

Quantum acoustics

jilawww.colorado.edu/~lehnertk



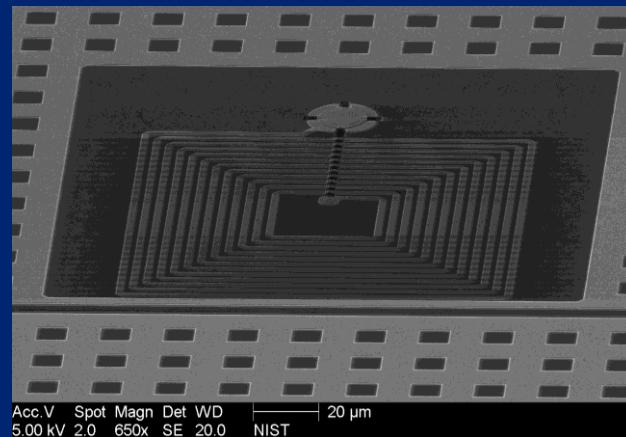
Konrad Lehnert

Post-docs

Tobias Donner
Francois Mallet
Tauno Palomaki

Collaborators

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Ray Simmonds
Kent Irwin
Cindy Regal

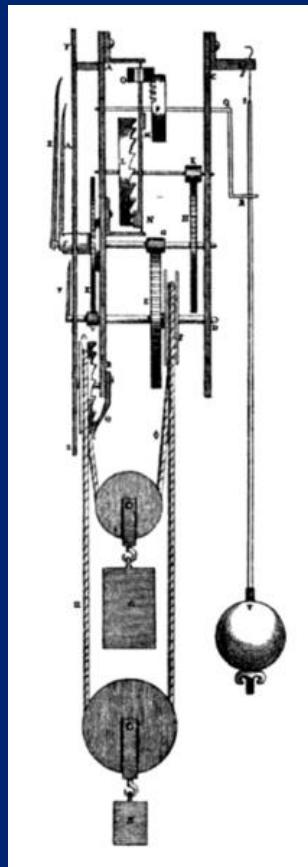


Graduate students

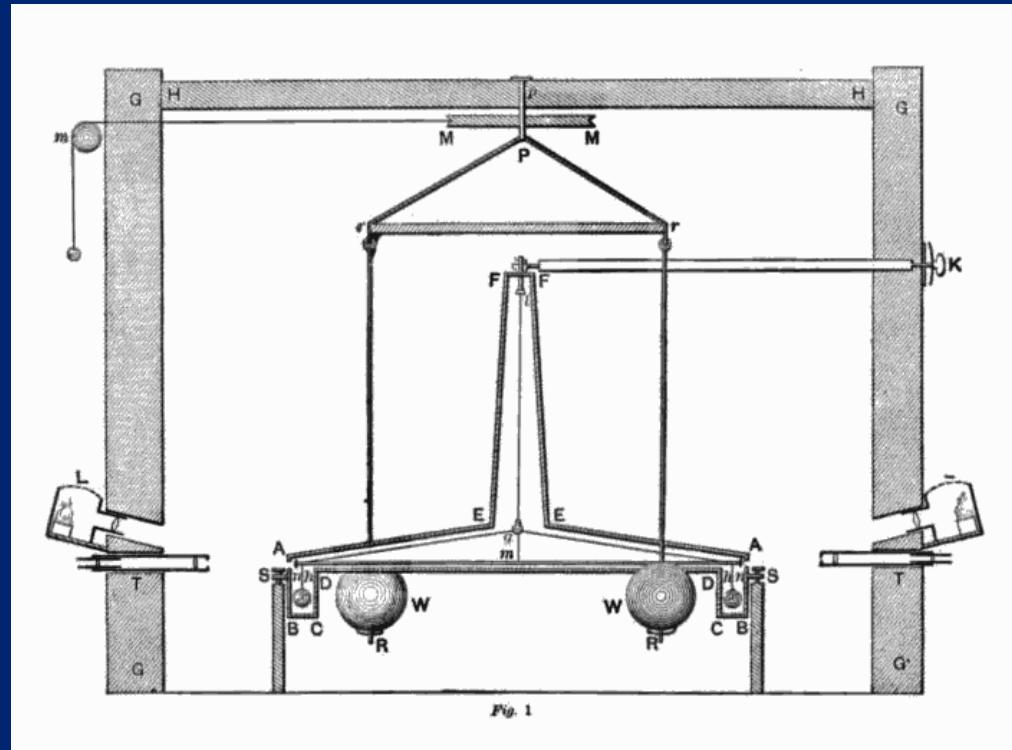
Jennifer Harlow
Reed Andrews
Hsiang-Shen Ku
William Kindel
Manuel Castellanos-Beltran
Nathan Flowers-Jacobs



Precision measurement tools were once mechanical oscillators



Huygens pendulum clock



The Cavendish balance
for weighing the earth

Modern measurement tools exploit optics and electronics, not mechanics

Laser light

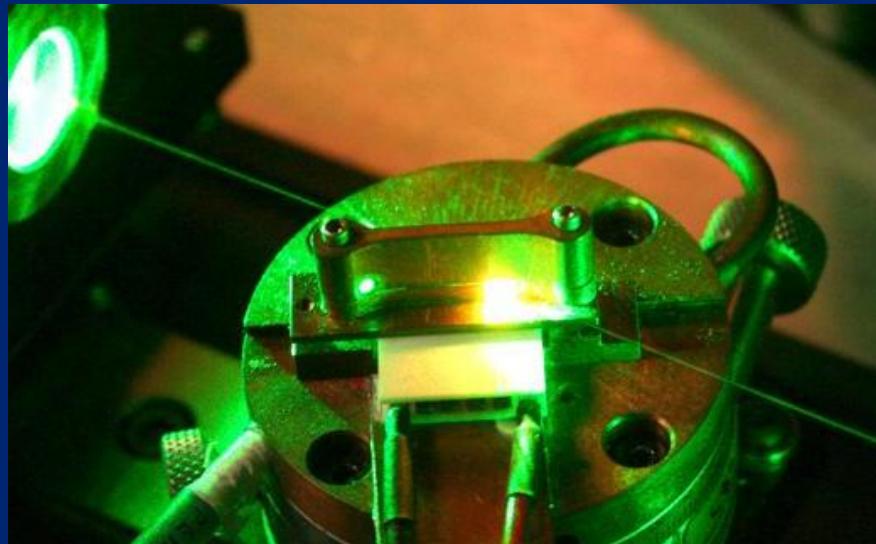


Image: Cundiff lab JILA

electricity



Described by:
Maxwell's equations

$$\nabla \times \vec{E} = -\mu \frac{\partial \vec{H}}{\partial r}$$

$$\nabla \times \vec{H} = \vec{J}_c + \epsilon \frac{\partial \vec{E}}{\partial r}$$

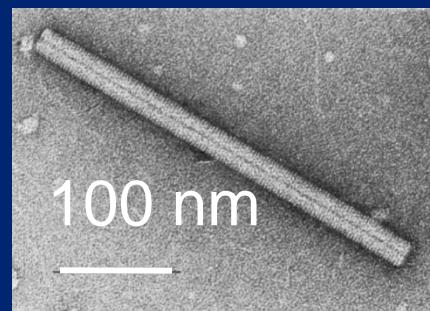
$$\nabla \cdot \vec{D} = p_v$$

$$\nabla \cdot \vec{B} = 0$$

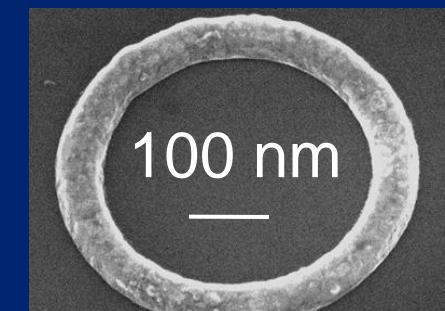
...AND THERE WAS LIGHT

Optical probes are ill-suited to directly measuring many interesting systems

non-atomic system



nuclear spins
in a virus



electrons in an
aluminum ring
(Harris lab, Yale)

Systems with:

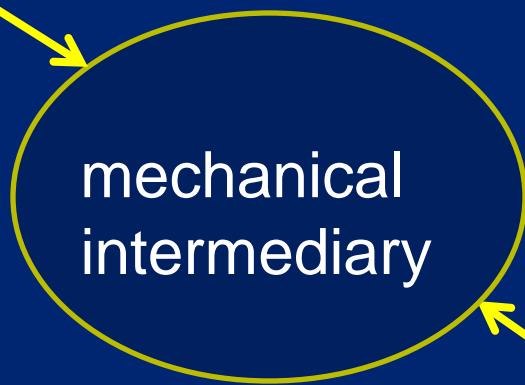
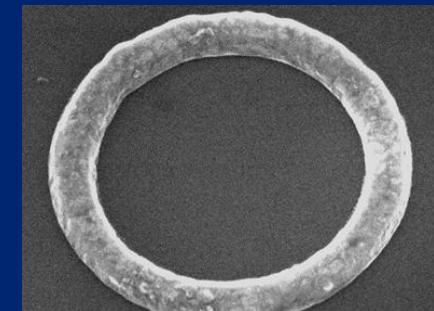
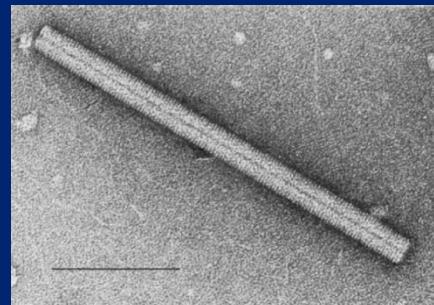
dense low-energy spectra

nanometer length scales

weak coupling to light

quantum
probe: light

Mechanical oscillators enable measurements of non-atomic systems

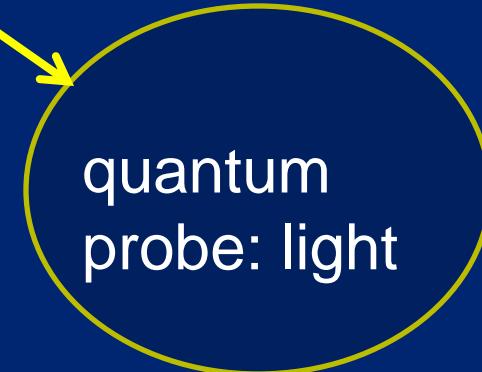


Systems with:

dense low-energy spectra

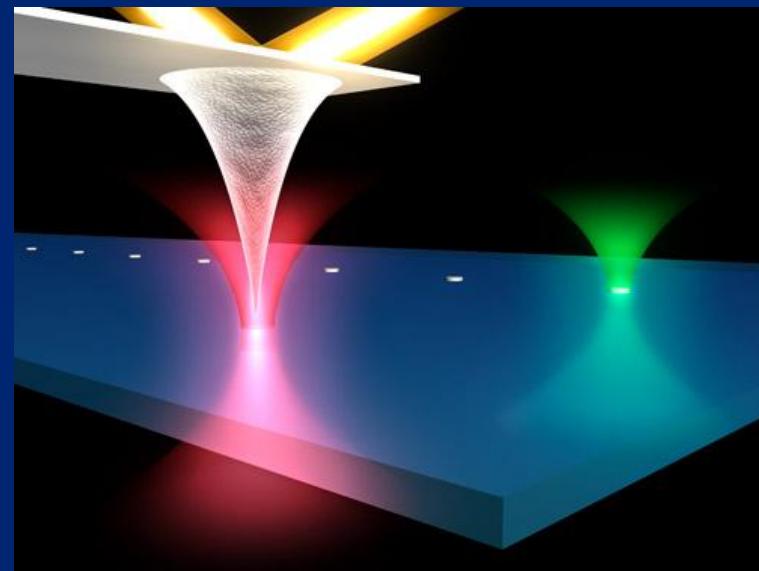
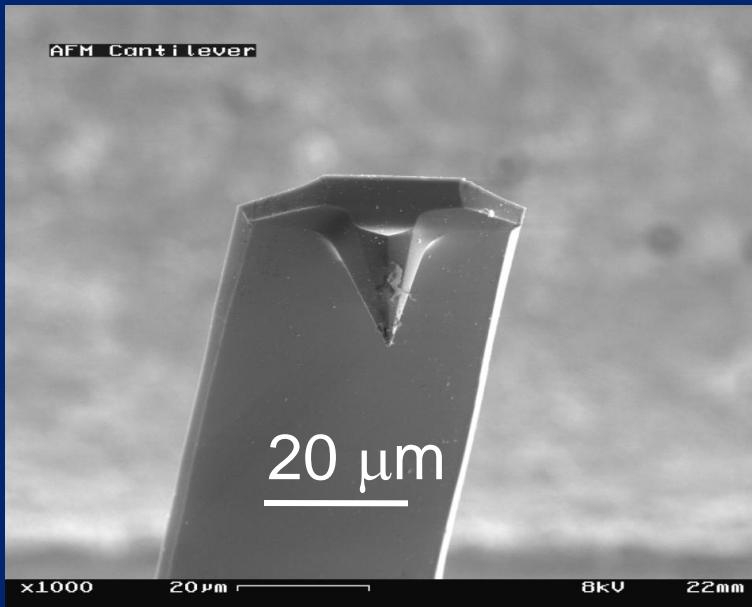
nanometer length scales

weak coupling to light



Mechanical oscillators are tools that access the nano-world

Atomic Force Microscope

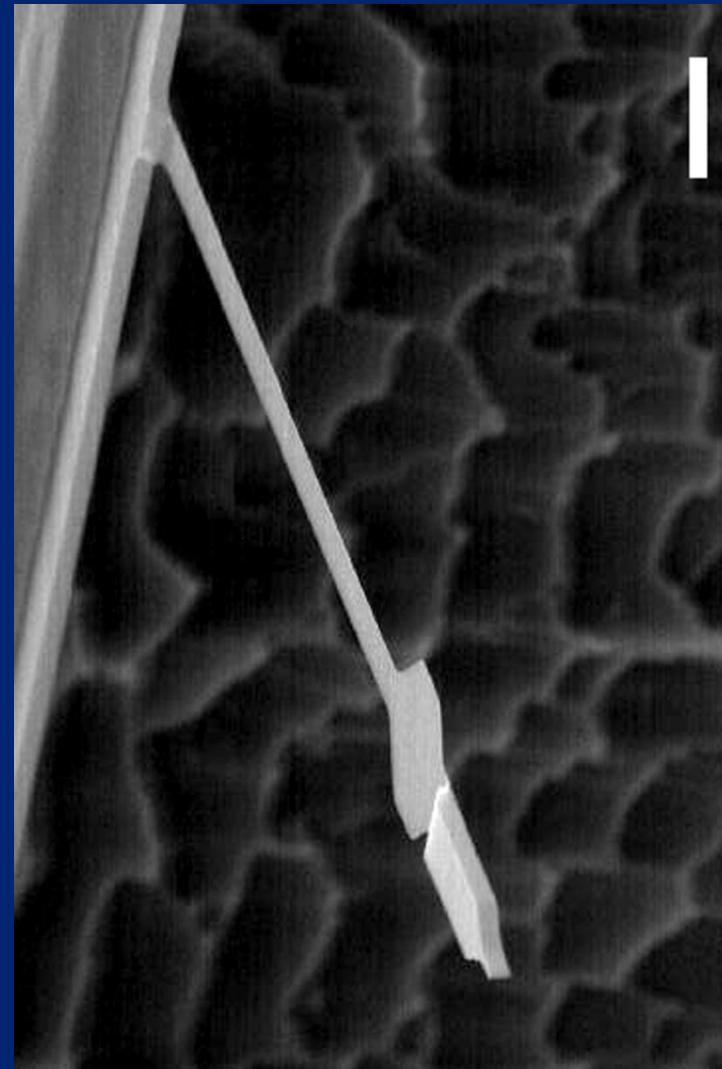
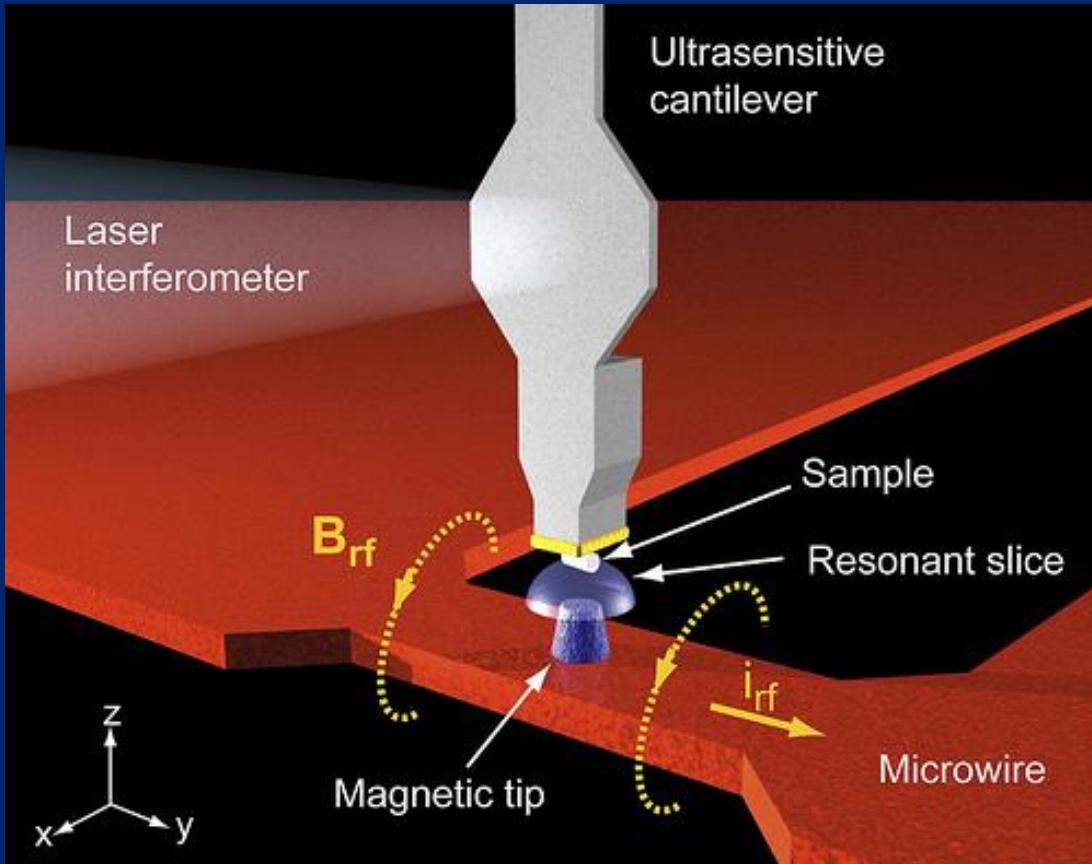


Perkins lab, JILA

Mechanical oscillator
nanometer probe
universal coupling (senses any force)

Optical interferometer detects oscillator motion

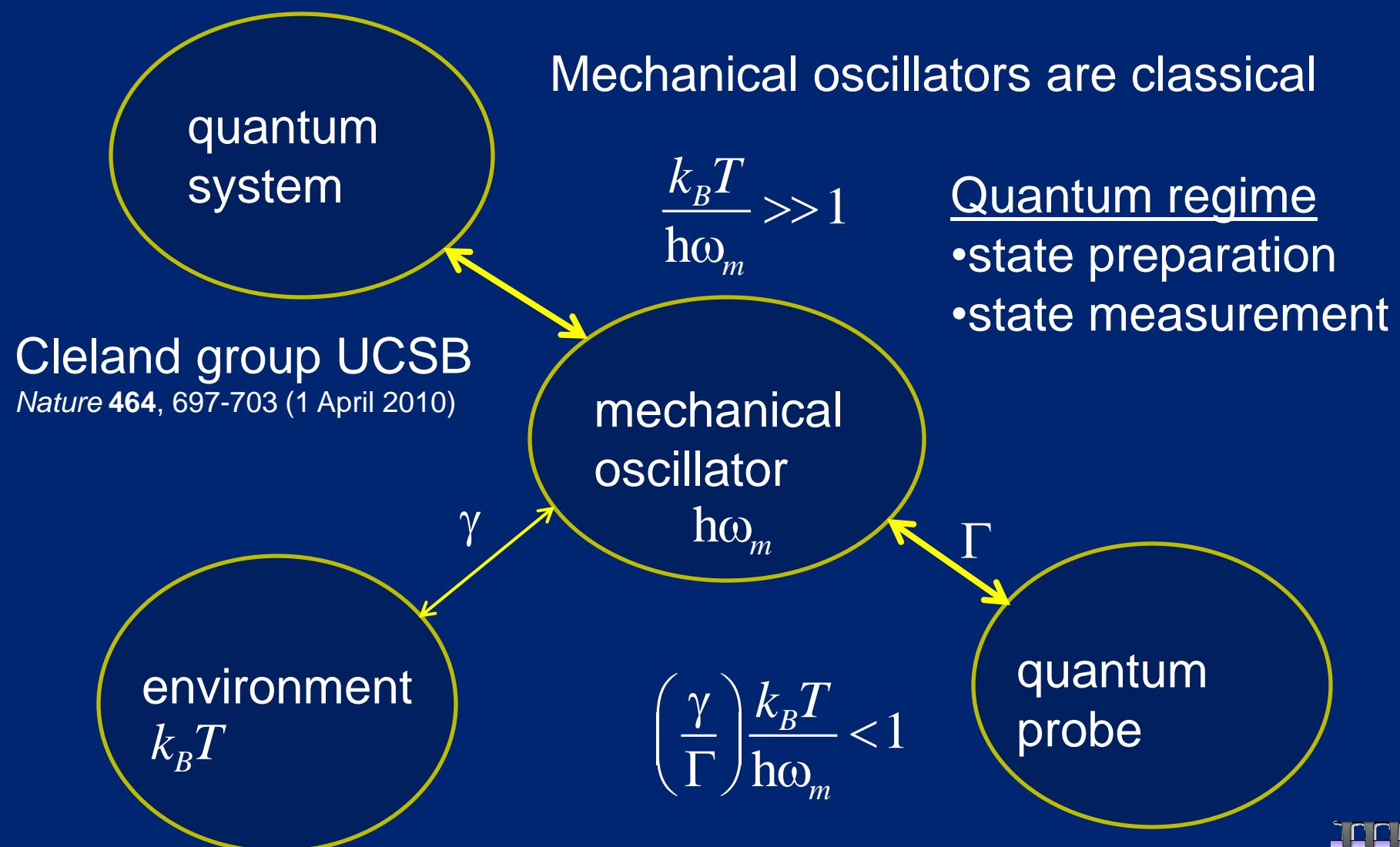
Mechanical oscillators form ultrasensitive, mesoscopic magnetometers



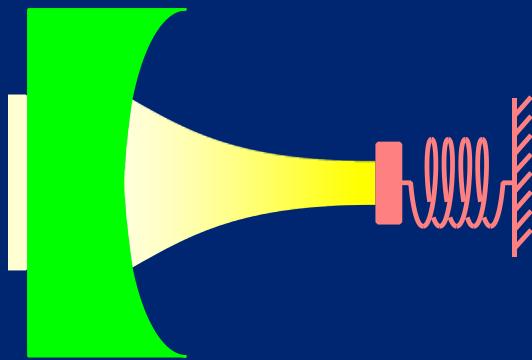
nanoscale MRI of a single virus

Rugar Lab, IBM $\sqrt{S_f^{\text{tot}}} = 0.8 \text{ aN/Hz}^{1/2}$

Mechanical oscillators as quantum coherent interfaces between incompatible systems



Cavity optomechanics: Use radiation pressure for state preparation and measurement



Fabry-Perot cavity with
oscillating mirror

$$\hat{H}_{\text{tot}} = \hbar\omega_c(a^\dagger a + \frac{1}{2}) + \hbar\omega_m(b^\dagger b + \frac{1}{2}) + H_I$$

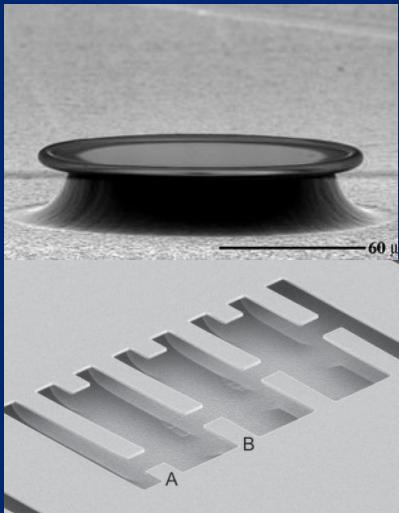
$$\hat{H}_I = \hat{F} \cdot \hat{x} = \hbar a^\dagger a g x_{zp} (b^\dagger + b)$$

Infer motion through optical phase

Cool with cavity-retarded radiation force

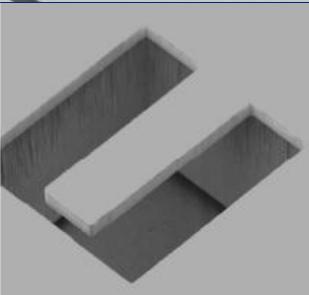
$$\Gamma \approx \frac{g^2 x_{zp}^2}{4\hbar\omega_c^2 \kappa} P_{\text{circ}} \quad g \sim 100 \text{ MHz/nm}$$

Images of cavity optomechanical systems

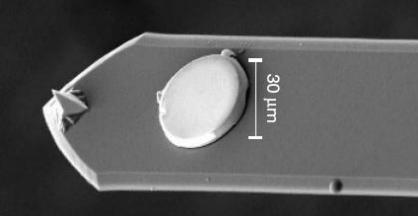


Caltech, Vahala
MPQ, Kippenberg

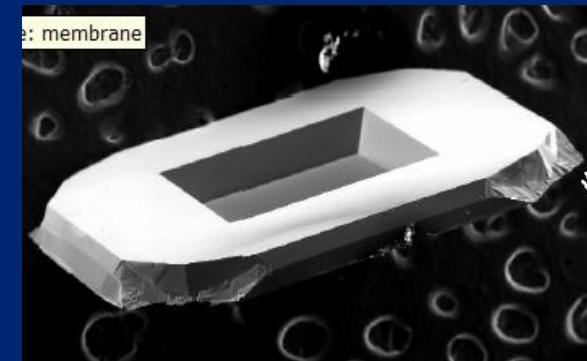
10 ng



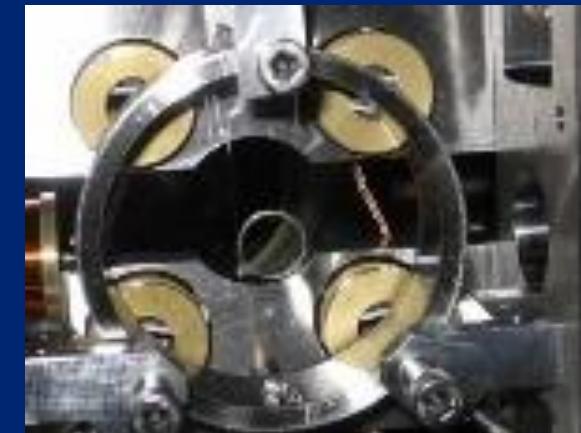
ENS: Pinard and Heidmann



UCSB: Bouwmeester



Yale, Harris



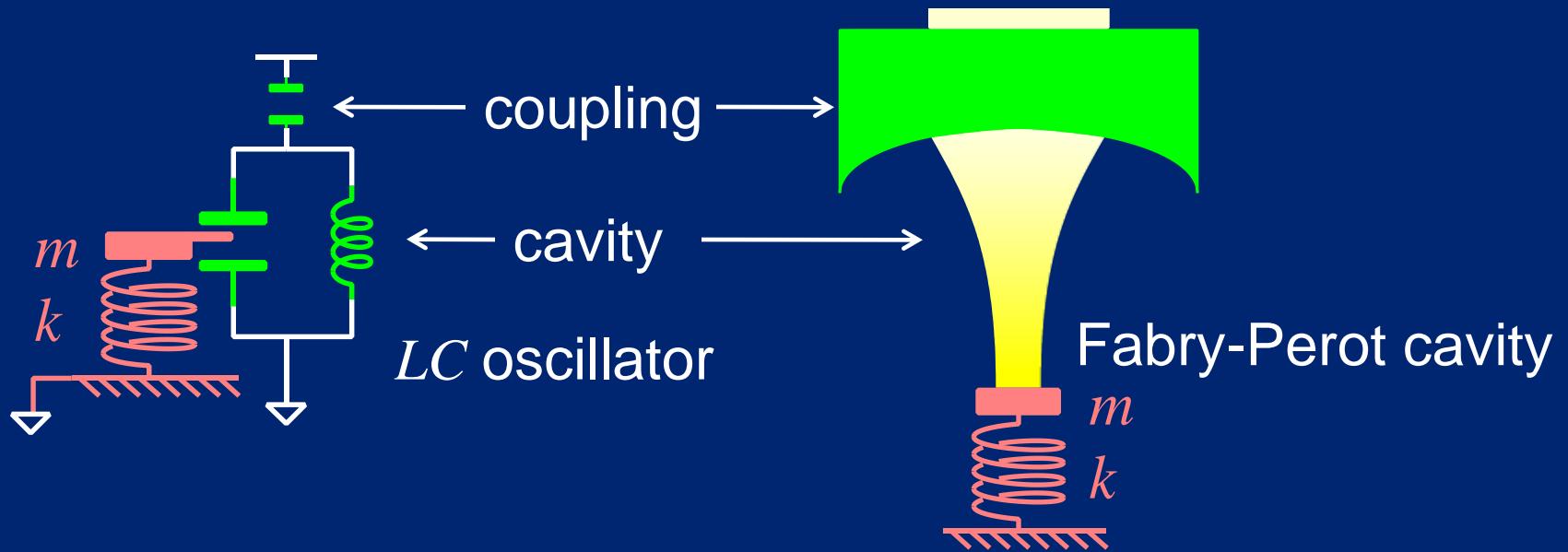
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MIT, Mavalvala

Microwave cavity optomechanics

Reduce coupling to the environment by lowering temperature: microwave optomechanics

Microwave “light” in ultralow temperature cryostat



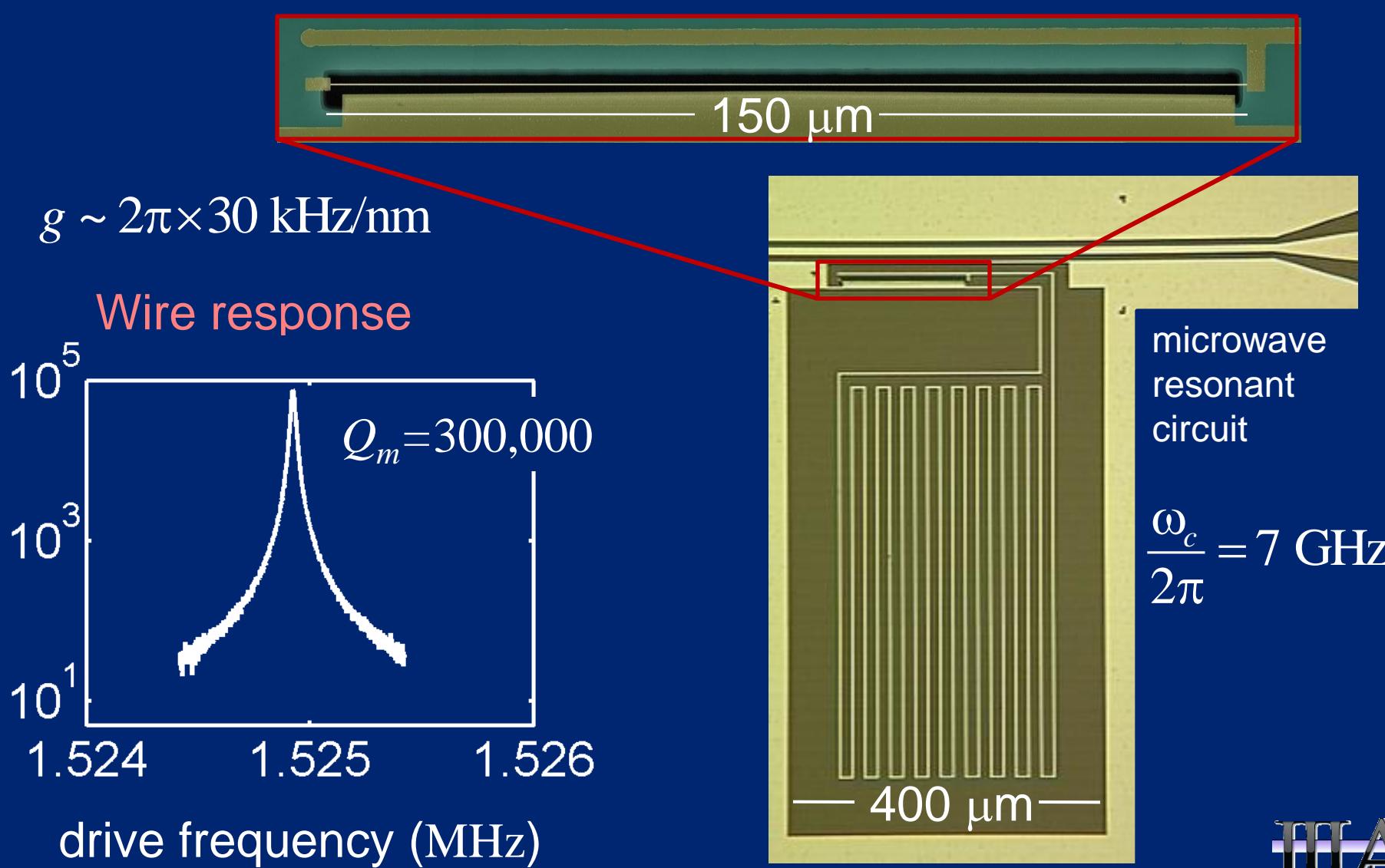
Strategy

Cool environment to $T \ll 1 \text{ K}$

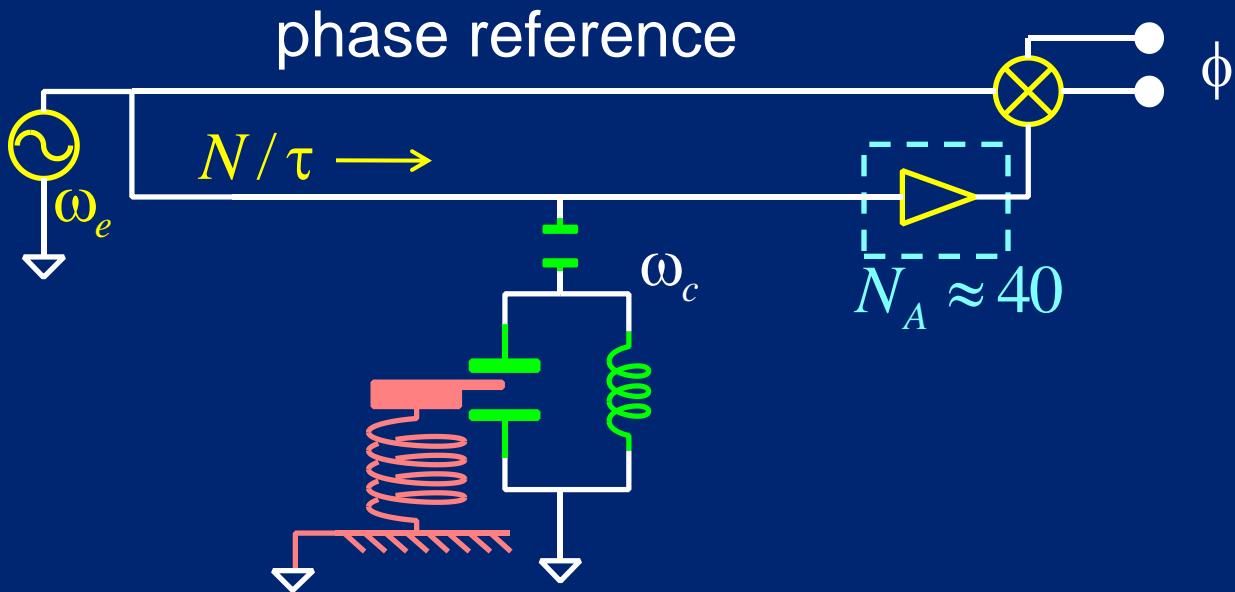
High Q mechanical oscillators

$$\left(\frac{\gamma}{\Gamma}\right) \frac{k_B T}{\hbar \omega_m} < 1$$

Cavity optomechanical system realized from a nanomechanical wire in a resonant circuit



Nanomechanical motion monitored with a microwave Mach-Zehnder interferometer

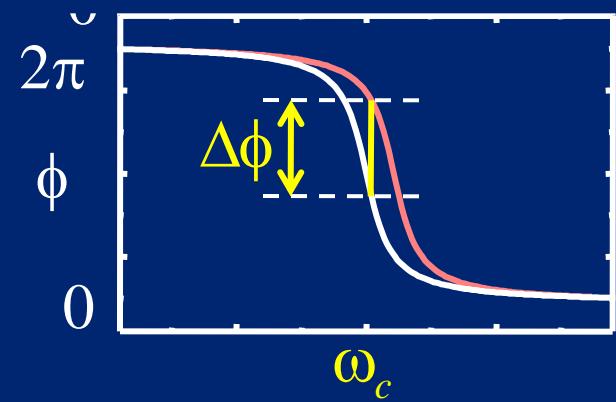


Infer wire motion from phase shift

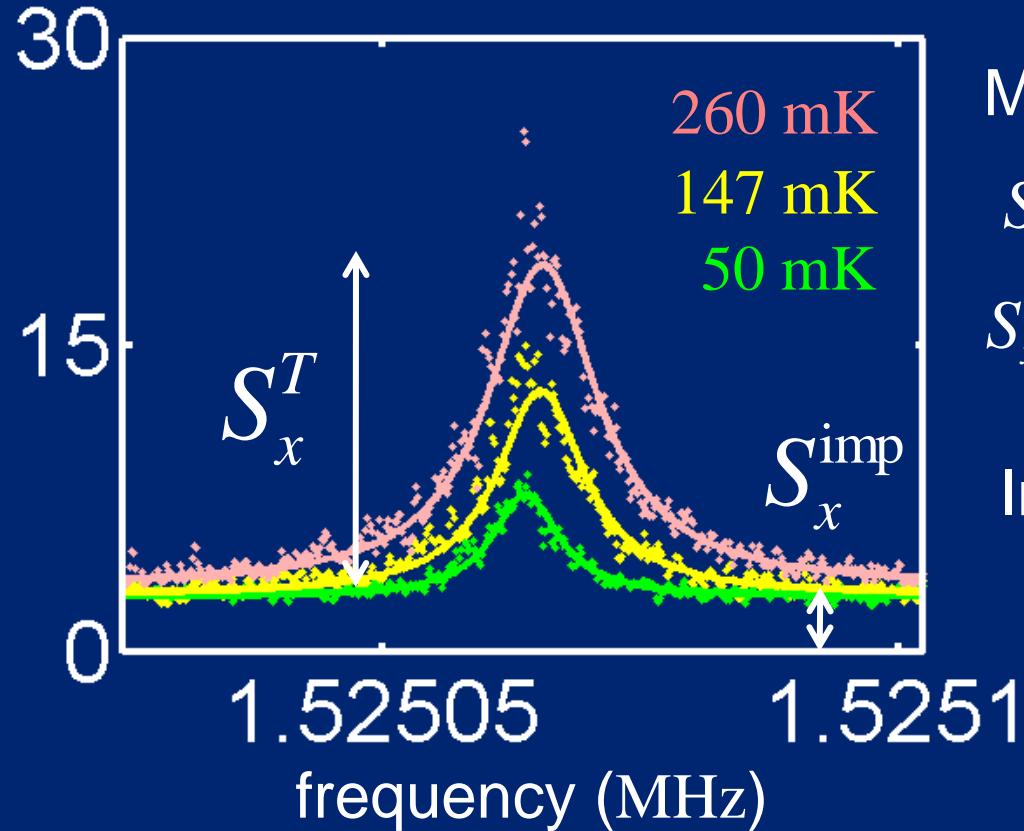
$$\Delta\phi = \frac{\Delta\omega_c}{\gamma_c} = \frac{gx}{\gamma_c}$$

Phase sensitivity limited by amplifier (HEMT)

$$S_\phi = \frac{N_A + \frac{1}{2}}{N/\tau} = \frac{\text{noise quanta}}{\text{photon flux}}$$



Thermal motion of beam calibrates interferometer noise (imprecision)



Minimum imprecision

$$S_x^{\text{imp}} = 145 \text{ ZPE}$$

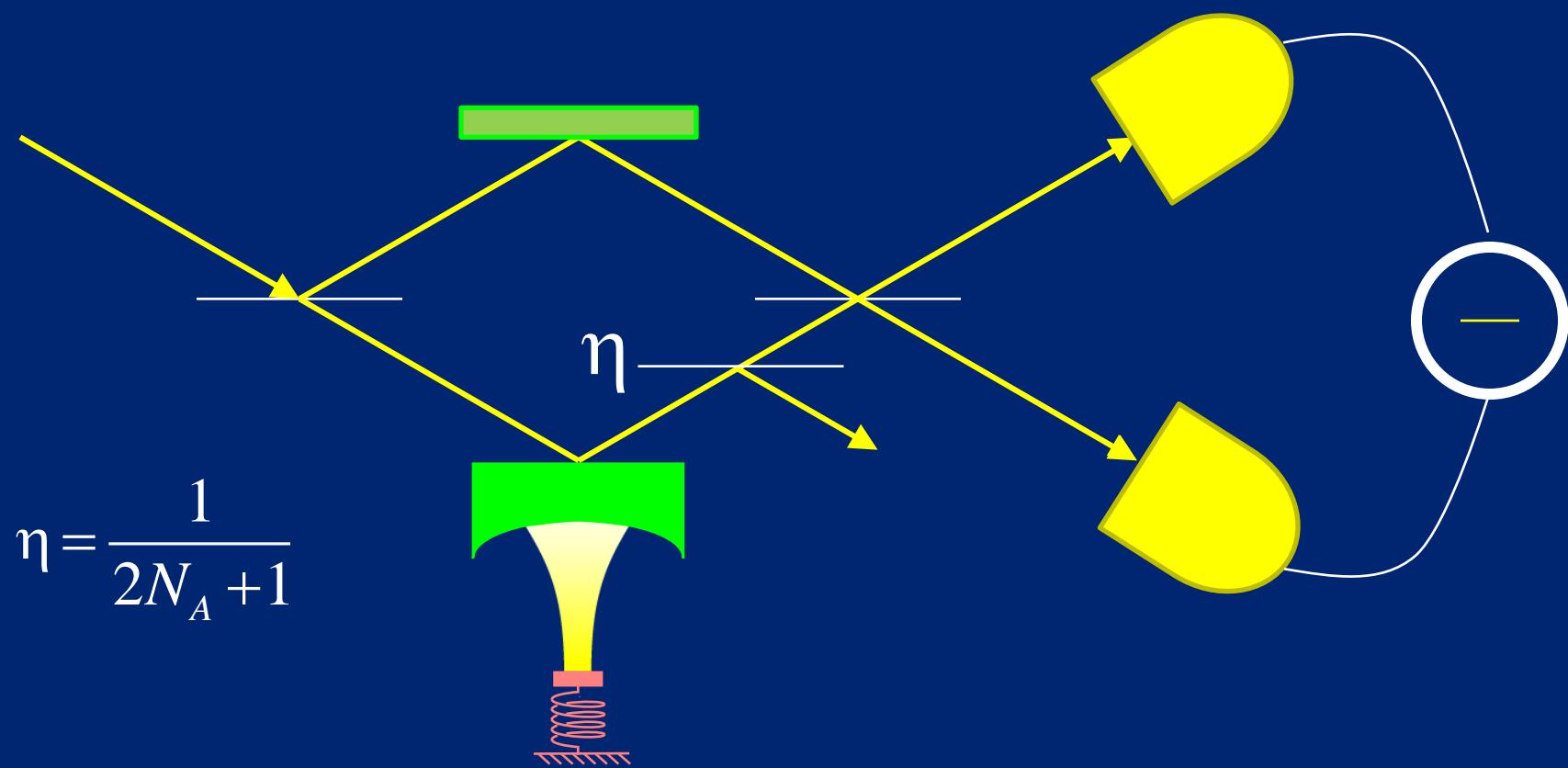
$$S_x^{\text{imp}} = 290 \times \text{SQL}$$

Imprecision at the SQL

$$S_x^{\text{sql}} = h/m\omega_m\gamma_m$$

Determine measurement imprecision S_x^{imp}

Amplifier added noise mimics quantum inefficiency



Excellent microwave amplifier:

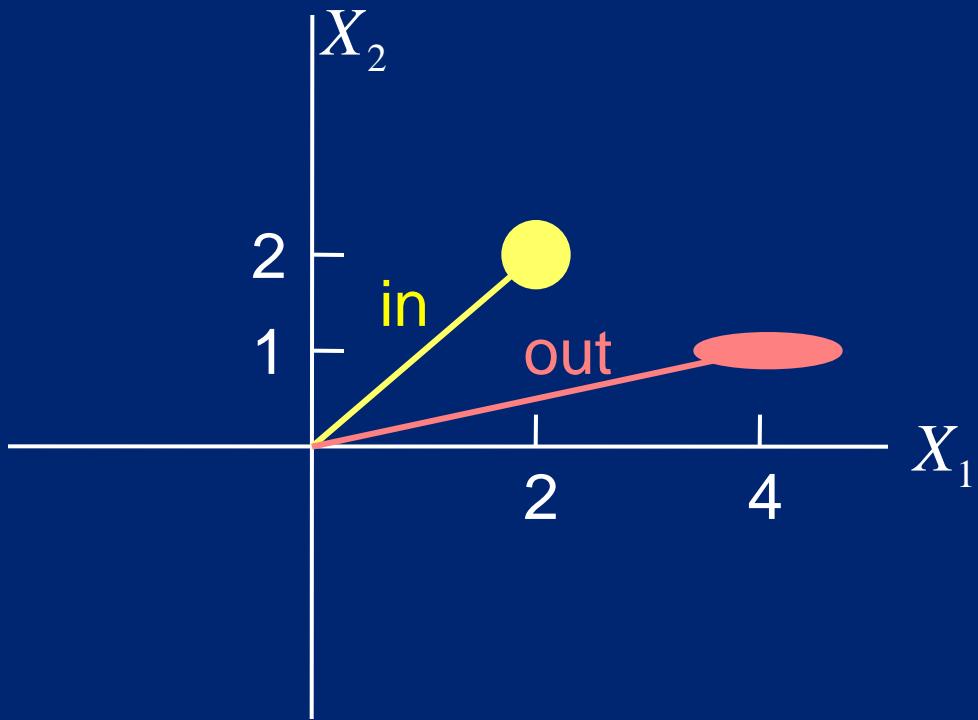
$$N_A = 40 \quad \eta = 1.2\%$$

Efficient quantum measurement

A single quadrature amplifier preserves entropy with photon number gain



$$V(t) = V_q \left(X_1^{\text{out}} \cos \omega t + X_2^{\text{out}} \sin \omega t \right)$$



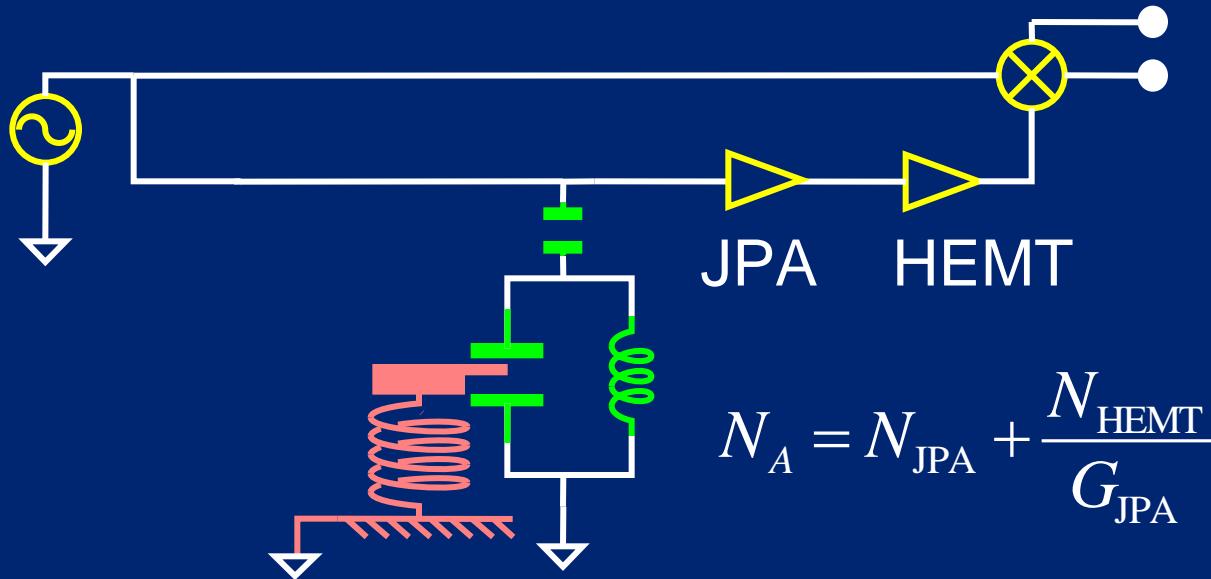
$$X_1^{\text{out}} = G X_1^{\text{in}}$$

$$X_2^{\text{out}} = \frac{1}{G} X_2^{\text{in}}$$

$$X_1^{\text{out}} X_2^{\text{out}} - X_2^{\text{out}} X_1^{\text{out}} = \frac{i}{2}$$

$$N_A \geq 0$$

Incorporate quantum pre-amplifier into the Mach-Zehnder interferometer

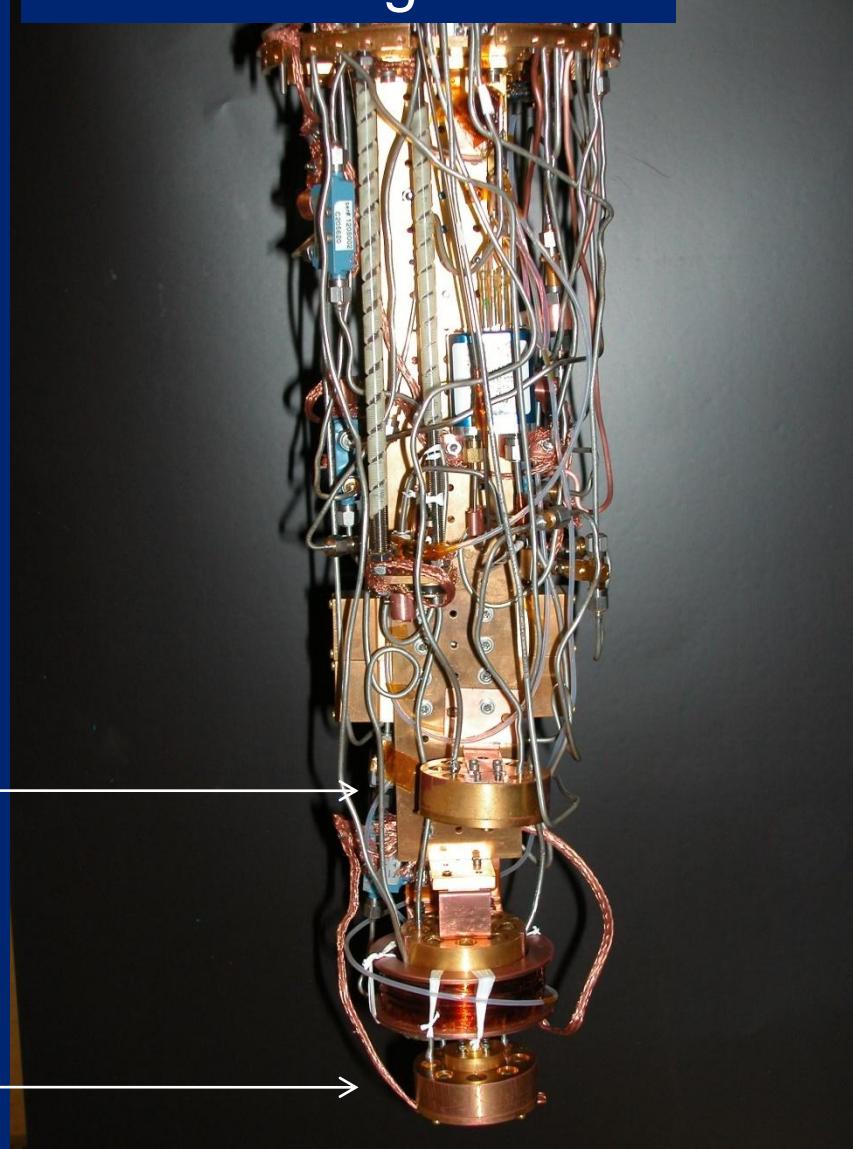


$$N_A = N_{\text{JPA}} + \frac{N_{\text{HEMT}}}{G_{\text{JPA}}}$$

Josephson parametric amplifier (JPA)
makes more photons without more entropy

Diagram conceals some complexity

Dilution refrigerator

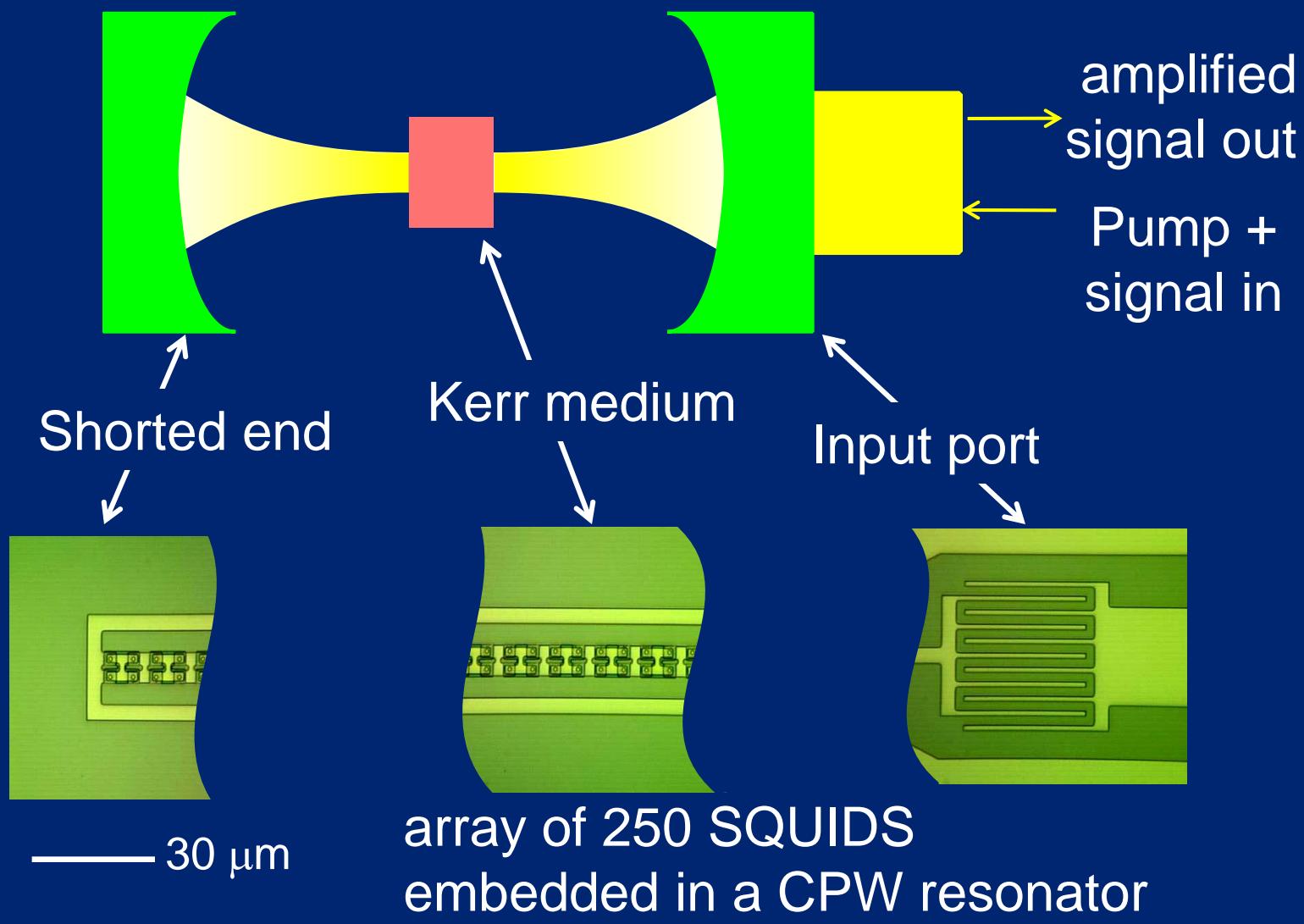


mechanics

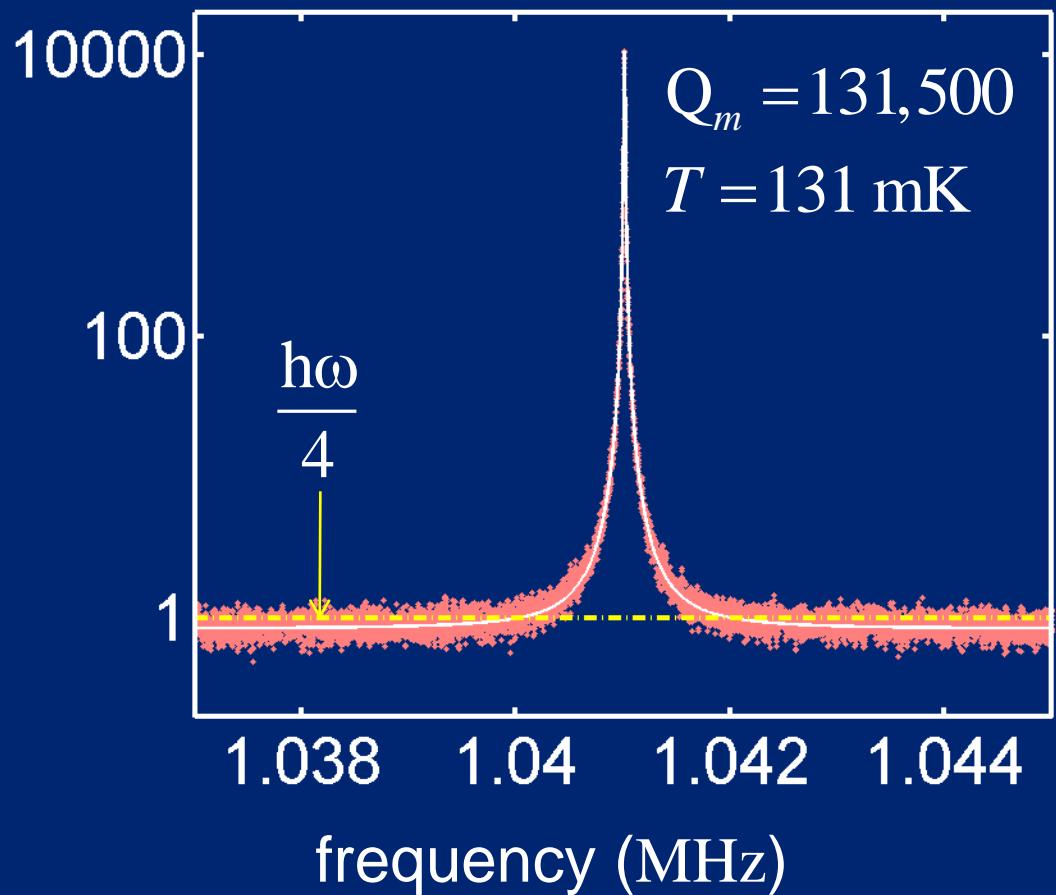
JPA

50 cm

SQUIDs embedded in a cavity form an optical parametric amplifier at microwave frequency

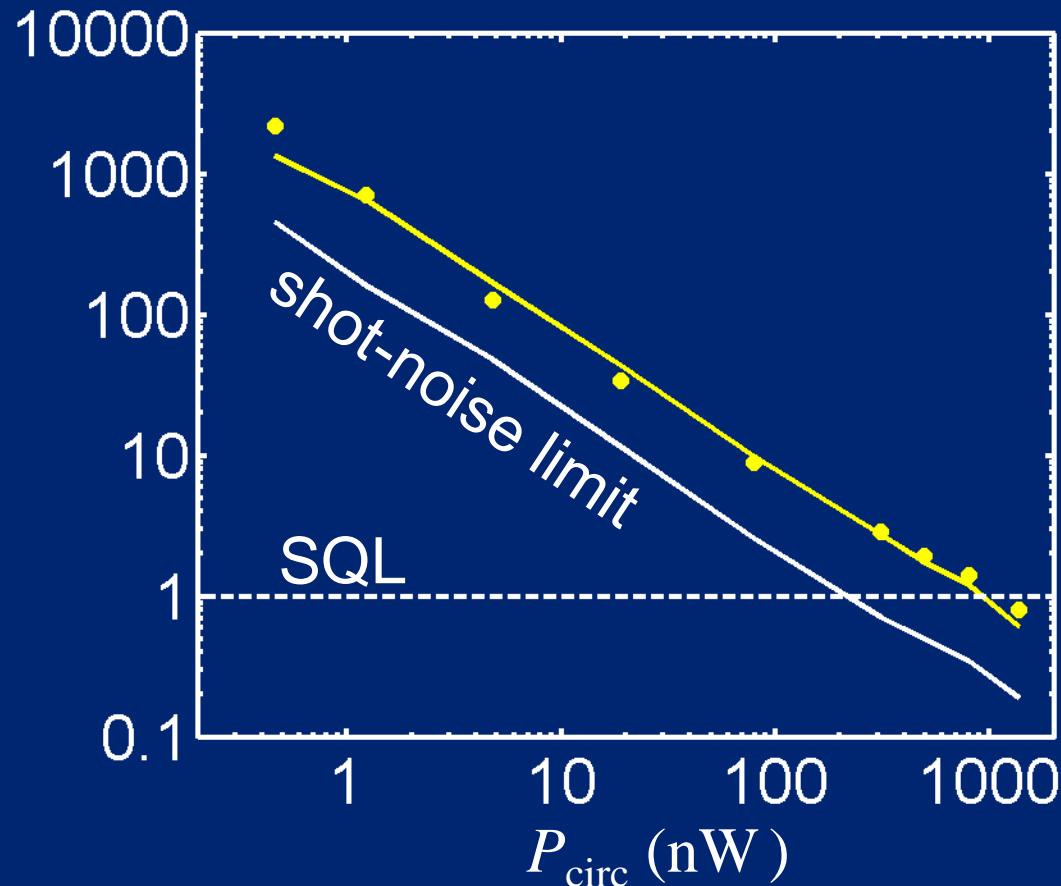


Imprecision noise is below the standard quantum limit with the JPA



$$S_x / S_x^{\text{sql}} = 0.83$$
$$\sqrt{S_x^{\text{sql}}} = 5.7 \text{ fm}/\sqrt{\text{Hz}}$$

All sources of loss and added noise yield an interferometer with 30% quantum efficiency



$$S_x^{\text{imp}} = \gamma_c^2 \frac{(N_A + \frac{1}{2})\hbar\omega_e}{g^2 P_{\text{circ}}}$$

Shot-noise limit
 $N_A = 0$

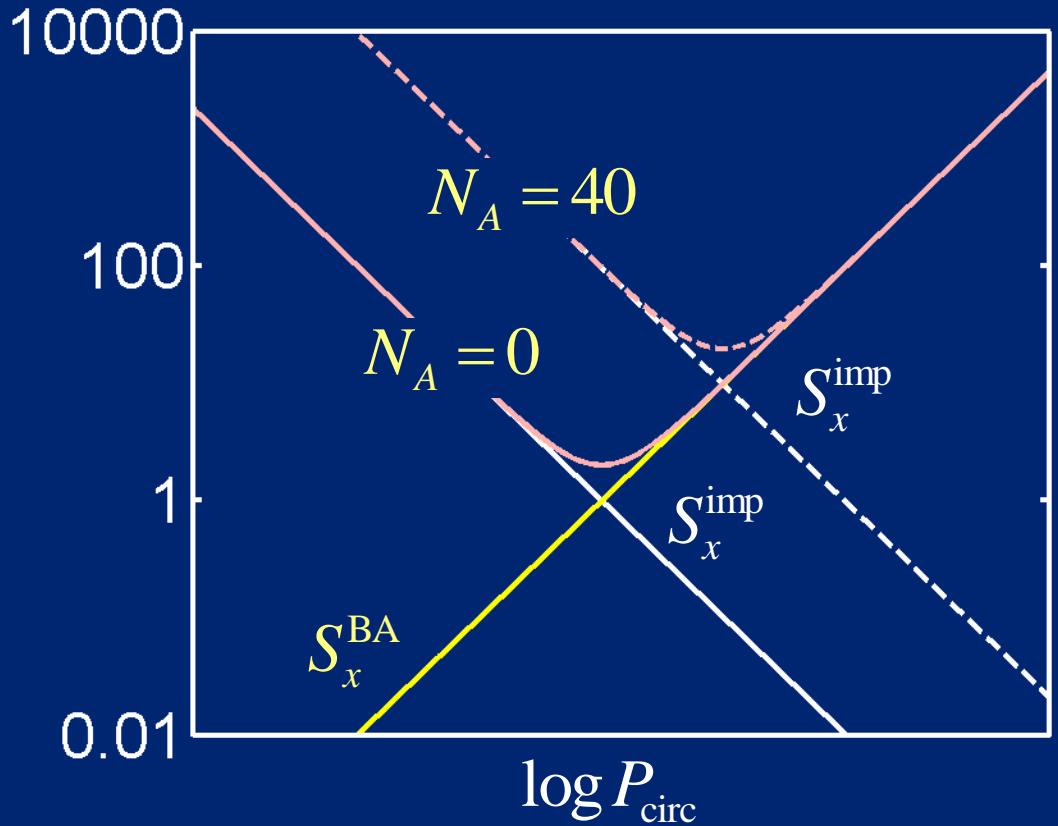
This interferometer
 $N_A = 1.3$ $\eta = 0.3$

Can we measure
zero-point motion?

YES!



The SQL is a compromise between imprecision and backaction



$$S_x^{\text{imp}} = \gamma_c^2 \frac{(N_A + \frac{1}{2})\hbar\omega_e}{g^2 P_{\text{circ}}}$$

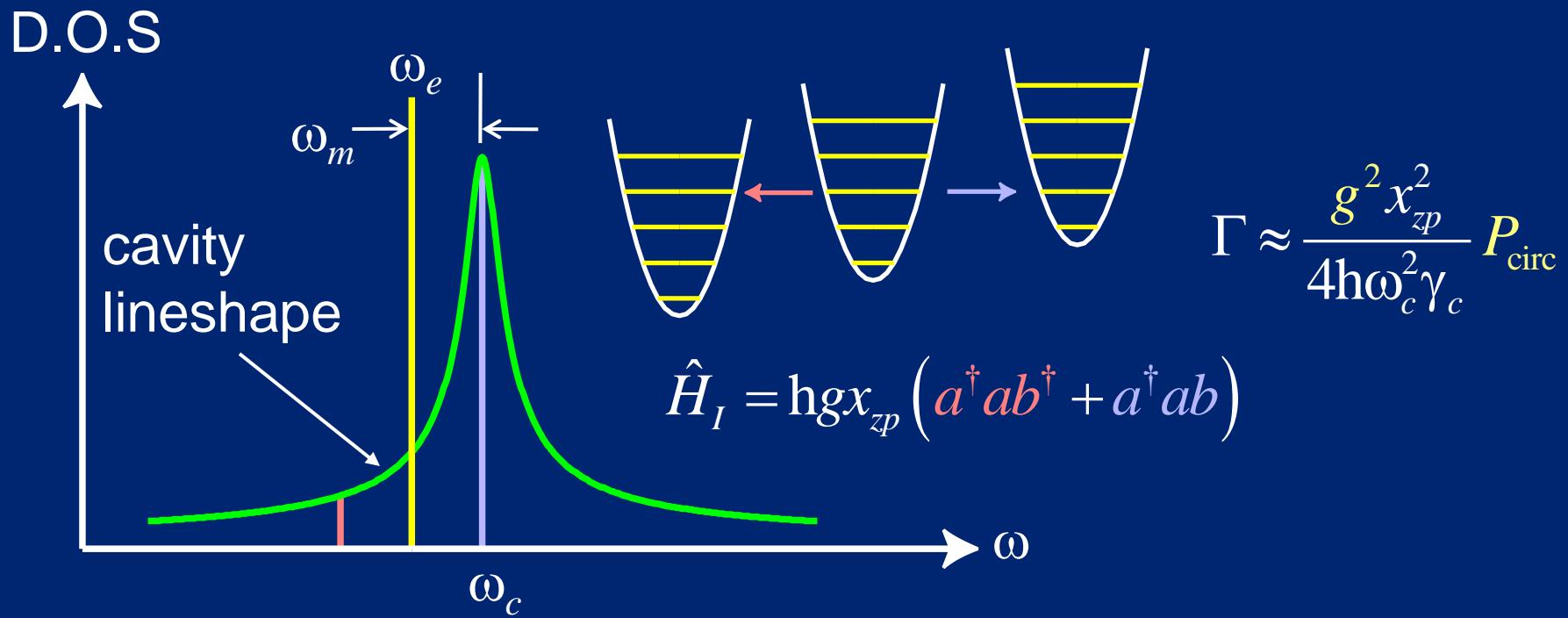
$$P_{\text{circ}}^{\text{sql}} = \gamma_c^2 \frac{(N_A + \frac{1}{2})\hbar\omega_e}{g^2 S_x^{\text{sql}}}$$

P_{circ} is constrained
-dumb heating
-superconductivity

Imprecision at the SQL $S_x^{\text{sql}} = \frac{\hbar}{m\omega_m \gamma_m}$

Radiation pressure cooling

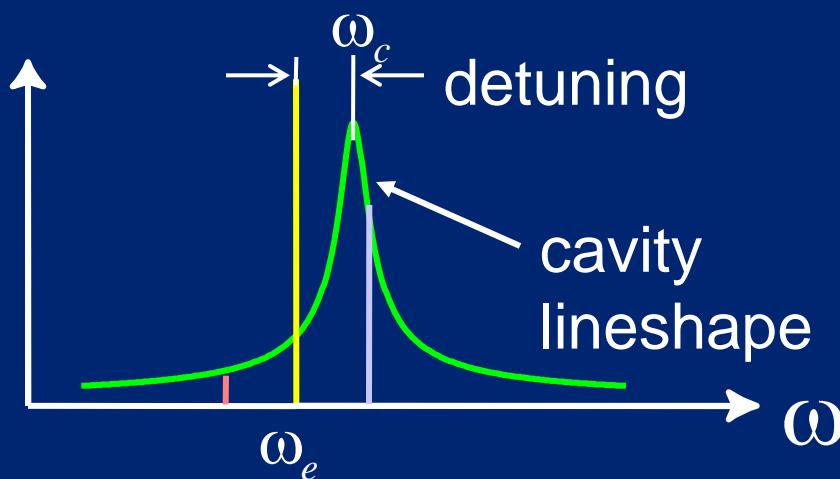
Radiation pressure can cool the beam to ground state in the resolved sideband limit



Requirements for sideband cooling to the ground state

- ✓ resolved sideband limit $\omega_m > \gamma_e$
- ✓ high frequency, high quality beams $Q_m > \frac{k_B T_{\text{bath}}}{\hbar\omega_m} = \langle n_{\text{bath}} \rangle$
- ✓ strong coupling $\Gamma > \langle n_{\text{bath}} \rangle \gamma_m^0$

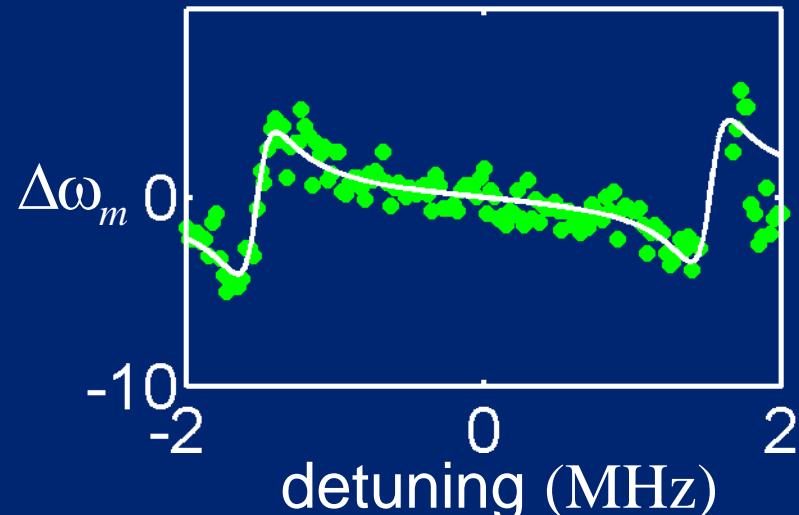
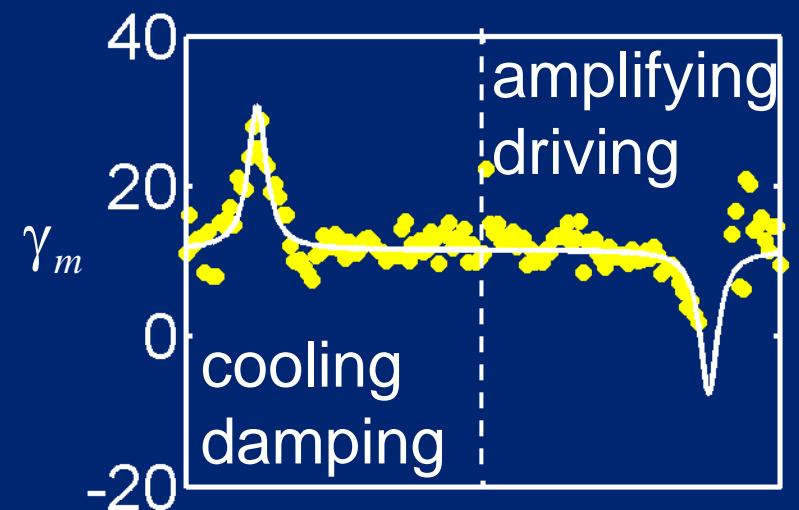
Radiation pressure changes the wire's damping rate and resonance frequency



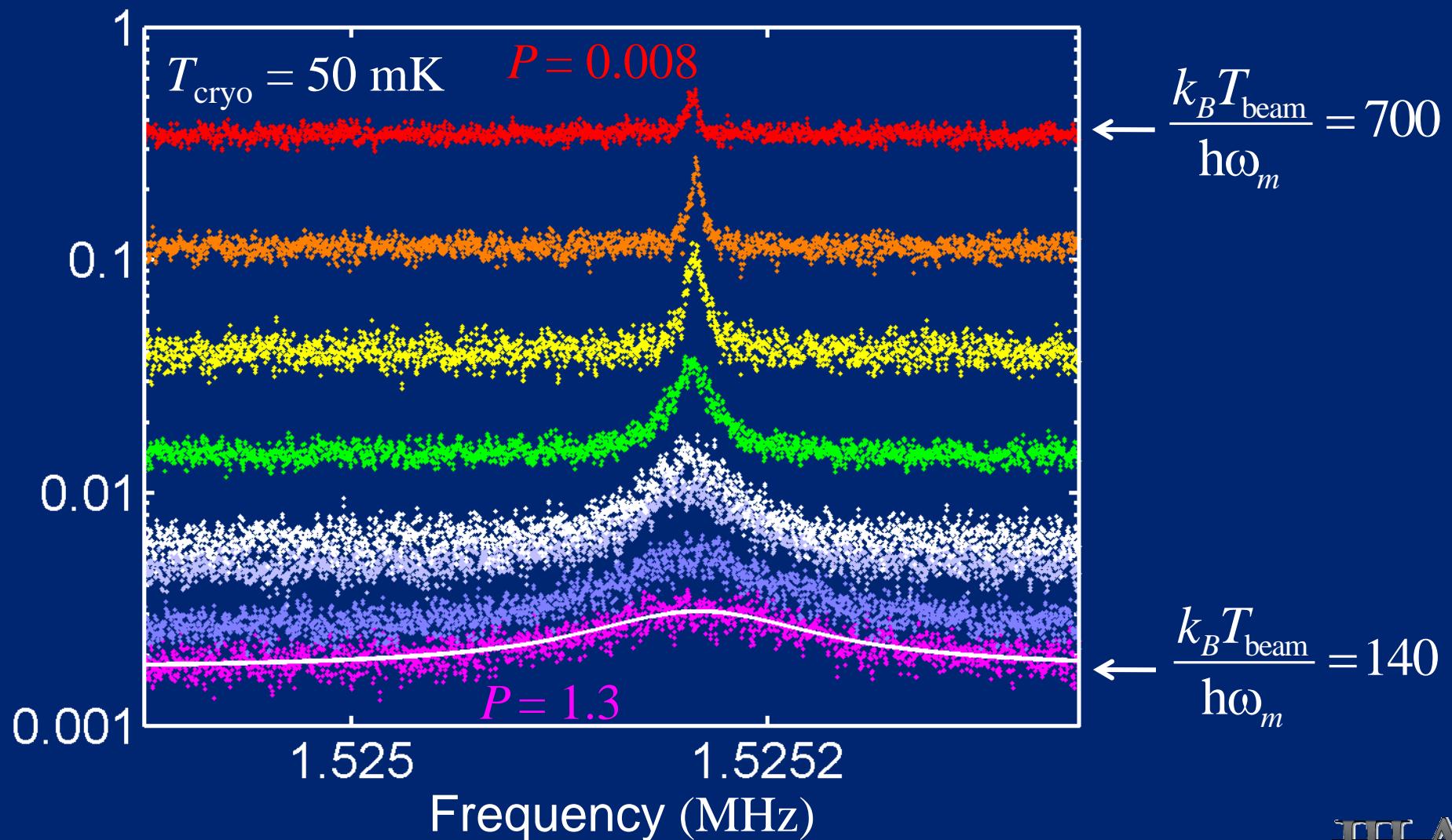
Mechanical response measures antisymmetric force noise

$$\Gamma = \frac{x_{zp}^2}{h^2} [S_f(\omega_m) - S_f(-\omega_m)]$$

$$f = \hbar g a^\dagger a \quad g = 6.4 \frac{\text{kHz}}{\text{nm}}$$

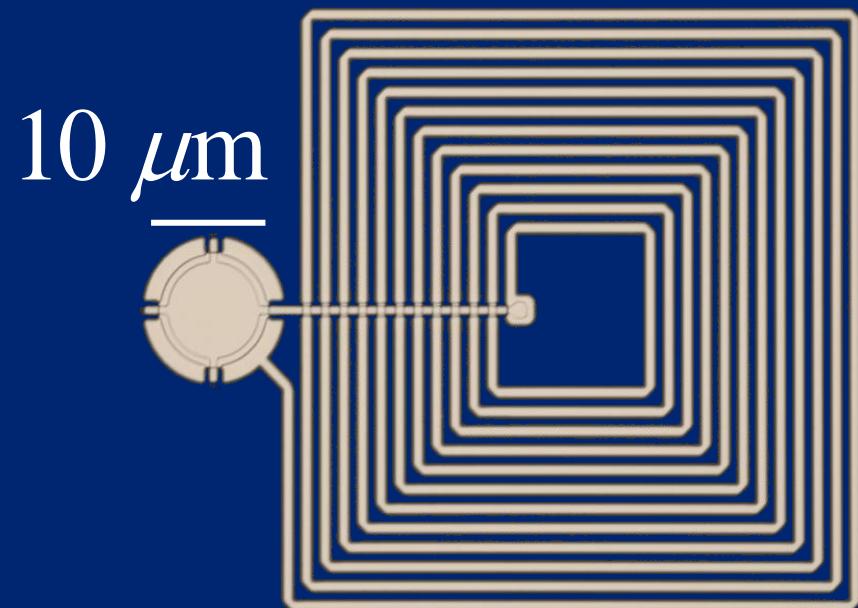
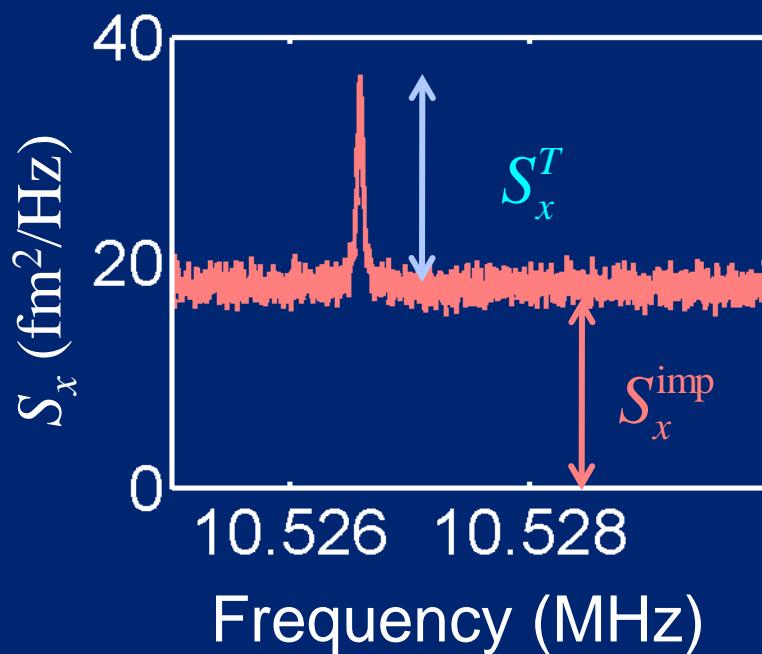
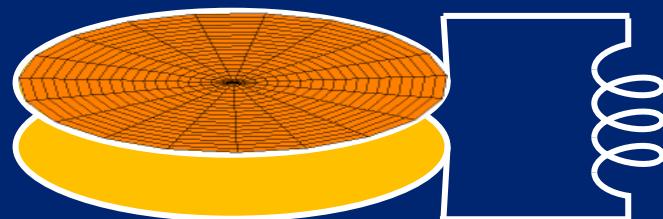


Coupling to radiation is too weak to cool wire to motional ground state



Cooling and measuring mechanical motion

capacitor built with suspended micromechanical membrane*



Electrical circuit
resonant at 7 GHz

John Teufel and
Ray Simmonds

Conclusions

- Measure and manipulate nanomechanical elements with microwaves
- Optomechanical performance: in quantum regime!
 - cooling: 0.35 phonons
 - imprecision: $0.83 \times \text{SQL}$
 - force: $0.5 \text{ aN/Hz}^{1/2}$
- Microwave Mach-Zehnder interferometer
 - quantum efficiency 30%



Graduate and post-doc positions available

Funding: NSF, NIST, DARPA, NASA