

Non-collider Detectors

NEWS-G -- continuing with gas chambers

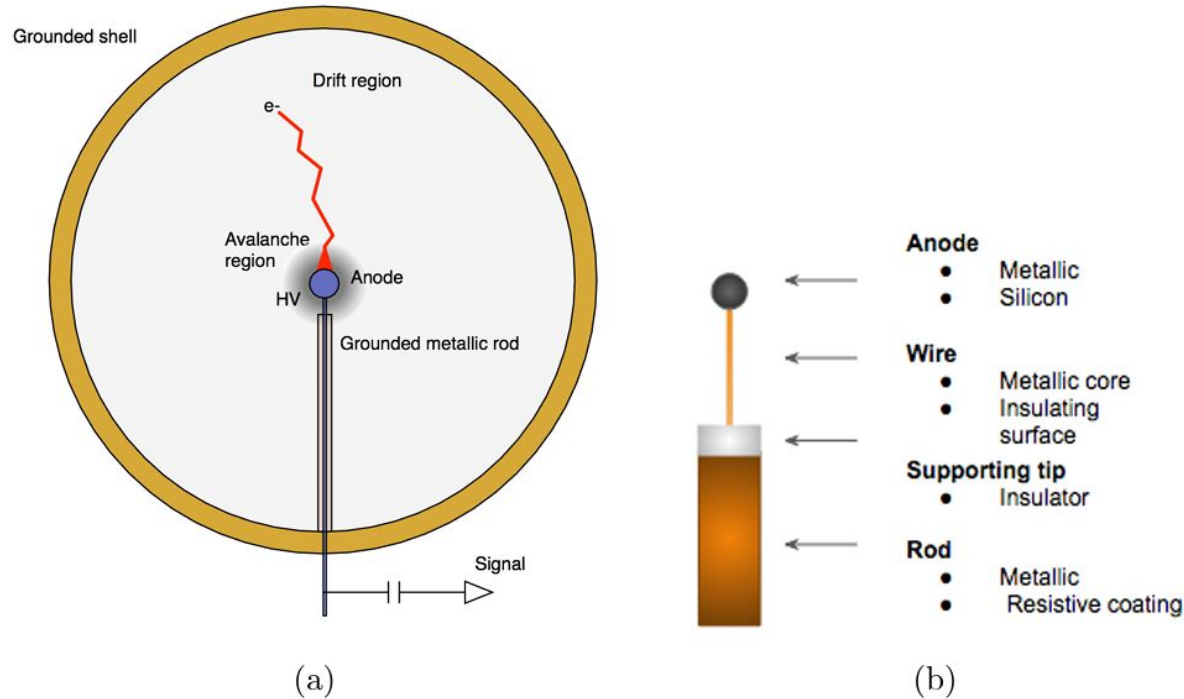
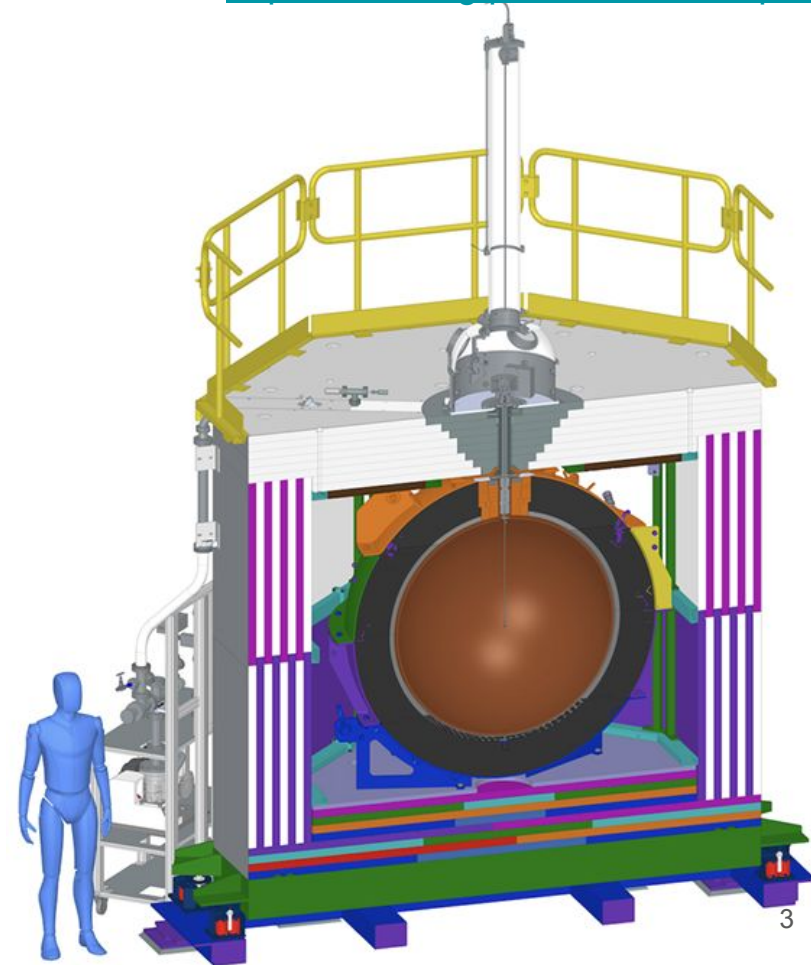
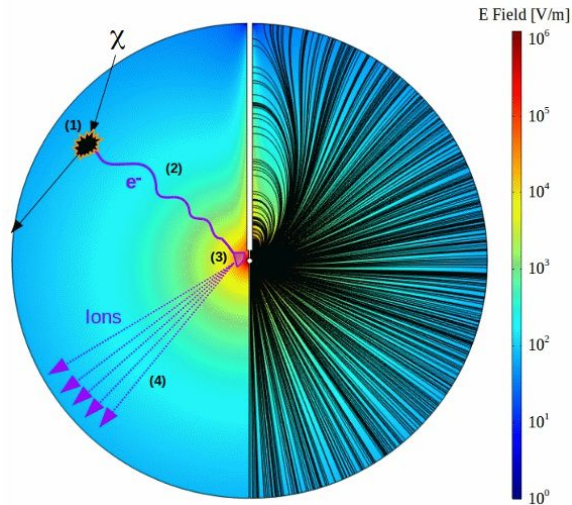


Figure 1: (a) SPC design and principle of operation and (b) illustration of the basic read out sensor⁵.

NEWS-G

- New Experiments With Spheres-Gas
- Spherical Proportional Counter
- Noble gases as targets
- Search for light dark matter
 - Down to sub-GeV mass region



NEWS-G -- Electrode

- From last lecture:
 - Mostly measuring drifting ions
 - Most ions generated near surface of the anode
 - Large E-field $E(r) \approx \frac{V}{r^2} r_A$
 - Tend to make anodes small
 - → larger 2nd ionization
- For small r_A , E field at large r becomes small as well
- Attachment & recombination becomes a problem
- → “Achinos” sensor:
 - Large r E-field determined by overall sensor shape
 - Local E-field determined by individual anode

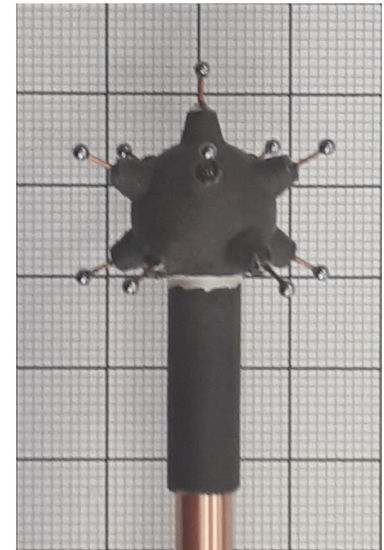
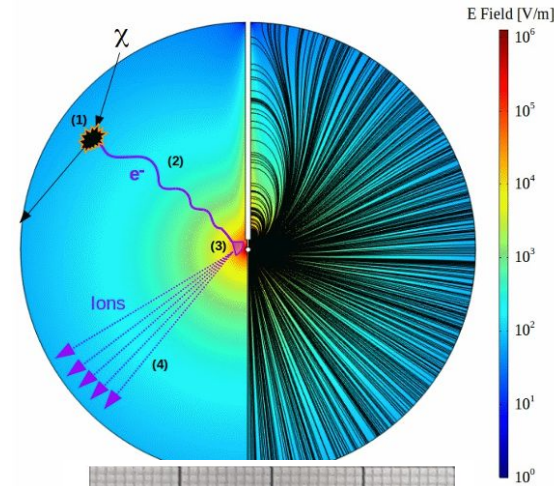


Figure 2. An 11-anode ACHINOS sensor.

Shockley–Ramo theorem

The Shockley–Ramo theorem states that the instantaneous current i induced on a given electrode due to the motion of a charge is given by:

$$i = E_v q v$$

where

q is the charge of the **particle**;

v is its instantaneous **velocity**; and

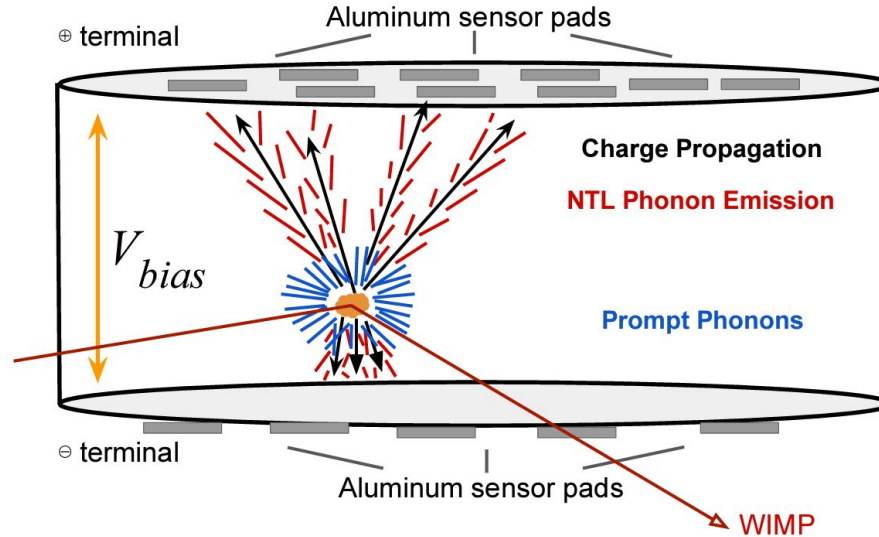
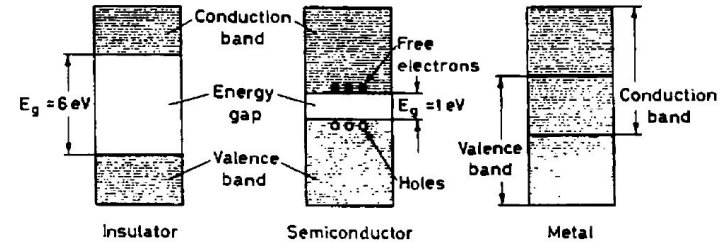
E_v is the component of the **electric field** in the direction of v at the charge's instantaneous position, under the following conditions: charge removed, given electrode raised to unit potential, and all other conductors grounded.

Can verify expression in last lecture

INDUCED SIGNAL	$dI = \frac{Q}{C \epsilon V_0} \frac{dV}{dr} dr.$
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Take away: The signal measured is charge-drifting induced signal...

Semiconductor detectors: Charge and phonon readouts

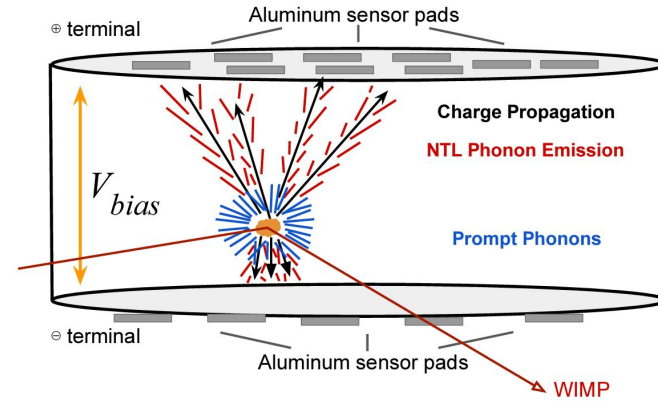


$$\begin{aligned}
 \text{Phonon energy} &= E_{\text{recoil}} + E_{\text{Luke}} \\
 &= E_{\text{recoil}} + n_{\text{eh}} e^- \Delta V
 \end{aligned}$$

Semiconductor detectors

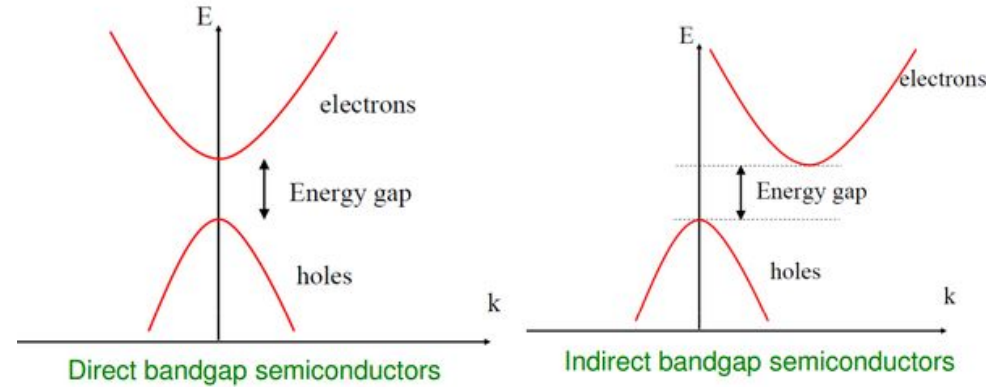
Microscopic interactions

- Detector crystal is voltage biased
- Initial recoil generates electron-hole pairs
- Electrons and holes gain initial kinetic energy
 - Immediately emit as phonons -- NTL effect
- Electrons and holes drift across semiconductor
 - Drift speed depends on E-field and temperature
 - ~ 5000 cm/s at low temperature, ~ 1 V/cm
 - Notice electrons don't follow electric field
 - Tensor mass of electron in semiconductor
 - All ionization happens at \sim initial recoil
 - Tend not to induce avalanche (except Avalanche PhotoDiodes/SiPM)
 - Emit phonons as they go
- Charge measured by Shockley–Ramo theorem, $E_{\text{Phonon}} = E_{\text{recoil}} + n_{\text{eh}} e^{-} \Delta V$



Direct band-gap vs indirect band-gap

- Energy and momentum both need to be conserved in a transition
- During absorption, momentum compensation can be via phonon absorption (room temp.) or phonon emission (low temp.)
- For direct band-gap materials, excited electron can recombine by **photon** emission
 - → Materials for scintillator detectors / LED (InAs, GaAs)
- Indirect band-gap materials don't emit photons
 - → Charge detectors (Si, Ge)



Charge amplifier

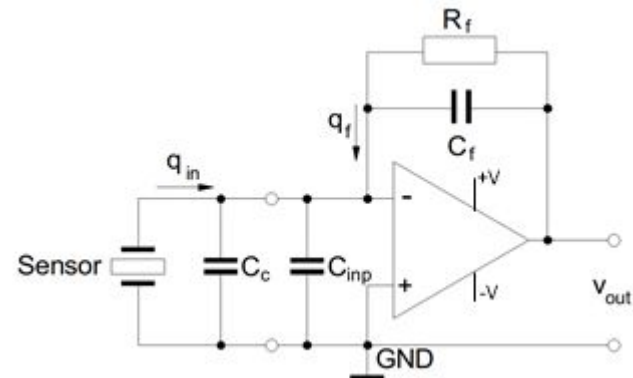
- Detector design is often integrated with readouts
- Charge detector:
 - Collecting charge on its capacitance
 - Amplifying via “charge amplifier”
 - Discharge via feedback resistor

$$q_{in} = q_f$$

$$V_{out} = q_f / C_f = q_{in} / C_f$$

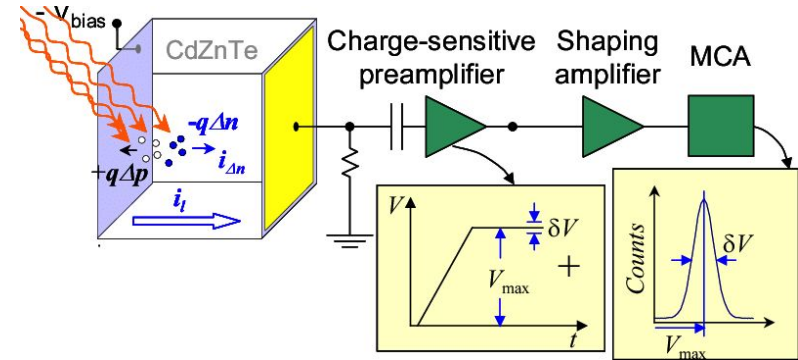
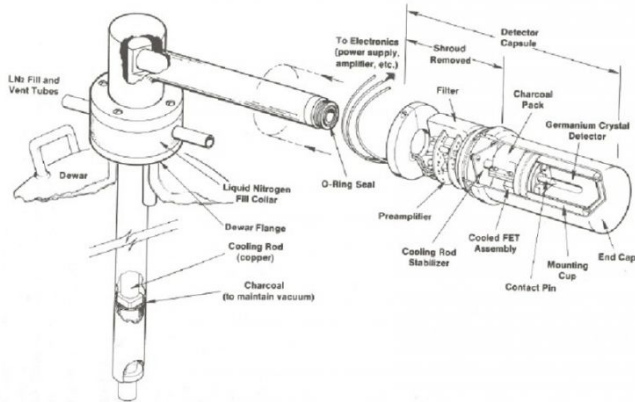
$$\text{Decay time: } R_f * C_f$$

- Detector capacitance critical for noise
 - Voltage noise (v_n) amplified by C_{in} / C_f



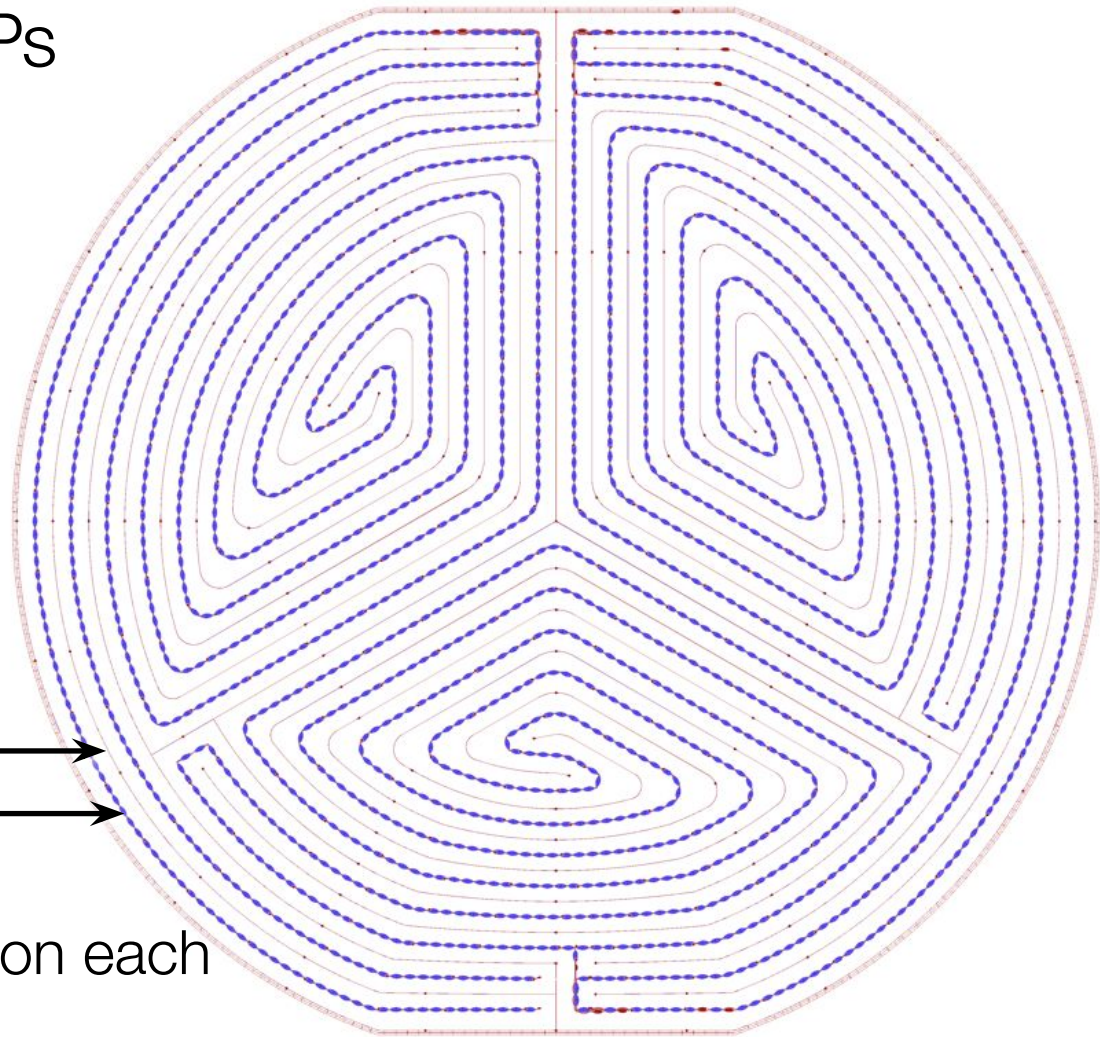
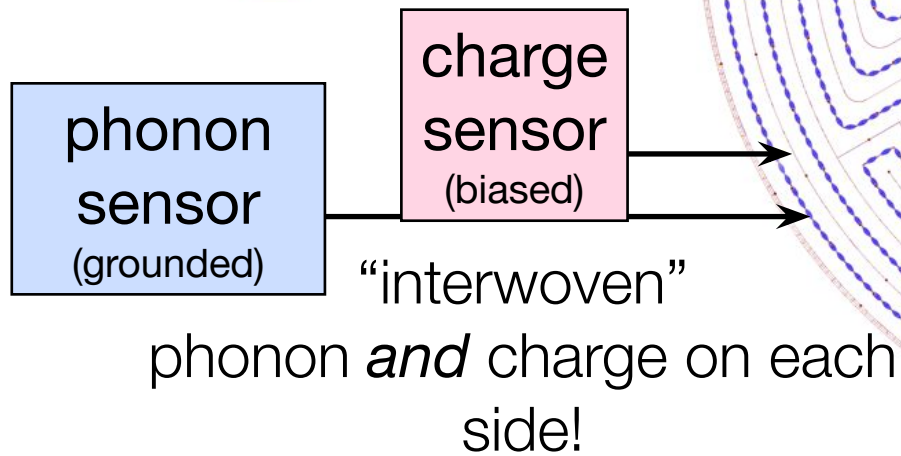
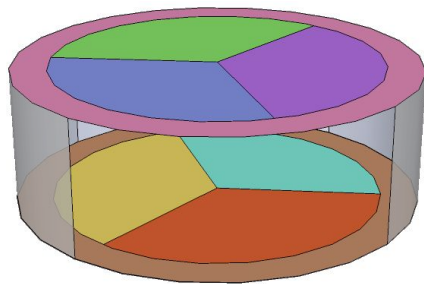
Semiconductor charge detector examples

- High purity germanium (HPGe)
 - Work at liquid Nitrogen temperature
 - Can cool down electronics
 - Resolution \sim keV
- CdTe/CdZnTe
 - Work at room temperature
 - Resolution \sim a few keV



SuperCDMS Soudan iZIPs

8 phonon channels +
4 charge sensors =
Lots of information per event!



SuperCDMS Soudan iZIP

- Minimize C_D and C_P
- Cool down first stage amplifier (JFET)
 - Next generation uses HEMT, operating at 1 Kelvin

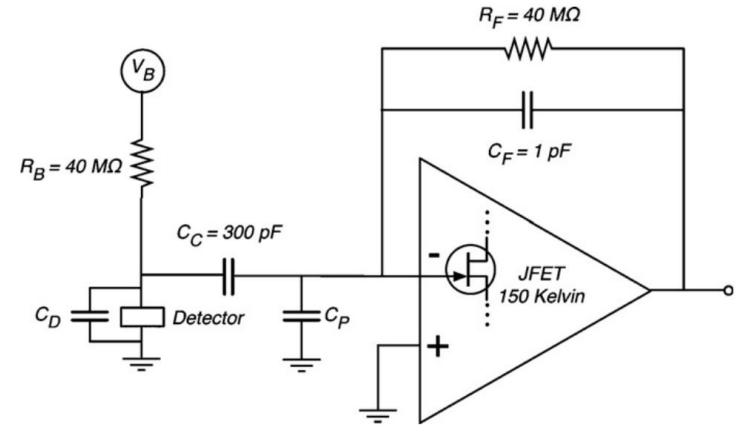
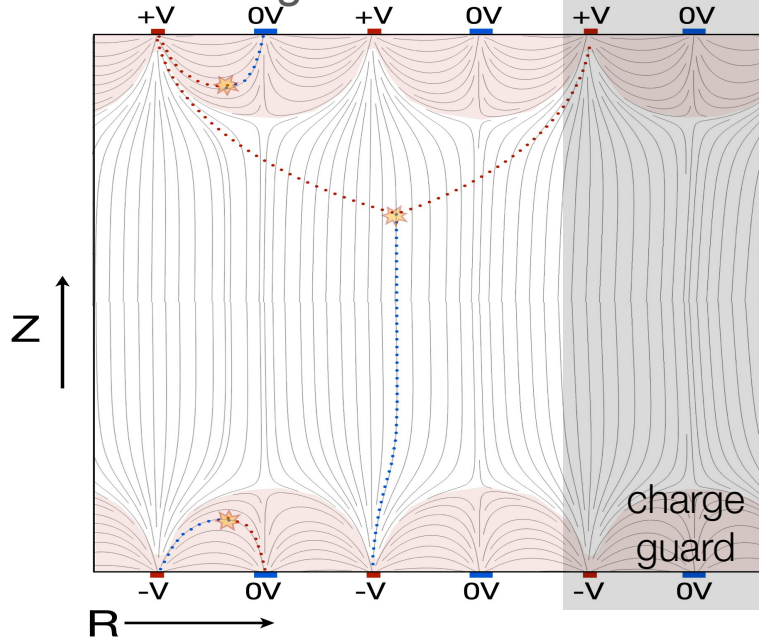


Fig. 5. Charge readout circuit. The heart of the amplifier is a JFET. The parasitic capacitance has been measured to be $C_P = 100$ pF and the detector capacitance is typically $C_D = 50$ pF. G denotes the gate of the JFET.

Ricochet HEMT readout

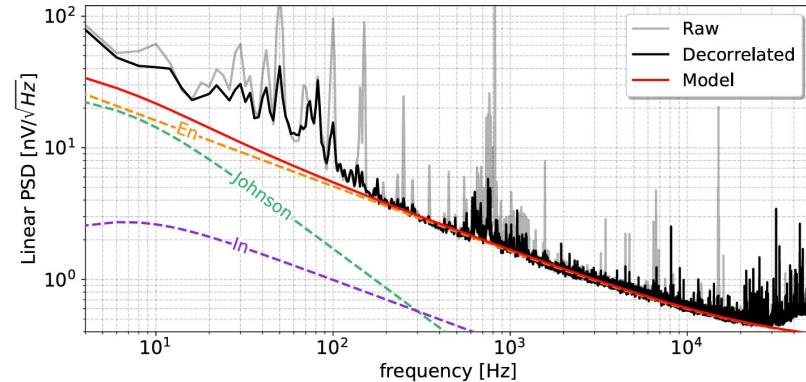
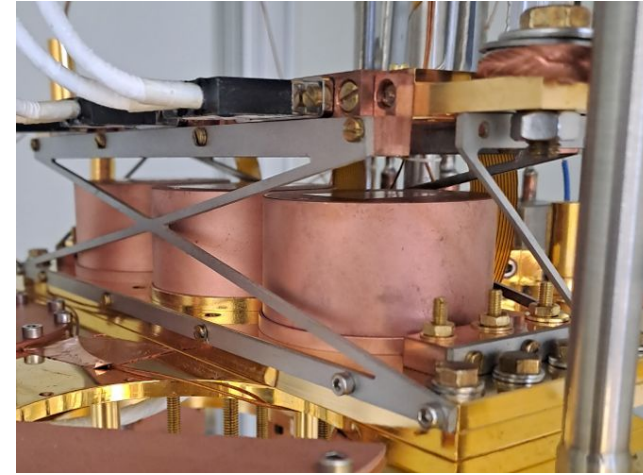
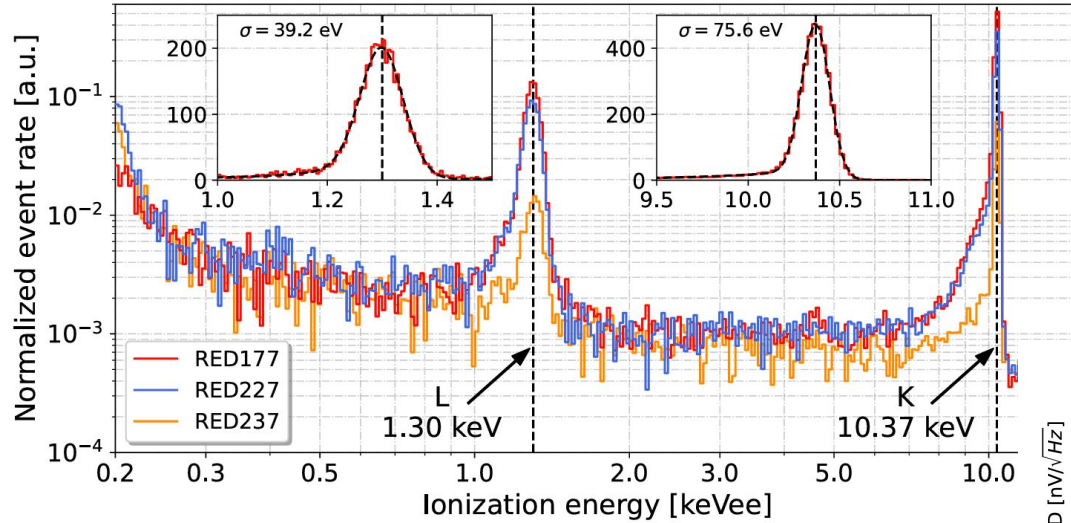
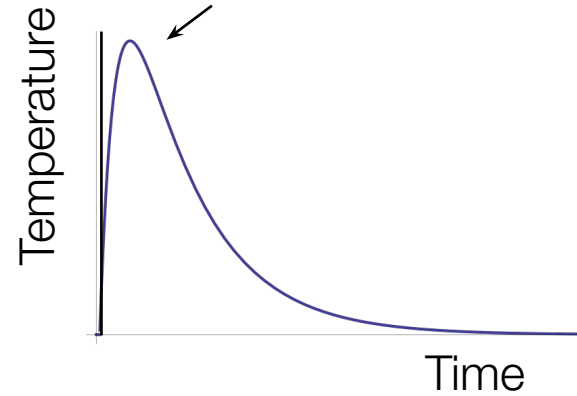
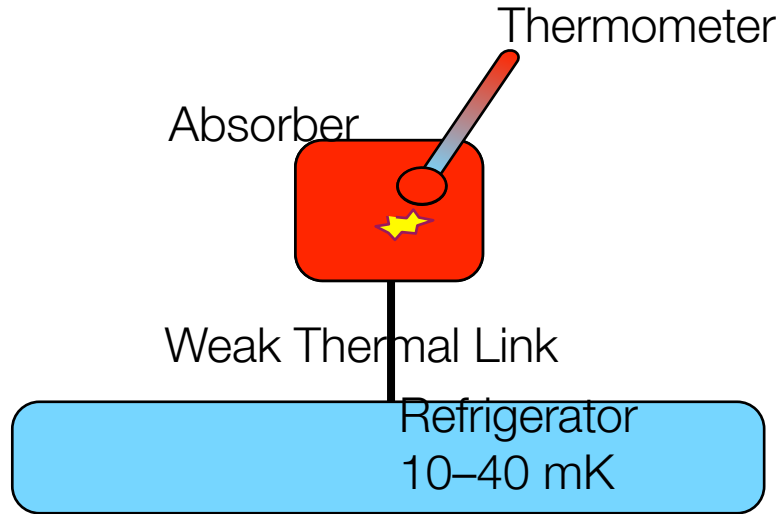


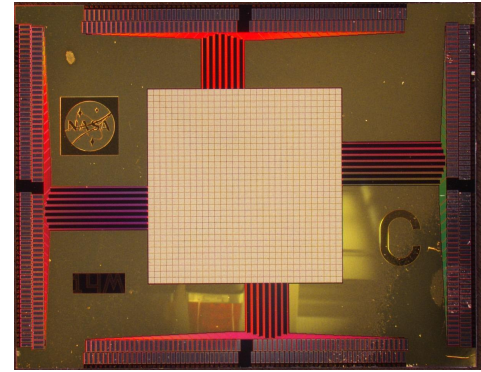
Fig. 2 Differential ($V_B - V_A$) noise power spectrum of RED227 with $V_{ds} = 100$ mV and $I_{ds} = 300$ μ A and the mixing chamber at 17 mK. The black, grey and red solid lines show the decorrelated and raw data, and our noise model considering a parasitic capacitance of 20 pF, respectively. The total contributions from the two ionization channels A and B of the current (I_n), voltage (E_n), and Johnson noise sources are also shown as dashed purple, orange, and green lines, respectively.

Measuring phonons: Cryogenic Crystal Detectors



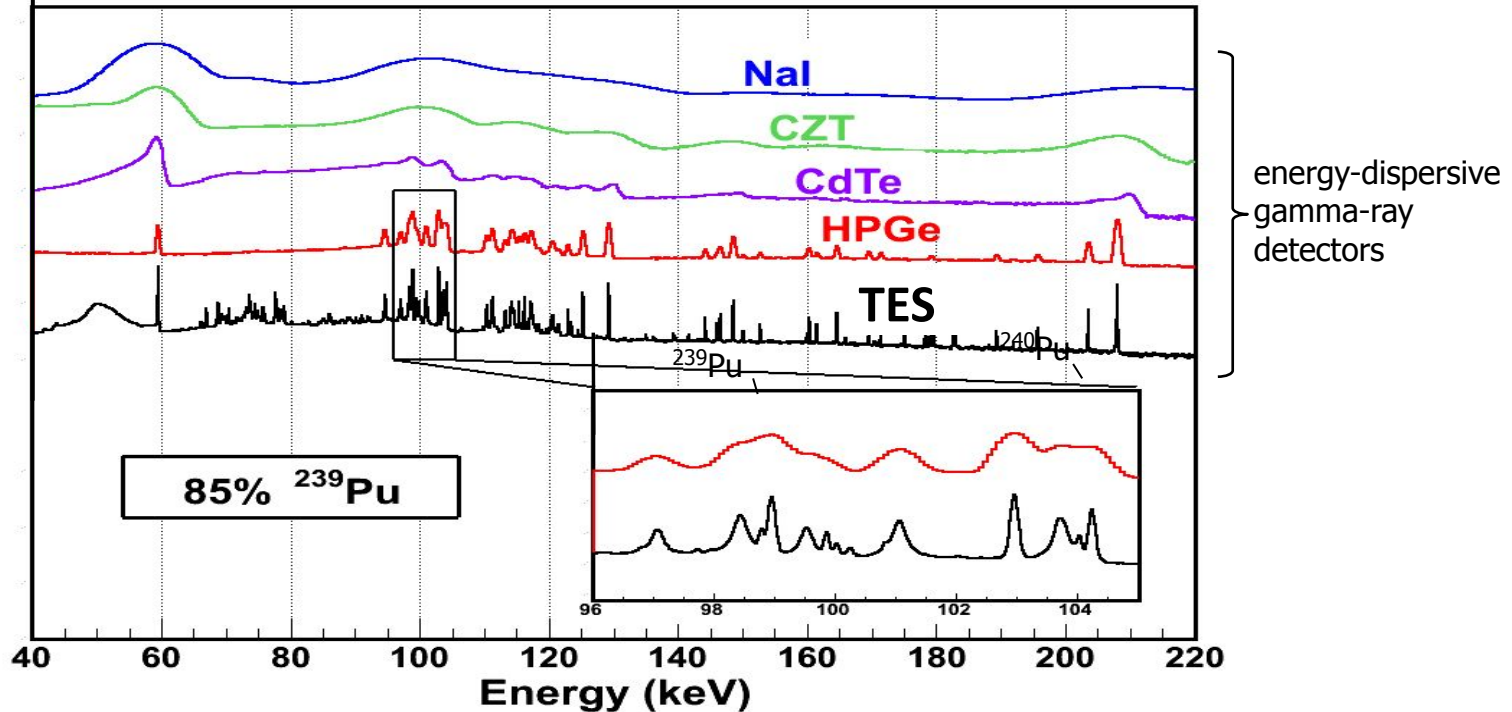
Cryogenic Crystal Detectors are used in...

- Astrophysics
 - mm to gamma-ray energies
- Particle Physics
 - Dark Matter Detectors
 - Neutrino Physics
- Materials-analysis
- Others!

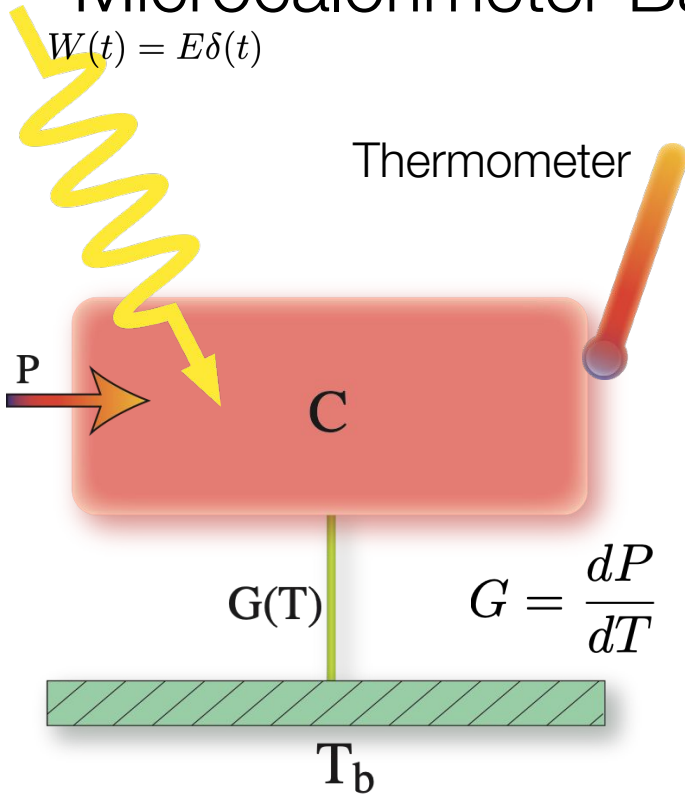


Why Use Cryogenic Detectors?

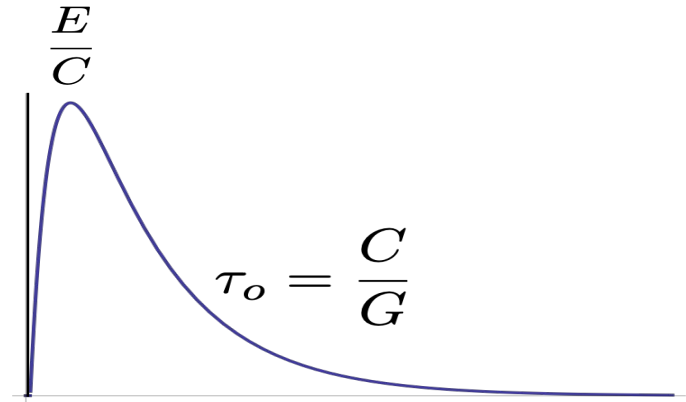
Cryogenic microcalorimeters can provide a unique combination of energy sensitivity and efficiency



Microcalorimeter Basics



$$C \frac{dT}{dt} = P - G(T - T_b) + W(t)$$



$$T(t) = \frac{E}{C} e^{t/t_o} + \left(\frac{P}{G} + T_b \right)$$

Some types of thermometers:

- resistive
- capacitive
- inductive
- paramagnetic
- thermoelectric

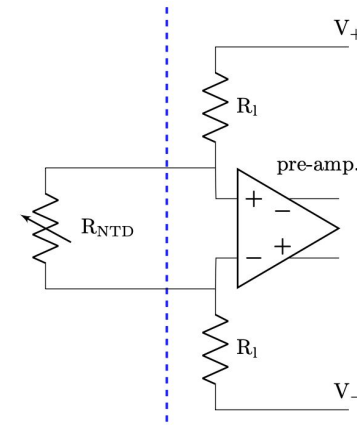
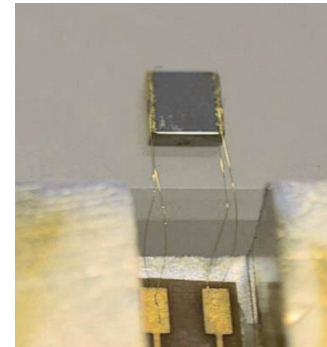
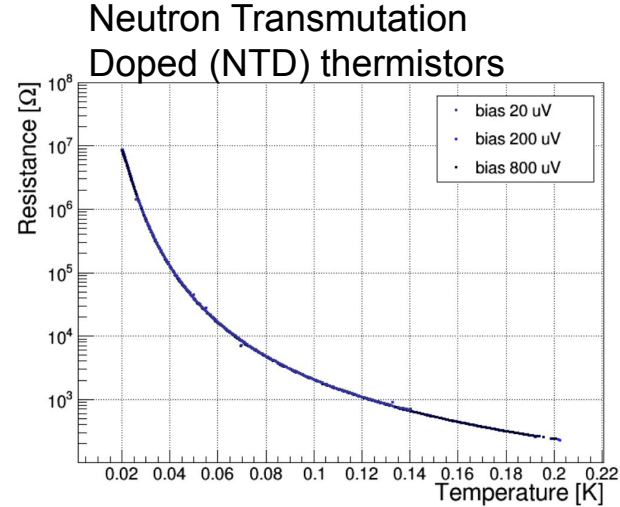


Neutron Transmutation Doped (NTD) thermistors

- Doped germanium/silicon chips
- Resistance follows Efros-Shklovskii law:

$$R = R_0 e^{\sqrt{T_0/T}}$$

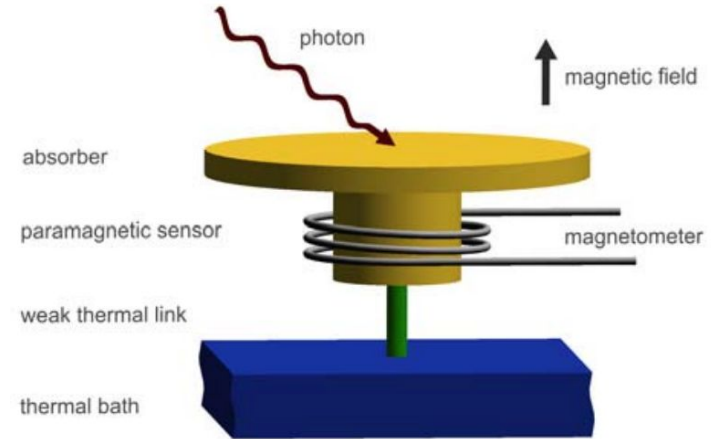
- Taking advantage of the steep slope at low temperature
- Also comes with high dynamic range
- Readout with FETs, operating at room temperature or in cold



Metallic magnetic calorimeter (MMC)

- Paramagnetic sensor positioned in weak magnetic field
- Heat changes its induced magnetic field
- Readout by SQUIDs as magnetometer

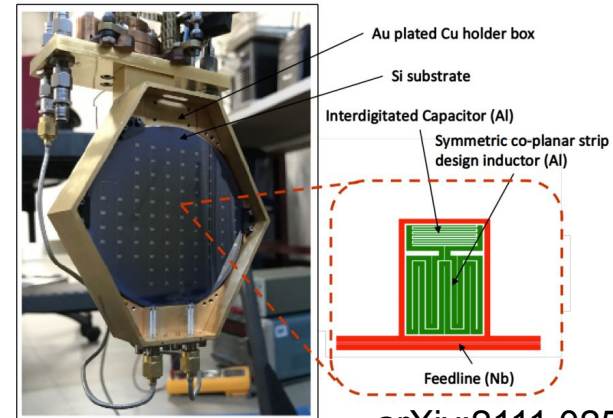
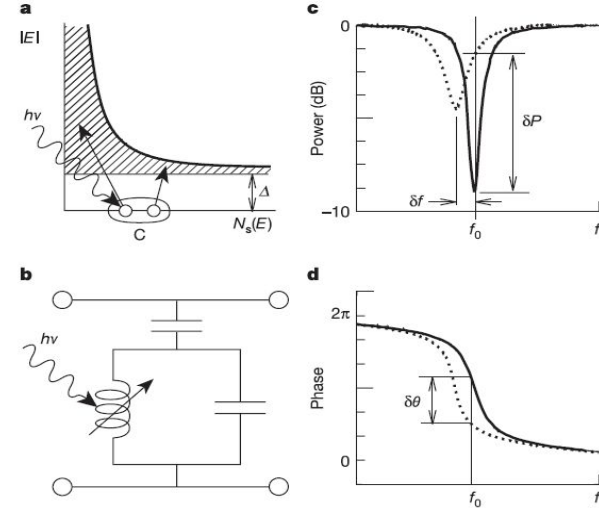
Metallic magnetic calorimeter (MMC)



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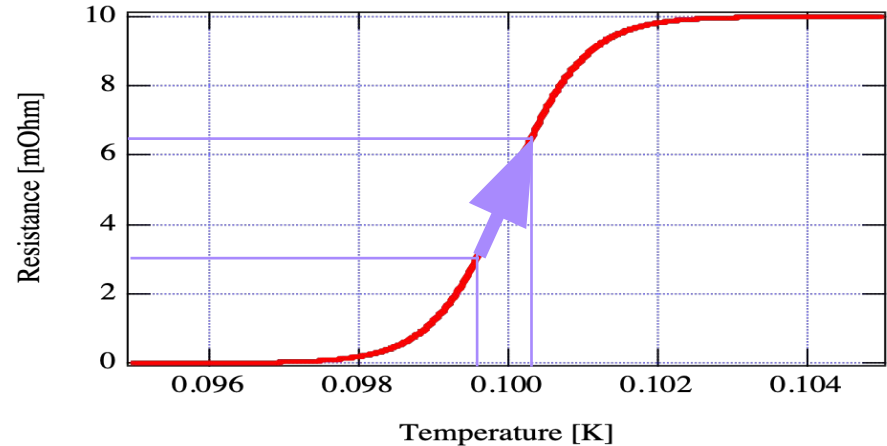
Microwave Kinetic Inductance Detectors (MKIDs)

- Resonators made of superconducting metal films
- Resonance frequency and phase response depending on its temperature
- Radio-Frequency (RF) Readout system
- Intrinsic capability for multiplexing



Transition-Edge Sensors

- Superconductor biased in its transition
- Several metal systems are used:
 - Elemental: W, Al, Re, Pb, etc.
 - Paramagnetic impurity doped: Al/Fe, Al/Mn, etc.
 - Bi-layers: Mo/Au, Mo/Cu, Ti/Al, etc.
- $50 < \alpha < 1000$
- Low resistance allows read out with SQUIDs



$$\alpha = \frac{T}{R} \frac{dR}{dT}$$

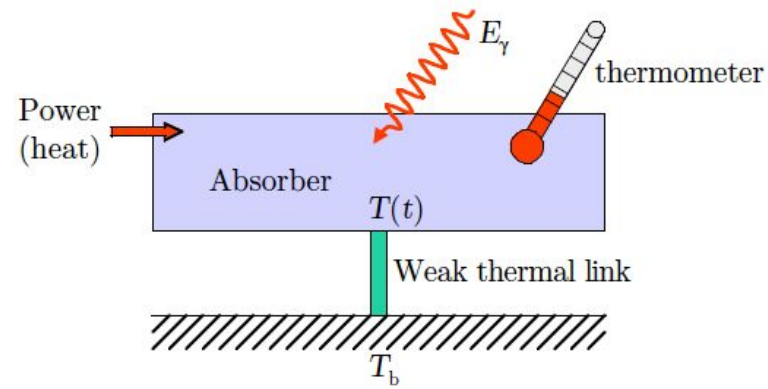
Balance the heat

- Sensor heat capacity C
- P is heat flowing into the sensor
 - For now assume as a constant
- Energy deposition E_γ at t_γ

$$C \frac{dT(t)}{dt} = P - P_{\text{link}}(T(t), T_b) + E_\gamma \delta(t - t_\gamma)$$

$$P_{\text{link}}(T(t), T_b) = G(T(t) - T_b)$$

$$T(t) = \frac{E_\gamma}{C} e^{-t/\tau_o} + \left(\frac{P}{G} + T_b \right) \quad \tau_o \equiv C/G$$

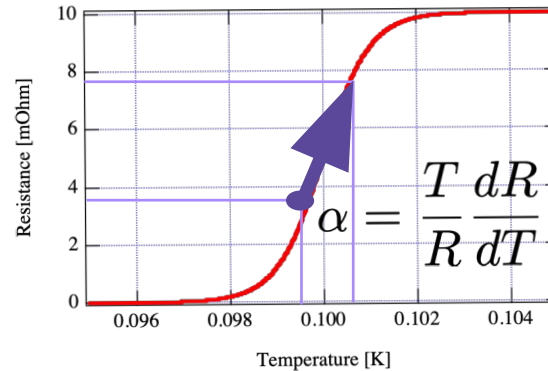
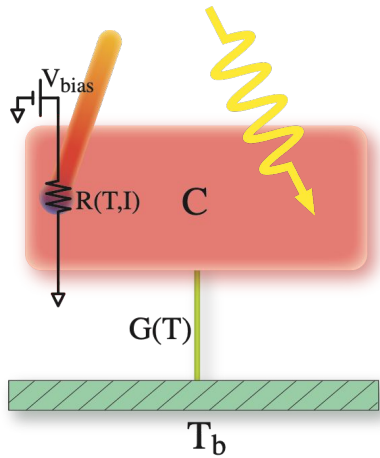


Electro-thermal Feedback (ETF)

$$P = \frac{V_{\text{bias}}^2}{R(T)}$$

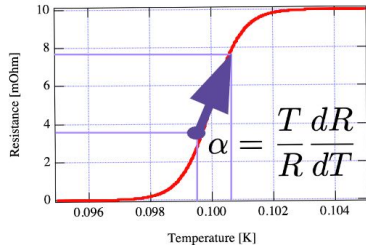
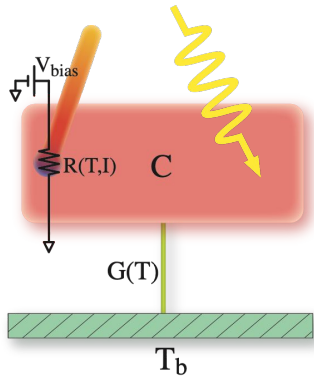
$$C \frac{dT(t)}{dt} = P - P_{\text{link}}(T(t), T_b) + E_{\gamma} \delta(t - t_{\gamma})$$

$$P_{\text{link}}(T(t), T_b) = K(T(t)^n - T_b^n)$$



Electro-thermal Feedback (ETF)

$$P = \frac{V_{\text{bias}}^2}{R(T)}$$



$$C \frac{dT(t)}{dt} = \frac{V^2}{R(T)} - K(T(t)^n - T_b^n) + E_\gamma \delta(t - t_\gamma)$$

- In quiescence:

$$P = \frac{V^2}{R} = K(T^n - T_b^n)$$

- Note, in quiescence $T \sim T_c$. Thus P is roughly constant
- Taylor expand with ΔT , note R_0 , T_0 are the quiescent values

$$R = R_0 + \alpha \frac{R_0}{T_0} \Delta T$$

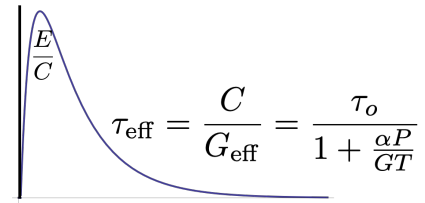
$$C \frac{d\Delta T}{dt} = \frac{V^2}{R_0} - K(T_0^n - T_b^n) - \frac{V^2}{R_0^2} \frac{dR}{dT} \Delta T - nKT_0^{n-1} \Delta T + E_\gamma \delta(t - t_\gamma)$$

- Define

$$G \equiv \frac{dP}{dT} = nKT^{n-1}$$

$$\Delta \dot{T}(t) = - \left(\frac{\alpha P}{TC} + \frac{G}{C} \right) \Delta T + \frac{E_\gamma}{C} \delta(t - t_\gamma)$$

Electro-thermal Feedback (ETF)



$$P = \frac{V_{\text{bias}}^2}{R(T)}$$

$$\Delta \dot{T}(t) = -\left(\frac{\alpha P}{TC} + \frac{G}{C}\right) \Delta T + \frac{E_{\gamma}}{C} \delta(t - t_{\gamma})$$

- Solution is a simple exponential, with

$$\tau_{\text{eff}} = \frac{\tau_0}{1 + \frac{\alpha P_0}{T_0 G}} = \frac{\tau_0}{1 + \frac{\alpha}{n} \left(1 - \left(\frac{T_b}{T_0}\right)^n\right)}$$

- Define

$$\phi = 1 - \left(\frac{T_b}{T_0}\right)^n \quad \tau_{\text{eff}} = \frac{\tau_0}{1 + \frac{\alpha \phi}{n}}$$

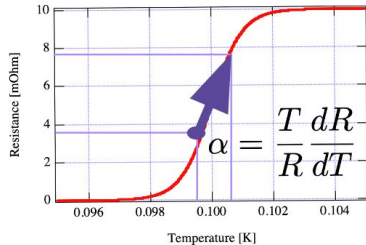
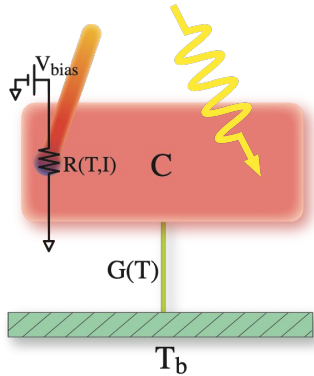
- In “extreme electrothermal feedback regime”

- Cold fridge $\rightarrow T_0^n \ll T_b^n$
- Excellent detector $\rightarrow \alpha/n \gg 1$

$$\tau_{\text{eff}} = \frac{n\tau_0}{\alpha}$$

- Higher $\alpha \rightarrow$ Faster detector

- Will see later that higher α also leads to better detector



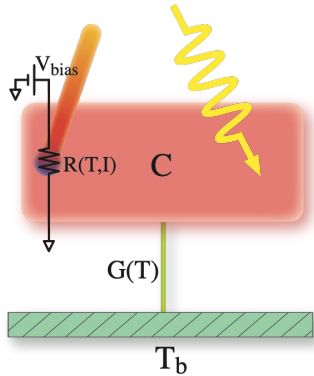
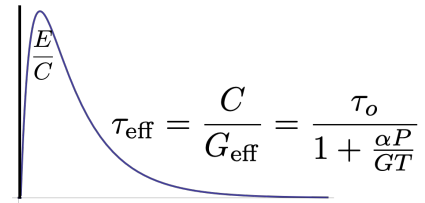
ETF as virtual conductance

$$P = \frac{V_{\text{bias}}^2}{R(T)}$$

$$G_{\text{ETF}} \equiv \frac{\alpha P}{T}$$

$$\tau_{\text{ETF}} \equiv \frac{C}{G_{\text{ETF}}} = \frac{C}{\frac{\alpha P}{T}} = \frac{\alpha P}{TG} \tau_o$$

$$\tau_{\text{eff}} = \frac{C}{G_{\text{eff}}} = \frac{C}{G + G_{\text{ETF}}} = \frac{C/G}{1 + \frac{\alpha P}{TG}} = \frac{1}{\frac{1}{\tau_o} + \frac{1}{\tau_{\text{ETF}}}}$$

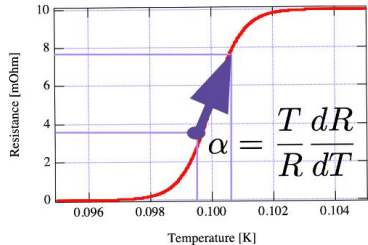


- ETF acts like another conductance, in parallel with the natural thermal conductance
- For ETF to be stable

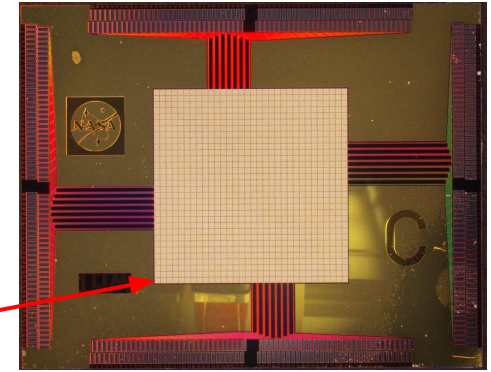
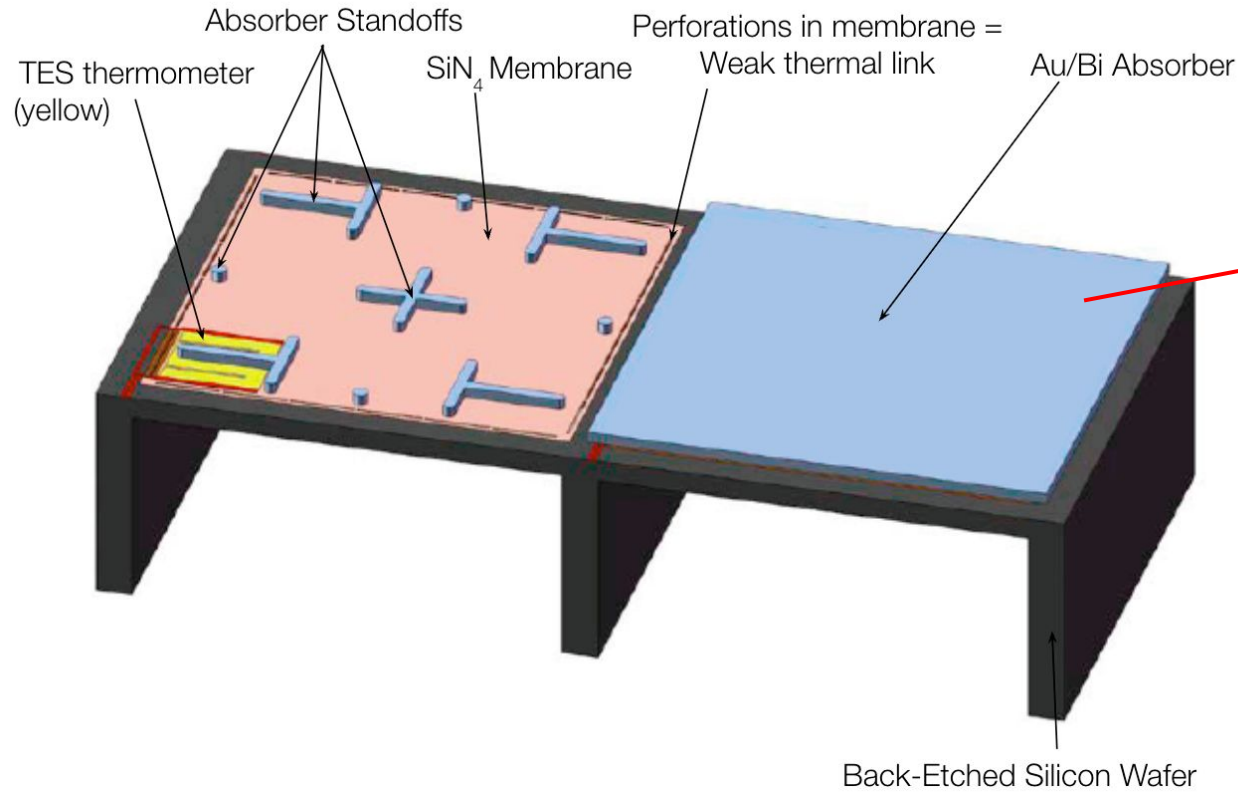
$$\frac{\alpha P}{TG} > -1$$

Naturally satisfied for TES, flipped for NTD

- → TES needs to be “Voltage biased”
NTD needs to be “current biased”



Example microcalorimeter -- MicroX



Example microcalorimeter -- MicroX

