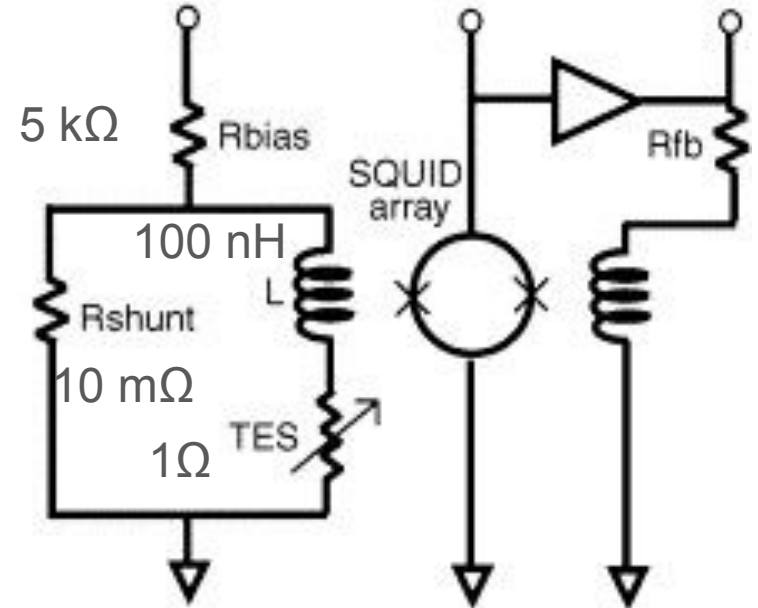
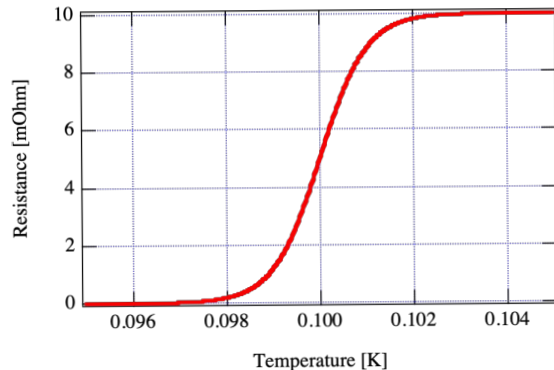


# Transition Edge Sensor Modeling

Following “Transition-Edge Sensors, K.D. Irwin and G. C. Hilton”

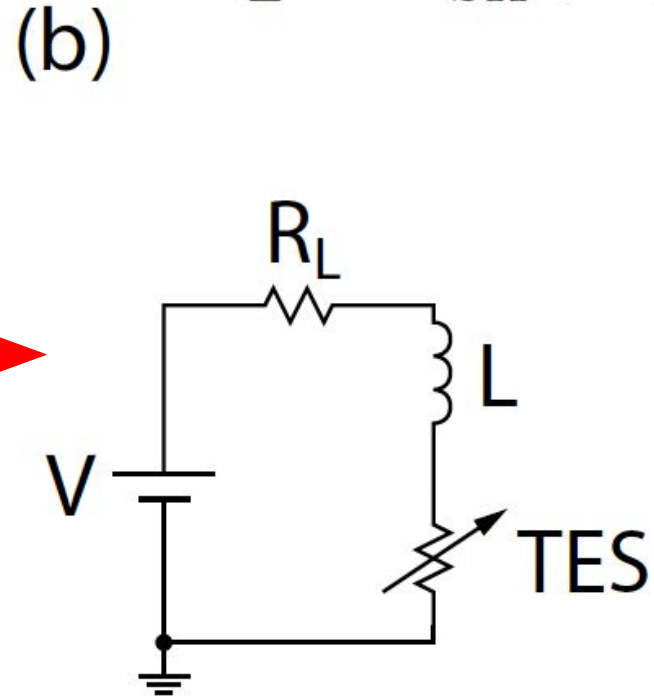
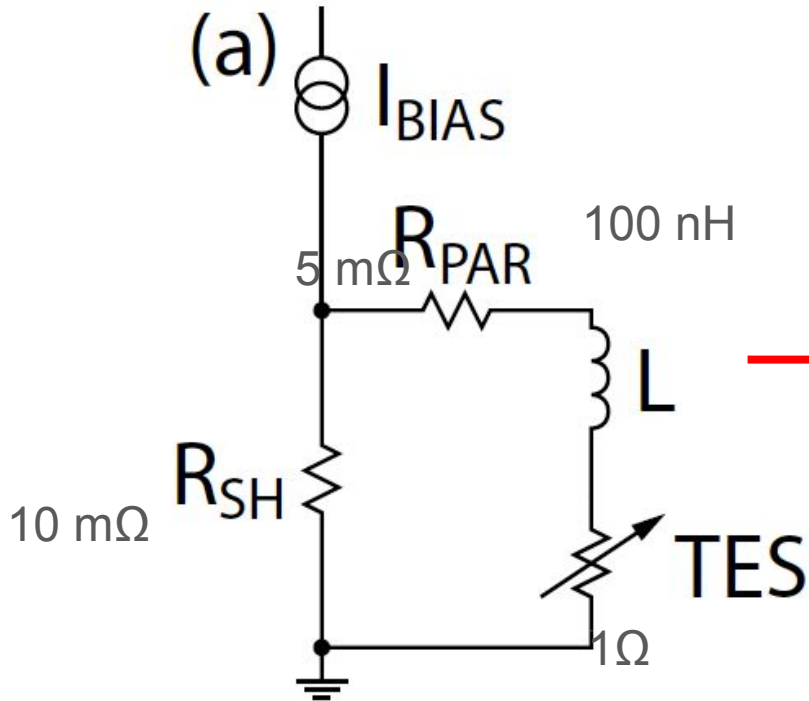
# TESS operation

- Set a **voltage** across the TES ( $V_{\text{TES}}$ )
- Put TES in series with an inductor
- Current in TES  $\rightarrow$  magnetic flux in inductor
- $\rightarrow$  SQUID measures magnetic flux
- $\rightarrow$  SQUID measures current  $I_{\text{TES}}$
- Calculate TES resistance with  $V_{\text{TES}}/I_{\text{TES}}$
- $R \rightarrow T$



# Thevenin Equivalent Circuit

$$V = I_{\text{BIAS}} R_{\text{SH}}$$
$$R_L = R_{\text{SH}} + R_{\text{PAR}}$$

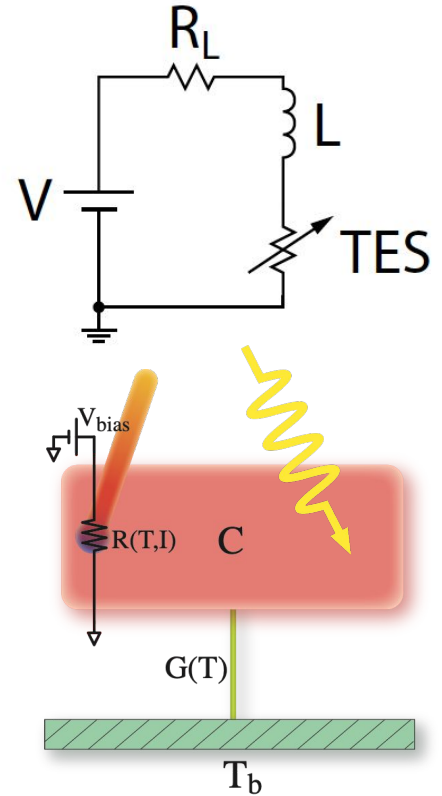


# Differential equations

$$L \frac{dI}{dt} = V - IR_L - IR(T, I)$$

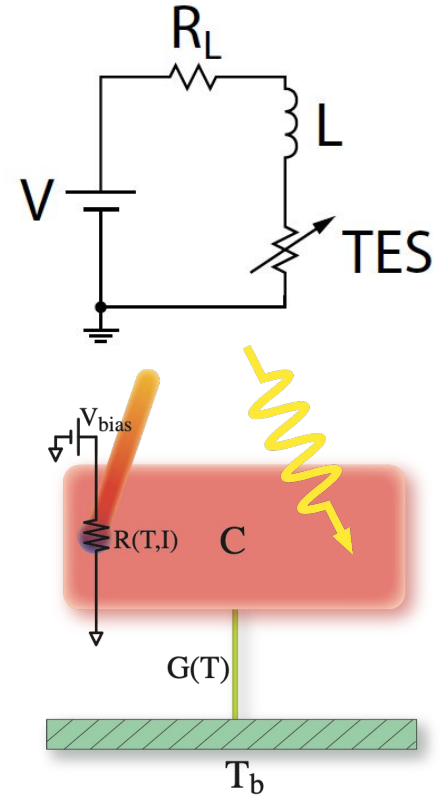
$$C \frac{dT}{dt} = -P_{\text{bath}} + P_J + P$$

$$P_{\text{bath}} = K (T^n - T_{\text{bath}}^n)$$



# Quiescent solution and Electro-thermal stability

- Typical scenario:
  - $T_b \sim 10$  mK
  - $T_{TES}$  (critical temperature  $T_c$ )  $\sim 50$  mK
- Joule heating keeps TES at elevated temperature
  - $P = V^2/R$
- In Quiescence, TES operates at
  - $T_0$ : Operating temperature, close to  $T_c$
  - $R_0$ : TES operating resistance
    - $\sim 30\%$  of  $R_n$  (normal Resistance)
  - $P_0$ : “Bias power”, power to heat TES from  $T_b$  to  $T_0$
- When particle hit  $\rightarrow$  TES heats up  $\rightarrow R$  goes up  $\rightarrow P$  decreases  $\rightarrow$  TES cools back down



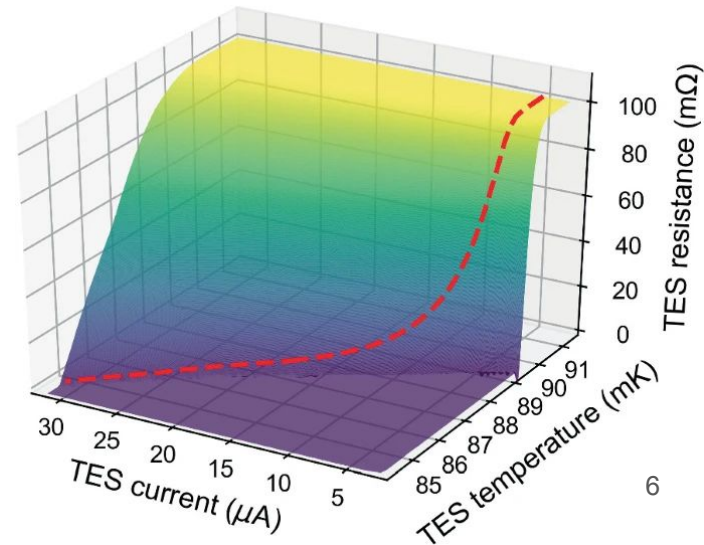
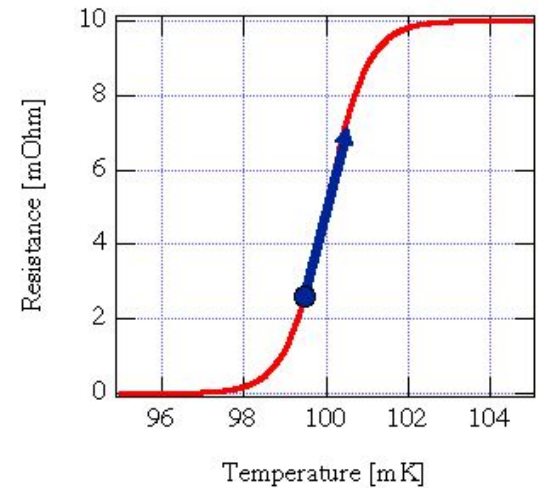
# TES model

$$R(T, I) \approx R_0 + \left. \frac{\partial R}{\partial T} \right|_{I_0} \delta T + \left. \frac{\partial R}{\partial I} \right|_{T_0} \delta I$$

$$\alpha_I \equiv \left. \frac{\partial \log R}{\partial \log T} \right|_{I_0} = \frac{T_0}{R_0} \left. \frac{\partial R}{\partial T} \right|_{I_0}$$

$$\beta_I \equiv \left. \frac{\partial \log R}{\partial \log I} \right|_{T_0} = \frac{I_0}{R_0} \left. \frac{\partial R}{\partial I} \right|_{T_0}$$

$$R(T, I) \approx R_0 + \alpha_I \frac{R_0}{T_0} \delta T + \beta_I \frac{R_0}{I_0} \delta I$$



## More “random” definitions (aka jargons)

- Constant temperature dynamic resistance  $R_{\text{dyn}} \equiv \left. \frac{\partial V}{\partial I} \right|_{T_0} = R_0 (1 + \beta_I)$
- Natural thermal time constant  $\tau \equiv \frac{C}{G}$
- Low frequency loop gain with constant current  $\mathcal{L}_I \equiv \frac{P_{J_0} \alpha_I}{GT_0}$
- Electrical time constant  $\tau_{\text{el}} = \frac{L}{R_L + R_0 (1 + \beta_I)} = \frac{L}{R_L + R_{\text{dyn}}}$
- Current biased thermal time constant  $\tau_I = \frac{\tau}{1 - \mathcal{L}_I}$

# Linearized differential equation

- $$L \frac{dI}{dt} = V - IR_L - IR(T, I)$$
$$C \frac{dT}{dt} = -P_{\text{bath}} + P_J + P$$

- Small signal limit, Taylor expansion, neglecting all higher order terms

$$\frac{d\delta I}{dt} = -\frac{R_L + R_0(1 + \beta_I)}{L} \delta I - \frac{\mathcal{L}_I G}{I_0 L} \delta T + \frac{\delta V}{L}$$

$$\frac{d\delta T}{dt} = \frac{I_0 R_0(2 + \beta_I)}{C} \delta I - \frac{(1 - \mathcal{L}_I)}{\tau} \delta T + \frac{\delta P}{C}$$

$\delta V$  and  $\delta P$  are small perturbations around bias voltage  $V$  and input power  $P$

- Put in for later analyses, eg. particle power injection, noise



## Matrix differential equation

$$\frac{d}{dt} \begin{pmatrix} \delta I \\ \delta T \end{pmatrix} = - \begin{pmatrix} \frac{1}{\tau_{el}} & \frac{\mathcal{L}_I G}{I_0 L} \\ -\frac{I_0 R_0 (2 + \beta_I)}{C} & \frac{1}{\tau_I} \end{pmatrix} \begin{pmatrix} \delta I \\ \delta T \end{pmatrix} + \begin{pmatrix} \frac{\delta V}{L} \\ \frac{\delta P}{C} \end{pmatrix}$$

- Solve for homogeneous equations first  $\rightarrow$  setting  $\delta V$  and  $\delta P$  to 0
- Diagonalize the matrix  $\rightarrow$  two eigenvalues with two eigenvectors

$$\frac{d}{dt} f_{\pm}(t) = -\lambda_{\pm} f_{\pm}(t)$$

## Differential equation solution

$$\frac{d}{dt} \begin{pmatrix} \delta I \\ \delta T \end{pmatrix} = - \begin{pmatrix} \frac{1}{\tau_{el}} & \frac{\mathcal{L}_I G}{I_0 L} \\ -\frac{I_0 R_0 (2 + \beta_I)}{C} & \frac{1}{\tau_I} \end{pmatrix} \begin{pmatrix} \delta I \\ \delta T \end{pmatrix} + \begin{pmatrix} \frac{\delta V}{L} \\ \frac{\delta P}{C} \end{pmatrix}$$

$$\frac{d}{dt} f_{\pm}(t) = -\lambda_{\pm} f_{\pm}(t)$$

- Solve for homogeneous equations first  $\rightarrow$  setting  $\delta V$  and  $\delta P$  to 0
- Diagonalize the matrix  $\rightarrow$  two eigenvalues with two eigenvectors

$$\begin{pmatrix} \delta I \\ \delta T \end{pmatrix} = A_+ e^{-\lambda_+ t} \vec{v}_+ + A_- e^{-\lambda_- t} \vec{v}_-$$

$$\frac{1}{\tau_{\pm}} \equiv \lambda_{\pm} = \frac{1}{2\tau_{el}} + \frac{1}{2\tau_I} \pm \frac{1}{2} \sqrt{\left(\frac{1}{\tau_{el}} - \frac{1}{\tau_I}\right)^2 - 4 \frac{R_0}{L} \frac{\mathcal{L}_I (2 + \beta_I)}{\tau}}$$

$$\vec{v}_{\pm} = \begin{pmatrix} \frac{1 - \mathcal{L}_I - \lambda_{\pm} \tau}{2 + \beta_I} \frac{G}{I_0 R_0} \\ 1 \end{pmatrix}$$

# Specific solution 1 -- Impulse energy input

- Particle incident and instantaneously thermalization

$$\delta T(0) = \delta T = E/C$$

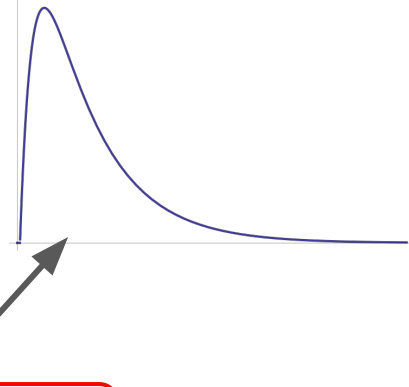
$$\delta I(0) = 0$$

- Solve for  $A_{\pm}$ , then...

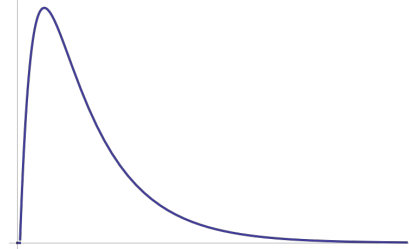
$$\delta I(t) = \left( \frac{\tau_I}{\tau_+} - 1 \right) \left( \frac{\tau_I}{\tau_-} - 1 \right) \frac{1}{(2 + \beta_I)} \frac{C \delta T}{I_0 R_0 \tau_I^2} \left( e^{-t/\tau_+} - e^{-t/\tau_-} \right)$$

$$\delta T(t) = \left( \left( \frac{1}{\tau_I} - \frac{1}{\tau_+} \right) e^{-t/\tau_-} + \left( \frac{1}{\tau_I} - \frac{1}{\tau_-} \right) e^{-t/\tau_+} \right) \frac{\delta T}{(1/\tau_+ - 1/\tau_-)}$$

- $\tau_+$ ,  $\tau_-$  : “rise time” and “fall time” of a pulse with impulse energy deposition



# Specific solution 1 -- Impulse energy input



- When L (inductance) is small,  $\tau_+ \ll \tau_-$

$$\tau_+ \rightarrow \tau_{el},$$

$$\tau_- \rightarrow \tau \frac{1 + \beta_I + R_L/R_0}{1 + \beta_I + R_L/R_0 + (1 - R_L/R_0)\mathcal{L}_I} = \tau_{\text{eff}}$$

$$\mathcal{L}_I \equiv \frac{P_{J_0} \alpha_I}{GT_0}$$

- Recall from last lecture, when  $R_L/R_0 \ll 1$ , “Stiff voltage bias limit”
- $\beta_I \rightarrow 0$  (not modelled in last lecture)

$$\tau_{\text{eff}} = \frac{\tau_0}{1 + \frac{\alpha P_0}{T_0 G}} = \frac{\tau_0}{1 + \frac{\alpha}{n} \left(1 - \left(\frac{T_b}{T_0}\right)^n\right)}$$

## Specific solution 2 -- small sinusoidal power input

$$\delta P = \text{Re}(\delta P_0 e^{i\omega t})$$

- Useful to derive **responsivity**

$$\frac{d}{dt} \begin{pmatrix} \delta I \\ \delta T \end{pmatrix} = - \begin{pmatrix} \frac{1}{\tau_{\text{el}}} & \frac{\mathcal{L}_I G}{I_0 L} \\ \frac{-I_0 R_0 (2 + \beta_I)}{C} & \frac{1}{\tau_I} \end{pmatrix} \begin{pmatrix} \delta I \\ \delta T \end{pmatrix} + \begin{pmatrix} 0 \\ \frac{\delta P_0}{C} \end{pmatrix} e^{i\omega t}$$

- After some math.....

$$\frac{dI}{dP}(\omega) = s_I(\omega) = -\frac{1}{I_0 R_0} \frac{1}{(2 + \beta_I)} \frac{(1 - \tau_+/\tau_I)(1 - \tau_-/\tau_I)}{(1 + i\omega\tau_+)(1 + i\omega\tau_-)}$$

$$\frac{dT}{dP}(\omega) = s_T(\omega) = \frac{1}{G} \frac{\tau_+ \tau_-}{\tau^2} \frac{(\tau/\tau_+ + \tau/\tau_- + \mathcal{L}_I - 1 + i\omega\tau)}{(1 + i\omega\tau_+)(1 + i\omega\tau_-)}$$

- Two low-pass filters

- For a review of one-pole filter and its response in Laplace domain, see

eg. <https://www.embeddedrelated.com/showarticle/590.php>

# Complex impedance

- Impedance of the system in a complex plane, as a function of frequency
  - Given Voltage excitation of  $\omega$  angular frequency, what's current and phase?

$$\frac{d}{dt} \begin{pmatrix} \delta I \\ \delta T \end{pmatrix} = - \begin{pmatrix} \frac{1}{\tau_{el}} & \frac{\mathcal{L}_I G}{I_0 L} \\ -\frac{I_0 R_0 (2 + \beta_I)}{C} & \frac{1}{\tau_I} \end{pmatrix} \begin{pmatrix} \delta I \\ \delta T \end{pmatrix} + \begin{pmatrix} \frac{\delta V}{L} \\ \frac{\delta P}{C} \end{pmatrix}$$

Fourier Transform  $\rightarrow$

$$\begin{pmatrix} \frac{1}{\tau_{el}} + i\omega & \frac{\mathcal{L}_I G}{I_0 L} \\ -\frac{I_0 R_0 (2 + \beta_I)}{C} & \frac{1}{\tau_I} + i\omega \end{pmatrix} \begin{pmatrix} I_\omega \\ T_\omega \end{pmatrix} = \begin{pmatrix} \frac{V_\omega}{L} \\ 0 \end{pmatrix}$$

Voltage excitation of  $V_\omega$

# Complex impedance

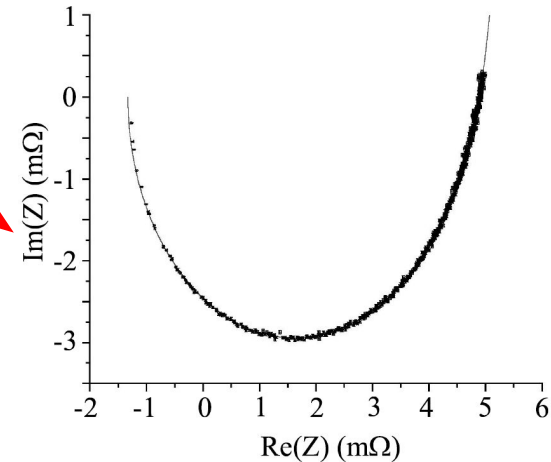
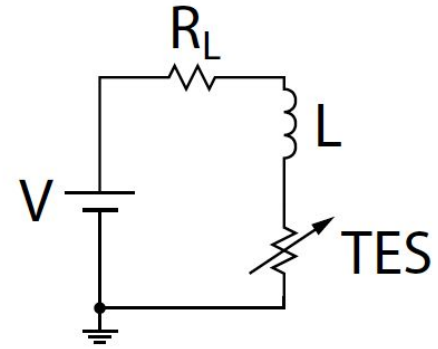
$$\begin{pmatrix} \frac{1}{\tau_{el}} + i\omega & \frac{\mathcal{L}_I G}{I_0 L} \\ \frac{-I_0 R_0 (2 + \beta_I)}{C} & \frac{1}{\tau_I} + i\omega \end{pmatrix} \begin{pmatrix} I_\omega \\ T_\omega \end{pmatrix} = \begin{pmatrix} \frac{V_\omega}{L} \\ 0 \end{pmatrix}$$

$$Z_\omega = V_\omega / I_\omega = R_L + i\omega L + Z_{TES}$$

$$Z_{TES} = R_0(1 + \beta_I) + \frac{R_0 \mathcal{L}_I}{1 - \mathcal{L}_I} \frac{2 + \beta_I}{1 + i\omega \tau_I}$$

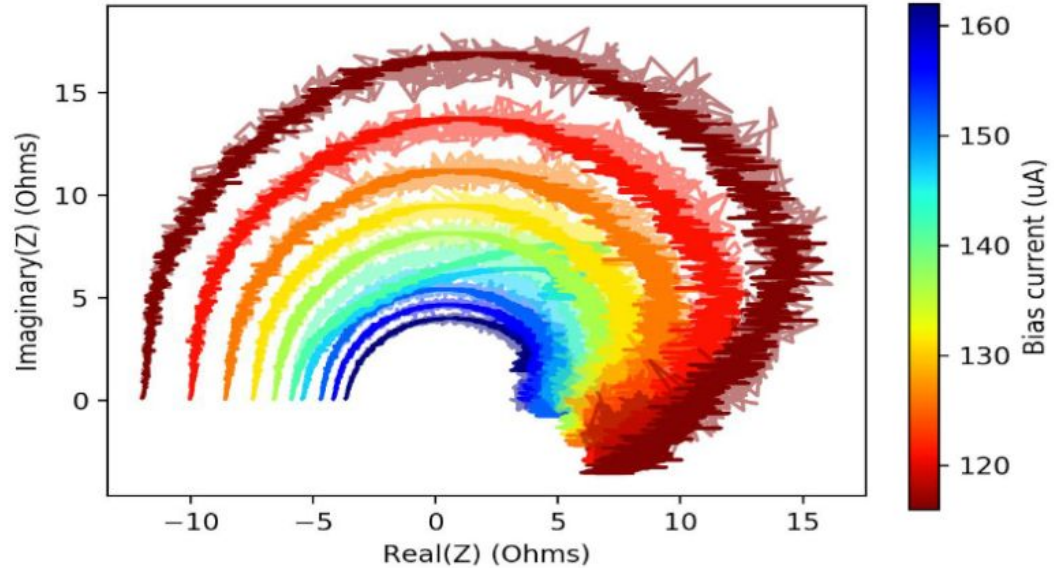
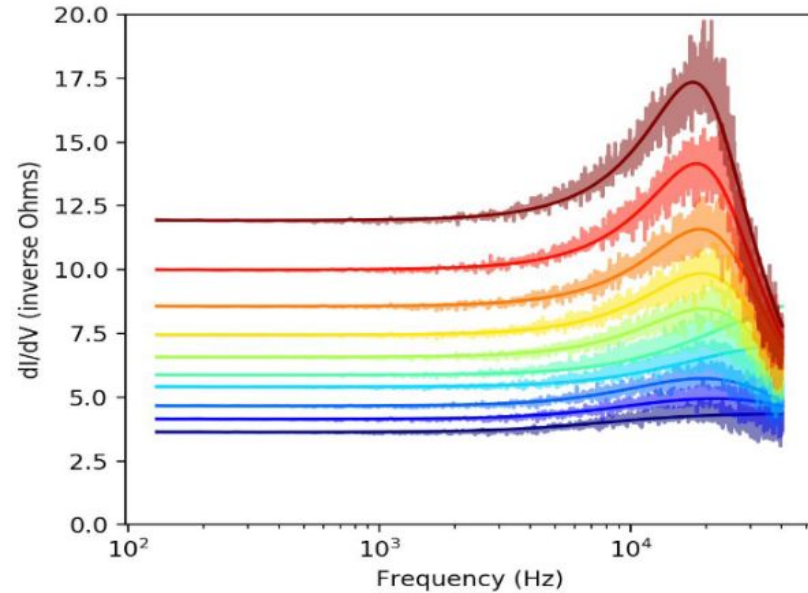
Or.. complex admittance  $Y=1/Z$

$$Y(\omega) = s_I(\omega) I_0 \frac{\mathcal{L}_I - 1}{\mathcal{L}_I} (1 + i\omega \tau_I)$$



# Complex impedance used to extract TES parameters

- Often takes data at different  $T_b$  and different  $R_0$





# TES Stability

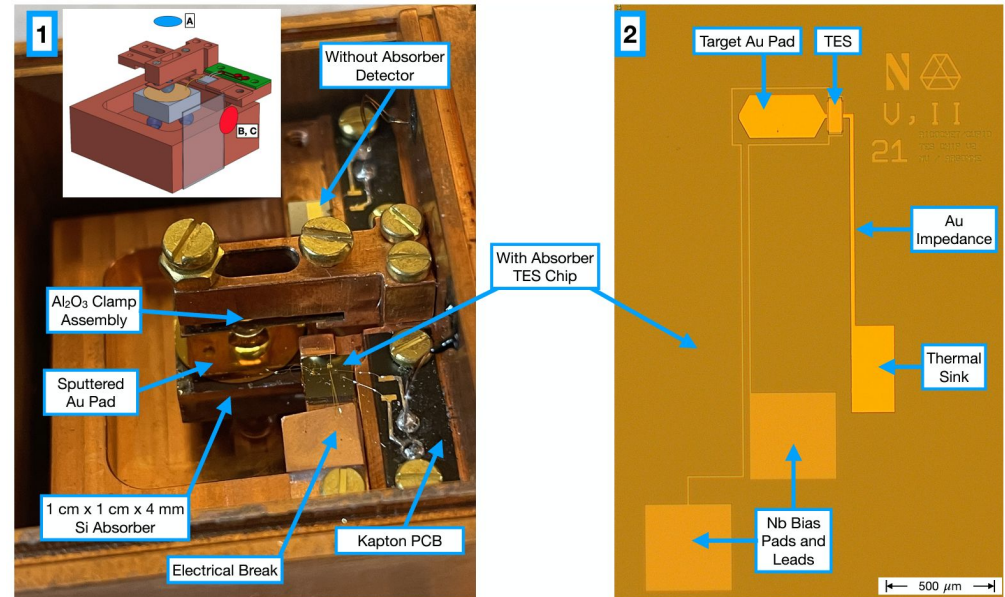
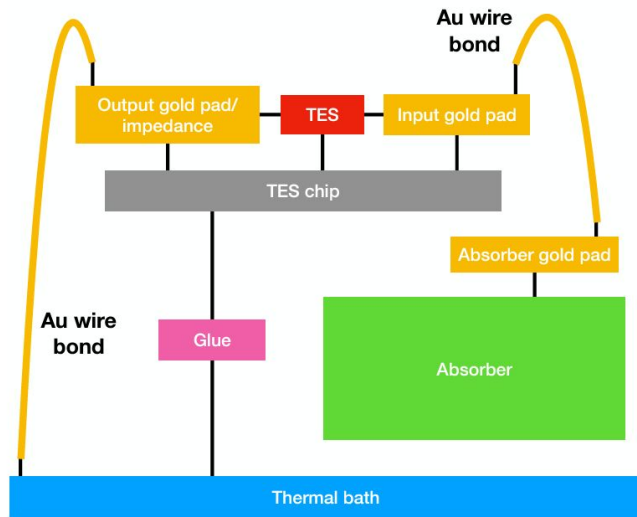
- For a TES to be stable:
  - $\tau_{\pm}$  need to be **real (not oscillating)**
  - $\text{Re}(\tau_{\pm})$  need to be **positive (converging back to quiescence)**

$$\frac{1}{\tau_{\pm}} \equiv \lambda_{\pm} = \frac{1}{2\tau_{\text{el}}} + \frac{1}{2\tau_I} \pm \frac{1}{2} \sqrt{\left(\frac{1}{\tau_{\text{el}}} - \frac{1}{\tau_I}\right)^2 - 4 \frac{R_0}{L} \frac{\mathcal{L}_I(2 + \beta_I)}{\tau}}$$

$$R_0 > \frac{(\mathcal{L}_I - 1)}{(\mathcal{L}_I + 1 + \beta_I)} R_L$$

- In large  $L_I$  limit (large amplification),  $R_0 > R_L$
- If  $\text{Re}(\tau_{\pm})$  is positive, but  $\tau_{\pm}$  isn't real  $\rightarrow$  TES is underdamped (oscillate) but will converge  $\rightarrow$  **Electrothermal oscillation**
- If  $\text{Re}(\tau_{\pm})$  is negative  $\rightarrow$  thermal runaway

# Complex detector modeling



$$L \frac{dI}{dt} = I_0 * R_0 - I * R_{tes}$$

$$C_a \frac{dT_a}{dt} = \frac{E}{c\sqrt{2\pi}} e^{-\frac{(t-t_0)^2}{2c^2}} - K_{a,b}(T_a^{n_{a,b}} - T_b^{n_{a,b}}) - K_{a,auq}(T_a^{n_{a,au1}} - T_{au1}^{n_{a,au1}}) - K_{a,g}(T_a^{n_{a,g}} - T_g^{n_{a,g}})$$

$$C_{au1} \frac{dT_{au1}}{dt} = K_{a,au1}(T_{au1}^{n_{a,au1}} - T_a^{n_{a,au1}}) - K_{au1,wb1}(T_{au1}^{n_{au1,wb1}} - T_{wb1}^{n_{au1,wb1}})$$

$$C_g \frac{dT_g}{dt} = K_{a,g}(T_g^{n_{a,g}} - T_a^{n_{a,g}}) - K_{g,si}(T_g^{n_{g,si}} - T_{si}^{n_{g,si}})$$

$$C_{wb1} \frac{dT_{wb1}}{dt} = K_{au1,wb1}(T_{wb1}^{n_{au1,wb1}} - T_{au1}^{n_{au1,wb1}}) - K_{wb1,au2}(T_{wb1}^{n_{wb1,au2}} - T_{au2}^{n_{wb1,au2}})$$

$$C_{au2} \frac{dT_{au2}}{dt} = K_{si,au2}(T_{au2}^{n_{si,au2}} - T_{si}^{n_{si,au2}}) - K_{wb1,au2}(T_{au2}^{n_{wb1,au2}} - T_{wb1}^{n_{wb1,au2}}) - K_{au2,tes}(T_{au2}^{n_{au2,tes}} - T_{tes}^{n_{au2,tes}})$$

$$C_{si} \frac{dT_{si}}{dt} = K_{g,si}(T_{si}^{n_{g,si}} - T_g^{n_{g,si}}) - K_{si,au2}(T_{si}^{n_{si,au2}} - T_{au2}^{n_{si,au2}}) - K_{si,tes}(T_{si}^{n_{si,tes}} - T_{tes}^{n_{si,tes}}) \\ - K_{si,m}(T_{si}^{n_{si,m}} - T_m^{n_{si,m}})$$

$$C_{tes} \frac{dT_{tes}}{dt} = I^2 R_{tes} - K_{au2,tes}(T_{tes}^{n_{au2,tes}} - T_{au2}^{n_{au2,tes}}) - K_{si,tes}(T_{tes}^{n_{si,tes}} - T_{si}^{n_{si,tes}}) - K_{tes,m}(T_{tes}^{n_{tes,m}} - T_m^{n_{tes,m}})$$

$$C_m \frac{dT_m}{dt} = K_{tes,m}(T_m^{n_{tes,m}} - T_{tes}^{n_{tes,m}}) - K_{si,m}(T_m^{n_{si,m}} - T_{si}^{n_{si,m}}) - K_{m,wb2}(T_m^{n_{m,wb2}} - T_{wb2}^{n_{m,wb2}})$$

$$C_{wb2} \frac{dT_{wb2}}{dt} = K_{m,wb2}(T_{wb2}^{n_{m,wb2}} - T_m^{n_{m,wb2}}) - K_{wb2,b}(T_{wb2}^{n_{wb2,b}} - T_b^{n_{wb2,b}})$$

# Differential equation $\rightarrow$ Matrices (time, frequency)

$$\frac{d}{dt} \begin{bmatrix} \Delta I \\ \Delta T_{tes} \\ \Delta T_a \\ \Delta T_{au1} \\ \Delta T_g \\ \Delta T_{wb1} \\ \Delta T_{au2} \\ \Delta T_{si} \\ \Delta T_m \\ \Delta T_{wb2} \end{bmatrix} = \begin{bmatrix} \frac{-R(1+\beta_I)}{L} & \frac{-\alpha V}{T_{tes}L} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \frac{(2+\beta_I)V}{C_{tes}} & \frac{-(G_{au2,tes}^{tes} + G_{tes,m}^{tes} + G_{si,tes}^{tes} - \frac{\alpha P}{T_{tes}})}{C_{tes}} & 0 & 0 & 0 & 0 & \frac{G_{au2,tes}^{au2}}{C_{tes}} & \frac{G_{si,tes}^{si}}{C_{tes}} & \frac{G_{tes,m}^m}{C_{tes}} & 0 \\ 0 & 0 & -(G_{a,b}^a + G_{a,au1}^a + G_{a,g}^a) & \frac{G_{a,au1}^{au1}}{C_a} & \frac{G_{a,g}^g}{C_a} & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & \frac{G_{a,au1}^a}{C_{au1}} & \frac{-(G_{a,au1}^{au1} + G_{au1,wb1}^{au1})}{C_{au1}} & 0 & \frac{G_{au1,wb1}^{wb1}}{C_{au1}} & 0 & 0 & 0 & 0 \\ 0 & 0 & \frac{G_{a,g}^g}{C_g} & 0 & \frac{-(G_{a,g}^g + G_{g,si}^g)}{C_g} & 0 & 0 & \frac{G_{g,si}^{si}}{C_g} & 0 & 0 \\ 0 & 0 & 0 & \frac{G_{au1,wb1}^{wb1}}{C_{wb1}} & 0 & \frac{-(G_{au1,wb1}^{wb1} + G_{wb1,au2}^{wb1})}{C_{wb1}} & \frac{G_{wb1,au2}^{au2}}{C_{wb1}} & 0 & 0 & 0 \\ 0 & \frac{G_{au2,tes}^{tes}}{C_{au2}} & 0 & 0 & 0 & \frac{G_{wb1,au2}^{wb1}}{C_{au2}} & \frac{-(G_{au2,tes}^{au2} + G_{wb1,au2}^{wb1} + G_{au2,si}^{au2})}{C_{au2}} & \frac{G_{au2,si}^{si}}{C_{au2}} & 0 & 0 \\ 0 & \frac{G_{si,tes}^{tes}}{C_{si}} & 0 & 0 & \frac{G_{g,si}^{si}}{C_{si}} & 0 & 0 & \frac{-(G_{g,si}^{si} + G_{si,au2}^{si} + G_{si,tes}^{si} + G_{si,m}^{si})}{C_{si}} & \frac{G_{si,m}^m}{C_{si}} & 0 \\ 0 & \frac{G_{tes,m}^{tes}}{C_m} & 0 & 0 & 0 & 0 & 0 & \frac{G_{si,m}^{si}}{C_m} & \frac{-(G_{tes,m}^m + G_{si,m}^m + G_{m,wb2}^m)}{C_m} & \frac{G_{m,wb2}^{wb2}}{C_m} \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{G_{m,wb2}^{wb2}}{C_{wb2}} & \frac{-(G_{m,wb2}^{wb2} + C_{wb2,b})}{C_{wb2}} \end{bmatrix} \begin{bmatrix} \Delta I \\ \Delta T_{tes} \\ \Delta T_a + w(t) \\ \Delta T_{au1} \\ \Delta T_g \\ \Delta T_{wb1} \\ \Delta T_{au2} \\ \Delta T_{si} \\ \Delta T_m \\ \Delta T_{wb2} \end{bmatrix}$$

$$\begin{bmatrix} R(1+\beta_I) + i\omega L & -1 & \frac{\alpha V}{T_c} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ R_l & 1 + i\omega R_l C_{cap} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -(2+\beta_I)V & 0 & (G_{au2,tes}^{tes} + G_{tes,m}^{tes} + G_{si,tes}^{tes} - \frac{\alpha P}{T_{tes}}) + i\omega C_{tes} & 0 & 0 & 0 & -G_{au2,tes}^{au2} & -G_{si,tes}^{si} & -G_{tes,m}^m & 0 & 0 \\ 0 & 0 & 0 & (G_{a,b}^a + G_{a,au1}^a + G_{a,g}^a) + i\omega C_a & -G_{a,au1}^{au1} & -G_{a,g}^g & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -C_{a,au1}^{au1} & (G_{a,au1}^{au1} + G_{au1,wb1}^{au1}) + i\omega C_{au1} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -C_{a,g}^g & 0 & (G_{a,g}^g + G_{g,si}^g) + i\omega C_g & 0 & 0 & -C_{g,si}^{si} & 0 & 0 \\ 0 & 0 & 0 & 0 & -C_{au1,wb1}^{wb1} & 0 & (G_{au1,wb1}^{wb1} + G_{wb1,au2}^{wb1}) + i\omega C_{wb1} & -G_{wb1,au2}^{au2} & 0 & 0 & 0 \\ 0 & 0 & -G_{au2,tes}^{tes} & 0 & 0 & 0 & -G_{wb1,au2}^{wb1} & (G_{au2,tes}^{au2} + G_{wb1,au2}^{wb1} + G_{au2,si}^{au2}) + i\omega C_{au2} & -C_{au2,si}^{si} & 0 & 0 \\ 0 & 0 & -G_{si,tes}^{tes} & 0 & 0 & 0 & -C_{si,tes}^{si} & 0 & 0 & 0 & 0 \\ 0 & 0 & -G_{tes,m}^{tes} & 0 & 0 & -C_{g,si}^{si} & 0 & -C_{si,m}^{si} & (G_{tes,m}^m + G_{si,m}^m + G_{m,wb2}^m) + i\omega C_m & -C_{m,wb2}^{wb2} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & (G_{m,wb2}^{wb2} + C_{wb2,b}) + i\omega C_{wb2} & -C_{wb2,b}^{wb2} \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -C_{wb2,b}^{wb2} & (C_{wb2,b}^{wb2} + C_{wb2,l}) + i\omega C_{wb2} \end{bmatrix} \begin{bmatrix} \Delta I \\ \Delta V \\ \Delta T_{tes} \\ \Delta T_a \\ \Delta T_g \\ \Delta T_{wb1} \\ \Delta T_{au1} \\ \Delta T_{au2} \\ \Delta T_{si} \\ \Delta T_m \\ \Delta T_{wb2} \end{bmatrix} = \begin{bmatrix} \epsilon_{int} \\ \epsilon_{ext} \\ P_{au2,tes} + P_{si,tes} - P_{tes,m} - I_{\epsilon_{int}} \\ -P_{a,g} - P_{a,au1} - P_{a,b} \\ P_{a,au1} - P_{au1,wb1} \\ P_{a,g} - P_{g,si} \\ P_{au1,wb1} - P_{wb1,au2} \\ P_{wb1,au2} + P_{si,au2} - P_{au2,tes} \\ P_{g,si} - P_{si,tes} - P_{si,au2} - P_{si,m} \\ P_{tes,m} + P_{si,m} - P_{m,wb2} \\ P_{m,wb2} - P_{wb2,b} \end{bmatrix}$$

# One note about noise and resolution

- With responsivity  $s_1$ , can derive theoretical resolution:
  - Model various noise
    - Johnson noise, thermal fluctuation noise, etc.
  - Convert to “Noise Equivalent Power”

- $$\Delta E_{\text{rms}} = \left( \sqrt{\int_0^\infty \frac{4df}{\text{NEP}(f)^2}} \right)^{-1}$$

- Irwin’s derivation shows

$$\Delta E_{\text{rms}} = \sqrt{4k_{\text{B}}T^2C \frac{1}{\alpha} \sqrt{\frac{n}{2}}}$$

- Will come back to this next class