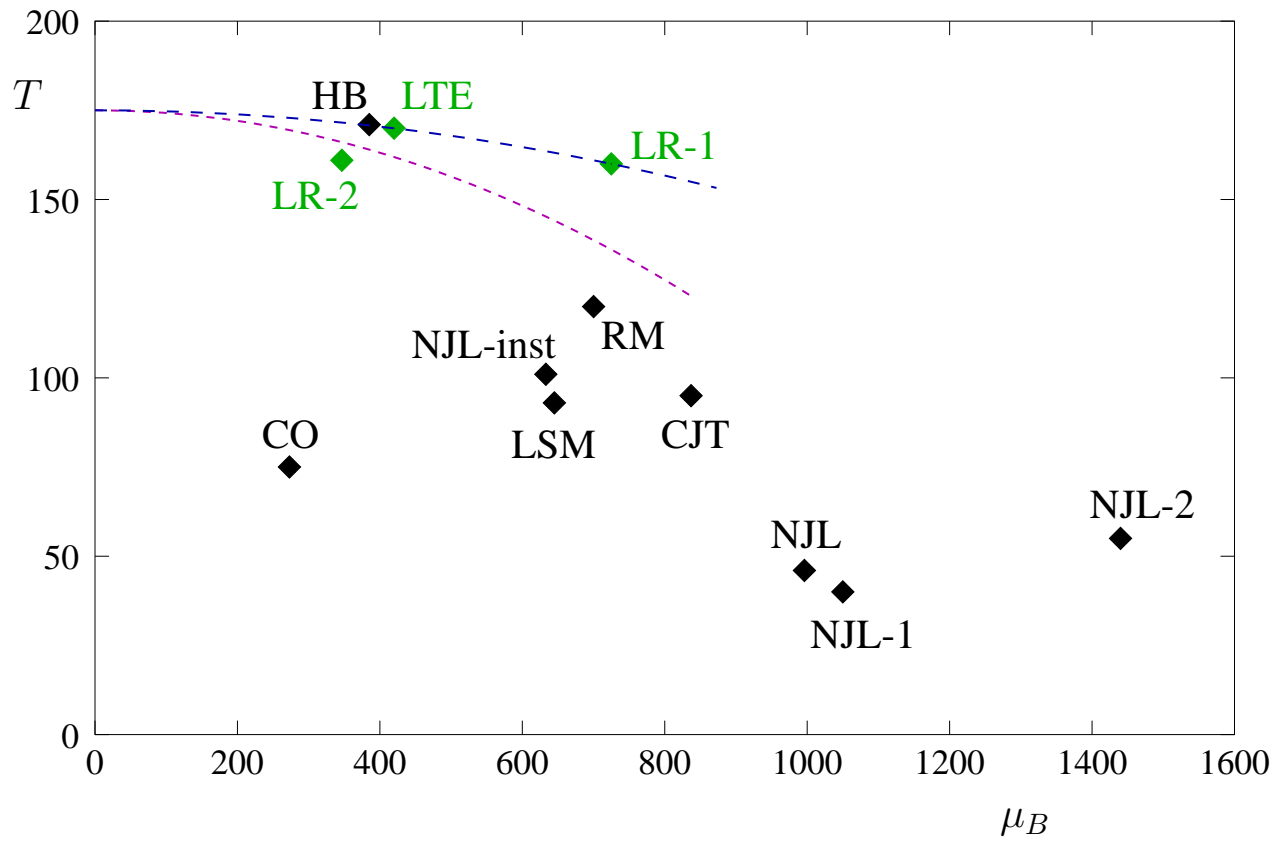


Ground state *locks color and flavor*. CFL (Alford, Rajagopal, Wilczek)

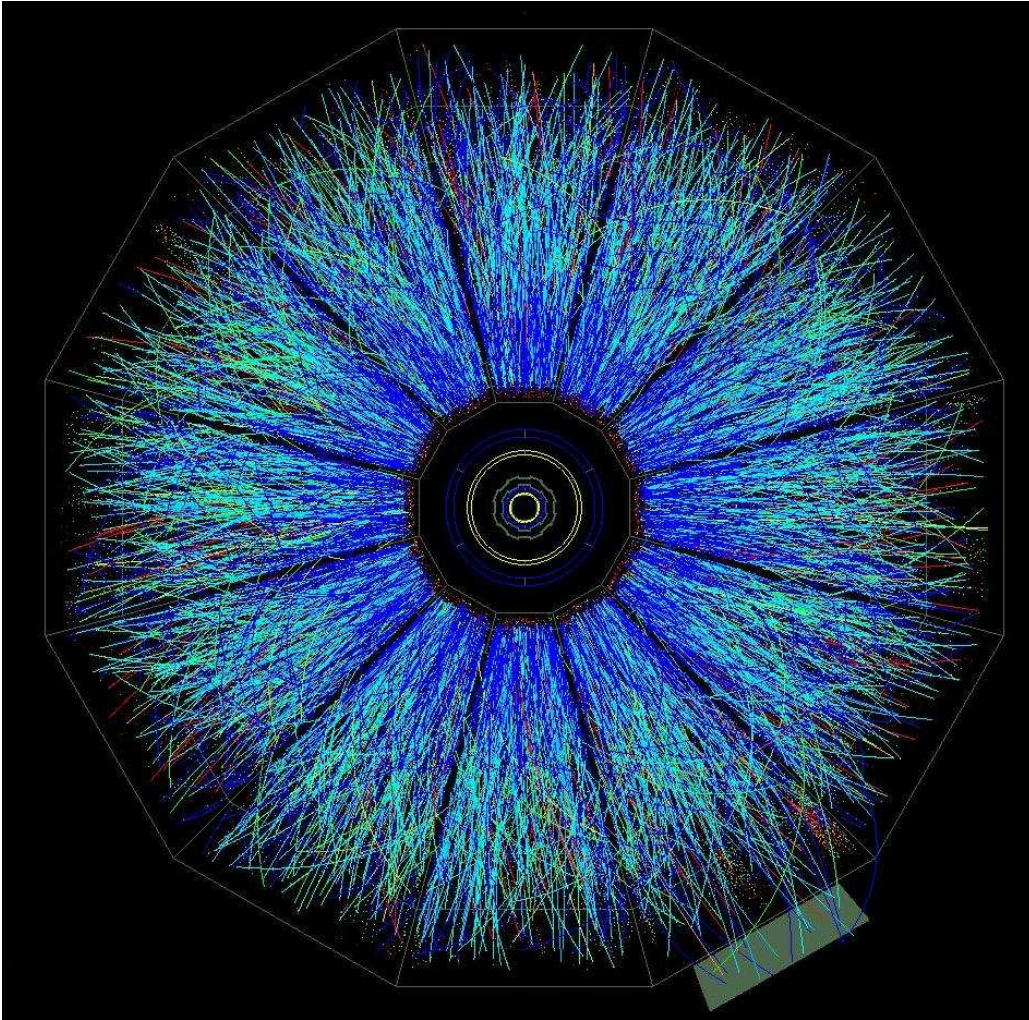
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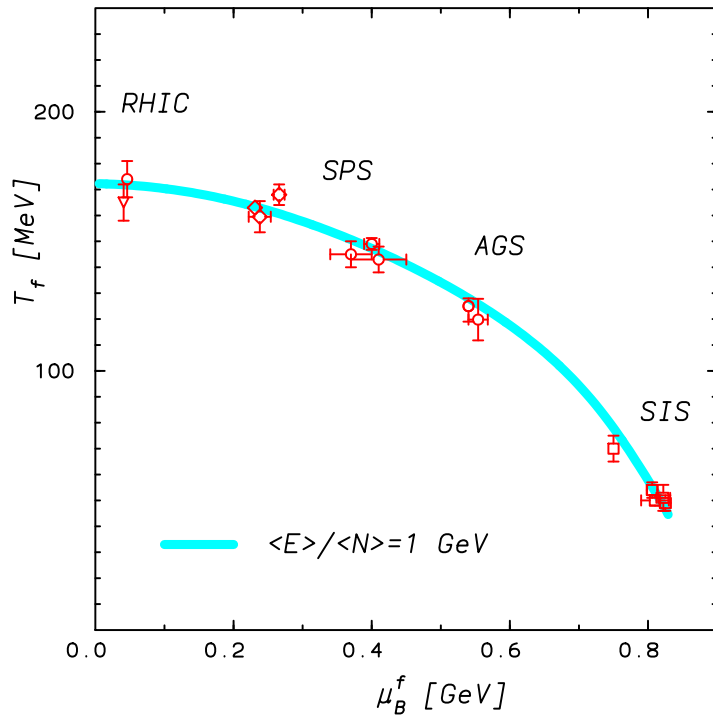
Location of the CP (theory)



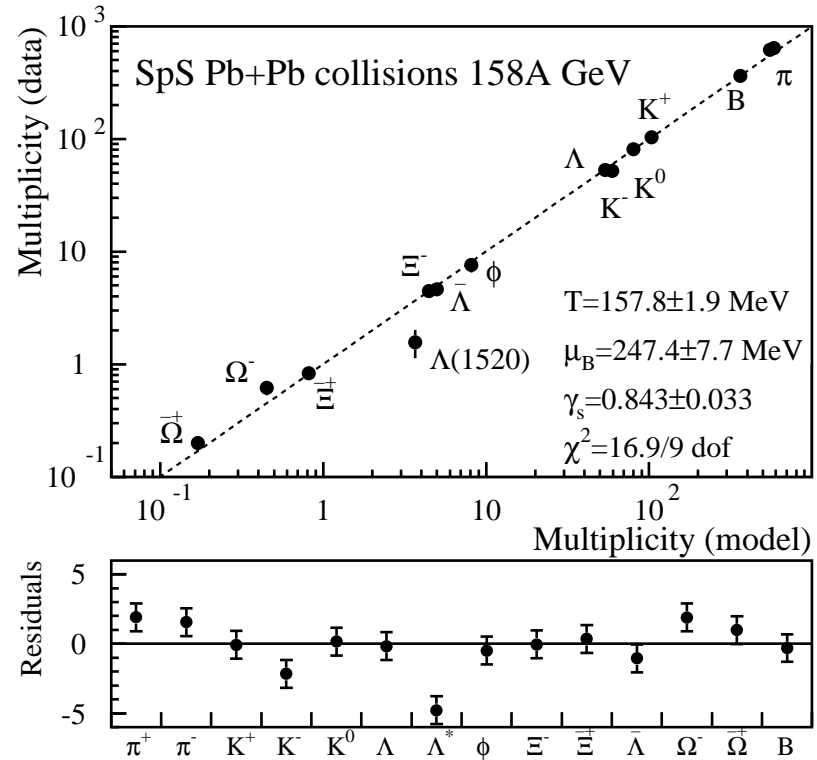
Event



Heavy-ion collisions and the phase diagram

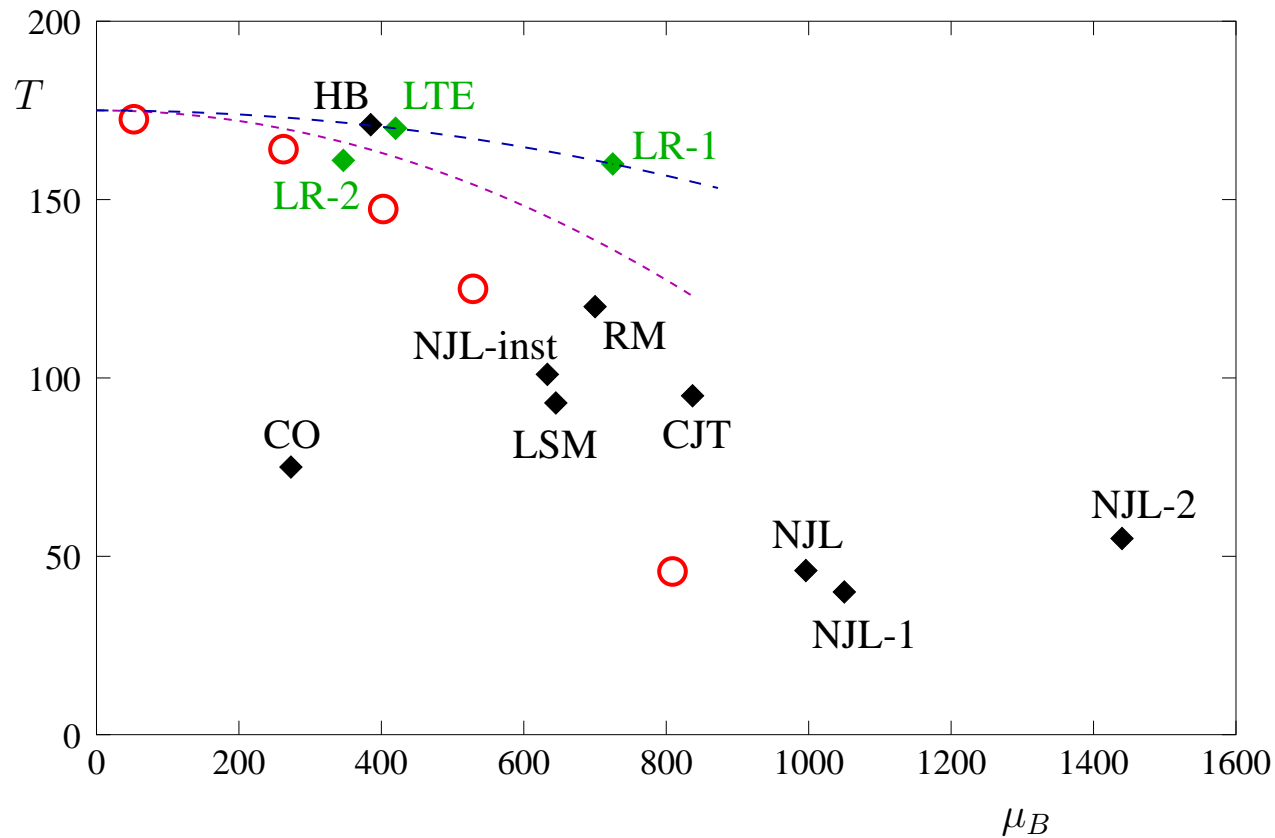


(from Braun-Munzinger, Redlich, Stachel)



(from Becattini et al)

Locating the QCD critical point experimentally

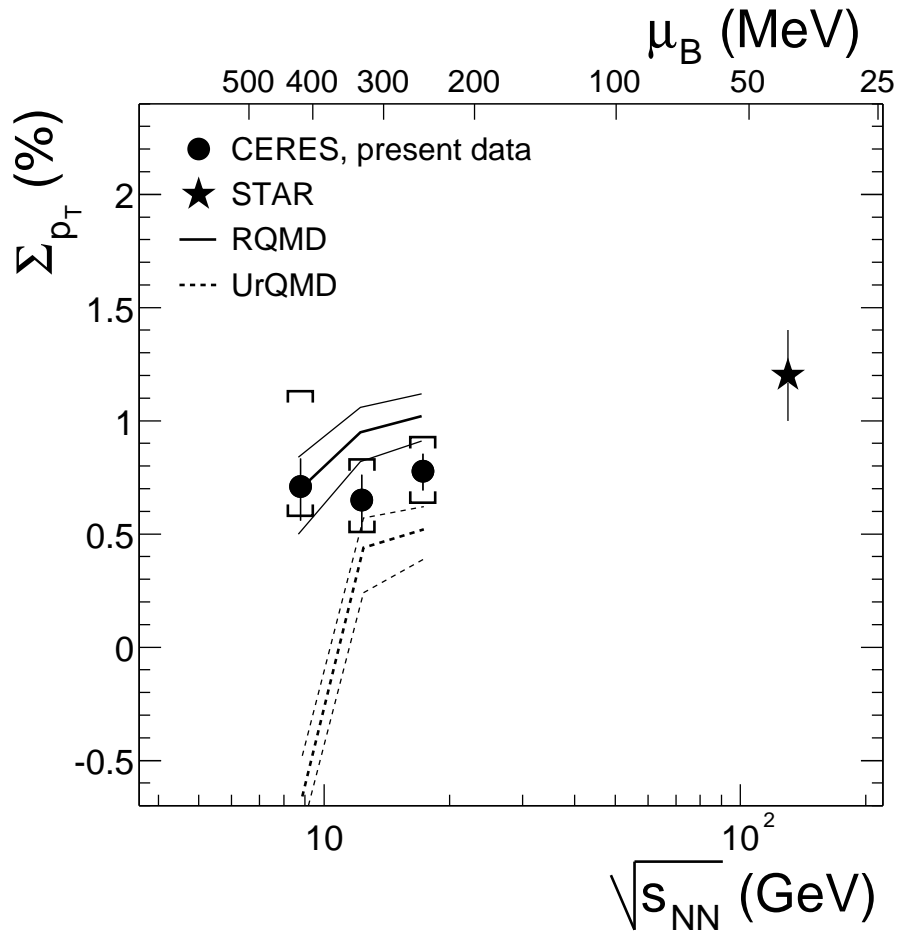


Energy scan.

Signatures: event-by-event fluctuations.

Susceptibilities diverge \Rightarrow fluctuations grow towards the critical point.

Data (example): p_T fluctuations (CERES)



Near the critical point (for CERES acceptance):

$$\sim 2\% \times \left(\frac{G}{300 \text{ MeV}} \right)^2 \left(\frac{\xi_\sigma}{3 \text{ fm}} \right)^2$$

$$(\xi_\sigma = 1/m_\sigma)$$

Signal one is looking for:

non-monotonic dependence on \sqrt{s} .

Summary and conclusions

- Phase diagram of QCD is a challenging, long-standing problem.
- $\alpha_s \ll 1$ calculations can address regions of $T, \mu \gg 1 \text{ fm}^{-1}$. Many interesting phenomena.
- The most interesting region is still out of reach of controllable calculations.
- Existence of the critical point on the QCD phase diagram is suggested by many theoretical approaches, including, recently, lattice QCD.
- Heavy ion experiments can discover the critical point – by measuring fluctuation observables as a function of collision parameters.
Search is underway.
- Needed:

Theory: a reliable (systematically improvable) method to study $T\mu$ plane and locate CP.
Experiment: [scan of the QCD phase diagram](#) — \sqrt{s} scan.